

A Simple Buckling Analysis of Aorta Artery

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Abstract

Aortas are the largest artery in the body and they carry the blood away which is pumped from the heart. Aorta artery is also the artery which is affected by the highest blood pressure. Its stability has a vital importance to humans and animals. The loss of stability in arteries may lead to arterial tortuosity and kinking. This situation causes to blackouts and serious permanent health problems. In this article, the buckling analysis of aorta artery is investigated by using Euler-Bernoulli beam theory for different boundary conditions. The aorta artery is modeled as a cylindrical tube with different average diameters. Results are presented in figures and table.

Keywords: Aorta artery, buckling, Euler-Bernoulli.

1. Introduction

Aorta artery is first mechanically modeled and its stability under blood pressure was studied to determine the critical buckling loads by Han in 2007 [1]. His researches showed that arteries may buckle and become tortuous due to reduced axial strain, hypertensive pressure, and a weakened wall. [2-9]. Han also has investigated the critical buckling by using nonlinear elastic thick-walled cylindrical model with residual stress which was developed by Han [3]. On the other hand, collapse of the vessel lumen and, the bent buckling of tubular arteries was recently reported [10-13]. The buckling of aorta artery causes additional wall stress and it leads to affect blood flow and cause extra loads to itself. In this study the simple classical Euler-Bernoulli buckling theory will be used in order to calculate the buckling loads of aorta artery for four different boundary conditions. In literature many researches about the buckling theory for small-scaled and micro-scaled structures [14-46] and rod, beam, plate, shell models are being used in order to determine the vibration of continuous systems. [47-52]

2. Buckling analysis of aorta artery

The demonstration of aorta artery and its continuum model are shown in Fig.(1) and Fig.(2). In order to calculate the critical buckling load of the model, Euler-Bernoulli beam theory is used. For modeling, L is the length of microtubules, R_i and R_o is inner and outer radius, D_{avg} average diameter, t thickness, E Young's modulus.



Fig. 1. Demonstration of aorta artery

Arteries are composed of three layers. These layers are intima, media and adventitia respectively from inside to outside of vessels (Fig.1). Intima is the intermost layer of the artery wich is covering the lumen side of vessels and it is composed of endothelial cells and lines the entire circulatory system, from the heart and the large arteries all the way down to the very tiny capillary beds. The intima layer also contain extracellular matrix and a supporting layer of collagenous tissue. Endothelial cells sorted in a single layer along the lumen side. Media is the muscular middle layer of the arteries and veins. It is composed of smooth muscle layers. Adventitia is outermost layer of vessels surrounding the media layer. It is mainly composed of collagen and, in arteries, is supported by external elastic lamina.



Fig. 2. Continuum model of aorta artery with variable boundary conditions

The aorta artery has a curved cylindrical shaped structure *in vivo*. In this study the structure is modeled as homogenous, straight cylindrical shaped tube. Critical buckling loads are investigated for simply supported (Fig. 2(a)), clamped (Fig. 2(b)), popped (Fig. 2(c)) and cantilever (Fig. 2(d)) boundary conditions.

3. Euler-Bernoulli formulation

The buckling equation of a beam is:

$$EI\frac{d^{4}y}{dx^{4}} + P\frac{d^{2}y}{dx^{2}} = 0$$
 (1)

If setting $\alpha^2 = \frac{P}{EI}$, Eq.(1) can be simplified as:

$$y^{\prime\nu} + \alpha^2 y^{\prime\prime} = 0 \tag{2}$$

If setting $y = e^{rx}$, Eq.(2) can be simplified as:

$$Br^{4}e^{rx} + \alpha^{2}Br^{2}e^{rx} = 0$$
(3)

By reducing Eq.(3), we can obtain:

$$r^4 + \alpha^2 r^2 = 0 \tag{4}$$

Solving Eq.(4), the result is:

$$r^{2} = -\alpha^{2}$$

$$r_{1,2} = 0 \qquad \text{and} \qquad r_{3,4} = \pm i\alpha$$
(5)

 $r_{1,2}$ and $r_{3,4}$ are two pairs of single complex root of Eq.(4).

By substitution roots into Eq.(5) and solving it, we obtain:

$$y = C_1 \sin \alpha x + C_2 \cos \alpha x + C_3 x + C_4 \tag{6}$$

 C_1, C_2, C_3, C_4 are constants which can be obtained from boundary conditions. The first order derivative of Eq.(6) is:

$$y' = C_1 \alpha \cos \alpha x - C_2 \alpha \sin \alpha x + C_3 \tag{7}$$

The second order derivative of Eq.(6) is:

$$y'' = -C_1 \alpha^2 \sin \alpha x - C_2 \alpha^2 \cos \alpha x \tag{8}$$

The third order derivative of Eq.(6) is:

$$y''' = -C_1 \alpha^3 \cos \alpha x + C_2 \alpha^3 \sin \alpha x \tag{9}$$

For a beam which is Clamped-Free supported, the boundary conditions would be as followed:

$$y(0) = y'(0) = 0$$
, $y''(l) = y'''(l) + \alpha^2 y'(l) = 0$ (10)

By substituting boundary conditions into Eq.(6), Eq.(7), Eq.(8) and Eq.(9) we obtain:

$$v(0) = C_2 + C_4 = 0 \tag{11}$$

$$y(0) = C_2 + C_4 = 0$$
(11)

$$y'(0) = C_1 \alpha + C_3 = 0$$
(12)
(12)

$$y''(l) = -C_1 \alpha^2 \sin \alpha l - C_2 \alpha^2 \cos \alpha l = 0$$
(13)

$$y'''(l) + \alpha^2 y'(l) = C_3 \alpha^2 = 0$$
(14)

As it is mentioned above C_1, C_2, C_3, C_4 are constants and we can obtain those constant by using Eq.(11), Eq.(12), Eq.(13) and Eq.(14). The solution is obtained as follow:

$$\alpha^5 \cos(\alpha l) = 0 \tag{15}$$

There are 2 possibilities which make the Eq.(15) equal to zero.

$$\alpha^5 = 0 \tag{16}$$
$$\cos(\alpha l) = 0 \tag{17}$$

By substituting $\alpha^2 = \frac{P}{EI}$ into Eq.(17) we can obtain:

$$\cos(\sqrt{\frac{P}{EI}}l) = 0, \text{ so } \sqrt{\frac{P}{EI}}l = n\frac{\pi}{2}$$
(18)

So the buckling load can be obtained via this formula:

$$P = \frac{n^2 \pi^2 EI}{4l^2} \tag{19}$$

For a beam which is simply supported at both ends, the boundary conditions would be as followed:

$$y(0) = y''(0) = 0, \qquad y(l) = y''(l) = 0$$
 (20)

By substituting boundary conditions into Eq.(6) and Eq.(8) we obtain:

$$y(0) = C_2 + C_4 = 0 \tag{21}$$

$$y''(0) = -C_2 \alpha^2 = 0 \tag{22}$$

$$y(l) = C_1 \sin \alpha l + C_2 \cos \alpha l + C_3 l + C_4 = 0$$
(23)

$$y''(l) = -C_1 \alpha^2 \sin(\alpha l) - C_2 \alpha^2 \cos(\alpha l) = 0$$
 (24)

As it is mentioned above C_1, C_2, C_3, C_4 are constants and we can obtain those constant by using Eq.(21), Eq.(22), Eq.(23) and Eq.(24). The solution is obtained as follow:

$$-\alpha^4 \sin(\alpha l) = 0 \tag{25}$$

There are 2 possibilities which make the Eq.(15) equal to zero.

$$-\alpha^4 = 0 \tag{26}$$

$$\sin(\alpha l) = 0 \tag{27}$$

By substituting $\alpha^2 = \frac{P}{EI}$ into Eq.(27) we can obtain:

$$\sin(\sqrt{\frac{P}{EI}}l) = 0, \text{ so } \sqrt{\frac{P}{EI}}l = n\pi$$
(28)

So the buckling load can be obtained via this formula:

$$P = \frac{n^2 \pi^2 EI}{l^2} \tag{29}$$

For a beam which is Clamped-Simple supported at ends, the boundary conditions would be as followed:

$$y(0) = y'(0) = 0,$$
 $y(l) = y''(l) = 0$ (30)

By substituting boundary conditions into Eq.(6) and Eq.(8) we obtain:

$$y(0) = C_2 + C_4 = 0 \tag{31}$$

$$y'(0) = C_1 \alpha + C_3 = 0 \tag{32}$$

$$y(l) = C_1 \sin \alpha l + C_2 \cos \alpha l + C_3 l + C_4 = 0$$
(33)

$$y''(l) = -C_1 \alpha^2 \sin(\alpha l) - C_2 \alpha^2 \cos(\alpha l) = 0$$
 (34)

As it is mentioned above C_1, C_2, C_3, C_4 are constants and we can obtain those constant by using Eq.(31), Eq.(32), Eq.(33) and Eq.(34). The solution is obtained as follow:

$$\alpha \left[\sin(\alpha l) - \alpha l \cos(\alpha l) \right] = 0 \tag{35}$$

By arranging the transcendent Eq.(35) we obtain:

$$\tan(\alpha l) = \alpha l \tag{36}$$

The result of the Eq.(36) is $\alpha l = 4.49$, by substituting $\alpha^2 = \frac{P}{EI}$ into Eq.(36) we can obtain:

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$$P = \alpha^2 E I \tag{37}$$

So the buckling load can be obtained via this formula:

$$P = 2.05 \frac{n^2 \pi^2 EI}{l^2}$$
(38)

For a beam which is fixed at both ends, the boundary conditions would be as followed:

$$y(0) = y'(0) = 0,$$
 $y(l) = y'(l) = 0$ (39)

By substituting boundary conditions into Eq.(6) and Eq.(8) we obtain:

$$y(0) = C_2 + C_4 = 0 \tag{40}$$

$$y'(0) = C_1 \alpha + C_3 = 0 \tag{41}$$

$$y(l) = C_1 \sin \alpha l + C_2 \cos \alpha l + C_3 l + C_4 = 0$$
(42)

$$y'(l) = C_1 \alpha \cos \alpha l - C_2 \alpha \sin \alpha l + C_3 \tag{43}$$

As it is mentioned above C_1, C_2, C_3, C_4 are constants and we can obtain those constant by using Eq.(31), Eq.(32), Eq.(33) and Eq.(34). The solution is obtained as follow:

$$\sin(\frac{1}{2}\alpha l) = 0 \tag{44}$$

By arranging the transcendent Eq.(44) we obtain:

$$\frac{1}{2}\alpha l = \pi \tag{45}$$

So the buckling load can be obtained via this formula:

$$P = 4 \frac{n^2 \pi^2 EI}{l^2} \tag{46}$$

4. Numerical Examples

In this study, the buckling of aorta artery with various boundary conditions is investigated via classical Euler-Bernoulli beam theory. Some of the results which are showing the buckling loads for simply supported aorta arteries are in Figure (3). These results show the decreasing of buckling load as the length of the artery increases. Three different average diameter is taken into account (D_{avg} =4.38mm, D_{avg} =5.5mm and D_{avg} =6mm). The elasticity modulus is E=200kPa [1], the thickness is t=1mm, the moment of inertia is I= π tR_{avg}³.(R_{avg}=D_{avg}/2)





Fig.3. Variation of buckling load of Aorta Artery for different average diameters for first 3 modes respectively.

Boundary conditions	15	25	<i>L/r</i> 45	65	85
C-F	0.0151	0.0054	0.0017	0.0008	0.0005
S-S	0.0604	0.0217	0.0067	0.0032	0.0019
C-S	0.1237	0.0445	0.0137	0.0066	0.0039
C-C	0.2414	0.0869	0.0268	0.0129	0.075

Table 1. The buckling load of aorta artery (N) for different boundary conditions and L/r ratio

The influences of the length on the buckling load for simply supported aorta artery for first three modes are illustrated in Figs. 3(a, b, c), respectively. In Fig. 3(a) the buckling loads are critical buckling load since it is for the first mode. As it can be seen in Figs. 2 (a,b,c,d), the buckling load is investigated for simply supported, clamped, propped and cantilever boundary conditions respectively. Fig.3(a, b, c) shows that the buckling load is decreasing dramatically with the increasing of length for all three modes. In Table 1. It can be seen that the highest buckling load is for Clamped boundary condition and lowest L/r ratio. As the L/r ratio increase, the artery is buckling under lower loads.

5. Concluding remarks

Buckling analysis of aorta artery is investigated for variable boundary conditions. Present equations from literature [14] are used in order to calculate the critical buckling loads. Results are presented in figures and table. The maximum buckling load is found for clamped case and lowest L/r ratio. On the other hand the minimum buckling load is for Clamped-Free case and highest L/r ratio.

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References

- [1] Han, H.C., A biomechanical model of artery buckling, *Journal of Biomechanics*, 40, 3672-3678, 2007.
- [2] Hayman, M. H., Zhang, J., Liu, Q., Xiao Y., Han H. C., Smooth muscle cell contraction increases the critical buckling pressure of arteries, *Journal of Biomechanics*, 46, 841-844, 2013.
- [3] Han, H.C., Nonlinear buckling of blood vessels: A theoretical study, *Journal of Biomechanics*, 41, 2708-2713, 2008.
- [4] Liu, Q., Han, H.C., Mechanical buckling of artery under pulsatile pressure, *Journal of Biomechanics*, 45, 1192-1198, 2012.
- [5] Han, H.C., Blood vessel buckling within soft surrounding tissue generates tortuosity, *Journal of Biomechanics*, 42, 2797-2801, 2009.
- [6] Lee, A. Y., Han, H.C., A nonlinear Thin-Wall Model for Vein Buckling, *Cardiovascular Engineering and Technology*, 1, 282-289, 2010.
- [7] Liu, Q., Han, H.C., Mechanical buckling of arterioles in collateral development, *Journal* of *Theoretical Biology*, 316, 42-48, 2013.
- [8] Lee, A. Y., Sanyal, A., Xiao, Y., Shadfan, R., Han, H.C., Mechanical instability of normal and aneurysmal arteries, *Journal of Biomechanics*, 47, 3868-3875, 2014.
- [9] Han, H.C., Chesnutt, J. K., Garcia, J. R., Liu, Q., Wen, Q., Artery Buckling: New Phenotypes, Models, and Applications, *Annals of Biomedical Engineering*, 41, 1399-1410, 2012.
- [10] Datir, P., Lee, A. Y., Lamm, S. D., Han, H.C., Effects of Geometric Variations on the Buckling of Arteries, *Int. J. Appl. Mech.*, 3(2), 385-406, 2011.
- [11] Ganguly, A., Simons, J., Schneider, A., Keck, B., Bennett, N. R., Fahrig, R., In-vitro Imaging of Femoral Artery Nitinol Stents for Deformation Analysis, *Journal of Vascular and Interventional Radiolog*, 22(2), 236-243, 2011.

- [12] Christensen, E. E., Landay, M. J., Dietz, G. W., Brinley, G., Buckling of the innominate artery simulating a right apical lung mass. *American Journal of Roentgenology*, 131, 119-123, 1978.
- [13] Arani, A. T., Arani, A. G., Kolahchi, R., Non-Newtonian pulsating blood flow-induced dynamic instability of visco-carotid artery within soft surrounding visco-tissue using differential cubature method, *Proceedings of the Institution of Mechanical Engineers, Part C: Journal of Mechanical Engineering Science*, 0954406214566038, 2015.
- [14] Chen, W.F., Lui, E.M., Structural Stability, *Elsevier*, New York/Amsterdam/London, 1987.
- [15] Ajori, S. and Ansari, R., Torsional buckling behavior of boron-nitride nanotubes using molecular dynamics simulations, *Curr. Appl. Phys.*, 14, 1072-1077, 2014.
- [16] Arani, A. G., Jamali, S. A., Amir, S., Maboudi, M. J., Electro-thermo-mechanical nonlinear buckling of Pasternak coupled DWBNNTs based on nonlocal piezoelasticity theory, *Turkish J. Eng. Env. Sci.*, 37, 231-246, 2013.
- [17] Shokuhfar, A., Ebrahimi-Nejad, S., Effects of structural defects on the compressive buckling of boron nitride nanotubes, *Physica E*, 48, 53-60, 2013.
- [18] Ansari, R. and Ajori, S., Molecular dynamics study of the torsional vibration characteristics of boron-nitride nanotubes, *Physics Letters A*, 378, 2876-2880, 2014.
- [19] Chowdhury, R., Wang, C. Y., Adhikari, S., Scarpa, F., Vibration and symmetrybreaking of boron nitride nanotubes, *Nanotechnology*, 21, 365702-365711, 2010.
- [20] Arani, A. G. and Roudbari, M. A., Nonlocal piezoelastic surface effect on the vibration of visco-Pasternak coupled boron nitride nanotube system under a moving nanoparticle, *Thin Solid Films*, 542, 232-241, 2013.
- [21] Panchal, M. B. and Upadhyay, S. H., Cantilevered single walled boron nitride nanotube based nanomechanical resonators of zigzag and armchair forms, *Physica E*, 50, 73-82, 2013.
- [22] Yan, Z. and Jiang, L. Y., The vibrational and buckling behaviors of piezoelectric nanobeams with surface effects, *Nanotechnology*, 22, 245703-245710, 2011.
- [23] H. L. Lee, R. P. Chang, W.-J. Chang, Buckling analysis of nonuniform nanowires under axial compression, Proceedings of the world congrees on engineering, 3 (2011).
- [24] Samaei, A. T., Bakhtiari, M., Wang ,G. F., Timoshenko beam model for buckling of piezoelectric nanowires with surface effects, *Nanoscale research letters*, 7, 1-6, 2012.
- [25] Lee, H. L. and Chang, W. J, Surface effects on axial buckling of nonuniform nanowires using non-local elasticity theory, *Micro & nano letters*, 6(1), 19-21, 2011.
- [26] Chiu, M. S. and Chen, T., Higher-order surface stress effects on buckling of nanowires under uniaxial compression, *Procedia engineering*, 10, 397-402,2011.
- [27] Yao, H. and Yun, G., The effect of nonuniform surface surface elasticity on buckling of ZnO nanowires, *Physica E*, 44, 1916-1919, 2012.
- [28] Reddy, J. N., Nonlocal continuum theories of beams for the analysis of carbon nanotubes, *Journal of applied physics*, 103, 023511, 2008.
- [29] Arda, M. and Aydogdu, M., Analysis of free torsional vibration in carbon nanotubes embedded in a viscoelastic medium, *Advances in science and technology research journal*, 9(26), 28-33, 2015.

- [30] Mohammadimehr, M., Mohandes, M., Moradi, M., Size dependent effect on the buckling and vibration analysis of double-bonded nanocomposite piezoelectric plate reinforced by boron nitride nanotube based on modified couple stress theory, *Journal of vibration and control*, 1077546314544513, 2014.
- [31] Sahmani,, S. and Ansari, R., Nonlocal beam models for buckling of nanobeams using state-space method regarding different boundary conditions, *Journal of mechanical science and technology*, 25(9), 2365-2375, 2011.
- [32] Kumar, D., Heinrich, C., Waas, A. M., Buckling analysis of carbon nanotubes modeled using nonlocal continuum theories, *Journal of applied physics*, 103(7), 073521, 2008.
- [33] Sudak, L. J., Column buckling of multiwalled carbon nanotubes using nonlocal continuum mechanics, *Journal of applied physics*, 94(11), 7281-7287, 2003.
- [34] Juntarasaid, C., Pulngern, T., Chucheepsakul, S., Bending and buckling of nanowires including the effects of surface stress and nonlocal elasticity, *Physica E*, 46, 68-76, 2012.
- [35] Wang, Q., Varadan, V. K., Quek, S. T., Small scale effect on elastic buckling of carbon nanotubes with nonlocal continuum models, *Physics letters A*, 357, 130-135, 2006.
- [36] Khademolhosseini, F., Rajapakse, R. K. N. D., Nojeh, A., Torsional buckling of carbon nanotubes based on nonlocal elasticity shell models, *Computational materials science*, 48, 736-742, 2010.
- [37] Civalek, Ö., Demir, Ç., Buckling and bending analyses of cantilever carbon nanotubes using the euler-bernoulli beam theory based on non-local continuum model, *Asian Journal of Civil Engineering*, 12(5), 651-661, 2011.
- [38] Akgöz, B., Civalek, Ö., Strain gradient elasticity and modified couple stress models for buckling analysis of axially loaded micro-scaled beams, *International Journal of Engineering Science*, 49(11), 1268-1280, 2011.
- [39] Akgöz, B., Civalek, Ö., Buckling analysis of linearly tapered micro-columns based on strain gradient elasticity, *Structural Engineering and Mechanics*, 48(2), 195-205, 2013.
- [40] Akgöz, B., Civalek, Ö., Buckling analysis of cantilever carbon nanotubes using the strain gradient elasticity and modified couple stress theories, *Journal of Computational and Theoretical Nanoscience*, 8(9), 1821-1827, 2011.
- [41] Akgöz, B., Civalek, Ö., A new trigonometric beam model for buckling of strain gradient microbeams, *International Journal of Mechanical Sciences*, 81, 88-94, 2014.
- [42] Civalek, Ö., Akgöz, B., Vibration analysis of micro-scaled sector shaped graphene surrounded by an elastic matrix, *Computational Materials Science*, 77, 295-303, 2013.
- [43] Akgöz, B., Civalek, Ö., A novel microstructure-dependent shear deformable beam model, *International Journal of Mechanical Sciences*, 99, 10-20, 2015.
- [44] Emsen, E., Mercan, K., Akgöz, B., Civalek, Ö., Modal Analysis Of Tapered Beam-Column Embedded In Winkler Elastic Foundation, *International Journal of Engineering* & *Applied Sciences*, 7(1), 25-35, 2015.
- [45] Mercan, K., Demir, Ç., Akgöz, B., Civalek, Ö., Coordinate Transformation for Sector and Annular Sector Shaped Graphene Sheets on Silicone Matrix, *International Journal* of Engineering & Applied Sciences, 7(2), 56-73, 2015.
- [46] Akgöz, B., Civalek, Ö., Thermo-mechanical buckling behavior of functionally graded microbeams embedded in elastic medium, *International Journal of Engineering Science* 85, 90-104, 2014.

- [47] Civalek, Ö., The determination of frequencies of laminated conical shells via the discrete singular convolution method, *Journal of Mechanics of Materials and Structures*, 1(1), 163-182, 2006.
- [48] Civalek, Ö., Gürses, M., Free vibration analysis of rotating cylindrical shells using discrete singular convolution technique, *International Journal of Pressure Vessels and Piping*, 86(10), 677-683, 2009.
- [49] Civalek, Ö., Fundamental frequency of isotropic and orthotropic rectangular plates with linearly varying thickness by discrete singular convolution method, *Applied Mathematical Modelling*, 33(10), 3825-3835, 2009.
- [50] Akgöz, B., Civalek, Ö., Longitudinal vibration analysis for microbars based on strain gradient elasticity theory, *Journal of Vibration and Control*, 20(4), 606-616, 2014.
- [51] Challamel, N., Wang, C.M., Elishakoff, I., Discrete systems behave as nonlocal structural elements; bending buckling and vibration analysis, *Euro. J. Mech. A/Solids*, 44, 125-135, 2014.
- [52] Peddieson, J., Buchanan, G.R., McNitt, R.P., Application of nonlocal continuum models to nanotechnology. Int. J. Eng. Sci. 41, 305–312, 2003.