A Circuit Model-Based Analysis of Magnetically Coupled Resonant Loops in Wireless Power Transfer Systems

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ABSTRACT

Magnetically coupled resonant loops can be represented by a lumped element circuit model. Each parameter in the lumped element model can be expressed as a function of loop geometry and the separation between the loops; therefore, the geometry can be systematically changed, and the power transfer efficiency of the coupled loops can be predicted. This paper presents a simulation-based efficiency analysis for wireless power transfer systems utilizing magnetically coupled resonant loops. The behavior of power transfer efficiency is studied for various loop geometry parameters, and the simulation results are presented in detail. These results clearly show that there is a trade-off between peak efficiency and critical coupling distance, both of which depend on the loop size, frequency of operation, and source-load impedances. To verify the accuracy model, two identical circular loops are fabricated and measured, and the measurement results agree well with the model.

Keywords: Wireless power transfer systems, magnetically coupled loops, power transfer efficiency

Introduction

Wireless power transfer (WPT) systems can be categorized into radiative and non-radiative systems [1, 2]. Radiative systems utilize highly directed receiver and transmitter antennas and work at microwave frequencies. Their efficiency is usually lower than that of non-radiative systems and a line of sight propagation is needed; however, power can be transferred to very long distances. On the other hand, non-radiative systems usually exhibit greater efficiency, but they are limited to very short power transfer distances. In these systems, power transfer is carried out via near-field coupling. Coupling is achieved through either electric field (capacitive coupling) or magnetic field (inductive coupling).

Today, due to very broad coverage of existing electrical grids, long-range wireless power transfer systems are not needed in all but several very specific applications. However, widespread use of mobile devices and battery-powered vehicles open up a new usage for wireless power transfer at medium range (midrange) distances, from several centimetres to a couple of meters. For mid-range applications, radiative systems are not convenient as they are based on far-field radiation. Capacitive coupling based non-radiative systems are not safe for daily use applications as the electric field near the transmitter and receiver devices interacts strongly with surrounding objects including humans. Magnetic field interacts weakly with surrounding objects, yet the power transfer distances in traditional inductive coupling systems are quite low (a couple of centimetres) for charging mobile devices and electric vehicles compared with several tens of centimetres.

The magnetically coupled loops, when resonated out, was proposed for use in midrange power transfer [1, 3]. It has been shown that one can transfer power at a distance several times the radius or the length of the loops in magnetically coupled resonant loops [4, 5]. Moreover, when the loops operate in a strongly coupled regime, it is almost possible to achieve a distance-independent power transfer efficiency.
Magnetically coupled resonant loops can be analysed either by a coupled mode theory (CMT) [6,7] or by using an equivalent circuit model. It has been shown that an equivalent circuit model is simpler and quite accurate for the analysis of resonant magnetically coupled loops [7]. Numerous works, based on equivalent circuit models, have been reported [7-15].

This paper presents an equivalent circuit model based analysis for WPTs using magnetically coupled resonant loops. As compared to previously reported works on circuit model based analysis of WPTs, this work presents comprehensive simulation results by systematically changing the geometry parameters and source & load resistances. Using these simulation results, the effect of loop size and source & load resistances on efficiency response in both weakly and strongly coupled regimes is presented. The conditions for operating in each of these regimes are examined in terms of loop size and load & source resistances. Using these simulation results, two identical circular loops are fabricated and measured. Measurement results exhibit good agreement with the model.

An Equivalent Circuit Model for Inductively Coupled Resonant Loops

The magnetically coupled resonant loops, as in Figure 1a, can be represented by a circuit model shown in Figure 1b. In this model, the magnetic coupling between the loops is represented by a mutual inductance (M). The components $L_1$, $R_c$, and $C_1$ are the self-inductance, parasitic resistance and the resonance capacitance for the transmitting loop, respectively. Similarly, $L_2$, $R_r$ and $C_2$ are the same parameters for the receiving loop. Each component in the circuit is dependent on the geometry parameters of the loops.

The self-inductance ($L_{1,2}$) for a circular loop as a function of loop radius ($r$), radius of the wire ($a$) and number of turns ($N$) can be found as follows [4]:

$$L_{1,2} = N^2 \mu_0 \sqrt{(r-a)(r+a)} \left[ \left( \frac{2}{k} \right) K(k) - \left( \frac{2}{k} \right) E(k) \right]$$ (1.a)

where $K$ and $E$ are the complete elliptic integrals of the first and second kind as follows

$$K(k) = \int_0^{\pi/2} \frac{d\beta}{\sqrt{1-k^2 \sin^2 \beta}}$$ (1.b)

$$E(k) = \int_0^{\pi/2} d\beta \sqrt{1-k^2 \sin^2 \beta}$$ (1.c)

$$k = \sqrt{\frac{4(r-a)(r+a)}{(2r)^2 + (2a)^2}}.$$ (1.d)

where $\mu_0$ is the magnetic permeability of the air. The parasitic resistance ($R_{1,2}$) includes the effect of both conductor loss ($R_c$) and the radiation loss ($R_r$) and can be calculated as follows [16]:

$$R_c = \frac{2\pi r N}{2\pi \rho} \sqrt{\mu_0 \omega}$$ (2.a)

$$R_r = \frac{2\pi r N^2}{2\pi \rho} \sqrt{\mu_0}$$ (2.b)

$$R = R_c + R_r$$ (2.c)

where $\omega$ is the radial frequency, $\lambda$ is the wavelength at the operating frequency and $\sigma$ is the conductivity of the wire. The mutual inductance between the loops can be calculated as function of loop geometry and separation between the loops as follows [4]:

$$M = N^2 \mu_0 \sqrt{r_1 r_2} \left( \frac{2}{l} - l F(l) - \frac{2}{l} E(l) \right)$$ (3a)

where $K$ and $E$ are the complete elliptic integrals of the first and second kind as follows:

$$F(l) = \int_0^{\pi/2} \frac{d\beta}{\sqrt{1-l^2 \sin^2 \beta}}$$ (3b)

$$E(l) = \int_0^{\pi/2} d\beta \sqrt{1-l^2 \sin^2 \beta}$$ (3c)

$$l = \sqrt{\frac{4rr}{(r_1 + r_2)^2 + d^2}}.$$ (3d)

Analysis of Circuit Model for Efficiency Calculations

For efficiency calculations, an RF source with an internal resistance ($R_s$) and a load resistance ($R_l$) are connected to the input and output of the model, respectively, as shown in Figure 1.
A mathematical expression for the power transfer efficiency can be achieved as (4) by analyzing the circuit model in Figure 2. Usually, the load resistance ($R_l$) is fixed and its value is dependent on the device to be powered by wireless transfer. The input impedance seen towards the transmitting resonant loop ($Z_{in}$), as shown in Figure 2, is the function of loop geometry ($C_1, L_1, R_1, C_2, L_2, R_2$ and $M$), the separation between the loops ($M_s$), and the load resistance ($R_l$). A mathematical expression for $Z_{in}$ can be obtained as given in (5).

The power transfer efficiency ($\eta$) is degraded by the mismatch between $Z_{in}$ and $R_s$ and is proportional to the value of $1-|\Gamma_{in}|^2$. This mismatch effect is accounted in the efficiency expressions (4) with the inclusion of $R_s$. The above expressions in (1)-(4) constitute the whole set of equations to predict the efficiency of magnetically coupled resonant loop pairs.

**Simulation Results**

Coupled loops can operate in strongly coupled regimes or weakly coupled regimes depending on the $M$ between the loops, losses of each loop and the load resistance ($R_l$) which represents the device to be powered or charged. The $\eta$ vs frequency response exhibits different characteristics depending on whether the coupled loops operate either one these regimes. A typical $\eta$ response of two coupled identical resonant loops, for various ($M_s$), is shown in Figure 3.

Each loop is resonant at $\omega_0$ in order to achieve a maximum efficiency value. As seen in Figure 3, when $M_s$ are lower than $M_s$ (e.g. $M_3, M_s$), peak efficiency is observed at $\omega_0$ and increases with increasing $M$. This is the typical efficiency characteristic of coupled resonant loops operating at the weakly coupled regime. When $M_s$ are larger than $M_s$ (e.g. $M_5, M_s$), two efficiency peaks are observed at two resonance frequencies of $\omega_{odd}$ and $\omega_{even}$ which are called the odd mode resonance frequency and the even mode resonance frequency, respectively. The efficiency is maximum at $\omega_{odd}$ and $\omega_{even}$ and is independent of $M$. This is the typical efficiency response characteristic of coupled resonant loops operating in strongly coupled regime. The $\omega_{odd}$ and $\omega_{even}$ get separated from the $\omega_0$ with increasing $M$. In other words, as $M$ increases, coupled resonant loops get in to strongly coupled regime from the weakly coupled regime and $\omega_0$ is split into two more resonance frequencies of $\omega_{odd}$ and $\omega_{even}$ in the strongly coupled regime (Figure 3). The $M$ at which coupled resonant loops change their operation from weakly coupled regime to strongly coupled regime is called the critically coupled point (e.g. $M_3$ in Figure 3). Strongly coupled regime is very advantageous in wireless power transfer; because, $M$-independent efficiency means that one can transfer a constant power even if the distance between the loops varies. To realize that, one needs to set up a frequency-tuned system like those that have recently been proposed in several works [18-21]. However, the distance range over which efficiency is constant is limited to a critically coupling distance ($d_{critical}$). The efficiency response, in particular the peak efficiency and the critically coupled distance ($d_{critical}$), are dependent on the loop size and the source and load resistances.

Figure 4 a-d, show four different graphs for coupled identical single turn ($N=1$) circular loops with a radius of 25 cm, 50 cm,
The loops are resonant with series connected capacitances at 5 MHz. Each graph depicts a simulated efficiency ($\eta$) vs frequency data for various separation distances ($d$) between the loops. Here the simulations are performed for the same load and source resistance values of 2 $\Omega$ ($R_s = R_L = 2 \Omega$).

Figure 4, 5-7, show similar efficiency graphs for 10 $\Omega$, 25 $\Omega$ and 50 $\Omega$ source & load resistance values, respectively.

In Figure 4, where source and load resistances are 2 $\Omega$, peak efficiencies at strongly coupled regime are 86 $\%$, 77 $\%$, 68 $\%$ and 59 $\%$, for loops with radius of 25 cm, 50 cm, 75 cm and 100 cm, respectively. It can clearly be seen that the peak efficiency is significantly reduced with increasing loop sizes. The critical coupling distance, on the other hand, increases from 35 cm to 250 cm as loop size increases from 25 cm to 100 cm. Such behaviour of peak efficiency and critical coupling distance can be explained as follows:

- As the size of the loops increases, larger $M$ values are obtained for the same loop separations; therefore, loops can stay in strongly coupled regime for larger separations between the loops ($d_{\text{critical}}$ increases). This is advantageous because one can operate with almost distance-independent efficiency in larger loop separations. However, due to frequency splitting phenomena in strongly coupled operation, one needs a frequency tuned wireless power transfer system.
- As the size of the loops increases, the radiation loss resistance ($R_{\text{rad}}$) and conductor loss resistance ($R_{\text{c}}$) increases; hence the peak efficiency decreases due to these larger loss resistances.

It should be remembered that the efficiency plots in Figure 4 are for source and load resistances of 2 $\Omega$. Figure 5-7, show efficiency graphs for source & load resistances ($R_s$ and $R_L$) of 10 $\Omega$, 25 $\Omega$ and 50 $\Omega$, respectively. To analyze the effect of $R_s$ and $R_L$ on peak efficiency and critical coupling distance, one can focus on the efficiency response for the same loop size in each figure. For example, in Figure 4 d, Figure 5 d, Figure 6 d and Figure 7 d, where loop radius is 100 cm, it can clearly be seen that the peak efficiency increases as $R_s$ and $R_L$ increase. However, the critical coupling distance decreases with increasing $R_s$ and $R_L$.

A major conclusion from these simulation results is that one can increase critical coupling distance by either increasing loop size or by decreasing $R_s$ and $R_L$ values. In either case, any increase in critical coupling distance would be at the cost of reduced peak efficiency. Therefore, loop sizes and input & output matching circuits should be carefully designed considering this tradeoff. Figure 8 a, b show peak efficiency ($\eta_{\text{peak}}$) and critical coupling distance ($d_{\text{critical}}$) as a function of loop radius for various $R_s$ and $R_L$. 

Figure 4. a-d. Efficiency ($\eta$) vs Frequency graphs for circular loops of radius of 25 cm (a); 50 cm (b); 75 cm (c) and 100 cm. $R_s = R_L = 2 \Omega$ (d)
Figure 5. a-d. Efficiency ($\eta$) vs Frequency graphs for circular loops of radius of 25 cm (a); 50 cm (b); 75 cm (c) and 100 cm (d) $R_s = R = 10 \, \Omega$.

Figure 6. a-d. Efficiency ($\eta$) vs Frequency graphs for circular loops of radius of 25 cm (a); 50 cm (b); 75 cm (c) and 100 cm (d) $R_s = R = 25 \, \Omega$. 
In the previous section, a significant analysis was performed using circuit model and geometry-dependent analytic equations. In this section, we fabricated two multiturn circular loops to verify the accuracy of the model utilized in simulations. The loops are fabricated by winding solid copper wires on an epoxyglass. The radius of the loops is 25 cm and the number of turns is 5. Each loop’s self-inductance $L$ is measured using an RLC meter. Then each loop is resonated out using a series connected capacitor at 1 MHz. Simulations were performed at 5 MHz. Here, we intentionally resonated out the loops at 1 MHz to reduce the inter-winding parasitic capacitance in multiturn loops. Table 1 shows the theoretical and measured self-inductance values of each 5 turn loops. The required values of series connected capacitors to form a resonance at 1 MHz are also shown in Table 1.

The S-parameters of magnetically coupled resonant loops are measured from 100 kHz to 5 MHz using Rohde & Schwarz FSH8 Spectrum +Network analyzer as shown in Figure 9.
The efficiency is extracted from the measured S-parameters by taking the square of $S_{21}$ as follows:

$$\eta = \left| S_{21} \right|^2 \quad (7)$$

Measured efficiency as a function of frequency for various loop separations are shown in Figure 10 a-f. $R_S$ and $R_L$ resistances in these measurements are 20 Ω and 10 Ω, respectively. Figure 11 shows the plot of peak efficiency vs distance for both simulation and measurement results. As seen in Figure 10, 11; the measurement results agree well with the simulation results specifically in strongly coupled regime. The critical coupling distance $d_{\text{critical}}$ is approximately 30 cm. The difference between the simulations and the measurement results could be attributed to some losses that are not considered in the model, such as the proximity effect of multi-turn loops and the loss of series connected resonance capacitors.

**Conclusion**

In this paper we have presented a circuit model based analysis of magnetically coupled resonant loops for use in wireless power transfer applications. The theoretical and measured values of self inductances for fabricated loops are given in Table 1.

<table>
<thead>
<tr>
<th></th>
<th>Theoretical</th>
<th>Measured</th>
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<tbody>
<tr>
<td>Loop 1</td>
<td>$L_1 = 35.8 \mu H$</td>
<td>$L_1 = 29.1 \mu H$</td>
</tr>
<tr>
<td></td>
<td>$C_1 = 707 \text{ pF}$</td>
<td>$C_1 = 875 \text{ pF}$</td>
</tr>
<tr>
<td>Loop 2</td>
<td>$L_2 = 35.8 \mu H$</td>
<td>$L_2 = 28.8 \mu H$</td>
</tr>
<tr>
<td></td>
<td>$C_2 = 707 \text{ pF}$</td>
<td>$C_2 = 875 \text{ pF}$</td>
</tr>
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</table>

Figure 9. S-parameter measurement setup for magnetically coupled loops

Figure 10. Comparison of simulation and measurement results for magnetically coupled resonant loops with $R_S = 20 \Omega$ and $R_L = 10 \Omega$.
power transfer systems. The model employs analytical equations for loop’s inductive parameters such as self and mutual inductances, and loss parameters such as conductor and radiation loss resistances. Hence overall power transfer efficiency was accurately predicted using this model. A thorough analysis was also performed showing how the peak efficiency and the critical coupling distance changes as loop size and load & source resistances vary. The trade off between peak efficiency and critical coupling distance is discussed in detail. To validate the model’s accuracy, two 5 turn identical loops are fabricated and their efficiencies are measured. Future work would be to employ the proximity loss effects into the model and to conduct a study on reducing the electrical losses in the loops.

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Conflict of Interest: The authors have no conflicts of interest to declare.

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