DECOMPOSITION OF CURVATURE TENSOR FIELD IN A RECURRENT KAEHLERIAN SPACE

ومعاورة المتعام وترجعا

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In the present paper the decomposition of the curvature tensor field R^{h}_{ijk} of a recurrent Kaehlerian space in terms of a vector field and two tensor fields has been considered and several theorems about this decomposition have been investigated.

1. Introduction. An n(=2m) dimensional Kaehletian space K^n is a Riemannian space which admits a tensor field ϕ_i^h satisfying the conditions

$$\phi_i^{\ h} \phi_h^{\ j} = -\delta_{ij}^{\ 1}, \tag{1.1}$$

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With the second second

$$\phi_{ij} = -\phi_i^{\ j} \quad (\phi_{ij} = \phi_i^{\ a} g_{aj}) , \qquad (1.2)$$

and

$$\phi^h_{i,j} = 0 \tag{1.3}$$

where the comma followed by an index denotes the operator of covariant differentiation with respect to the metric tensor g_{ij} of the Riemannian space.

The Riemannian curvature tensor field is defined by

$$R^{h}_{ijk} \stackrel{\text{def.}}{=} \partial_{i} \left< {}^{h}_{jk} \right> - \partial_{j} \left< {}^{h}_{ik} \right> + \left< {}^{h}_{il} \right> \left< {}^{l}_{jk} \right> - \left< {}^{h}_{jl} \right> \left< {}^{l}_{ik} \right> {}^{2}; \qquad (1.4)$$

the Ricci tensor and scalar curvature are given by $R_{ij} = R^a_{aij}$ and $R = g^{ij} R_{ij}$ respectively.

It is well known that these tensors satisfy the identity $[2]^{3}$

$$R^{a}_{\ ijk,a} = R_{jk,i} - R_{ik,j} \ . \tag{1.5}$$

The holomorphically projective curvature tensor P^{h}_{ijk} is defined by

⁾ All Latin indices run over the range from 1 to n.

²⁾ $\partial_i = \partial/\partial x^i$, where $\{x^i\}$ denotes real local coordinates.

³ Numbers in square brackets refer to the references at the end of the paper.

$$P^{h}_{ijk} \stackrel{\text{def.}}{=} R^{h}_{ijk} + \frac{1}{n+2} \left(R_{ik'} \delta_{j}^{h} - R_{jk} \delta_{i}^{h} + S_{ik} \phi_{j}^{h} - S_{jk} \phi_{l}^{h} + 2 S_{ij} \phi_{k}^{h} \right) \quad (1.6)$$

where $S_{ij} = \phi_i^{\ a} \ R_{aj}$.

The Bianchi identities in K^n are given by

$$R^{h}_{ijk} + R^{h}_{jki} + R^{h}_{kil} = 0 (1.7)$$

and

$$R^{h}_{ijk,a} + R^{h}_{ika,j} + R^{h}_{iaj,k} = 0.$$
(1.8)

A Kaehlerian space K^n is said to be Kaehlerian recurrent if its curvature tensor field satisfies the condition ([³])

$$R^{h}{}_{ijk,a} = \lambda_{a} R^{h}{}_{ijk} \tag{1.9}$$

where λ_a is a non-zero recurrence vector field.

The following relations follow immediately from (1.9):

$$R_{ij,a} = \lambda_a R_{ij} \tag{1.10}$$

and

 $R_{a} = \lambda_a R$.

In the present paper, we have considered the decomposition of curvature tensor field R^{h}_{ijk} in terms of a vector field and two tensor fields and several theorems have been investigated.

2. Decomposition of Curvature Tensor Fields \mathbf{R}^{h}_{ijk} . We consider the decomposition of recurrent curvature tensor field R^{h}_{ijk} in the following form:

$$R^{h}_{ijk} = X^{h}_{i} v^{l} \psi_{ijk} , \qquad (2.1)$$

where v^{l} is a non-zero vector field and X_{l}^{h} , ψ_{iJk} are two non-zero tensor fields such that

$$X_i^h \lambda_h = p_i \tag{2.2}$$

and

$$\lambda_h v^h = 1 . \tag{2.3}$$

 p_l is called decomposed vector field and this is a non-zero vector field.

We shall prove the following :

Theorem (2.1). Under the decomposition (2.1), the Bianchi identities for R^{h}_{ijk} take the form

$$\Psi_{ijk} + \Psi_{jkl} + \Psi_{klj} = 0$$
, $\Psi_{ijk} = - \Psi_{ikj}$ (2.4)

and

$$\lambda_a \,\psi_{ijk} + \lambda_j \,\psi_{ika} + \lambda_k \,\psi_{iaj} = 0 \,. \tag{2.5}$$

Proof. In view of equations (1.10), (1.11), (1.9) and (2.1), we obtain

$$X_{l}^{h} v^{l} (\psi_{ijk} + \psi_{jki} + \psi_{kij}) = 0$$
(2.6)

and

$$X_l^h v^l \left(\lambda_a \psi_{ijk} + \lambda_j \psi_{ika} + \lambda_k \psi_{iaj} \right) = 0 . \qquad (2.7)$$

The identities (2.4) and (2.5) follow immediately from these equations and the fact $X_l^h v^l \neq 0$.

The following main theorems may be proved in the same way as it is proved in the recent paper [5]:

Theorem (2.2). The vector field λ_a and the tensor field X_l^h given by equations (1.9) and (2.1) behave like recurrent vector and recurrent tensor fields and their recurrent forms are given by

$$\lambda_{a,m} = \mu_m \,\lambda_a \tag{2.8}$$

and

$$X^h_{l,m} = v_m X_l^h , \qquad (2.9)$$

where μ_m and ν_m are non-zero recurrence vector fields.

Theorem (2.3). Under the decomposition (2.1), the decomposed vector field p_i behaves like a recurrent vector field and its recurrent form is given by

$$p_{l,m} = (\mu_m + \nu_m) p_l \,. \tag{2.10}$$

Now, we prove the following :

Theorem (2.4). Under the decomposition (2.1), the vector field v^{t} and the tensor field ψ_{ijk} behave like recurrent vector and recurrent tensor fields.

Proof. Multiplying equation (2.5) by v^{α} and using relation (2.3), we obtain

$$\Psi_{ijk} = \lambda_k \, \Psi_{ij} - \lambda_j \, \Psi_{ik} \,, \tag{2.11}$$

where $\psi_{ij_k} v^k = \psi_{ij}$ is a tensor field.



Therefore, the relation (2.1) takes the form whether the state of the

$$R^{h}{}_{ijk} = X_{l}{}^{h} \operatorname{v}^{l} \left(\lambda_{k} \, \psi_{ij} - \lambda_{j} \, \psi_{ik} \right). \tag{2.12}$$

Differentiating equation (2.12) covariantly with respect to x^m and using equations (1.9), (2.8), (2.9), (2.12), we get

$$(\lambda_k \psi_{ij} - \lambda_j \psi_{ik}) \mathbf{v}^i_{,m} = \mathbf{v}^i \{ \mathbf{v}_m (\lambda_j \psi_{jk} - \lambda_k \psi_{ij}) + \mu_m (\lambda_j \psi_{ik} - \lambda_k \psi_{ij}) \}.$$
(2.13)

Multiplying this equation by v^a , we obtain

$$(\lambda_k \psi_{ij} - \lambda_j \psi_{ik}) v^a v^l_{,m} = v^l v^a \{ \nu_m (\lambda_j \psi_{ik} - \lambda_k \psi_{ij}) + \mu_m (\lambda_j \psi_{ik} - \lambda_k \psi_{ij}) \}, \quad (2.14)$$

which yields

$$v^{a} v^{l}_{m} = v^{l} v^{a}_{m}.$$
 (2.15)

Since $v^{l} \neq 0$, there exists a proportional non-zero vector field π_{m} such that

$$\mathbf{v}^l_{,m} = \boldsymbol{\pi}_m \, \mathbf{v}^l \,. \tag{2.16}$$

Therefore, v^l is recurrent vector field.

Further, differentiating equation (2.1) covariantly with respect to x^m and using equations (1.9), (2.1), (2.8), (2.9) and (2.11), we obtain

$$\Psi_{ijk,m} = (\lambda_m - \nu_m - \pi_m) \Psi_{ijk} . \qquad (2.17)$$

Hence, ψ_{ijk} is a recurrent tensor field.

If $v_m + \pi_m \neq 0$, we have

Corollary (2.1). Under the decomposition (2.1), the vector field $X_l^h v^l$ is recurrent with the recurrence vector field $(v_m + \pi_m)$.

Proof. Differentiating the vector field $X_l^h v^l$ covariantly with respect to x^m and using equations (2.9) and (2.16), we get the proof.

On the other hand, if $v_m + \pi_m = 0$, we have

Corollary (2.2). Under the decomposition (2.1), ψ_{ijk} will be recurrent with the same recurrence vector λ_m as the curvature tensor field R^{h}_{ijk} .

Proof. The proof follows immediately from equation (2.17).

Theorem (2.5). Under the decomposition (2.1), the vector field v^t and tensor fields R^h_{ijk} , R_{ij} , ψ_{ijk} satisfy the relations

DECOMPOSITION OF CURVATURE TENSOR FIELD

$$\lambda_h R^h_{ijk} = \lambda_i R_{jk} - \lambda_j R_{ik} = p_l v^l \psi_{ijk} . \qquad (2.18)$$

Proof. With the help of equations (1.5), (1.9) and (1.10) we obtain

$$\lambda_h R^h{}_{ijk} = \lambda_i R_{jk} - \lambda_j R_{ik} . \qquad (2.19)$$

Multiplying equation (2.1) by λ_h and using relation (2.2), we obtain

$$\lambda_h R^h{}_{ijk} = p_l \, \mathbf{v}^l \, \psi_{ijk} \,. \tag{2.20}$$

From equations (2.19) and (2.20), we get the relations (2.18).

Theorem (2.6). Under the decomposition (2.1), the curvature tensor R^{h}_{ijk} and holomorphically projective curvature tensor fields are equal iff

$$\delta_{j}{}^{h} \psi_{ik} - \delta_{i}{}^{h} \psi_{jk} + \psi_{ak} \left(\phi_{j}{}^{h} \phi_{i}{}^{a} - \phi_{i}{}^{h} \phi_{j}{}^{a} \right) + 2 \phi_{k}{}^{h} \phi_{i}{}^{a} \psi_{aj} = 0.$$
 (2.21)

Proof. Equation (1.6) may be expressed in the form

$$P^{h}{}_{ijk} = R^{h}{}_{ijk} + D^{h}{}_{ijk} , \qquad (2.22)$$

where

$$D^{h}_{ijk} = \frac{1}{n+2} \left(R_{ik} \, \delta_{j}^{h} - R_{jk} \, \delta_{i}^{h} + S_{ik} \, \phi_{j}^{h} - S_{jk} \, \phi_{i}^{h} + 2S_{ij} \, \phi_{k}^{h} \right). \quad (2.22a)$$

Contracting indices h and k in (2.1), we obtain

$$R_{ij} = X_l^k \,\mathbf{v}^l \,\psi_{ijk} \,. \tag{2.23}$$

With the help of equation (2.23), we have

$$S_{ij} = \phi_i^{\ a} R_{aj} = \phi_i^{\ a} X_i^r \, v^l \, \psi_{ajr} \,. \tag{2.24}$$

Making use of equations (2.23) and (2.24) in (2.22a), we obtain

$$D^{h}_{ijk} = \frac{X_{i}^{r} \nabla^{l}}{n+2} \{ \Psi_{ikr} \, \delta_{j}^{h} - \Psi_{jkr} \, \delta_{i}^{h} + \Psi_{akr} \, (\phi_{j}^{h} \, \phi_{i}^{a} - \phi_{i}^{h} \, \phi_{j}^{a}) + 2 \, \phi_{k}^{h} \, \phi_{i}^{a} \, \Psi_{ajr} \} \,.$$

$$(2.25)$$

From equation (2.22a), it is clear that $P_{ijk}^{h} = R_{ijk}^{h}$, iff $D_{ijk}^{h} = 0$, which in view of equation (2.25) becomes

$$\psi_{ikr}\,\delta_{j}{}^{h} - \psi_{jkr}\,\delta_{i}{}^{h} + \psi_{akr}\,(\phi_{j}{}^{h}\,\phi_{i}{}^{a} - \phi_{i}{}^{h}\,\phi_{i}{}^{a}) + 2\,\phi_{k}{}^{h}\,\phi_{i}{}^{a}\,\psi_{ajr} = 0\,.$$
(2.26)

Multiplying this equation by v^r and using the relation

$$\Psi_{ijk} \, \mathrm{v}^k = \Psi_{ij}$$
 ,

we have the required equation.

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ÖZET

Bu çalışmada bir tekrarlamalı Kaehler uzayının R^{h}_{ijk} eğrilik tansörü alanının bir vektör alanı ile iki tansör alanı cinsinden ayrılış formülü ele alınarak bu ayrılışa dair bazı teoremler araştırılmıştır.