



Limit-point Classification for Singular Conformable Fractional Sturm-Liouville Operators

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ABSTRACT. In this work, we study the following conformable fractional Sturm–Liouville problem

$$l[y] = -T_\alpha(p(t)T_\alpha y(t)) + q(t)y(t),$$

where $t \in [0, \infty)$, the real-valued functions p and q satisfy the following conditions:

- (i) $q \in L_\alpha^2[0, \infty)$,
- (ii) p is absolutely continuous on $[0, \infty)$,
- (iii) $p(t) > 0$ for all $t \in [0, \infty)$.

The conformable fractional Sturm–Liouville problem is of the limit-point case if the number of linearly independent α –square integrable solutions of the equation $l[y] = \lambda y$ is less than 2. We give a criterion for the limit point classification of conformable fractional Sturm-Liouville operators in singular case.

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1. INTRODUCTION

Differential equations; it is used as a mathematical model in many different sciences such as physics, chemistry, engineering, biology, and economics. Newton and Leibnitz studied differential and integral calculus in the second half of the 17th century. Due to the limited application of fractional derivatives and problem-solving, it was not possible to make sufficient studies on this subject in the 17th century. With the discovery of the fractional calculus technique; many problems such as chaos, propagation and wave motions, filtering and irreversibility can be modeled and explained more accurately using fractional analysis. This interesting topic is used in fields such as diffusion, Schrödinger equation, materials science, transmission lines theory, chemical analysis of liquids, heat transfer, fluids, biology, biophysics, bioengineering, electromagnetic theory, mechanics, physics and control theory, analytical and numerical methods. The day comes out with new and interesting applications. Besides, the use of fractional calculation to model the physical states of viscoelastic (adhesive and flexible) materials (cartilage, skin, muscle) in which hysteresis, memory and tension factors occur naturally occurs spontaneously [10–13].

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In recent years, conformable fractional derivative and conformable fractional integral were defined by Khalil and his friends. In their study, they presented a linearity requirement for the new definition of fractional derivative, multiplication rule, division rule, fractional Rolle theorem and fractional average value theorem for conformable fractional derivative [7]. Later in [1], Abdeljawad defined the right and left conformable fractional derivatives, the fractional chain rule and fractional integrals of higher orders.

Conformable fractional derivative aims to expand the derivative definition as known by providing the natural characteristics of classical derivative and to gain new perspectives for differential equation theory with the help of conformable differential equations obtained as using this derivative definition [8].

The Levinson criterion is considered to be one of the most important limit-point criteria for the differential operator, since it provides a more general boundary point criterion [9]. Baleanu et al. [3] gave the Levinson criterion for 2α -order conformable fractional Sturm-Liouville operator. Zhaowen et al. studied the 2α -order Sturm-Liouville operator with conformable fractional derivatives. They have obtained two limit-point case criteria for this operator [15].

In this study, we shall give a criterion for the limit-point classification of conformable fractional Sturm-Liouville operators in singular case. While proving our results, we use the machinery and methods of [4–6].

Now, some preliminary concepts and theorems are presented for the convenience of the reader.

Definition 1.1 ([1]). Let $0 < \alpha < 1$. For a function $f : (0, \infty) \rightarrow \mathbb{R} := (-\infty, \infty)$, the conformable derivative of order α of f at $t > 0$ is defined by

$$T_\alpha f(t) = \lim_{\varepsilon \rightarrow 0} \frac{f(t + \varepsilon t^{1-\alpha}) - f(t)}{\varepsilon},$$

and the conformable derivative at 0 is defined by

$$(T_\alpha f)(0) = \lim_{t \rightarrow 0^+} (T_\alpha f(t)).$$

Definition 1.2 ([1]). The left conformable derivative of order α of a function $f : [a, \infty) \rightarrow \mathbb{R}$ is defined by

$$(T_\alpha^a f)(t) = \lim_{\varepsilon \rightarrow 0} \frac{f(t + \varepsilon(t-a)^{1-\alpha}) - f(t)}{\varepsilon}, \quad 0 < \alpha \leq 1.$$

When $a = 0$ we write T_α . If $(T_\alpha f)(t)$ exists on (a, b) then

$$(T_\alpha^a f)(a) = \lim_{t \rightarrow a^+} (T_\alpha^a f)(t).$$

Definition 1.3 ([1]). The right conformable derivative of order α of $f : (-\infty, b]$ is defined by

$$({}^b T_\alpha f)(t) = -\lim_{\varepsilon \rightarrow 0} \frac{f(t + \varepsilon(b-t)^{1-\alpha}) - f(t)}{\varepsilon},$$

where $0 < \alpha \leq 1$ and $({}^b T_\alpha f)(b) = \lim_{t \rightarrow b^-} ({}^b T_\alpha f)(t)$.

Theorem 1.4 ([1]). Let $x, y : [0, b] \rightarrow \mathbb{R}$ be two functions such that x and y are conformable fractional differentiable. Then, we have

$$\int_0^b y(t) T_\alpha(x)(t) d_\alpha t + \int_0^b x(t) T_\alpha(y)(t) d_\alpha t = x(b)y(b) - x(0)y(0). \quad (1.1)$$

Definition 1.5 ([2]). Let $L_\alpha^2[0, \infty)$ be the space of all complex-valued functions defined on $[0, \infty)$ such that

$$\|y\| := \sqrt{\int_0^\infty |y(t)|^2 d_\alpha t} = \sqrt{\int_0^\infty |y(t)|^2 t^{\alpha-1} d_\alpha t} < \infty.$$

The space $L_\alpha^2[0, \infty)$ is a Hilbert space with the inner product

$$\langle y, z \rangle := \int_0^\infty y(t) \overline{z(t)} d_\alpha t, \quad \text{where } y, z \in L_\alpha^2[0, \infty).$$

Definition 1.6 ([6]). Two functions f, g will be said to be effectively proportional if there are constants β_1, β_2 , not both zero, such that $\beta_1 f \equiv \beta_2 g$. A null function is effectively proportional to any function.

Theorem 1.7. *i) If $\beta_1, \beta_2, \dots, \beta_n$ are positive and $\beta_1 + \beta_2 + \dots + \beta_n = 1$, then*

$$\int |f^{\beta_1} g^{\beta_2} \dots l^{\beta_n}| d_\alpha t < \left(\int |f| d_\alpha t \right)^{\beta_1} \left(\int |g| d_\alpha t \right)^{\beta_2} \dots \left(\int |l| d_\alpha t \right)^{\beta_n},$$

unless one of the functions is null or all are effectively proportional.

ii) If $k > 1$ and $\frac{1}{k} + \frac{1}{k'} = 1$, then

$$\int |fg| d_\alpha t < \left(\int |f|^k d_\alpha t \right)^{\frac{1}{k}} \left(\int |g|^{k'} d_\alpha t \right)^{\frac{1}{k'}}$$

unless either (a) f^k and $g^{k'}$ are effectively proportional or (b) fg is null.

Proof. See [6], as the proof is unchanged except for minor notational modifications. □

Theorem 1.8. *If $p > 1$, $f \in L_\alpha^p(0, a)$, and*

$$F(x) = \int_0^x |f(t)| d_\alpha t,$$

then we have

$$F(x) = o\left(\left[\frac{x}{\alpha}\right]^{\frac{\alpha}{p}}\right),$$

for small x .

Proof. It follows from Theorem 1.7 that

$$F^p \leq \int_0^x |f|^p d_\alpha t \left(\int_0^x d_\alpha t \right)^{p-1} = \frac{x^{\alpha(p-1)}}{\alpha} \int_0^x |f|^p d_\alpha t,$$

and the second factor tends to 0. □

Theorem 1.9. *If $p > 1$ and $f \in L_\alpha^p(0, \infty)$, then $F(x) = o\left(\left[\frac{x}{\alpha}\right]^{\frac{\alpha}{p}}\right)$ both for small and for large x .*

Proof. For small x , the proof is obtained by Theorem 1.8. Now, we prove the result for large x . Choose η so that

$$\int_\eta^\infty |f|^p d_\alpha x < \varepsilon^p$$

and suppose $x > \eta$. Then, we obtain

$$\begin{aligned} (F(x) - F(\eta))^p &= \left(\int_\eta^x |f| d_\alpha t \right)^p \leq \frac{(x - \eta)^{\alpha(p-1)}}{\alpha} \int_\eta^x |f|^p d_\alpha t < \frac{\varepsilon^p x^{\alpha(p-1)}}{\alpha}, \\ F(x) &< F(\eta) + \varepsilon \frac{x^{\frac{\alpha}{p}}}{\alpha} < 2\varepsilon \frac{x^{\frac{\alpha}{p}}}{\alpha} \end{aligned}$$

for sufficiently large x . □

2. MAIN RESULTS

We consider the following conformable fractional singular Sturm–Liouville expression

$$l(y) = -T_\alpha(p(t)T_\alpha y(t)) + q(t)y(t) \text{ on } [0, \infty), \tag{2.1}$$

where the real-valued functions p and q satisfy the following conditions:

- (i) $q \in L_\alpha^2[0, b]$ for all $b > 0$,
 - (ii) p is absolutely continuous on $[0, b]$ for all $b > 0$,
 - (iii) $p(t) > 0$ for all $t \in [0, \infty)$.
- (2.2)

The set D is defined by: $y \in D$ if

$$\begin{aligned} & \text{(i)} \quad y \in L_\alpha^2[0, \infty), \\ & \text{(ii)} \quad T_\alpha y \text{ is absolutely continuous on } [0, b] \text{ for all } b > 0, \\ & \text{(iii)} \quad l(y) \in L_\alpha^2[0, \infty), \\ & \text{(iv)} \quad y(0) = 0. \end{aligned} \tag{2.3}$$

For $y_1, y_2 \in D$, we have the following Green's formula [2]

$$\int_0^\infty l(y_1)(t)\overline{y_2(t)}d_\alpha t - \int_0^\infty y_1(t)\overline{l(y_2)(t)}d_\alpha t = [y_1, y_2](\infty) - [y_1, y_2](0),$$

where

$$[y_1, y_2](t) = p(t) \left\{ y_1(t)\overline{T_\alpha y_2(t)} - T_\alpha y_1(t)\overline{y_2(t)} \right\}, t \in [0, \infty)$$

and

$$[y_1, y_2](\infty) := \lim_{t \rightarrow \infty} [y_1, y_2](t).$$

Theorem 2.1. *If the function q is bounded below on $[0, \infty)$ and*

$$\int_0^\infty \{p(t)\}^{-\frac{1}{2}} d_\alpha t < \infty, \tag{2.4}$$

then the differential operator l defined by (2.1) is in the limit-point case at infinity.

Proof. It is known, see ([4, 14]), that l is limit-point at infinity if and only if

$$[y_1, y_2](\infty) = 0.$$

for all $y_1, y_2 \in D$. Hence it is sufficient to prove that, for all $y_1, y_2 \in D$ (where y_1 and y_2 are real-valued functions),

$$\lim_{b \rightarrow \infty} p(b)y_1(b)T_\alpha y_2(b) = 0. \tag{2.5}$$

Without loss of generality we may assume that there is a positive constant k such that

$$q(t) \geq k > 0 \text{ for all } t \in [0, \infty). \tag{2.6}$$

From (2.6), for all $b > 0$, we get,

$$Q(b) = \int_0^b q(t)d_\alpha t \geq k \frac{b^\alpha}{\alpha}. \tag{2.7}$$

(2.7) implies that

$$[Q(b)]^{-1/2} \leq \left[k \frac{b^\alpha}{\alpha} \right]^{-1/2}, \tag{2.8}$$

for all $b > 0$.

By integration by parts (1.1), we obtain

$$\int_0^b \{pT_\alpha y_1 T_\alpha y_2 + qy_1 y_2\} d_\alpha t = [py_1 T_\alpha y_2] \Big|_0^b + \int_0^b l(y_2)y_1 d_\alpha t, \tag{2.9}$$

for all $y_1, y_2 \in D$ and for all $b > 0$.

If we take $y_1 = y_2 \in D$ in (2.9), then we get

$$\int_0^b \{p(T_\alpha y_1)^2 + qy_1^2\} d_\alpha t = p(b)y_1(b)T_\alpha y_1(b) - \int_0^b l(y_1)y_1 d_\alpha t.$$

It follows from (2.2) and (2.6) that the integrand on the left is non-negative. From (2.3), the integrand on the right is $L_\alpha^2[0, \infty)$. If $p(T_\alpha y_1)^2 + qy_1^2 \notin L_\alpha^1[0, \infty)$, then $p(b)y_1(b)T_\alpha y_1(b) \rightarrow \infty$, as $b \rightarrow \infty$. This is impossible since then both $y_1(b)$ and $T_\alpha y_1(b)$ would have the same sign for all large b and y_1 could not then belong to $L_\alpha^2[0, \infty)$. This contradicts our assumption $y_1 \in D \subset L_\alpha^2[0, \infty)$. Thus, for all $y_1 \in D$, we conclude that

$$p^{\frac{1}{2}} T_\alpha y_1 \in L_\alpha^2[0, \infty), \quad q^{\frac{1}{2}} y_1 \in L_\alpha^2[0, \infty). \tag{2.10}$$

From (2.9), we deduce that, for all $y_1, y_2 \in D$,

$$\lim_{b \rightarrow \infty} p(b)y_1(b)T_{\alpha}y_2(b) \tag{2.11}$$

exists and is finite.

If $\Psi \in L^2_{\alpha}[0, \infty)$, then $\Psi \in L^1_{\alpha}[0, b]$. It follows from Theorem 1.9 that

$$\lim_{b \rightarrow \infty} \left(\frac{b}{\alpha}\right)^{-\frac{1}{2}\alpha} \int_0^b |\Psi(t)| d_{\alpha}t = 0 \tag{2.12}$$

For all $b \geq 0$, define

$$P(b) := \int_0^b \{p(t)\}^{-\frac{1}{2}} d_{\alpha}t.$$

By (2.4), we get

$$\lim_{b \rightarrow \infty} P(b) = K, \text{ where } 0 < K < \infty. \tag{2.13}$$

From (1.1), we conclude that

$$\begin{aligned} \int_0^b p(t)^{\frac{1}{2}} T_{\alpha}y_2(t) d_{\alpha}t &= P(b)p(b)T_{\alpha}y_2(b) \\ &+ \int_0^b P(t)l(y_2(t)) d_{\alpha}t - \int_0^b P(t)q(t)y_2(t) d_{\alpha}t. \end{aligned}$$

Multiply this result by $\{Q(b)\}^{-\frac{1}{2}}$ and consider the separate terms. it follows from (2.8), (2.10) and (2.12) that

$$\begin{aligned} \{Q(b)\}^{-\frac{1}{2}} \int_0^b p(t)^{\frac{1}{2}} T_{\alpha}y_2(t) d_{\alpha}t &= \\ O\left(\left(\frac{b}{\alpha}\right)^{-\frac{\alpha}{2}} \int_0^b p(t)^{\frac{1}{2}} |T_{\alpha}y_2(t)| d_{\alpha}t\right) &= o(1) \text{ as } b \rightarrow \infty. \end{aligned}$$

Similarly, using (2.8), (2.10) and (2.12) as $b \rightarrow \infty$, we get

$$\begin{aligned} \{Q(b)\}^{-\frac{1}{2}} \int_0^b P(t)l(y_2(t)) d_{\alpha}t &= \\ = O\left(K\left(\frac{b}{\alpha}\right)^{-\frac{\alpha}{2}} \int_0^b |l(y_2)| d_{\alpha}t\right) &= o(1). \end{aligned}$$

Let $b' > 0$ be fixed. Then, for all $b > b'$, we have

$$\begin{aligned} &Q(b)^{-\frac{1}{2}} \int_0^b P(t)q(t)y_2(t) d_{\alpha}t \\ &= Q(b)^{-\frac{1}{2}} \left\{ \int_0^{b'} + \int_{b'}^b \right\} (Pqy_2) d_{\alpha}t = O(b^{-\frac{1}{2}}) \\ &+ O\left(K \left\{ Q(b)^{-1} \int_{b'}^b q d_{\alpha}t \int_{b'}^b qy_2^2 d_{\alpha}t \right\}^{\frac{1}{2}}\right) \\ &= o(1) + O\left(K \left\{ \int_{b'}^b qy_2^2 d_{\alpha}t \right\}^{\frac{1}{2}}\right). \end{aligned}$$

It follows from (2.10) that the left-hand side tends to zero as $b \rightarrow \infty$.

Hence by (2.13), we get, for all $y_2 \in D$,

$$\lim_{b \rightarrow \infty} \{Q(b)\}^{-\frac{1}{2}} p(b) T_{\alpha} y_2(b) = 0. \quad (2.14)$$

Let us consider $\{Q\}^{\frac{1}{2}} y_1$ where $y_1 \in D$; assume that

$$\lim_{b \rightarrow \infty} \inf \{Q(b)\}^{\frac{1}{2}} |y_1(b)| > 0.$$

There is a constant S , $0 < S < \infty$, such that for all $t > b_0$ (say) we have $|y_1(t)|^2 \geq S^2 \{Q(t)\}^{-1}$. Multiply this inequality by the positive number $q(t)$ and integrate over $[b_0, b]$ to give, using (2.7),

$$\begin{aligned} \int_{b_0}^b q(t) |y_1(t)|^2 d_{\alpha} t &\geq S^2 \int_{b_0}^b \frac{q(t)}{Q(t)} d_{\alpha} t \\ &= S^2 \left[\ln Q\left(\frac{t^{\alpha}}{\alpha}\right) \right] \Big|_{b_0}^b \geq BS^2 \ln\left(k \frac{b^{\alpha}}{\alpha}\right), \end{aligned}$$

where B is a positive constant depending on b_0 and $k > 0$. This implies that $q^{\frac{1}{2}} y_1 \notin L_{\alpha}^2[0, \infty)$. This contradicts (2.10). Then there exists a sequence $\{b_i; i \geq 1\}$, such that $b_i \rightarrow \infty$ as $i \rightarrow \infty$, and for which

$$\lim_{i \rightarrow \infty} Q(b_i)^{\frac{1}{2}} y_1(b_i) = 0 \quad (2.15)$$

From (2.15) and (2.14), we conclude that there is a sequence $\{b_i; i \geq 1\}$ such that

$$\lim_{i \rightarrow \infty} p(b_i) y_1(b_i) T_{\alpha} y_2(b_i) = 0,$$

for every pair $y_1, y_2 \in D$. It now follows from (2.11) that (2.5) is satisfied for all $y_1, y_2 \in D$. The proof is complete. \square

CONFLICTS OF INTEREST

The authors declare that there are no conflicts of interest regarding the publication of this article.

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