Commun. Fac. Sci. Univ. Ank. Series A₁ V. 41. pp. 79-84 (1992)

THREE LORENTZIAN PLANES MOVING WITH RESPECT TO ONE ANOTHER AND POLE POINTS

ABDULLAH AZİZ ERGİN

Ground Military School, Ankara.

(Received May 4, 1992; Accepted May 29, 1992)

ABSTRACT

In this paper we have considered a Lorentzian plane A moving with respect to the planes L and L'. We have discussed here on the three Lorentzian planes A, L and L' instead of the two planes in [1]. Also, Lorentzian pole points and pole line of the mentioned planes are evaluated at a time "t".

I. INTRODUCTION

Let A and L be two moving Lorentzian planes with coordinate systems $\{B; \vec{a}_1, \vec{a}_2\}$ and $\{O; \vec{l}_1, \vec{l}_2\}$ respectively and L' be a fixed Lorentzian plane with coordinate system $\{O'; \vec{l}_1' \vec{l}_2\}$.

At first let's evaluate some resembling equations given in [1]. By using Figure. 1 we find

$$\vec{\mathbf{a}}_{1} = \vec{\mathbf{l}}_{1} \operatorname{ch} \varnothing + \vec{\mathbf{l}}_{2} \operatorname{sh} \varnothing$$

$$\vec{\mathbf{a}}_{2} = \vec{\mathbf{l}}_{1} \operatorname{sh} \varnothing + \vec{\mathbf{l}}_{2} \operatorname{ch} \varnothing$$

$$\vec{\mathbf{OB}} = \vec{\mathbf{b}} = \vec{\mathbf{a}}_{1} \mathbf{b}_{1} + \vec{\mathbf{a}}_{2} \mathbf{b}_{2}$$
(2)

And we obtain equations of LM motion A/L as

$$\overrightarrow{da}_{1} = \overrightarrow{a}_{2} d \varnothing$$

$$\overrightarrow{da}_{2} = \overrightarrow{a}_{1} d \varnothing$$

$$\overrightarrow{db} = \overrightarrow{a}_{1} (db_{1} + b_{2} d \varnothing) + \overrightarrow{a}_{2} (db_{2} + b_{1} d \varnothing)$$
(2)

Likewise, equations of LM motion A/L' are obtained as follows

$$\overrightarrow{d'a}_{1} = \overrightarrow{a}_{2} d \varnothing'
\overrightarrow{d'a}_{2} = \overrightarrow{a}_{1} d \varnothing'
\overrightarrow{d'b'} = \overrightarrow{a}_{1} (db_{1}' + b_{2}'d \varnothing') + \overrightarrow{a}_{2} (db_{2}' + b_{1}'d \varnothing')$$
(3)

For the sake of shortness let's use

$$d \varnothing = \tau$$
, $d \varnothing' = \tau'$, $db_1 + b_2 d \varnothing = \sigma_1$,

 $db_2 + b_1 \, d \, \varnothing = \sigma_2, \, db_1' + b_2' \, d \, \varnothing' = \sigma_1', \, db_2' + b_1' \, d \, \varnothing' = \sigma_2'.$ And $\sigma_j, \, \sigma_j', \, \tau$ and τ' will be called the Lorentzian Pfaffian forms of the 1-parameter Lorentzian motion, where $1 \leq j \leq 2$. So we obtain (2) and (3) again as follows, respectively,

$$\vec{da}_{1} = \vec{a}_{2} \tau$$

$$\vec{da}_{2} = \vec{a}_{1} \tau$$

$$\vec{db} = \vec{a}_{1} \sigma_{1} + \vec{a}_{2} \sigma_{2}$$

$$(4)$$

and

$$\vec{\mathbf{d}'} \vec{\mathbf{a}}_{1} = \vec{\mathbf{a}}_{2} \, \tau'
\vec{\mathbf{d}'} \vec{\mathbf{a}}_{2} = \vec{\mathbf{a}}_{1} \, \tau'
\vec{\mathbf{d}'} \vec{\mathbf{b}} = \vec{\mathbf{a}}_{1} \, \sigma_{1}' + \mathbf{a}_{2} \, \sigma_{2}'$$
(5)

Now let us consider a point $X = (x_1, x_2)$ on the moving plane A. Then the following vectors are obtained.

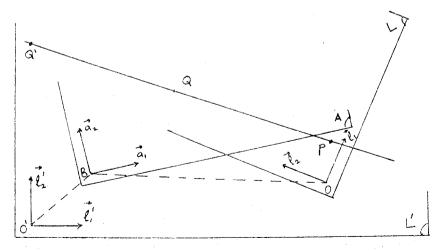


Figure.1.

$$\overrightarrow{BK} = \overrightarrow{a}_1 x_1 + \overrightarrow{a}_2 x_2$$

$$\overrightarrow{x} = \overrightarrow{OX} = \overrightarrow{OB} + \overrightarrow{BK} = \overrightarrow{b} + \overrightarrow{a}_1 x_1 + \overrightarrow{a}_2 x_2$$

$$\overrightarrow{x'} = \overrightarrow{O'X} = \overrightarrow{O'B} + \overrightarrow{BX} = \overrightarrow{b'} + \overrightarrow{a}_1 x_1 + \overrightarrow{a}_2 x_2$$

From (4) we find differentiation of X in L as

$$\vec{dx} = \vec{a}_1 (dx_1 + \sigma_1 + x_2 \tau) + \vec{a}_2 (dx_2 + \sigma_2 + x_1 \tau)$$
 (6)

And the relative velocity vector is $\vec{V}_r = \frac{\overrightarrow{dx}}{dt}$. If $\vec{V}_r = 0$, then X

is a fixed point on L. That is, we can find the conditions that X is a fixed point on L as follows:

and so

$$dx_1 = -(\sigma_1 + x_2 \tau)$$

$$dx_2 = -(\sigma_2 + x_1 \tau)$$

Similarly, we find the conditions of X as a constant point on the fixed Lorentzian plane L' as follows:

$$\vec{d'x'} = \vec{a}_1 (dx_1 + \sigma_1' + x_2 \tau') + \vec{a}_2 (dx_2 + \sigma_2' + x_1 \tau') \quad .. \quad (7)$$

and $\vec{V}_a = \frac{\overrightarrow{d'x'}}{dt}$ is the absolute velocity vector. If X is a fixed point

on L', then $\overset{\rightarrow}{V_a}=\vec{o}.$ So $\overset{\rightarrow}{d'x'}=\vec{o}$ and we obtain

$$\begin{array}{l} dx_1 = - \left(\sigma_1 + x_2 \; \tau' \right) \\ dx_2 = - \left(\sigma_2 + x_1 \; \tau' \right) \end{array} \right\}$$

Hence we can give the following theorem.

I.1. Theorem: Let A and L be moving and L' be a fixed Lorentzian planes and X be a point on A. If X is a fixed point on L, then

$$d\mathbf{x}_1 = -(\sigma_1 + \mathbf{x}_2 \tau)$$

$$d\mathbf{x}_2 = -(\sigma_2 + \mathbf{x}_1 \tau)$$
(8)

and if X is a fixed point on L,, then

$$\frac{d\mathbf{x}_1 = -(\sigma_1' + \mathbf{x}_2 \, \tau')}{d\mathbf{x}_2 = -(\sigma_2' + \mathbf{x}_1 \, \tau')}$$
(9)

Now let X be a fixed point on L' then the sliding velocity vector

of X according to L' is $\overrightarrow{V}_f = \frac{\overrightarrow{d_f x}}{dt}$, which corresponds to the differential $\overrightarrow{d_f x}$ of X with respect to L'. Therefore by using (8), (7) changes into

$$\vec{\mathbf{d}_f} \mathbf{x} = \vec{\mathbf{a}}_1 \left(-\sigma_1 - \mathbf{x}_2 \ \tau + \sigma'_1 + \mathbf{x}_2 \ \tau' \right) + \vec{\mathbf{a}}_2 \left(-\sigma_2 - \mathbf{x}_1 \ \tau + \sigma_2' + \mathbf{x}_1 \ \tau' \right)$$
and so

$$\overrightarrow{\mathbf{d}_{\mathbf{f}}\mathbf{x}} = \overrightarrow{\mathbf{a}}_1 \left\{ (\sigma_1' - \sigma_1) - \mathbf{x}_2 (\tau - \tau') \right\} + \overrightarrow{\mathbf{a}}_2 \left\{ (\sigma'_2 - \sigma_2) - \mathbf{x}_1 (\tau - \tau') \right\} \dots (10)$$

I.2. Theorem: If X is a fixed point on L, then

$$\overrightarrow{d'x'} = \overrightarrow{d_fx} + \overrightarrow{dx}.$$

II. MOVING PLANES WITH RESPECT TO ONE ANOTHER AND ROTATION POLES

In the 1-parameter Lorentzian motion the rotation pole is characterized by vanishing sliding velocity. That is, at the pole point the relative velocity equals to the absolute velocity at a t time. And so, if we take $d_f \mathbf{x} = 0$ and use (10), the pole point $P = (p_1, p_2)$ of the 1-parameter Lorentzian motion L/L' is obtained as

$$p_{1} = x_{1} = \frac{\sigma_{2}' - \sigma_{2}}{\tau - \tau'}$$

$$p_{2} = x_{2} = \frac{\sigma_{1}' - \sigma_{1}}{\tau - \tau'}$$

$$(11)$$

where $\overrightarrow{BP} = \vec{a}_1 p_1 + \vec{a}_2 p_2$.

Now let us consider 1-parameter Lorentzian motions A/L and A/L'. Each of them has an exact pole point at a "t" time.

The Lorentzian A/L motion was known by (4), where $\tau = d \varnothing$

was the infinitesimal rotation angle and $\frac{\tau}{dt}$ was the angular velocity.

Differentiating N on A we find the vector

$$\overrightarrow{dBX} = \overrightarrow{a}_1 \, dx_1 + \overrightarrow{a}_2 \, dx_2$$

which corresponds to the relative velocity vector of X on A. We have given by (6) the differential of X with respect to L. If we write (6) in the following form

 $\overrightarrow{dx} = [\overrightarrow{a}_1 \ dx_1 + \overrightarrow{a}_2 \ dx_2] + [\overrightarrow{a}_1 \ (\sigma_1 + x_2 \ \tau) + \overrightarrow{a}_2 \ (\sigma_2 + x_1 \ \tau)].$ The first paranthesis corresponds to the differential of the relative velocity vector of X on A, while the second paranthesis corresponds to the sliding velocity vector of X in A/L.

The rotation pole $Q=(q_1,\,q_2),$ at any t time, of the Lorentzian motion $A/_L$ is obtained by the vanishing of the sliding velocity. So

$$\left. \begin{array}{l} \sigma_1 + x_2 \, \tau = 0 \\ \sigma_2 + x_1 \, \tau = 0 \end{array} \right\}$$

and from that we find

$$x_1 = q_1 = -\frac{\sigma_2}{\tau}$$

$$x_2 = q_2 = -\frac{\sigma_1}{\tau}$$
(12)

Similarly, the pole point $Q'=(q_{1}{}',\,q_{2}{}')$ of the 1-parameter Lorentzian motion $A\,/\,_{L},$ is found as

$$q_{1}' = -\frac{\sigma_{2}'}{\tau'}$$

$$q_{2}' = -\frac{\sigma_{1}'}{\tau'}$$
(13)

II.1. Theorem: Let P, Q and Q' be the pole points of the Lorentzian motions $L/_{L'}$, $A/_{L}$ and $A/_{L'}$ respectively, then P, Q and Q' are collinear.

Proof: The slopes of [PQ], [PQ'] and [QQ'] are all equal to

$$\frac{\sigma_1 \ \tau' - \sigma_1' \ \tau}{\sigma_2 \ \tau' - \sigma_2' \ \tau}.$$

That completes the proof of the theorem.

II.1. Definition: The straight line, indicated by the points P, Q and Q' is called the Lorentzian pole line of the Lorentzian motions $L/_{L'}$, $A/_{L}$ and $A/_{L'}$.

ÖZET

Bu çalışmada L ve L' Lorentz düzlemlerine göre hareket eden bir A Lorentz düzlemi göz önüne alındı. [1] deki iki Lorentz düzlemine karşılık burada A, L ve L' gibi üç Lorentz düzlemi üzerinde çalışıldı. Ayrıca bahsedilen düzlemlerin bir "t" anındaki Lorentz anlamında pole noktaları ve pol doğrusu bulundu.

REFERENCES

- ERGÍN, A.A., "On The 1-Parameter Lorentzian Motion", Commun. Fac. Sci. Univ. Ank. Series A₁, V. 40, (1991).
- [2] HACISALİHOĞLU, H.H., "On the Geometry of Motion in The Euclidean n-Space", Commun. Fac. Sci. Univ. Ank. Series A, V. 23 pp. 95-108, (1974).
- [3] MÜLLER, H.R., "Kinematik Dersleri", Ank. Univ. Yayınları No: 27, (1963).
- [4] O'NEILL, B., "Semi-Riemannian Geometry", Academic Press, New York, London (1983).