# Note on the convergence of fractional conformable diffusion equation with linear source term 

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#### Abstract

In this paper, we study the diffusion equation with conformable derivative. The main goal is to prove the convergence of the mild solution to our problem when the order of fractional Laplacian tends to $1^{-}$. The principal techniques of our paper is based on some useful evaluations for exponential kernels.


Keywords: Conformable derivative, diffusion equation, convergent estimate.
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## 1. Introduction

In this paper, we are interested to study the following problem

$$
\left\{\begin{array}{l}
\frac{C \partial^{\alpha}}{\partial t^{\alpha}} v+(-\Delta)^{s} v(x, t)=F(x, t), \quad(x, t) \in \Omega \times(0, T),  \tag{1.1}\\
v(x, t)=0, \quad x \in \partial \Omega, \quad t \in(0, T),
\end{array}\right.
$$

with the initial condition

$$
\begin{equation*}
v(x, 0)=v_{0}(x), \quad x \in \Omega, \tag{1.2}
\end{equation*}
$$

where $v_{0}$ and $F$ are the input data. The symbol $(-\Delta)^{s}, s>0$ says that the fractional Laplacian which is defined later. Here $\Omega \subset \mathbb{R}^{N}(N \geq 1)$ is a bounded domain with the smooth boundary $\partial \Omega$, and $T>0$ is a given positive number. The above equation has various applications in areas such as the harmonic oscillator, the damped oscillator and the forced

[^0]oscillator (see [3]), electrical circuits (see [5]), chaotic systems in dynamics (see [6]), projectile motion (see [7]). Our paper is one of the braches of directions about fractional PDEs, see [15, 16, 17, [18, 20].

The symbol $\frac{C \partial^{\alpha}}{\partial t^{\alpha}}$ is understood as the conformable derivative. Let us now give a clear definition of conformable derivative on the Banach space. Let us given $B$ is a Banach space, and $f:[0, \infty) \rightarrow B$. Let $\frac{c_{\partial \partial^{\alpha}}}{\partial t^{\alpha}}$ be the conformable derivative of order $0<\alpha \leq 1$ which is given by

$$
\frac{c_{\partial^{\alpha}} f(t)}{\partial t^{\alpha}}:=\lim _{\epsilon \rightarrow 0} \frac{f\left(t+\epsilon t^{1-\alpha}\right)-f(t)}{\epsilon} \quad \text { in } B,
$$

for each $t>0$. Some more details on conformable derivative can be found in [1, 2, 8, 14, ,9, 18]. It is easy to see that $\alpha=s=1$, Problem (1.1)-(1.2) becomes the classical heat equation.

We mention now some previous results for conformable derivative. In [12], an inverse problem for second boundary with conformable diffusion is shown. In [13], Bayrak and his colleagues investigated approximate solution of the time-fractional Fisher equation with small delay. In [10], the authors studied some nonlinear partial differential equations with conformable derivative. In [14], the authors focused on a mild solution of conformable fractional abstract initial value problem.

The well-posedness of Problem (1.1)-(1.2) was established in [11]. Indeed, the paper [11] derived more clearly the existence and the regularity of the mild solution of Problem (1.1)-(1.2). One of the highlights of our problem is the occurrence the fractional Laplacian $(-\Delta)^{s}$ for any $0<s \leq 1$. Our main goal in this paper is to study the limit problem of the mild solution when $s \rightarrow 1^{-}$. Up to now, there has not been any literature surveying the mentioned issue. Our paper is the first result concerned with the limit problem for the fractional diffusion equation with conformable derivatives. In order to overcome some complicated evaluations, we use some new techniques for the computations for exponential functions.

## 2. Initial value problem

### 2.1. Premilinaries

Let us recall that the spectral problem

$$
\begin{cases}(-\Delta)^{s} e_{n}(x)=\lambda_{n}^{s} e_{n}(x), & x \in \mathcal{D}, \\ e_{n}(x)=0, & x \in \partial D,\end{cases}
$$

admits the eigenvalues $0<\lambda_{1} \leq \lambda_{2} \leq \cdots \leq \lambda_{n} \leq \ldots$ with $\lambda_{n} \rightarrow \infty$ as $n \rightarrow \infty$. The corresponding eigenfunctions are $e_{n} \in H_{0}^{1}(\Omega)$.

Next, Let a given positive number $\sigma \geq 0$. Let us also define the Hilbert scale space as follows

$$
\begin{equation*}
\mathbb{H}^{\sigma}(\Omega)=\left\{\psi \in L^{2}(\Omega): \sum_{n=1}^{\infty} \lambda_{n}^{2 \sigma}\left\langle\psi, e_{n}\right\rangle^{2}<+\infty\right\}, \tag{2.1}
\end{equation*}
$$

with the following norm $\|\psi\|_{\mathbb{H}^{\sigma}(\Omega)}=\left(\sum_{n=1}^{\infty} \lambda_{n}^{2 \sigma}\left\langle\psi, e_{n}\right\rangle^{2}\right)^{\frac{1}{2}}$.

### 2.2. The linear case

In this section, we focus the following initial value problem (1.1) under the linear case with the initial condition (1.2). Here $v_{0}$ and source function $F$ are defined later.

In order to find a precise formulation for solutions, we consider the spectral decomposition

$$
v_{s}(x, t)=\sum_{n=1}^{\infty}\left(\int_{\Omega} v_{s}(x, t) e_{n}(x) d x\right) e_{n}(x) .
$$

Thanks for the work [11], we get the following equality

$$
\begin{align*}
\int_{\Omega} v_{s}(x, t) e_{n}(x) d x & =\exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}^{s}\right) \int_{\Omega} v_{0}(x) e_{n}(x) d x \\
& +\int_{0}^{t} \theta^{\alpha-1} \exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}\right)\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right) d \theta \tag{2.2}
\end{align*}
$$

Then the mild solution to problem (1.1)-(1.2) is defined by

$$
\begin{align*}
v_{s}(x, t) & =\sum_{n=1}^{\infty} \exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}^{s}\right)\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right) e_{n}(x) \\
& +\sum_{n=1}^{\infty}\left(\int_{0}^{t} \theta^{\alpha-1} \exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}\right)\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right) d \theta\right) e_{n}(x) \tag{2.3}
\end{align*}
$$

Here $0<s<1$.
Lemma 1. Let $0<s<1$. Then we have the following inequality

$$
\begin{equation*}
\left|\lambda_{n}^{s}-\lambda_{n}\right| \leq C(m) \lambda_{n}^{s-m}(1-s)^{m}, \quad \lambda_{n} \leq 1 \tag{2.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\lambda_{n}^{s}-\lambda_{n}\right| \leq C(m) \lambda_{n}^{1-m}(1-s)^{m}, \quad \lambda_{n}>1 \tag{2.5}
\end{equation*}
$$

Proof. If $\lambda_{n} \leq 1$ then we get that the following bound

$$
\begin{align*}
\left|\lambda_{n}^{s}-\lambda_{n}\right|=\lambda_{n}^{s}-\lambda_{n} & =\lambda_{n}^{s}\left(1-\exp \left(-(1-s) \log \left(\frac{1}{\lambda_{n}}\right)\right)\right) \\
& \leq C(m) \lambda_{n}^{s}(1-s)^{m} \log ^{m}\left(\frac{1}{\lambda_{n}}\right) \leq C(m) \lambda_{n}^{s-m}(1-s)^{m} \tag{2.6}
\end{align*}
$$

where we have used the inequality $1-e^{-z} \leq C(m) z^{m}$.
If $\lambda_{n} \geq 1$ then we get that the following bound

$$
\begin{align*}
\left|\lambda_{n}^{s}-\lambda_{n}\right|=\lambda_{n}-\lambda_{n}^{s} & =\lambda_{n}\left(1-\exp \left(-(1-s) \log \left(\frac{1}{\lambda_{n}}\right)\right)\right) \\
& \leq C(m) \lambda_{n}(1-s)^{m} \log ^{m}\left(\frac{1}{\lambda_{n}}\right) \leq C(m) \lambda_{n}^{1-m}(1-s)^{m} \tag{2.7}
\end{align*}
$$

Theorem 1. Let $v_{0} \in \mathbb{H}^{1-m+p}(\Omega)$ and $F \in L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)$ for any $p>0$. Then we get

$$
\begin{equation*}
\left\|v_{s}-v^{* *}\right\|_{L^{\infty}\left(0, T ; \mathbb{H}^{p}(\Omega)\right)} \leq C(m, T, \alpha)(1-s)^{m}\left(\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)}+\|F\|_{L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)}\right) \tag{2.8}
\end{equation*}
$$

for any $p \geq 0$.
Proof. The mild solution of Problem (1.1)-(1.2) with $s=1$ is given by

$$
\begin{align*}
v^{* *}(x, t) & =\sum_{n=1}^{\infty} \exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}\right)\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right) e_{n}(x) \\
& +\sum_{n=1}^{\infty}\left(\int_{0}^{t} \theta^{\alpha-1} \exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}\right)\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right) d \theta\right) e_{n}(x) \tag{2.9}
\end{align*}
$$

By subtracting both sides of the two expressions above, we get the following difference

$$
\begin{align*}
& v_{s}(x, t)-v^{* *}(x, t) \\
& =\sum_{n=1}^{\infty}\left[\exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}^{s}\right)-\exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}\right)\right]\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right) e_{n}(x) \\
& +\sum_{n=1}^{\infty}\left(\int_{0}^{t} \theta^{\alpha-1}\left[\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}-\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}\right)\right]\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right) d \theta\right) e_{n}(x)\right. \\
& =B_{1}(x, t)+B_{2}(x, t) \tag{2.10}
\end{align*}
$$

Let us first treat the quantity $B_{1}$. Using Parseval's equality and the inequality $\left|e^{-a}-e^{-b}\right| \leq|a-b|$ for any $a, b>0$, we find that

$$
\begin{align*}
\left\|B_{1}(., t)\right\|_{\mathbb{H} p(\Omega)}^{2} & =\sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left[\exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}^{s}\right)-\exp \left(-\alpha^{-1} t^{\alpha} \lambda_{n}\right)\right]^{2}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} \\
& \leq \alpha^{-2} t^{2 \alpha} \sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} \\
& =\alpha^{-2} t^{2 \alpha} \sum_{\lambda_{n} \leq 1} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} \\
& +\alpha^{-2} t^{2 \alpha} \sum_{\lambda_{n}>1} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} . \tag{2.11}
\end{align*}
$$

In view of Lemma 1, we derive that

$$
\begin{align*}
\left\|B_{1}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}^{2} & \leq C(m) \alpha^{-2} t^{2 \alpha}(1-s)^{2 m} \sum_{\lambda_{n} \leq 1} \lambda_{n}^{2 s-2 m+2 p}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} \\
& +C(m) \alpha^{-2} t^{2 \alpha}(1-s)^{2 m} \sum_{\lambda_{n}>1} \lambda_{n}^{2-2 m+2 p}\left(\int_{\Omega} v_{0}(x) e_{n}(x) d x\right)^{2} \\
& \leq C(m) \alpha^{-2} t^{2 \alpha}(1-s)^{2 m}\left(\left\|v_{0}\right\|_{\mathbb{H}^{s-m+p}(\Omega)}^{2}+\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)}^{2}\right) \tag{2.12}
\end{align*}
$$

Since the fact that

$$
\left\|v_{0}\right\|_{\mathbb{H}^{s-m+p}(\Omega)} \leq C(s, m, p)\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)}
$$

we know that the following estimate

$$
\begin{equation*}
\left\|B_{1}(., t)\right\|_{\mathbb{H}^{p}(\Omega)} \leq C(m, s, p) \alpha^{-1} t^{\alpha}(1-s)^{m}\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)} \tag{2.13}
\end{equation*}
$$

Let us to study the second term $\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}$. Indeed, we get that

$$
\begin{align*}
& \left\|B_{2}(., t)\right\|_{\mathbb{H} p}^{2}(\Omega) \\
& =\sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left(\int_{0}^{t} \theta^{\alpha-1}\left[\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}-\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}\right)\right]\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right) d \theta\right)^{2}\right. \\
& \leq \sum_{n=1}^{\infty} \lambda_{n}^{2 p} \int_{0}^{t} \theta^{\alpha-1}\left[\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}-\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}\right)\right]^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} d \theta\right. \tag{2.14}
\end{align*}
$$

By applying the inequality $\left|e^{-a}-e^{-b}\right| \leq|a-b|, a, b \geq 0$ we derive that

$$
\begin{equation*}
\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}^{s}-\exp \left(-\alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right) \lambda_{n}\right) \leq \alpha^{-1}\left(t^{\alpha}-\theta^{\alpha}\right)\left|\lambda_{n}^{s}-\lambda_{n}\right|\right. \tag{2.15}
\end{equation*}
$$

Hence, we get that

$$
\begin{equation*}
\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}^{2} \leq \frac{\alpha^{-2} T^{\alpha}}{\alpha}\left(\int_{0}^{t} \theta^{\alpha-1}\left(t^{\alpha}-\theta^{\alpha}\right)^{2} \sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} d \theta\right) \tag{2.16}
\end{equation*}
$$

By looking at Lemma 1, we know that

$$
\begin{align*}
& \sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} \\
& =\sum_{\lambda_{n} \leq 1} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} \\
& +\sum_{\lambda_{n}>1} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} \\
& \leq C(m)(1-s)^{2 m} \sum_{\lambda_{n} \leq 1} \lambda_{n}^{2 p+2 s-2 m}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} \\
& +C(m)(1-s)^{2 m} \sum_{\lambda_{n}>1} \lambda_{n}^{2 p+2-2 m}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} . \tag{2.17}
\end{align*}
$$

Since the fact that $s<1$, we follows from the latter estimate that

$$
\begin{equation*}
\sum_{n=1}^{\infty} \lambda_{n}^{2 p}\left|\lambda_{n}^{s}-\lambda_{n}\right|^{2}\left(\int_{\Omega} F(x, \theta) e_{n}(x) d x\right)^{2} \leq C(m)(1-s)^{2 m}\|F(., \theta)\|_{\mathbb{H}^{1+p-m}(\Omega)}^{2} \tag{2.18}
\end{equation*}
$$

Combining 2.16, 2.18), we obtain that the following bound

$$
\begin{align*}
\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}^{2} & \leq C(m, T, \alpha)(1-s)^{2 m} \int_{0}^{t} \theta^{\alpha-1}\left(t^{\alpha}-\theta^{\alpha}\right)^{2}\|F(., \theta)\|_{\mathbb{H}^{1+p-m}(\Omega)}^{2} d \theta \\
& \leq C(m, T, \alpha)(1-s)^{2 m}\|F\|_{L^{\infty}\left(0, T ; ; \mathbb{H}^{1+p-m}(\Omega)\right)}^{2}\left(\int_{0}^{t} \theta^{\alpha-1}\left(t^{\alpha}-\theta^{\alpha}\right)^{2} d \theta\right) \tag{2.19}
\end{align*}
$$

Set $z=\theta^{\alpha}$ then $d z=\alpha \theta^{\alpha-1} d \theta$. Then we obtain that

$$
\begin{equation*}
\int_{0}^{t} \theta^{\alpha-1}\left(t^{\alpha}-\theta^{\alpha}\right)^{2} d \theta=\int_{0}^{t^{\alpha}}\left(t^{\alpha}-z\right)^{2} d z=\frac{t^{3 \alpha}}{3} \tag{2.20}
\end{equation*}
$$

This implies that

$$
\begin{equation*}
\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}^{2} \leq C(m, T, \alpha)(1-s)^{2 m}\|F\|_{L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)}^{2} \tag{2.21}
\end{equation*}
$$

which implies that

$$
\begin{equation*}
\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)} \leq C(m, T, \alpha)(1-s)^{m}\|F\|_{L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)} \tag{2.22}
\end{equation*}
$$

Combining (2.13) and 2.22), we derive that

$$
\begin{align*}
\left\|v_{s}(., t)-v^{* *}(., t)\right\|_{\mathbb{H}^{p}(\Omega)} & \leq\left\|B_{1}(., t)\right\|_{\mathbb{H}^{p}(\Omega)}+\left\|B_{2}(., t)\right\|_{\mathbb{H}^{p}(\Omega)} \\
& \leq C(m, T, \alpha)(1-s)^{m}\left(\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)}+\|F\|_{L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)}\right) \tag{2.23}
\end{align*}
$$

Since the right-hand side of 2.23 is independent of $t$, we derive that

$$
\begin{equation*}
\left\|v_{s}-v^{* *}\right\|_{L^{\infty}\left(0, T ; \mathbb{H}^{p}(\Omega)\right)} \leq C(m, T, \alpha)(1-s)^{m}\left(\left\|v_{0}\right\|_{\mathbb{H}^{1-m+p}(\Omega)}+\|F\|_{L^{\infty}\left(0, T ; \mathbb{H}^{1+p-m}(\Omega)\right)}\right) \tag{2.24}
\end{equation*}
$$

## 3. Conclusion

In this work, the diffusion equation with conformable derivative, and fractional Laplacian tends to $1^{-}$. We proved the convergence of the mild solution, while the principle techniques is exponential kernels.

## References

[1] R. Khalil, M. Al Horani, A. Yousef, M. Sababheh, A new definition of fractional derivative, J. Comput. Appl. Math. 264 (2014), 65-70.
[2] A.A. Abdelhakim, J.A. T. Machado, A critical analysis of the conformable derivative, Nonlinear Dynamics, Volume 95, Issue 4, (2019), 3063-3073.
[3] W.S. Chung, Fractional Newton mechanics with conformable fractional derivative, Journal of Computational and Applied Mathematics, Volume 290 (2015), Pages 150-158.
[4] A. Jaiswal, D. Bahuguna, Semilinear Conformable Fractional Differential Equations in Banach Spaces, Differ. Equ. Dyn. Syst. 27 , no. 1-3, (2019), 313-325.
[5] V.F. Morales-Delgado, J.F. Gómez-Aguilar, R.F. Escobar-Jiménez, M.A. Taneco-Hernández, Fractional conformable derivatives of Liouville-Caputo type with low-fractionality, Physica A: Statistical Mechanics and its Applications, Volume 503 (2018), 424-438.
[6] S. He, K. Sun, X. Mei, B. Yan, S. Xu, Numerical analysis of a fractional-order chaotic system based on conformable fractional-order derivative, Eur. Phys. J. Plus, (2017) 132: 36. https://doi.org/10.1140/epjp/i2017-11306-3.
[7] F.M. Alharbi, D. Baleanu, A. Ebaid, Physical properties of the projectile motion using the conformable derivative, Chinese Journal of Physics, Volume 58, (2019), Pages 18-28.
[8] T. Abdeljawad, On conformable fractional calculus, J. Comput. Appl. Math., 279 (2015), 57-66
[9] A. Atangana, D. Baleanu, A. Alsaedi, New properties of conformable derivative, Open Math., 13 (2015), 889-898
[10] Y. Çenesiz, D. Baleanu, A. Kurt, O. Tasbozan, New exact solutions of Burgers' type equations with conformable derivative, Waves Random Complex Media, 27 (2017), no. 1, 103-116.
[11] N.H. Tuan, T.B. Ngoc, D. Baleanu, D. O'Regan, On well-posedness of the sub-diffusion equation with conformable derivative model, Communications in Nonlinear Science and Numerical Simulation Volume 89, October 2020, 105332.
[12] Y. Çakmak, Inverse nodal problem for a conformable fractional diffusion operator, Inverse Probl. Sci. Eng. 29 (2021), no. 9, $1308-1322$.
[13] A.M. Bayrak, A. Demir, E. Ozbilge, On the numerical solution of conformable fractional diffusion problem with small delay, Numer. Methods Partial Differential Equations 38 (2022), no. 2, 177-189.
[14] A. Jaiswal, D. Bahuguna, Semilinear Conformable Fractional Differential Equations in Banach Spaces, Differ. Equ. Dyn. Syst. 27 (2019), no. 1-3, 313-325.
[15] E. Karapinar, A.Fulga,M. Rashid, L.Shahid, H. Aydi, Large Contractions on Quasi-Metrics Spaces with a Application to Nonlinear Fractional Differential-Equations, Mathematics 2019, 7, 444.
[16] E.Karapinar, Ho Duy Binh, Nguyen Hoang Luc, and Nguyen Huu Can, On continuity of the fractional derivative of the time-fractional semilinear pseudo-parabolic systems, Advances in Difference Equations (2021) 2021:70.
[17] J. E. Lazreg, S. Abbas, M. Benchohra, and E. Karapinar, Impulsive Caputo-Fabrizio fractional differential equations in b-metric spaces , Open Mathematics 2021; 19: 363-372.
[18] J.E. Lazreg, S. Abbas, M. Benchohra, and E.Karapinar, Impulsive Caputo-Fabrizio fractional differential equations in b-metric spaces, Open Mathematics 2021; 19: 363-372.
[19] N. D. Phuong, Note on a Allen-Cahn equation with Caputo-Fabrizio derivative, Results in Nonlinear Analysis 4 (2021), 179-185.
[20] N. D. Phuong, N. H. Luc and L. D. Long, Modified Quasi Boundary Value method for inverse source problem of the bi-parabolic equation , Advances in the Theory of Nonlinear Analysis and its Applications 4 (2020), 132-142.


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