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# THE FLOW-CURVATURE OF PLANE PARAMETRIZED CURVES

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ABSTRACT. We introduce and study a new frame and a new curvature function for a fixed parametrization of a plane curve. This new frame is called *flow* since it involves the time-dependent rotation of the usual Frenet flow; the angle of rotation is exactly the current parameter. The flow-curvature is calculated for several examples obtaining the logarithmic spirals (and the circle as limit case) and the Grim Reaper as flat-flow curves. A main result is that the scaling with  $\frac{1}{\sqrt{2}}$  of both Frenet and flow-frame belong to the same fiber of the Hopf bundle. Moreover, the flow-Fermi-Walker derivative is defined and studied.

# 1. INTRODUCTION

The theory of geometric flows is a new and fascinating field of research in geometric analysis. The most simple of them is *the curve shortening flow* and already the excellent survey [3] is twenty years old. Recall that the main geometric tool in this last flow is the well-known curvature of plane curves. Hence, to give a re-start to this problem seams to search for variants of the curvature, or in terms of [9], *deformations* of the usual curvature. The goal of this short note is to propose such a deformation which in turn defines a Fermi-Walker type derivative.

Fix an open interval  $I \subseteq \mathbb{R}$  and consider  $C \subset \mathbb{R}^2$  a regular parametrized curve of equation:

$$C: r(t) = (x(t), y(t)) = x(t)\overline{i} + y(t)\overline{j}, \quad ||r'(t)|| > 0, \quad t \in I.'$$

The ambient setting  $\mathbb{R}^2$  is an Euclidean vector space with respect to the canonical inner product:

$$\langle u, v \rangle = x^1 y^1 + x^2 y^2, u = (x^1, x^2) \in \mathbb{R}^2, v = (y^1, y^2) \in \mathbb{R}^2, \quad 0 \le ||u||^2 = \langle u, u \rangle.$$

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The infinitesimal generator of the rotations in  $\mathbb{R}^2 = \mathbb{C}$  is the linear vector field, called *angular*:

$$\xi(u) := -x^2 \frac{\partial}{\partial x^1} + x^1 \frac{\partial}{\partial x^2}, \quad \xi(u) = i \cdot u = i \cdot (x^1 + ix^2), \quad i = \sqrt{-1}.$$

It is a complete vector field with integral curves the circles  $\mathcal{C}(O, r)$ :

$$\begin{cases} \gamma_{u_0}^{\xi}(t) = (u_0^1 \cos t - u_0^2 \sin t, u_0^1 \sin t + u_0^2 \cos t) = R(t) \cdot \begin{pmatrix} u_0^1 \\ u_0^2 \end{pmatrix}, & t \in \mathbb{R}, \\ r = \|u_0\| = \|(u_0^1, u_0^2)\|, & R(t) := \begin{pmatrix} \cos t & -\sin t \\ \sin t & \cos t \end{pmatrix} \in SO(2) = S^1 \end{cases}$$

and since the rotations R(t) are isometries of the Riemannian metric  $g_{can} = dx^2 + dy^2 = |dz|^2$  it follows that  $\xi$  is a Killing vector field of the Riemannian manifold  $(\mathbb{R}^2, g_{can})$ . The first integrals of  $\xi$  are the Gaussian functions i.e. multiples of the square norm:  $f_{\alpha}(x, y) = \alpha(x^2 + y^2), \ \alpha \in \mathbb{R}$ . For an arbitrary vector field  $X = A(x, y)\frac{\partial}{\partial x} + B(x, y)\frac{\partial}{\partial y}$  its Lie bracket with  $\xi$  is:

$$[X,\xi] = (yA_x - xA_y - B)\frac{\partial}{\partial x} + (A + yB_x - xB_y)\frac{\partial}{\partial y},$$

where the subscript denotes the variable corresponding to the partial derivative. For example,  $\xi$  commutes with the radial (or Euler) vector field  $E(x, y) = x \frac{\partial}{\partial x} + y \frac{\partial}{\partial y}$ , which is also a complete vector field having as integral curves the homotheties  $\gamma_{u_0}^E(t) = e^t u_0$  for all  $t \in \mathbb{R}$ .

The Frenet apparatus of the curve C is provided by:

$$\begin{cases} T(t) = \frac{r'(t)}{\|r'(t)\|}, \quad N(t) = i \cdot T(t) = \frac{1}{\|r'(t)\|} (-y'(t), x'(t)), \\ k(t) = \frac{1}{\|r'(t)\|} \langle T'(t), N(t) \rangle = \frac{1}{\|r'(t)\|^3} \langle r''(t), ir'(t) \rangle = \frac{1}{\|r'(t)\|^3} [x'(t)y''(t) - y'(t)x''(t)] \end{cases}$$

Hence, if C is naturally parametrized (or parametrized by arc-length) i.e. ||r'(s)|| = 1 for all  $s \in I$  then r''(s) = k(s)ir'(s). In a complex approach based on  $z(t) = x(t) + iy(t) \in \mathbb{C} = \mathbb{R}^2$  we have:

$$\begin{cases} k(t) = \frac{1}{|z'(t)|^3} Im(\bar{z}'(t) \cdot z''(t)) = \frac{1}{|z'(t)|} Im\left(\frac{z''(t)}{z'(t)}\right) = \frac{1}{|z'(t)|} Im\left[\frac{d}{dt}\left(\ln z'(t)\right)\right] \in \mathbb{R},\\ Re(\bar{z}'(t) \cdot z''(t)) = \frac{1}{2} \frac{d}{dt} \|r'(t)\|^2, \quad f_{\alpha}(z) = \alpha |z|^2. \end{cases}$$

The multiplication with the complex unit *i* corresponds to the rotation  $R\left(\frac{\pi}{2}\right)$ ; we have also:

$$\frac{d}{dt}R(t) = R\left(t + \frac{\pi}{2}\right) = R(t)R\left(\frac{\pi}{2}\right) = R\left(\frac{\pi}{2}\right)R(t),$$

and the Frenet equations can be unified by means of the column matrix  $\mathcal{F}(t) = \begin{pmatrix} T \\ N \end{pmatrix}(t)$  as:

$$\frac{d}{dt}\mathcal{F}(t) = \|r'(t)\|k(t)R\left(-\frac{\pi}{2}\right)\mathcal{F}(t).$$

It is an amazing fact that if the general rotation R(t) belongs to the Lie group  $SO(2) = S^1$  its particular values  $R(\pm \frac{\pi}{2})$  are elements of its Lie algebra so(2) of skew-symmetric  $2 \times 2$  matrices. In fact,  $\{R(\frac{\pi}{2})\}$  is exactly the basis of so(2).

### 2. Main Results

This short note defines a new frame and correspondingly a new curvature function for C:

**Definition 1.** The flow-frame of C consists in the pair of unit vectors  $(E_1^f(t), E_2^f(t)) \in T^2 := S^1 \times S^1$  given by:

$$\mathcal{E}(t) := \begin{pmatrix} E_1^f \\ E_2^f \end{pmatrix}(t) = R(t)\mathcal{F}(t) = \begin{pmatrix} \cos tT(t) - \sin tN(t) \\ \sin tT(t) + \cos tN(t) \end{pmatrix}$$
(1)

the letter f being the initial of the word "flow". The flow-curvature of C is the smooth function  $k_f: I \to \mathbb{R}$  given by the flow-equations:

$$\frac{d}{dt}\mathcal{E}(t) = \|r'(t)\|k_f(t)R\left(-\frac{\pi}{2}\right)\mathcal{E}(t).$$
(2)

Before starting its study we point out that this work is dedicated the memory of Academician Radu Miron (1927-2022). He was always interested in the geometry of curves and besides his theory of *Myller configurations* ([11]) he generalized also a type of curvature for space curves in [10]. We remark also that this note follows the idea of Bishop in his delightful note [2] and that the flow-curvature of spacelike parametrized curves in the Lorentz plane was introduced by the author in [4]. The hyperbolic curves are studied also by the author in [5].

Returning to our subject we note as a first main result:

**Proposition 1.** The expression of the flow-curvature is:

$$k_f(t) = k(t) - \frac{1}{\|r'(t)\|} < k(t).$$
(3)

As a consequence, the curve C and its trigonometrical rotation iC share the same flow-curvature.

**Proof** We have directly in the flow-frame:

$$\|r'(t)\|k_f(t)R\left(-\frac{\pi}{2}\right) = R\left(t + \frac{\pi}{2}\right)R(-t) + \|r'(t)\|k(t)R(t)R\left(-\frac{\pi}{2}\right)R(-t)$$
(4)

and the conclusion follows. Concerning the consequence it is obvious that C and  $iC: t \to (-y(t), x(t))$  share the same curvature k and the same second term from (3).  $\Box$ 

**Example 1.** i) If C is the line  $r_0 + tu, t \in \mathbb{R}$  with the vector  $u \neq \overline{0} = (0,0)$  then  $k_f$  is constant:

$$k_f(t) = -\frac{1}{\|u\|} = constant < 0.$$
(5)

In particular, if u is an unit vector then  $k_f(t) = -1$ .

ii) The circle C(O, R) with the usual parametrization  $r(t) = Re^{it}$  is a flat-flow curve i.e.  $k_f = 0$ . Indeed, the flow-frame is constant and universal for the families of concentric circles i.e. it does not depend on the radius R (exactly as the Frenet frame):

$$E_1^f = (0,1) = \overline{j}, \quad E_2^f = (-1,0) = -\overline{i}.$$
 (6)

More generally, if C is expressed in polar coordinates as  $C: \rho = \rho(t)$  for  $t \in I$ then C is a flat-flow curve if and only if C is a logarithmic spiral  $\rho_{R,\alpha}(t) = Re^{\alpha t}$ ,  $R, \alpha > 0$  and  $t \in \mathbb{R}$ . The limit case  $\alpha \to 0$  gives the circle  $\mathcal{C}(O, R)$  and the flow-frame of the logarithmic spiral is:  $E_1^f = \frac{1}{\sqrt{\alpha^2+1}}(\alpha, 1), E_2^f = \frac{1}{\sqrt{\alpha^2+1}}(-1, \alpha);$  if  $\alpha = \cot \varphi$  then  $E_1^f = e^{\varphi i}, E_2^f = e^{i(\frac{\pi}{2}+\varphi)}.$ 

iii) Fix  $R \in (0, +\infty)$  and the plane curve  $C : w = F(Re^{it})$  with t as an increasing parameter and F = F(z) a holomorphic function. Then the curvatures are:

$$k(t) = \frac{1}{|zF'(z)|} Re\left(1 + \frac{zF''(z)}{F'(z)}\right), \quad k_f(t) = \frac{1}{|zF'(z)|} Re\left(\frac{zF''(z)}{F'(z)}\right).$$
(7)

For the circle example of  $F(z) = z^2$  it results  $k = \frac{1}{R^2} = \text{constant}$  and  $k_f = \frac{1}{2R^2} = \text{constant}$  which proves the proper dependence of  $k_f$  on the parametrizations of C.  $\Box$ 

**Remark 1.** *i)* Suppose that I is symmetric with respect to  $0 \in \mathbb{R}$  and that C is positively oriented in the terms of Definition 1.14 from [14, p. 17]. Suppose also the C is convex; then applying the Theorem 1.18 of page 19 from the same book it results for the usual curvature the inequality  $k \ge 0$ . Hence the opposite curve  $C^-: t \in I \rightarrow r(-t)$  has the flow-curvature  $k_f < 0$ .

ii) An important tool in dynamics is the Fermi-Walker derivative. Let  $\mathcal{X}_C$  be the set of vector fields along the curve C. Then the Fermi-Walker derivative is the map ([6])  $\nabla_C^{FW} : \mathcal{X}_C \to \mathcal{X}_C$ :

$$\nabla_C^{FW}(X) := \frac{d}{dt} X + \|r'(\cdot)\|k[\langle X, N\rangle T - \langle X, T\rangle N] = \frac{d}{dt} X + \|r'(\cdot)\|k[X^{\flat}(N)T - X^{\flat}(T)N]$$
(8)

with  $X^{\flat}$  the differential 1-form dual to X with respect to the Euclidean metric. In a matrix form we can express this as follows:

$$\nabla_C^{FW} = \frac{d}{dt} - \|r'\|k \begin{vmatrix} (\cdot)^{\flat}(T) & (\cdot)^{\flat}(N) \\ T & N \end{vmatrix} = \frac{d}{dt} + \|r'\|k \begin{vmatrix} T & (\cdot)^{\flat}(T) \\ N & (\cdot)^{\flat}(N) \end{vmatrix}.$$
(9)

It is natural to make here a remark concerning rotation-minimizing fields  $X \in \mathcal{X}_C$  i.e. fields satisfying:

$$\frac{d}{dt}X(t) = \lambda(t)T(t), \quad \langle X(t), T(t) \rangle = 0$$

for a smooth function  $\lambda = \lambda(t)$ . Then the Fermi-Walker derivative of such X is also parallel with the tangent T:

 $\nabla_C^{FW} X(t) = [\lambda(t) + \|r'(t)\|k(t)\langle X(t), N(t)\rangle]T(t).$ 

Calculating the Fermi-Walker derivative on our frames we get:

$$\nabla_C^{FW}(T) = \nabla_C^{FW}(N) = 0, \quad \nabla_C^{FW}(E_1^f) = -E_2^f, \quad \nabla_C^{FW}(E_2^f) = E_1^f. \tag{10}$$

With the matrix notation we can express these relations as:

$$\nabla_C^{FW}(\mathcal{F}) = \begin{pmatrix} 0\\0 \end{pmatrix}, \quad \nabla_C^{FW}(\mathcal{E}) = R\left(\frac{\pi}{2}\right)\mathcal{E}$$
(11)

and the Fermi-Walker derivative can be expressed in terms of  $k_f$  as:

$$\nabla_C^{FW}(X) = \frac{d}{dt}X + (1 + ||r'||k_f)[X^{\flat}(N)T - X^{\flat}(T)N].$$
(12)

Also, we can define the flow-Fermi-Walker derivative as:

$$\nabla_C^{fFW}(X) := \frac{d}{dt}X + \|r'(\cdot)\|k_f[X^{\flat}(N)T - X^{\flat}(T)N] = \nabla_C^{FW}(X) + T \wedge N(X)$$
(13)

with the skew-symmetric endomorphism  $\wedge \in so(2)$  defined by:

$$X \wedge Y := \langle X, \cdot \rangle Y - \langle Y, \cdot \rangle X = (X^1 Y^2 - X^2 Y^1) R\left(\frac{\pi}{2}\right), X = (X^1, X^2), Y = (Y^1, Y^2) X = (Y^1, Y^1) X = (Y^1) X = (Y^1, Y^1) X = (Y^1) X = (Y^1) X = (Y^1)$$

Then:

$$\nabla_C^{fFW}(\mathcal{F}) = R\left(-\frac{\pi}{2}\right)\mathcal{F}, \quad \nabla_C^{fFW}(\mathcal{E}) = \begin{pmatrix} 0\\ 0 \end{pmatrix}.$$
 (14)

As in the usual case, if  $V, W \in \mathcal{X}_C$  are flow-Fermi-Walker fields i.e. with zero flow-Fermi-Walker derivative then the value  $\langle V, W \rangle \in \mathbb{R}$  is constant along C. iii) Remark that the 4-dimensional vectors  $\frac{1}{\sqrt{2}}\mathcal{F}$  and  $\frac{1}{\sqrt{2}}\mathcal{E}$  belong to the Clifford torus  $\frac{1}{\sqrt{2}}T^2 \subset S^3$ . A remarkable Riemannian submersion is the Hopf map H:  $S^3 \subset \mathbb{C}^2 \to S^2(\frac{1}{2}) \subset \mathbb{R} \times \mathbb{C}$ :

$$H(z,w) = \left(\frac{1}{2}(|z|^2 - |w|^2), z\bar{w}\right).$$
(15)

It follows:

$$H\left(\frac{1}{\sqrt{2}}\mathcal{F}(t)\right) = \left(0, \frac{1}{2}T(t)\bar{N}(t)\right) = \left(0, -\frac{i}{2}\right) = H\left(\frac{1}{\sqrt{2}}\mathcal{E}(t)\right).$$
(16)

Hence, considering H as a projection map of the  $S^1$ -principal bundle  $S^3 \to S^2(\frac{1}{2})$ we have that  $\frac{1}{\sqrt{2}}\mathcal{F}$  and  $\frac{1}{\sqrt{2}}\mathcal{E}$  belong to the same fiber, namely that over the South pole of the sphere  $S^2(\frac{1}{2})$ .

iv) Suppose now that our curve C belongs to the plane xOz of the physical space  $\mathbb{R}^3$  as C: r(t) = (f(t), 0, F(t)) with f > 0 on I and consider the rotational surface generated by C as:

$$\Sigma: \bar{r}(t,\varphi) := (f(t)\cos\varphi, f(t)\sin\varphi, F(t)), \quad \varphi \in S^1.$$

Its principal curvatures depend only on t, [8, p. 85]:

$$k_1 = k, \quad k_2 = \frac{F'}{\|r'\|f}$$
 (17)

and then for F' = f we have that  $k_f$  of C is exactly the difference  $k_1 - k_2$  of the principal curvatures of  $\Sigma$ ; consequently the umbilic circles of  $\Sigma$  are provided by the zeros of  $k_f$  and are parametrized by  $\varphi \in S^1$ .

For F' = f the curvatures of C are expressed only through the function F as:

$$k(t) = \frac{[F''(t)]^2 - F'(t)F'''(t)}{[F'(t)^2 + F''(t)^2]^{\frac{3}{2}}}, \quad k_f(t) = \frac{-F'(t)F'''(t) - [F'(t)]^2}{[F'(t)^2 + F''(t)^2]^{\frac{3}{2}}}$$
(18)

and due to the presence of the third derivative of F we recall its Schwarzian derivative:

$$S_F = \frac{F'''}{F'} - \frac{3}{2} \left(\frac{F''}{F'}\right)^2$$
(19)

which implies the new formulae:

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$$k = \frac{(F'')^2 - 2(F')^2 S_F}{2[(F')^2 + (F'')^2]^{\frac{3}{2}}}, \quad k_f = \frac{-3(F'')^2 - 2(F')^2 S_F - 2(F')^2}{2[(F')^2 + (F'')^2]^{\frac{3}{2}}}.$$
 (20)

In conclusion, a smooth F with negative Schwarzian derivative will give a positive curvature k for C while a positive Schwarzian derivative  $S_F$  produces a negative flow-curvature  $k_f$ .

v) The nature and the relationship between our frames can be put in the framework of moving frames of [8, p. 32]. Recall that the set of all orientation-preserving Euclidean isometries forms a Lie group,  $E(2) := \mathbb{R}^2 \times SO(2)$ , with the standard projection  $\pi_1$  on the first factor making  $E(2) \to \mathbb{R}^2$  an  $S^1$ -principal bundle. A moving frame along C is a map  $F: I \to E(2)$  such that  $\pi_1 \circ F = r$ . But C defines also a 1-parameter family of bijections of SO(2):

$$\begin{split} L^C: I &\to Bijections(SO(2)), t \to L^C(t): SO(2) \to SO(2), A \to R(t)A, (L^C(t))^{-1} = L^C(-t).\\ Then \ our \ frames \ are \ \mathcal{F}: I \to E(2) \ as \ \mathcal{F}(t) = (r(t), T(t), N(t)) \ and \ \mathcal{E}: I \to E(2) \\ as \ \mathcal{E}(t) = (r(t), (L^C(t) \circ \pi_2 \circ \mathcal{F})(t)). \end{split}$$

vi) Suppose now that the curve C is in the space  $\mathbb{R}^3$  and is bi-regular; hence it has the Frenet frame (T, N, B) and the pair (curvature, torsion)= $(k, \tau)$ . We define its flow-frame as:

$$\begin{pmatrix} T \\ E_2^f \\ E_3^f \end{pmatrix}(t) := \begin{pmatrix} 1 & 0_2(h) \\ 0_2(v) & R(t) \end{pmatrix} \begin{pmatrix} T \\ N \\ B \end{pmatrix}, \quad 0_2(h) := (0,0), \quad 0_2(v) := \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

and then, its matrix moving equation is:

$$\frac{d}{dt} \begin{pmatrix} T\\ E_2^f\\ E_3^f \end{pmatrix} (t) = \|r'(t)\| \begin{pmatrix} 0 & k_f^2(t) & k_f^3(t)\\ -k_f^2(t) & 0 & \tau_f(t)\\ -k_f^3(t) & -\tau_f(t) & 0 \end{pmatrix} \begin{pmatrix} T\\ E_2^f\\ E_3^f \end{pmatrix} (t).$$

A similar computation yields:

$$k_f^2(t) = k(t)\cos t, \quad k_f^3(t) = k(t)\sin t, \quad \tau_f(t) = \tau(t) - \frac{1}{\|r'(t)\|} < \tau(t).$$

We point out the formal similarity with the Darboux equations of a curve on a given surface and then a curve C with vanishing  $\tau_f$  will be called flow-geodesic in  $\mathbb{R}^3$ . Hence, if C is naturally parametrized then C is a flow-geodesic if and only if its torsion has the constant value 1; for this class of space curves and examples see [1]. In order to express the above moving equation in the compact form as in the theory of space curves:

$$\omega_f(t) \times T(t) = T'(t), \quad \omega_f(t) \times E_2^f(t) = (E_2^f)'(t), \quad \omega_f(t) \times E_3^f(t) = (E_3^f)'(t)$$

we associate a vector field along C, called flow-Darboux:

$$\omega_f(t) := \|\gamma'(t)\| [\tau_f(t)T(t) - k_f^3(t)E_f^2(t) + k_f^2(t)E_3^f(t)].$$

Something similar but with the rotation with respect to an angle  $\theta = \theta(s)$  appears in [13] under the name of quasi frame for C. Our choice corresponds to the angle  $\theta(s) = -s$ .

vii) Suppose that the curvature function  $t \to k(t)$  is always strictly positive (or strictly negative). Then the evolute of C is the curve:

$$C_e: r_e(t) := r(t) + \frac{1}{k(t)}N(t).$$

With this model in mind, for a non-flat-flow curve we associate its flow-evolute as being the curve:

$$C_{fe}: r_{fe}(t) := r(t) + \frac{1}{k_f(t)} E_2^f(t)$$

We will obtain this curve for some examples below. So, the line C discussed in the example 1 has the flow-evolute

$$C_{fe}: r_{fe}(t) = r_0 + (t - \sin t)u - \cos t(iu)$$

and for  $r_0 = (0,1) = iu$  this last curve is exactly the cycloid of radius R = 1 according to the example 3 below.  $\Box$ 

Returning to the plane curves let  $J \subseteq \mathbb{R}$  be another open interval and fix the diffeomorphism  $\varphi : s \in J \to t \in I$  with the smooth inverse  $\varphi^{-1} : t \in I \to s \in J$ . Since

 $r'(s) = \varphi'(s)r'(t(s))$  we restrict our study to the class  $Diff_+(J, I)$  of orientationpreserving diffeomorphisms:  $\varphi'(s) > 0$ , for all  $s \in J$ . The transformation of the flow-curvature under the action of  $\varphi$  is:

$$k_f(s) = k(t) - \frac{1}{\varphi'(s) \|r'(t)\|}$$
(21)

and then:

$$k_f(s) - k_f(t) = \frac{1}{\|r'(t)\|} \left[ 1 - \frac{1}{\varphi'(s)} \right].$$
 (22)

**Proposition 2.** (the rigidity of the flow-curvature) The only orientation-preserving diffeomorphism  $\varphi$  which preserves also the flow-curvature of C is an interval shift on the real line  $\varphi(s) = s + s_0$ ,  $s_0 \in (0, +\infty)$ .

A natural important problem is the class of curves with prescribed flow-curvature. For example, if we ask the vanishing of the flow-curvature for a graphic curve  $C_F: r(t) = (t, F(t))$  then it follows the differential equation:

$$\frac{F''(t)}{\left[1 + (F'(t))^2\right]^{\frac{3}{2}}} = \frac{1}{\left[1 + (F'(t))^2\right]^{\frac{1}{2}}}.$$
(23)

Since this equation reads:

$$\frac{F''(t)}{1 + (F'(t))^2} = 1 \tag{24}$$

we have exactly the Grim Reaper solution, [3, p. 28], a famous solution of the curve shortening flow:

$$F_u(t) = u - \ln(\cos t), \quad t \in \left(-\frac{\pi}{2}, \frac{\pi}{2}\right), \quad u \in \mathbb{R}$$
 (25)

with the usual curvature  $k(t) = \cos t$  and the frames:

$$\mathcal{F}(t) = \begin{pmatrix} e^{it} \\ e^{i\left(t+\frac{\pi}{2}\right)} \end{pmatrix}, \quad \mathcal{E} = \begin{pmatrix} (1,0) = \overline{i} \\ (0,1) = \overline{j} \end{pmatrix} = constant.$$
(26)

Another formalism is that of [15, p. 2] if  $r : S^1 \simeq [0, 2\pi) \to \mathbb{R}^2$  is naturally parametrized then there exists the smooth function  $\theta : S^1 \to \mathbb{R}$ , called *normal angle*, such that:

$$N(s) = e^{i\theta(s)} = (\cos\theta(s), \sin\theta(s)), \quad T(s) = -iN(s) = -ie^{i\theta(s)} = e^{i(\theta(s) - \frac{\pi}{2})}$$
(27)

and then the Frenet equations yield:

$$\frac{d\theta}{ds}(s) = k(s). \tag{28}$$

In conclusion, the constant value  $\beta \in \mathbb{R}$  of the flow-curvature of a closed convex curve means  $\theta(s) = (\beta + 1)s + \alpha$  for all  $s \in S^1$  with  $\alpha \in \mathbb{R}$  an arbitrary constant. The flow-frame corresponding to the equations (27) is:

$$E_1^f(s) = (\sin(\theta(s) - t(s)), -\cos(\theta(s) - t(s))), E_2^f(s) = (\cos(\theta(s) - t(s)), \sin(\theta(s) - t(s)))$$
(29)

which, in turn, is the Frenet frame of a new curve with the same natural parameter s but having the normal angle  $\tilde{\theta}(s) := \theta(s) - t(s)$ .

The formula (28) can be replaced with  $\frac{d(\theta - \pi/2)}{ds}(s) = k(s)$  which expresses the curvature k as the derivative of the angle between  $T \in \mathcal{X}_C$  and the unit vector  $\bar{i}$ . Following this approach the paper [7] generalizes k to a curvature-type function  $k_V$  defined with respect to an arbitrary  $V \in \mathcal{X}_C$ . A main result of the cited work is that  $k_V = k_W$  if and only if the angle between V and W is constant along C. Hence, we can apply the last statement of the Remark ii) and then two flow-Fermi-Walker unit vectors  $V, W \in \mathcal{X}_C$  yield the same curvature-type function.

In the following we present a couple of examples in order to remark the computational aspects of our approach.

**Example 2.** The involute of the unit circle  $S^1$  is:

$$C: r(t) = (\cos t + t\sin t, \sin t - t\cos t) = (1 - it)e^{it}, \quad t \in (0, +\infty).$$
(30)

A direct computation gives:

$$r'(t) = (t\cos t, t\sin t) = te^{it}, \quad \|r'(t)\| = t, \quad k(t) = \frac{1}{t} > 0,$$
(31)

and then this curve is also a flat-flow one and having the same flow-frame as the Grim Reaper. This example can be treated also with respect to a natural parameter  $s \in (0, +\infty)$  which is provided by  $t := \sqrt{2s}$ . For example, the normal angle function is  $\theta(s) = \frac{\pi}{2} + \sqrt{2s}$  since then  $r'(s) = e^{i\sqrt{2s}}$ . Comparing with the approach above it results the constants  $\alpha = \frac{\pi}{2}$  and  $\beta = \sqrt{2} - 1$ .  $\Box$ 

**Example 3.** Recall that for R > 0 the cycloid of radius R has the equation:

$$C: r(t) = R(t - \sin t, 1 - \cos t) = R[(t, 1) - e^{i(\frac{\pi}{2} - t)}], \quad t \in \mathbb{R}.$$
(32)

Remark that here we have a twisted situation of the Remark iv) namely the derivative of the first component of the vector r(t) is exactly the second component. The Schwarzian derivative is:

$$S_{t-\sin t}(t) = \frac{\cos t}{\sin t} - \frac{3}{2} \left( \frac{\cos \frac{t}{2}}{\sin \frac{t}{2}} \right)^2, \quad t \in \mathbb{R} \setminus \mathbb{Z}\pi.$$
(33)

We have immediately:

$$r'(t) = R(1 - \cos t, \sin t) = R[(1, 0) - e^{it}], ||r'(t)|| = 2R|\sin\frac{t}{2}|, k(t) = -\frac{1}{4R|\sin\frac{t}{2}|},$$
(34)

and then we restrict our definition domain to  $(0,\pi)$ . It follows:

$$\begin{cases} k_f(t) = -\frac{3}{4R\sin\frac{t}{2}} < 0, \\ E_1^f(t) = (\sin\frac{3t}{2}, \cos\frac{3t}{2}) = e^{i\left(\frac{\pi}{2} - \frac{3t}{2}\right)}, E_2^f(t) = (-\cos\frac{3t}{2}, \sin\frac{3t}{2}) = e^{i\left(\pi - \frac{3t}{2}\right)}. \end{cases}$$
(35)

Again a natural parameter s is provided by:  $t = 2 \arccos \left(1 - \frac{s}{4R}\right)$  and the flowevolute of C is the curve:

$$C_{fe}: r_{fe}(t) = R(t - \sin t, 1 - \cos t) + \frac{4}{3}R\sin\frac{t}{2}(\cos t, -\sin t), \quad t \in (0, \pi).$$

**Example 4.** The derivative curve r' from (31) is an Archimedes' spiral. This spiral is given in polar coordinates as:

$$A(spiral): \rho(t) = Rt, \quad R > 0 \tag{36}$$

and hence:

$$k_f(t) = \frac{1}{R(t^2 + 1)^{\frac{3}{2}}} > 0 \tag{37}$$

while its flow-evolute is the curve:

$$C_{fe}: r_{fe}(t) = R(t\cos t, t\sin t) + R(1+t^2)(-\sin t - t\cos t, \cos t - t\sin t).$$

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**Example 5.** Fix  $\alpha \in \mathbb{R}^*$  and the naturally parametrized curve C. Then the  $\alpha$ -parallel curve of C is the new curve:

$$C_{\alpha}: \tilde{r}(t) := r(t) + \alpha N(t), \quad t \in I$$
(38)

with:

$$\tilde{T}(t) = \frac{1 - \alpha k(t)}{|1 - \alpha k(t)|} T(t), \quad \tilde{N}(t) = \frac{1 - \alpha k(t)}{|1 - \alpha k(t)|} N(t), \quad \tilde{k}(t) = k(t).$$
(39)

Hence, we consider that  $\alpha$  does not belongs to the range of the function  $\frac{1}{k}$  and the new flow-curvature is:

$$\tilde{k}_f(t) = k(t) - \frac{1}{|1 - \alpha k(t)|}.$$
(40)

We finish this note with the problem raised in the beginning, namely the possible variants of the curve shortening flow. Recall that the setting of this question consists in a 1-parameter family of plane curves  $C_u : r = r_u(t) = r(t, u)$  satisfying:

$$\frac{\partial r(t,u)}{\partial u} = k(t,u)N(t,u). \tag{41}$$

It follows immediately an expression in terms of flow-apparatus:

$$\frac{\partial r(t,u)}{\partial u} = \left(k_f(t,u) + \frac{1}{\|r'(t,u)\|}\right) \left[-\sin t E_1^f(t,u) + \cos t E_2^f(t,u)\right].$$
(42)

The first variant which we propose as an open problem is to study the flow-variant of (41):

$$\frac{\partial r(t,u)}{\partial u} = k_f(t,u)E_2^f(t,u). \tag{43}$$

The second variant is to generalize all this study through a general smooth function  $\Omega \in C^{\infty}(\mathbb{R})$ . More precisely, we use the equation (1) with R replaced by  $R \circ \Omega$  to define the notion of  $\Omega$ -frame for the plane curve C; we note that for a particular choice of  $\Omega$  the 3-dimensional variant of the remark vi) is called *positional adapted* frame in [12]. Then the  $\Omega$ -curvature of the plane curve C is:

$$k_{\Omega}(t) = k(t) - \frac{\Omega'(t)}{\|r'(t)\|}$$
(44)

and the curves in polar coordinates with vanishing  $\Omega$ -curvature are provided by:

$$\rho(t) = Re^{\int_{t_0}^t \cot[\Omega(u) - u + C]du}, \quad R > 0, \quad C \in \mathbb{R}.$$
(45)

The flow-curvature corresponds to the identity map  $\Omega = 1_{\mathbb{R}}$ . Moreover, if C is naturally parametrized then  $k_{\Omega} = (\theta - \Omega)'$  which means that the case  $\Omega = \theta + constant$ provides a zero  $\Omega$ -curvature.

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# References

- Bates, L. M., Melko, O. M., On curves of constant torsion I, J. Geom., 104 (2) (2013), 213–227. https://doi.org/10.1007/s00022-013-0166-2
- Bishop, R. L., There is more than one way to frame a curve, Am. Math. Mon., 82 (1975), 246-251. https://doi.org/10.2307/2319846
- [3] Chou, K.-S., Zhu, X.-P., The Curve Shortening Problem, Boca Raton, FL: Chapman & Hall/CRC, 2001. Zbl 1061.53045

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- [4] Crasmareanu, M., The flow-curvature of spacelike parametrized curves in the Lorentz plane, Proceedings of the International Geometry Center, 15 (2) (2022), 100–108. https://doi.org/10.15673/tmgc.v15i2.2281
- [5] Crasmareanu, M., The flow-geodesic curvature and the flow-evolute of hyperbolic plane curves, *Int. Electron. J. Geom.*, 16 (2023), no. 1, 225–231. https://doi.org/10.36890/iejg.1229215
- [6] Crasmareanu, M., Frigioiu, C. Unitary vector fields are Fermi-Walker transported along Rytov-Legendre curves, Int. J. Geom. Methods Mod. Phys., 12 (10) (2015), Article ID 1550111. https://doi.org/10.1142/S021988781550111X
- [7] Góźdź, S., Curvature type functions for plane curves, An. Ştiinţ. Univ. Al. I. Cuza Iaşi Mat., 39 (3) (1993), 295–303. Zbl 0851.53001
- [8] Jensen, G. R., Musso, E., Nicolodi, L., Surfaces in Classical Geometries. A Treatment by Moving Frames, Universitext, Springer, 2016. Zbl 1347.53001
- [9] Mazur, B., Perturbations, deformations, and variations (and "near-misses") in geometry, physics, and number theory, Bull. Am. Math. Soc., 41 (3) (2004), 307–336. https://doi.org/10.1090/S0273-0979-04-01024-9
- [10] Miron, R., Une généralisation de la notion de courbure de parallélisme, Gaz. Mat. Fiz., Bucureşti, Ser. A 10 (63) (1958), 705–708. Zbl 0087.36101
- [11] Miron, R., The geometry of Myller configurations. Applications to theory of surfaces and nonholonomic manifolds, Bucharest: Editura Academiei Române, 2010. Zbl 1206.53003
- [12] Özen, K. E., Tosun, M., A new moving frame for trajectories with non-vanishing angular momentum, J. of Mathematical Sciences and Modelling, 4 (1) (2021), 7–18. https://doi.org/10.33187/jmsm.869698
- [13] Soliman, M. A., Nassar, H.A.-A., Hussien, R. A., Youssef, T., Evolutions of the ruled surfaces via the evolution of their directrix using quasi frame along a space curve, J. of Applied Mathematics and Physics, 6 (2018), 1748–1756. https://doi.org/10.4236/jamp.2018.68149
- [14] Younes, L., Shapes and Diffeomorphisms, 2nd Updated Edition, Applied Mathematical Sciences 171, Berlin, Springer, 2019. Zbl 1423.53002
- [15] Zhu, X.-P., Lectures on Mean Curvature Flows, AMS/IP Studies in Advanced Mathematics vol. 32, Providence, RI: American Mathematical Society, 2002. Zbl 1197.53087