



Calculation of Dynamic Properties of Drug-Added Aqueous Solutions with T_1 and T_2 Relaxation Times

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Abstract

The development of modern spectroscopic methods has facilitated and accelerated structure analysis. The NMR method is the most popular way to perform structural analysis of compounds with very complex structures. D_2O is a solvent that is frequently used in NMR analysis of both chemical molecules and many biological molecules such as drugs, proteins, and enzymes. In this paper, the study of residual water in proton drug-added protein solutions was carried out via NMR relaxation. The spin-lattice (T_1) and the spin-spin relaxation (T_2) times of residual water in drug-added protein solutions were studied depending on temperature by Avance Bruker 400 MHz 1H -NMR Spectrometer, and activation energies (E_a) and rotational correlation times (τ_0 and τ_c) have been determined for T_1 and T_2 relaxation times.

Keywords: D_2O ; Proton NMR; Relaxation times; Activation energy; Flurouacil.

İlaç Katkılı Sulu Çözeltilerin Dinamik Özelliklerinin T_1 ve T_2 Rölaksasyon Zamanları ile Hesaplanması

Öz



Modern spektroskopik yöntemlerin gelişimi, yapı analizini kolaylaştırmış ve hızlandırmıştır. NMR yöntemi, çok karmaşık yapılara sahip bileşiklerin yapısal analizini gerçekleştirmenin en popüler yoludur. D₂O, hem kimyasal moleküllerin hem de ilaçlar, proteinler, enzimler gibi birçok biyolojik molekülün NMR analizinde sıklıkla kullanılan bir çözücüdür. Bu çalışmada, ilaç katkılı protein çözeltilerinde residual su çalışması, proton NMR rölaksasyonu yoluyla gerçekleştirilmiştir. Avance Bruker 400 MHz ¹H-NMR Spektrometresi ile ilaç katkılı protein çözeltilerinde residual suyun spin-örgü (T₁) ve spin-spin rölaksasyon (T₂) süreleri sıcaklığa bağlı olarak ölçülmüş, T₁ ve T₂ rölaksasyon süreleri için aktivasyon enerjileri (E_a) ile τ₀ ve τ_c dönme korelasyon zamanları belirlenmiştir.

Anahtar Kelimeler: D₂O; Proton NMR; Rölaksasyon zamanları; Aktivasyon enerjisi; Flurouacil.

1. Introduction

The contribution of magnetic resonance to science and technology is one of the greatest discoveries of the century in the field of spectroscopy. Nuclear magnetic resonance technique (NMR) is a spectroscopic technique commonly used by researchers in many different fields to reveal the structures of chemicals, especially organic molecules, in liquid or solid form. After successful experiments by Purcell, Pound, Torrey at Harvard and Bloch, Hansen, Packard at Stanford in 1946, the field of application of magnetic resonance developed rapidly. The theoretical and technical development of this phenomenon, whose basic formulation is summarized by Abragam [1], is still ongoing.

The application of NMR relaxation times to understand the moving mechanisms of molecules in aqueous solutions has become an indispensable and frequently used method. The relaxation rates of water molecules are influenced by the local magnetic interactions between the nuclei of the water and the rate of molecular motion and proton exchange [2]. There are two mechanisms for the relaxation rates of water molecules: spin-lattice and spin-spin.

The relaxation time T₁, which defines the return to equilibrium in the magnetic field direction, is also known as spin-lattice relaxation or longitudinal relaxation and can be measured from the magnetization creation along the applied static magnetic field. T₁ relaxation can also be measured by the decay of a signal that creates a rotating magnetic field around the resonance frequency perpendicular to the static magnetic field. T₁ longitudinal or spin-lattice relaxation may be more sensitive to intramolecular mobility on flexible substrates, as it is associated with general rotational of the molecule in solution. The standard method for measuring T₁ is known as reversal recovery.

The relaxation time T_2 , which describes the decay of induced magnetization perpendicular to the applied magnetic field, is also called spin-spin relaxation or transverse relaxation. It is known that the observed spectral line width is associated with the T_2 relaxation time, that is, spin-spin relaxation, and is also affected by magnetic inhomogeneity. The relaxation time T_2 is used to determine the decay rate of the magnetization in the xy plane. T_1 and T_2 relaxation times are long in rapidly rotating and small molecules such as free water. As in the case of proteins, as molecular motion slows down, T_2 relaxation gets shorter while T_1 increases. As a result, T_2 relaxation is faster than T_1 relaxation, and however, the T_1 relaxation time always takes a value longer than or equal to T_2 [3].

The rotational correlation time (τ_c), used in molecular-scale viscosity measurements and protein characterization, refers to the average time it takes for a molecule to rotate one radian, and its value in solutions is on the order of picoseconds and is known to be 1.7 ps for water [4]. The rotational correlation time can be determined for a molecule by T_1 and T_2 relaxation times and can be measured by methods such as microwave, dielectric, and nuclear magnetic resonance NMR spectroscopy [5, 6].

The most widely known mechanism is the magnetic dipole-dipole interaction. In this interaction, the mechanism between a nucleus and another nucleus or the magnetic moment of another environment such as the electron, atom, ion, or molecule is considered. Besides magnetic dipole-dipole interaction, there are different mechanisms for relaxation such as Chemical shift anisotropy (CSA) relaxation, Spin rotation (SR) relaxation mechanism, and quadrupole relaxation mechanism. In addition to these mechanisms, molecular reorientation or tumbling and the electrostatic interaction between nuclei can also lead to spin transitions and relaxation.

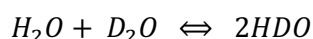
By analyzing aqueous solutions with NMR relaxation techniques, important data about rotational motion, molecular interactions and structure can be obtained. It was mentioned before that T_2 is smaller than T_1 . The T_1 and T_2 curves overlap in the region called the extreme narrowing limit ($\omega^2 \tau_c^2 \ll 1$) which includes small rotational correlation times [7-14]. The expression “extreme narrowing” describes rapid molecular rotation on the order of picoseconds and simplifies the equations of rotational correlation [15].

Using proton resonance, the water dynamics has been extensively investigated for low frequencies by NMR relaxation methods, and it has been thought that the use of proton resonance is a major disadvantage [2-7]. However, several biological molecules, especially

proteins and enzymes have been studied in D₂O to decrease the water in the environment and increase the effectiveness of the bound water on the protein [16-25].

D₂O is one of the most important and widely preferred solvents used in NMR studies of chemical molecules. The standard water (H₂O) ratio in this solvent, which is tried to be purified but cannot be 100% pure, contains some residual hydrogen, which is called residual water.

The hydrogen-deuterium exchange reaction (H-D or H/D Exchange) between residual water molecules and D₂O is written as follows.



This expression indicates that the residual water in D₂O is in the HDO form. Therefore, it is very important to investigate HDO relaxation for information about the relaxation mechanism in drug-protein containing solutions.

Since the residual water in D₂O is in the form of HDO, investigation of HDO relaxation is valuable for understanding the mechanism of relaxation in several protein solutions. [26-30]. To easily calculate the input parameters of the chemical change-exchange formula, the proton relaxation of residual water must be known.

In this paper, to investigate residual water in drug-added protein solutions, proton NMR relaxation methods have been used. T₁ and the T₂ relaxation times of residual water in drug-added protein solutions were studied depending on temperature by Avance Bruker 400 MHz ¹H-NMR Spectrometer, and activation energies (E_a) and rotational correlation times (τ_0 and τ_c) have been calculated for T₁ and T₂ relaxation times.

2. Theoretical Background

T₁ relaxation time is also called spin-lattice or longitudinal relaxation time. This relaxation time is a measure of the energy transfer from the core spin system to neighboring molecules or lattices. This relaxation takes place in the z-axis and causes the Boltzmann equilibrium to change.

The states of a dipole pair in a water molecule exposed to the external field B₀ applied in the z direction are as follows.

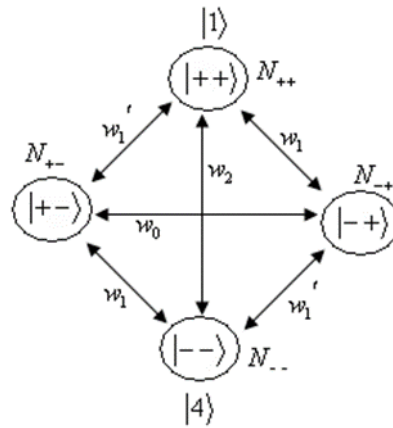


Figure 1: States of a dipole pair in a water molecule exposed to an external field in the z-direction

Figure 1 showing the eigenstates of I_z with respect to a spin pair are valid for all pairs in the example. Therefore, the population of each condition is shown above. If an RF field is applied to the water sample, there will be transitions between the above levels and saturation will occur. In other words, the magnetization in the z direction lies in the y direction. After the RF is cut off, the dipolar hamiltonian, which occurs due to the interaction of a couple among themselves, leads to transitions between states and eventually reaches the Boltzmann equilibrium. The dipolar hamiltonian leading to transitions between states can be written as;

$$H_D(t) = \left[I_z^{(1)} I_z^{(2)} - \frac{1}{4} (I_+^{(1)} I_-^{(2)} + I_-^{(1)} I_+^{(2)}) \right] F_0(t) + [I_-^{(1)} I_z^{(2)} + I_z^{(1)} I_-^{(2)}] F_1(t) + [I_+^{(1)} I_z^{(2)} + I_z^{(1)} I_+^{(2)}] F_{-1}(t) + [I_-^{(1)} I_-^{(2)}] F_2(t) + [I_+^{(1)} I_+^{(2)}] F_{-2}(t) \tag{1}$$

Here, the terms in square brackets denote the operators, while the terms abbreviated as F denote the time-dependent parts of the hamiltonian. It is known that the F terms expressing the time-dependent parts of the Hamiltonian are written as follows.

$$F_0(t) = d (1 - 3 \cos^2 \theta) \qquad d = \frac{\mu_0 \gamma_1 \gamma_2 \hbar^2}{4\pi b^3}$$

$$F_1(t) = -\frac{3}{2} d \sin \theta \cos \theta \exp (i\phi) \qquad F_{-1}(t) = F_1(t)^*$$

$$F_2(t) = -\frac{3}{4} d \sin^2 \theta \exp (2i\phi) \qquad F_{-2}(t) = F_2(t)^*$$

To examine the transitions between the $|4\rangle$ and $|1\rangle$ states shown in Figure 1, it will be useful to start with Perturbation theory by considering the perturbation calculation between paired spins that Solomon’s work in the past [13]. Accordingly, the transition probability between states $|4\rangle$ and $|1\rangle$ can be written as,

$$a_{41} = \frac{1}{i\hbar} \int_0^t \langle 1 | H'(t') | 4 \rangle \exp [i(\omega_1 + \omega_2) t'] dt' \quad (3)$$

Continuing the quantum mechanical calculations and considering the $R_{22}(\tau)$ autocorrelation function w_{41} transition probability per unit time is written as,

$$w_2 = \frac{1}{\hbar^2} \int_{-t}^t R_{22}(\tau) \exp[-i(\omega_1 + \omega_2)\tau] d\tau = \frac{2}{\hbar^2} \langle |F_2(t)|^2 \rangle \frac{\tau_c}{1+(\omega_1+\omega_2)^2\tau_c^2} = \frac{1}{\hbar^2} J_2(\omega_1) \quad (4)$$

As can be seen from Eqn. (4), only F_2 is included in the process. That's why w_{41} is called w_2 . If the transitions between other levels are calculated in a similar way, the following equations are obtained.

$$w_0 = \frac{1}{8\hbar^2} \langle |F_0(t)|^2 \rangle \frac{\tau_c}{1+(\omega_1-\omega_2)^2\tau_c^2} \quad (5)$$

$$w_1 = \frac{1}{2\hbar^2} \langle |F_1(t)|^2 \rangle \frac{\tau_c}{1+\omega_1^2\tau_c^2} \quad (6)$$

$$w_1' = \frac{1}{2\hbar^2} \langle |F_1(t)|^2 \rangle \frac{\tau_c}{1+\omega_2^2\tau_c^2} \quad (7)$$

The importance of dipolar interaction in NMR relaxation has been proven in the past work of Bloembergen, Purcell and Pound [14]. By substituting these equations in $\frac{1}{T_1} = 2(w_1 + w_2)$ and $\frac{1}{T_1} = w_0 + 2(w_1 + w_2)$ written for the relaxation time T_1 in similar spin and different spin states, $\frac{1}{T_1}$ relaxation rates are obtained for similar spin and different spin states, respectively.

$$\frac{1}{T_1} = \frac{6}{20} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma^4 \hbar^2}{b^6} \left(\frac{\tau_c}{1+\omega^2\tau_c^2} + \frac{4\tau_c}{1+4\omega^2\tau_c^2} \right) \quad (8)$$

$$\frac{1}{T_1} = \frac{1}{10} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma_1^2 \gamma_2^2 \hbar^2}{b^6} \left(\frac{\tau_c}{1+(\omega_1-\omega_2)^2\tau_c^2} + \frac{3\tau_c}{1+\omega_2^2\tau_c^2} + \frac{6\tau_c}{1+(\omega_1+\omega_2)^2\tau_c^2} \right) \quad (9)$$

It is known that after an RF pulse is applied to a system, the magnetization lies in the y direction and this magnetization goes to zero by dephasing, and the process of this magnetization to zero is equal to the T_2 time. Therefore, T_2 is also called dephasing time. The same process followed to calculate the T_1 time in a dipole pair sample in a water molecule under the external field effect can also be followed to calculate the T_2 time. For similar spins and different spins, the equations for T_2 are written as follows, respectively.

$$\frac{1}{T_2} = 2(u_1 + u_2) \quad (10)$$

$$\frac{1}{T_2} = (u_0 + 2 u_1 + u_2) \tag{11}$$

Considering Fig.1, the diagram showing all states and the transition probabilities between these states, by using the transition probabilities u_0 , u_1 and u_2 between states calculated for similar spins and different spins; $\frac{1}{T_2}$ relaxation rate for the two different situations are respectively written as follows.

$$\frac{1}{T_2} = 2 (u_1 + u_2) = \frac{3}{20} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma^4 \hbar^2}{b^6} \left(3 \tau_c + \frac{5 \tau_c}{1+\omega^2 \tau_c^2} + \frac{2\tau_c}{1+4 \omega^2 \tau_c^2} \right) \tag{12}$$

$$\frac{1}{T_2} = u_0 + 2 u_1 + u_2 = \frac{1}{20} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma_1^2 \gamma_2^2 \hbar^2}{b^6} \left(4 \tau_c + \frac{\tau_c}{1+(\omega_1-\omega_2)^2 \tau_c^2} + \frac{3 \tau_c}{1+\omega_1^2 \tau_c^2} \right. \\ \left. + \frac{6 \tau_c}{1+\omega_2^2 \tau_c^2} + \frac{6\tau_c}{1+(\omega_1+\omega_2)^2 \tau_c^2} \right) \tag{13}$$

Considering Eqns. (9) and (13), T_1 and T_2 relaxation rates because of dipolar interaction between hydrogen (H) and deuterium (D) were determined by Solomon-Bloembergen.

$$\frac{1}{T_1} = \frac{1}{10} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma_D^2 \gamma_H^2 \hbar^2}{b^6} \left(\frac{\tau_c}{1+(\omega_D-\omega_H)^2 \tau_c^2} + \frac{3 \tau_c}{1+\omega_D^2 \tau_c^2} + \frac{6\tau_c}{1+(\omega_D+\omega_H)^2 \tau_c^2} \right) \tag{14}$$

$$\frac{1}{T_2} = \frac{1}{20} \left(\frac{\mu_0}{4\pi}\right)^2 \frac{\gamma_D^2 \gamma_H^2 \hbar^2}{b^6} \left(4 \tau_c + \frac{\tau_c}{1+(\omega_D-\omega_H)^2 \tau_c^2} + \frac{3 \tau_c}{1+\omega_D^2 \tau_c^2} + \frac{6 \tau_c}{1+\omega_H^2 \tau_c^2} \right. \\ \left. + \frac{6\tau_c}{1+(\omega_D+\omega_H)^2 \tau_c^2} \right) \tag{15}$$

In Eqns. (14) and (15), γ_D and γ_H are the deuteron and proton gyromagnetic ratio, respectively, b is the distance between the H and D spins, \hbar is Planck's constant divided by 2π , τ_c is rotational correlation time, ω_H and ω_D are Larmor frequencies of proton and deuteron, respectively and μ is the magnetic permeability.

The rotational correlation time τ_c in Eqns. (14) and (15) is given by the Arrhenius equation as follows.

$$\tau_c = \tau_0 \exp(E_a/RT) \tag{16}$$

Here, E_a is the activation energy for molecular motions, R is the gas constant, and T is the temperature, τ_0 and τ_c are rotational correlation time constant and rotational correlation time, respectively.

It is known that the T_1 and T_2 times in Eqns. (14) and (15) are equal to each other in the case of extreme narrowing, expressed as $\omega^2 \tau_c^2 \ll 1$. In this case, Eqns. (14) and (15) turn into Eqn. (17).

$$\frac{1}{T_1} = \frac{1}{T_2} = \left(\frac{\mu_0}{4\pi}\right)^2 \left(\frac{\gamma_D^2 \gamma_H^2 \hbar^2}{b^6}\right) \tau_c \quad (17)$$

Using Eqns. (16) and (17), Eqn. (18) is obtained.

$$\frac{1}{T_1} = \frac{1}{T_2} = \left(\frac{\mu_0}{4\pi}\right)^2 \left(\frac{\gamma_D^2 \gamma_H^2 \hbar^2 \tau_0}{b^6}\right) \exp(E_a/RT) \quad (18)$$

Here, assuming that $A = \left(\frac{\mu_0}{4\pi}\right)^2 \left(\frac{\gamma_D^2 \gamma_H^2 \hbar^2 \tau_0}{b^6}\right)$ the Eqn. (18) is written as follows.

$$\frac{1}{T_1} = \frac{1}{T_2} = A \exp(E_a/RT) \quad (19)$$

$$\ln T_1 = \ln T_2 = \ln \frac{1}{A} - \left(\frac{E_a}{R}\right) \left(\frac{1}{T}\right) \quad (20)$$

According to Eqn. (20), it is seen that there is a linear dependence between $\ln T_1$ and $\left(\frac{1}{T}\right)$ and it also has a slope with a negative sign.

3. Experimental

The most commonly used deuterated solvent in NMR experiments is D₂O. All chemicals used in this study were obtained from the Sigma-Aldrich catalog, together with the high purity D₂O (99.9%) used as the solvent in this study.

Albumin, the most abundant protein in human serum, is a monomeric structure with a molecular weight of 66 kD, consisting of three helical segments, each divided into two subregions, and has many ligand-binding abilities [31]. The most important factor affecting the distribution of intravenous drugs is the binding affinity of that drug to albumin in human blood serum, and therefore, the interaction between the drug and albumin must be optimal for the drug to reach the target organ and provide an effective treatment.

5-Fluorouracil, which is the active ingredient of the drug named Biosyn, is a fluorinated pyrimidine derivative and is a chemotherapy drug belonging to the group of antimetabolites.

Spin-Lattice Relaxation Times (T_1) and Spin-Spin Relaxation Times (T_2) were measured with a 400 MHz Bruker NMR Spectrometer. Measurements of Spin-Lattice Relaxation Times were performed with the Inversion Recovery technique using (180- τ -90) pulse steps.

Measurements of Spin-Spin Relaxation Times were carried out using CPMG (Carr-Purcell-Meiboom-Gill) technique with (90- τ -180- τ) pulse steps. After the solutions were prepared, they were placed in 5 mm NMR tubes to make NMR measurements. The delay times τ required for each measurement were chosen to match each peak observed in the spectrum. Temperature dependence measurements were made at 25, 30, 35, 40, 45 °C values by increasing the sample temperature by 5 °C.

4. Results and Discussion

The ^1H - NMR spectrum of FU solution for relaxation time measurements is given in Fig. 2. As seen in Fig. 2, the peak that appears as a single narrow line is the HDO peak at 4.7044 ppm and a doublet peak at 7.5978 ppm and 7.5845 ppm. The peak of the CH molecule seen in the chemical formula of FU was split in half due to the neighboring HN molecule and formed a doublet. The distance between the peaks of this doublet gives the coupling constant (j).

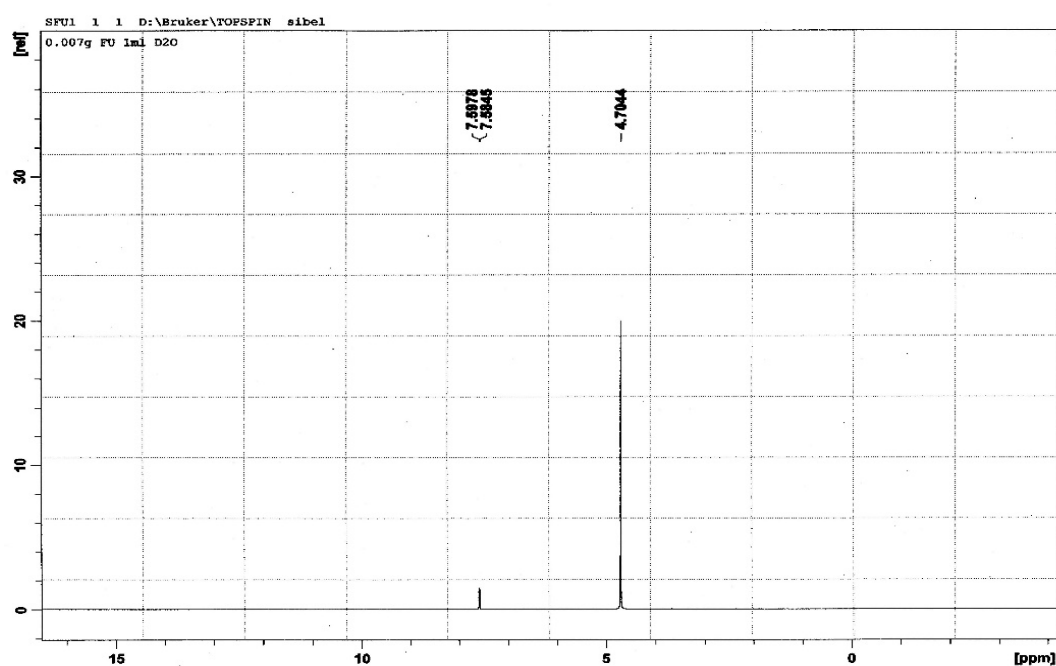


Figure 2: ^1H - NMR spectrum of FU in D_2O

T_1 and T_2 relaxation times graphs are shown in Fig. 3 a and b, respectively.

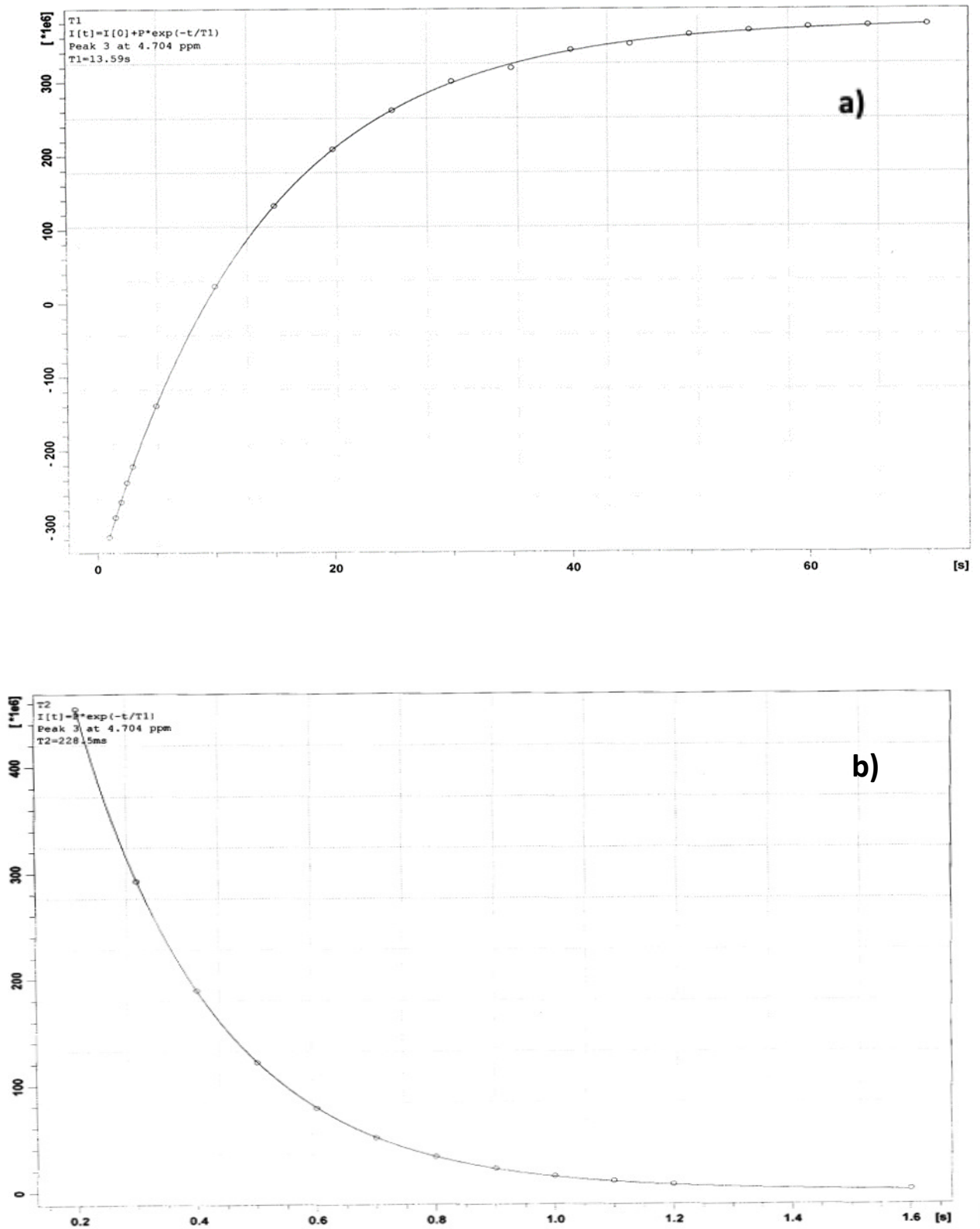


Figure 3: a) Inversion Recovery Spin-Lattice Relaxation Times (T_1), **b)** Spin Echo Spin-Spin Relaxation Times (T_2)

The measured T_1 and T_2 relaxation time values of HDO and doublet peaks in the spectrum shown in Fig. 2 for different temperatures are given in Table 1 and Table 2, respectively.

Table 1: T_1 and T_2 relaxation times of HDO peak in Fig. 2

T (K)	T_1 (s)	T_2 (ms)
298	7.815	118.6
303	7.991	111.0
308	8.486	109.8
313	8.885	108.6
318	8.994	107.3

Table 2: T_1 and T_2 relaxation times of Doublet peak in Fig. 2

T (K)	T_1 (s)	T_2 (ms)
298	1.728	114.2
303	2.063	147.4
308	2.297	193.4
313	2.573	261.4
318	2.591	340.1

As seen in Table 1 and Table 2, T_1 and T_2 values of Doublet peak and T_1 values of HDO peak increase depending on temperature, but T_2 values of HDO peak decrease depending on temperature. In addition, the case of $T_1 \gg T_2$ in both tables draws attention. This clearly shows that the dynamics of the water molecule are complex and that intramolecular movements are also possible in the system [8, 12].

Temperature dependence graphs of $\ln T_1$ and $\ln T_2$ values for both peaks in the spectrum, shown in Fig. 4 and Fig. 5, respectively, were created to find the activation energies and rotational correlation times.

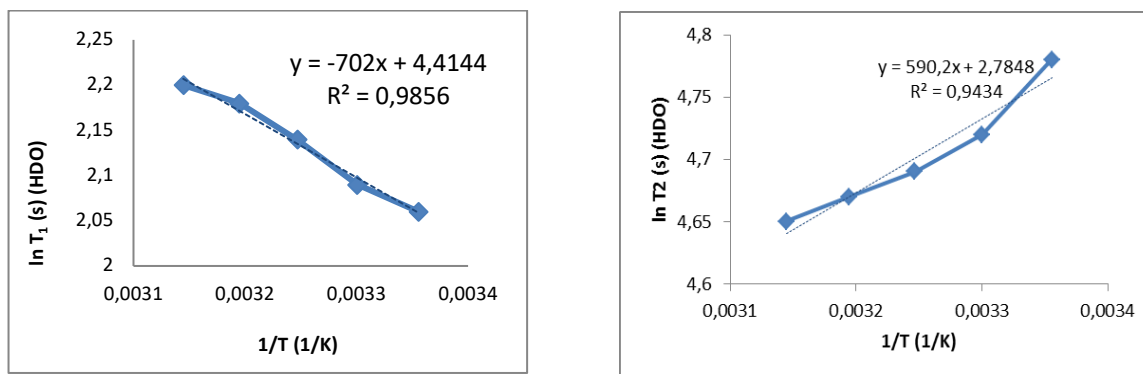


Figure 4: (a) $\ln T_1$ and (b) $\ln T_2$ vs. reciprocal temperature ($1/T$) for HDO peak measured at 400 MHz

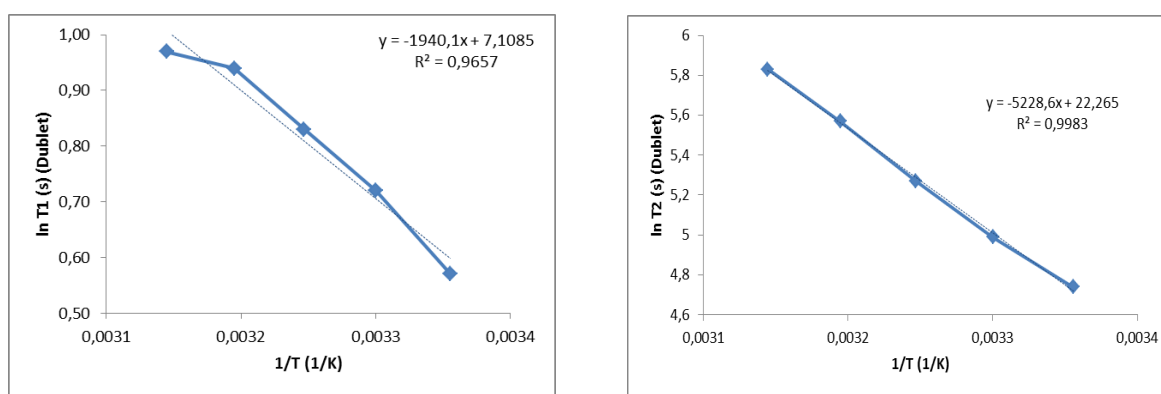


Figure 5: (a) $\ln T_1$ and (b) $\ln T_2$ vs. reciprocal temperature ($1/T$) for Doublet peak measured at 400 MHz

For the Doublet peak, both $\ln T_1$ and $\ln T_2$ decreased linearly with increasing $1/T$. For the HDO peak, while $\ln T_1$ decreased linearly with increasing $1/T$, $\ln T_2$ increased linearly. Therefore, in aqueous solutions, the T_1 and T_2 times are considered to be under high mobility. According to Eqn. (17), T_1 and T_2 relaxation times are inversely proportional to the rotational correlation time for molecular motion; According to Eqn. (19), it seems that T_1 and T_2 decrease as the temperature decreases.

In this case, the minimum T_1 relaxation appears to occur on the high-temperature side and τ_c provides the value $\omega \tau_c \ll 1$ of the correlation time characterizing the reorientation of the core. This approach should be applicable to some water molecules, or at least part of their dynamics. [11, 32, 37]. The intermolecular proton exchange affects the correlation time, and the spin-spin coupling causes the magnitude difference between T_1 and T_2 [32-36].

For the HDO peak, starting from the graph of T_1 and T_2 versus $1/T$ shown in Fig. 4, the following equations can be written, respectively.

$$\ln T_1 = -702 (1/T) + 4.4144 \quad R^2 = 0.9856 \quad (21)$$

$$\ln T_2 = 590.2 (1/T) + 2.7848 \quad R^2 = 0.9434 \quad (22)$$

Similarly, for the doublet peak, starting from the graph of T_1 and T_2 versus $1/T$ shown in Fig. 5, the following equations can be written, respectively.

$$\ln T_1 = -1940.1 (1/T) + 7.1085 \quad R^2 = 0.9657 \quad (23)$$

$$\ln T_2 = -5228.6 (1/T) + 22.265 \quad R^2 = 0.9983 \quad (24)$$

According to Eqn. (20), activation energies (E_a) were calculated from the slopes of the curves in Fig. 4 and Fig. 5. Eqns. (16) and (17) were used to calculate the rotational correlation times (τ_0 , τ_c). E_a and τ_0 , τ_c values for HDO peak and Doublet peak are given in Table 3 and Table 4, respectively, at 25 °C.

Table 3: E_a , and τ_0 , τ_c values of HDO peak at 25 °C

	T ₁	T ₂
E_a (kcal.mol ⁻¹)	1.39	1.17
τ_0 (s)	1.88×10^{-11}	1.038×10^{-9}
τ_c (s)	11.38×10^{-11}	7.49×10^{-9}

Table 4: E_a , and τ_0 , τ_c values of Doublet peak at 25 °C

	T ₁	T ₂
E_a (kcal.mol ⁻¹)	3.85	10.39
τ_0 (s)	8.697×10^{-10}	1.864×10^{-16}
τ_c (s)	5.15×10^{-10}	7.78×10^{-9}

5. Conclusion

For the HDO peak, the rotational correlation time and activation energy values were found as 1.14 ps, 1.39 kcal.mol⁻¹ for T₁ and 7.5 ns, 1.17 kcal.mol⁻¹ for T₂, respectively. For the Doublet peak, the rotational correlation time and activation energy values were found as 0.05 ps 3.85 kcal.mol⁻¹ for T₁ and 7.8 ns, 10.39 kcal.mol⁻¹ for T₂, respectively. As a result of the very fast molecular motion, the small τ_c value for T₁ was calculated and this result confirms the extreme narrowing condition mentioned above. However, the value of τ_c calculated for T₂ conforms to condition $\omega\tau_c > 1$. The reason for this situation can be considered as the effect of intermolecular proton exchange on shortening the correlation time and the spin-spin coupling that causes the difference in magnitude between T₁ and T₂. The calculated activation energies are compatible with the rotational activation energy of water given in the literature and are associated with the isotropic rotational motion of unbound water molecules. Consequently, the effective process for the spin-spin relaxation mechanism in an aqueous solution may be

chemical exchange and dipolar coupling between protons and deuterons, while spin-lattice relaxation may be caused by the molecular tumbling reorientation process. This indicates that two different molecular species exist in aqueous solutions with respect to reorientation dynamics.

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