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LP-Kenmotsu Manifolds Admitting Bach Almost Solitons

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Keywords: Bach almost solitons, LP- Kenmotsu manifolds, Perfect fluid, Weyl tensor 2010 AMS: 53C25, 53C44, 53C50 Keywords: Bach almost soliton (g, ζ, λ) , we explored the characteris Besides, we gave the necessary condition for $(LPK)_m$ $(m \ge 1)$	
2010 AMS: 53C25, 53C44, 53C50to be an η -Einstein manifold. Afterwards, we proved that steady when a Lorentzian para-Kenmotsu manifold of d soliton.Received: 27 February 2024steady when a Lorentzian para-Kenmotsu manifold of d soliton.Available online: 25 August 2024steady when a Lorentzian para-Kenmotsu manifold of d soliton.	cs of the norm of Ricci operator. 4) admitting Bach almost soliton Bach almost solitons are always

1. Introduction

In 1976, the concept of almost paracontact manifolds was proposed by Sato [1]. An almost paracontact structure on a semi-Riemannian manifold \mathcal{M} was established by Kaneyuki and Kozai in [2]. They created almost paracomplex shape on $\mathcal{M} \times R$. According to Kaneyuki et al. [3], the key variation among an almost paracontact manifold is the signature of metric. In 1995, the authors Sinha and Prasad described para-Kenmotsu as well as special para-Kenmotsu manifolds and found significant properties of para-Kenmotsu manifolds [4]. Afterwards, para-Kenmotsu manifolds drew huge attention and a number of mathematicians brought forward the significant characteristics of such manifolds [5–9].

Semi-Riemannian geometry, used in the relativity theory, was studied in [10]. About four decades ago, Kaigorodov has explored the curvature structure of the spacetime [11]. Raychaudhuri et al. [12] extended the above concepts of the general theory of spacetime. Recently, Haseeb and Rajendra introduced and studied the Lorentzian para-Kenmotsu manifolds [13, 14].

1921 was the year, when Bach initiated Bach tensor [15] to explore conformal geometry. He proved that the Bach tensor is a rank 2 trace-free tensor and is conformally invariant in dimension 4. So, in lieu of Hilbert-Einstein functional, the functional is taken in the following way

$$\mathscr{W}(g) = \int_{\mathscr{M}} \|\mathscr{C}\|_g^2 d\nu_g$$

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where, \mathcal{M} is a manifold of dimension-four and \mathcal{C} repersents the Weyl tensor of type (1,3) given by

$$\mathscr{C}(U,\mathscr{V})\mathscr{W} = \mathscr{R}(U,\mathscr{V})\mathscr{W} + \frac{1}{m-2}[\mathscr{S}(U,\mathscr{W})\mathscr{V} - \mathscr{S}(\mathscr{V},\mathscr{W})U + g(U,\mathscr{W})\mathscr{Q}\mathscr{V} - g(\mathscr{V},\mathscr{W})\mathscr{Q}U] - \frac{r}{(m-1)(m-2)}[g(U,\mathscr{W})\mathscr{V} - g(\mathscr{V},\mathscr{W})U],$$

$$(1.1)$$

here, \mathscr{R} represents the Riemannian curvature tensor, \mathscr{Q} is the Ricci operator and \mathscr{S} denotes the Ricci tensor, such that, $g(\mathscr{Q}U, \mathscr{V}) = \mathscr{S}(U, \mathscr{V})$, \forall differentiable vector fields $U, \mathscr{V}, \mathscr{W}$. Bach tensor of type (0,2) on a semi-Riemannian manifold (\mathscr{M}^m, g) of dimension $m(\geq 3)$ is given by

$$\mathscr{B}(U,\mathscr{V}) = \frac{1}{(m-3)} \sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j (\nabla_{\mathscr{E}_i} \nabla_{\mathscr{E}_j} \mathscr{C}') (U, \mathscr{E}_i, \mathscr{E}_j, \mathscr{V}) + \frac{1}{(m-2)} \sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j \mathscr{S}(\mathscr{E}_i, \mathscr{E}_j) \mathscr{C}' (U, \mathscr{E}_i, \mathscr{E}_j, \mathscr{V}),$$

$$(1.2)$$

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here, $g(\mathcal{E}_i, \mathcal{E}_i) = \varepsilon_i$, $g(\mathcal{C}(U, \mathcal{V})\mathcal{W}, \mathcal{Y}) = \mathcal{C}'(U, \mathcal{V}, \mathcal{W}, \mathcal{Y})$ and $\{\{\mathcal{E}_i\}_{i=1}^{m-1}, \mathcal{E}_m = \zeta\}$ is a local orthonormal frame at each point p of $T_p\mathcal{M}$. Relation (1.1), together with contracting Bianchi second identity, we obtain

$$div\mathscr{C} = \frac{(m-3)}{(m-2)}C_0,\tag{1.3}$$

where, C_0 is Cotton tensor [16] given by

$$C_0(U,\mathscr{V})\mathscr{W} = -(\nabla_{\mathscr{V}}\mathscr{S})(U,\mathscr{W}) + (\nabla_U\mathscr{S})(\mathscr{V},\mathscr{W}) + \frac{1}{2(m-1)}[(\mathscr{V}r)g(U,\mathscr{W}) - (Ur)g(\mathscr{V},\mathscr{W})].$$
(1.4)

In view of equation (1.3), together with equation (1.2), the Bach tensor takes the form,

$$\mathscr{B}(U,\mathscr{V}) = \frac{1}{(m-2)} \left[\sum_{i \in \{1,\dots,m\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i, U)\mathscr{V} + \sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j \mathscr{S}(\mathscr{E}_i, \mathscr{E}_j)\mathscr{C}'(U, \mathscr{E}_i, \mathscr{E}_j, \mathscr{V}) \right], \tag{1.5}$$

 \forall differentiable vector fields U, \mathscr{V} . For dimension three, the Weyl tensor vanishes. Therefore, Bach tensor given in equation (1.5) reduces to

$$\mathscr{B}(U,\mathscr{V}) = \sum_{i \in \{1,2,3\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i, U)\mathscr{V}.$$
(1.6)

For further study, the references [17–24] may be seen.

In 2012, Das and Kar [25] studied different characteristics of Bach flow on product manifolds and analysed their outcomes with the Ricci flow. Bach flow is suggested in [26] to specify the Harava-Lifschitz gravity in general relativity. In 2011, Bahuaud and Helliwell in [27] studied the presence of Bach flow for short time. Cao and Chen, in the year 2013, explored Bach flat Ricci solitons [28]. Subsequently, Ho [29] worked comprehensively on the solitons of Bach flow. He also studied the Bach flows on Lie group of dimension 4. In 2020, Helliwell specified Bach flow of dimension 4 on locally homogeneous product manifolds [30]. In recent times, Ghosh [31] investigated the Bach almost solitons (g, ζ, λ) in semi-Riemannian geometry and is given by

$$(\pounds_{\mathscr{X}}g + 2\mathscr{B} - 2\lambda g)(U, \mathscr{V}) = 0, \tag{1.7}$$

here, $\pounds_{\mathscr{X}}$ is the Lie derivative operator along \mathscr{X} ; \mathscr{X} is a potential vector field and $\lambda \in C^{\infty}(\mathscr{M}^m)$. The Bach almost solitons (g, ζ, λ) is said to be expanding, steady and shrinking according to $\lambda < 0$, $\lambda = 0$ and $\lambda > 0$, respectively.

This article is organized in the following manner: Section 1 contains introduction, based on development of almost paracontact manifold and other concepts. Preliminaries are given in Section 2, based on $(LPK)_m$. Section 3 contains the work on (g, ζ, λ) in $(LPK)_m$. In Section 4, we examine $(LPK)_m$ of dimension 3, which admits Bach almost solitons.

2. Preliminaries

An *m*-dimensional smooth manifold \mathcal{M}^m is called Lorentzian almost paracontact manifold, if it is equipped with a (1,1)-tensor field ϕ , a contravariant vector field ζ , a 1-form η and a Lorentzian metric *g* of type (0, 2). The following relations for an *m*-dimensional Lorentzian metric manifold hold [32],

$$\phi^2(U) = U + \eta(U)\zeta, \ \eta(\zeta) + 1 = 0, \tag{2.1}$$

$$g(U,\zeta) = \eta(U), \ g(\phi U, \phi \mathscr{V}) = \eta(U)\eta(\mathscr{V}) + g(U,\mathscr{V}), \tag{2.2}$$

 $\forall U, \mathscr{V} \text{ on } \mathscr{M}^m$, and the structure (ϕ, ζ, η, g) is named the Lorentzian almost paracontact structure. An \mathscr{M}^m endowed with (ϕ, ζ, η, g) is known as Lorentzian almost paracontact manifold and holding the following results:

$$\phi \zeta = 0, \ \eta(\phi U) = 0, \ \Omega(U, \mathscr{V}) = \Omega(\mathscr{V}, U), \tag{2.3}$$

here, $\Omega(U, \mathscr{V}) = g(U, \varphi \mathscr{V}).$

Definition 2.1. A Lorentzian almost paracontact manifold \mathcal{M}^m is known as $(LPK)_m$ if

$$(\nabla_U \phi)(\mathscr{V}) = -\eta(\mathscr{V})\phi U - g(\phi U, \mathscr{V})\zeta,$$

 $\forall U and \mathscr{V} on \mathscr{M}^m$.

Further, for $(LPK)_m$, following results hold good:

$$\nabla_U \zeta + U + \eta(U) \zeta = 0, \tag{2.4}$$

$$(\nabla_U \eta)(\mathscr{V}) + g(U, \mathscr{V}) + \eta(U)\eta(\mathscr{V}) = 0, \tag{2.5}$$

$$\mathscr{R}(U,\mathscr{V})\zeta = \eta(\mathscr{V})U - \eta(U)\mathscr{V},\tag{2.6}$$

$$\mathscr{R}(\zeta,\mathscr{V})U = g(U,\mathscr{V})\zeta - \eta(U)\mathscr{V},\tag{2.7}$$

$$\mathscr{R}(\zeta, U)\zeta = U + \eta(U)\zeta, \tag{2.8}$$

$$\mathscr{S}(U,\zeta) = (m-1)\eta(U), \tag{2.9}$$

$$\mathscr{Q}\zeta = (m-1)\zeta,\tag{2.10}$$

$$\mathscr{S}(\phi \mathscr{V}, \phi U) = \mathscr{S}(\mathscr{V}, U) + (m-1)\eta(\mathscr{V})\eta(U), \tag{2.11}$$

 $\forall U, \mathcal{V}, \mathcal{W}$ on $(LPK)_m$ [33, 34]. In the above results, ∇ represents the covariant differentiation operator w.r.t. g in semi-Riemannian manifolds.

Proposition 2.2. We assume \mathcal{M} to be an $(LPK)_m$. Subsequently, we have

$$\mathscr{S}(\phi U, \mathscr{V}) = \mathscr{S}(U, \phi \mathscr{V}), \tag{2.12}$$

 $\forall U, \mathscr{V} on (LPK)_m.$

Proof. Setting ϕU for U in (2.11), we get,

$$\mathscr{S}(\phi^2 U, \phi \mathscr{V}) = \mathscr{S}(\phi U, \mathscr{V}) + (m-1)\eta(\phi U)\eta(\mathscr{V}).$$

Using equations (2.1) and (2.3) in the foregoing equation, we yield

$$\mathscr{S}(U+\eta(U)\zeta,\phi\mathscr{V}) = \mathscr{S}(\phi U,\mathscr{V}).$$
(2.13)

From equation (2.13), the Proposition 2.2 follows.

3. Bach Almost Solitons and $(LPK)_m$

Definition 3.1. A semi-Riemannnian manifold is called Bach perfect fluid if Bach almost tensor is given by

$$\mathscr{B}(U,\mathscr{V}) = \beta \eta(U) \eta(\mathscr{V}) + \alpha g(U,\mathscr{V}), \quad \forall \mathscr{V}, U,$$

where, α and β are scalars.

Let $(LPK)_m$ admit (g, ζ, λ) . Then (1.7) holds and thus, we have

$$(\pounds_{\zeta}g)(U,\mathscr{V}) + 2\mathscr{B}(U,\mathscr{V}) = 2\lambda g(U,\mathscr{V}).$$
(3.1)

As we have

$$(\pounds_{\zeta}g)(U,\mathscr{V}) = g(\nabla_{U}\zeta,\mathscr{V}) + g(U,\nabla_{\mathscr{V}}\zeta).$$
(3.2)

The result (2.4), together with (3.2) yields

$$(\pounds_{\zeta}g)(U,\mathscr{V}) + 2[g(U,\mathscr{V}) + \eta(U)\eta(\mathscr{V})] = 0.$$
(3.3)

Putting the preceding result (3.3) in (3.1), we lead to

$$\mathscr{B}(\mathscr{V},U) = (1+\lambda)g(\mathscr{V},U) + \eta(\mathscr{V})\eta(U).$$
(3.4)

Result (3.4) shows the succeeding proposition:

Proposition 3.2. An $(LPK)_m$ admitting a Bach almost soliton (g, ζ, λ) is Bach perfect fluid.

Replacing \mathscr{W} by ζ in (1.1), we have

$$\mathscr{C}(U,\mathscr{V})\zeta = \mathscr{R}(U,\mathscr{V})\zeta + \frac{1}{(m-2)}[\mathscr{S}(U,\zeta)\mathscr{V} - \mathscr{S}(\mathscr{V},\zeta)U + g(U,\zeta)\mathscr{2}\mathscr{V} - g(\mathscr{V},\zeta)\mathscr{2}U] - \frac{r}{(m-1)(m-2)}[g(U,\zeta)\mathscr{V} - g(\mathscr{V},\zeta)U],$$
(3.5)

 \forall differentiable vector fields U, \mathscr{V} . Operating \mathscr{Q} in (3.5) and using relations (2.2), (2.6), (2.7) and (2.10), we get

$$\mathscr{Q}(\mathscr{C}(U,\mathscr{V})\zeta) = \frac{(r-m+1)}{(m-1)(m-2)} [-\eta(U)\mathscr{Q}\mathscr{V} + \eta(\mathscr{V})\mathscr{Q}U] - \frac{1}{(m-2)} [\eta(\mathscr{V})\mathscr{Q}^2U - \eta(U)\mathscr{Q}^2\mathscr{V}].$$
(3.6)

The inner product of (3.6) with \mathscr{X} leads to

$$g(\mathscr{Q}(\mathscr{C}(U,\mathscr{V})\zeta),\mathscr{X}) = \frac{(r-m+1)}{(m-1)(m-2)} [\eta(\mathscr{V})g(\mathscr{Q}U,\mathscr{X}) - \eta(U)g(\mathscr{Q}\mathscr{V},\mathscr{X})] - \frac{1}{(m-2)} [\eta(\mathscr{V})g(\mathscr{Q}^{2}U,\mathscr{X}) - \eta(U)g(\mathscr{Q}^{2}\mathscr{V},\mathscr{X})].$$

$$(3.7)$$

Let $\{\{\mathscr{E}_i\}_{i=1}^{m-1}, \mathscr{E}_m = \zeta\}$ be an orthonormal frame at each point p of $T_p\mathscr{M}$. Now, setting $\mathscr{V} = \mathscr{X} = \mathscr{E}_i$ in (3.7) with summation i = 1 to m and on evaluation, we get

$$\sum_{i \in \{1,...,m\}} \varepsilon_i g(\mathscr{Q}(\mathscr{C}(U,\mathscr{E}_i)\zeta),\mathscr{E}_i) = -\frac{(r-m+1)^2}{(m-1)(m-2)} \eta(U) + \frac{1}{(m-2)} [|\mathscr{Q}|^2 - (m-1)^2] \eta(U).$$
(3.8)

Setting ζ in place of \mathscr{W} in relation (1.4) gives

$$C_0(U,\mathscr{V})\zeta = g((\nabla_U\mathscr{Q})\mathscr{V},\zeta) - g((\nabla_{\mathscr{V}}\mathscr{Q})U,\zeta) - \frac{1}{2(m-1)}[U(r)\eta(\mathscr{V}) - \mathscr{V}(r)\eta(U)].$$
(3.9)

From equation (2.12), we have the relation

$$\phi \mathcal{Q}U = \mathcal{Q}\phi U. \tag{3.10}$$

From the equation (3.10), we also have

$$g((\nabla_U \mathscr{Q})\mathscr{V}, \zeta) = g(\mathscr{Q}U, \mathscr{V}) - (m-1)g(U, \mathscr{V}).$$
(3.11)

Applying above equation (3.11) in (3.9), it gives

$$C_0(U,\mathscr{V})\zeta = -\frac{1}{2(m-1)}[U(r)\eta(\mathscr{V}) - \mathscr{V}(r)\eta(U)].$$
(3.12)

After differentiating covariantly the above relation w.r.t. \mathcal{W} and using the relation (2.5), we obtain

$$(\nabla_{\mathscr{W}}C_{0})(U,\mathscr{V})\zeta = -(\nabla_{\mathscr{V}}\mathscr{S})(U,\mathscr{W}) + (\nabla_{U}\mathscr{S})(\mathscr{V},\mathscr{W}) - \frac{1}{2(m-1)}[g(\nabla_{\mathscr{W}}\mathscr{D}r,U)\eta(\mathscr{V}) - g(\nabla_{\mathscr{W}}\mathscr{D}r,\mathscr{V})\eta(U)],$$

$$(3.13)$$

here \mathscr{D} represents the gradient operator. Let $\{\{\mathscr{E}_i\}_{i=1}^{m-1}, \mathscr{E}_m = \zeta\}$ be the orthonormal frame at each point p of $T_p \mathscr{M}$. Replacing $U = \mathscr{W} = \mathscr{E}_i$ with summation over i = 1 to m in equation (3.13), this gives

$$\sum_{i \in \{1,\dots,m\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i,\mathscr{V})\zeta = -\frac{1}{2(m-1)} [(div\mathscr{D}r)\eta(\mathscr{V}) - g(\nabla_{\zeta}\mathscr{D}r,\mathscr{V})] - \frac{\mathscr{V}(r)}{2}.$$
(3.14)

Now, by rewriting the equation (1.5), we have

$$\mathscr{B}(U,\mathscr{V}) = \frac{1}{(m-2)} \left[\sum_{i \in \{1,\dots,m\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i, U)\mathscr{V} + \sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j \mathscr{S}(\mathscr{E}_i, \mathscr{E}_j)\mathscr{C}'(U, \mathscr{E}_i, \mathscr{E}_j, \mathscr{V}) \right].$$
(3.15)

After evaluation, the second term of the above equation takes the form

$$\begin{split} \sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j \mathscr{S}(\mathscr{E}_i,\mathscr{E}_j) \mathscr{C}'(U,\mathscr{E}_i,\mathscr{E}_j,\mathscr{V}) &= -\sum_{i \in \{1,\dots,m\}} \sum_{j \in \{1,\dots,m\}} \varepsilon_i \varepsilon_j g(\mathscr{Q}\mathscr{E}_i,\mathscr{E}_j) g(\mathscr{C}(U,\mathscr{E}_i)\mathscr{V},\mathscr{E}_j), \\ &= -\sum_{i \in \{1,\dots,m\}} \varepsilon_i g(\mathscr{Q}(\mathscr{C}(U,\mathscr{E}_i)\mathscr{V}),\mathscr{E}_i). \end{split}$$

Taking the above equation and equation (3.15) together, we obtain

$$\mathscr{B}(U,\mathscr{V}) = \frac{1}{(m-2)} \left[\sum_{i \in \{1,\dots,m\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i, U) \mathscr{V} - \sum_{i \in \{1,\dots,m\}} \varepsilon_i g(\mathscr{Q}(\mathscr{C}(U, \mathscr{E}_i) \mathscr{V}), \mathscr{E}_i)) \right].$$
(3.16)

Replacing \mathscr{V} for ζ in the above relation (3.16), it gives

$$\mathscr{B}(U,\zeta) = \frac{1}{(m-2)} \left[\sum_{i \in \{1,\dots,m\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_0)(\mathscr{E}_i, U)\zeta - \sum_{i \in \{1,\dots,m\}} \varepsilon_i g(\mathscr{Q}(\mathscr{C}(U,\mathscr{E}_i)\zeta), \mathscr{E}_i)].$$
(3.17)

Equations (3.8), (3.14) and (3.17) taken together give

$$\mathscr{B}(U,\zeta) = \frac{1}{(m-2)} \left[-\frac{U(r)}{2} - \frac{1}{2(m-1)} \left\{ (div\mathscr{D}r)\eta(U) - g(\nabla_{\zeta}\mathscr{D}r, U) \right\} + \frac{(r-m+1)^2}{(m-1)(m-2)} \eta(U) - \frac{1}{(m-2)} \left\{ |\mathscr{D}|^2 - (m-1)^2 \right\} \eta(U) \right].$$
(3.18)

Setting \mathscr{V} for ζ in equation (3.4), we get

$$\mathscr{B}(U,\zeta) = \lambda \eta(U). \tag{3.19}$$

Relation (3.18) and (3.19), taken together give

$$\lambda \eta(U) = \frac{1}{(m-2)} \left[-\frac{U(r)}{2} - \frac{1}{2(m-1)} \left\{ (div \mathscr{D}r) \eta(U) - g(\nabla_{\zeta} \mathscr{D}r, U) \right\} + \frac{(r-m+1)^2}{(m-1)(m-2)} \eta(U) - \frac{1}{(m-2)} \left\{ |\mathscr{Q}|^2 - (m-1)^2 \right\} \eta(U) \right].$$
(3.20)

Setting U for ϕU in relation (3.20), we obtain

$$\frac{1}{(m-2)} \left[-\frac{\phi U(r)}{2} + \frac{1}{2(m-1)} g(\nabla_{\zeta} \mathscr{D} r, \phi U) \right] = 0$$

This implies that

 $g(\nabla_{\zeta} \mathscr{D}r, \phi U) = (m-1)g(\mathscr{D}r, \phi U).$

This gives

$$\phi \nabla_{\zeta} \mathscr{D} r = (m-1)\phi \mathscr{D} r. \tag{3.21}$$

Taking covariant differentiation of equation (2.10) w.r.t. U and using the relations (2.3) and (2.4), we get

$$(\nabla_U \mathscr{D})\zeta = \mathscr{D}U - (m-1)U. \tag{3.22}$$

Contracting the preceding equation w.r.t. U, we have

$$\sum_{i \in \{1, \dots, m\}} \varepsilon_i g(\nabla_{\mathscr{E}_i} \mathscr{Q}) \zeta, \mathscr{E}_i) = \sum_{i=1}^m \varepsilon_i [g(\mathscr{Q} \mathscr{E}_i, \mathscr{E}_i) - (m-1)g(\mathscr{E}_i, \mathscr{E}_i)].$$

or,

$$(div\mathscr{Q})\zeta = r - (m-1)m,$$

or,

$$\zeta(r) = 2[r - m(m-1)], \tag{3.23}$$

which can be written as

 $\pounds_{\zeta} r = 2r - 2m(m-1).$

Applying the exterior derivative in the above relation, we have

 $d\pounds_{\zeta}r = 2dr.$

Since, d and the Lie derivative commutes, therefore, we have

$$\pounds_{\zeta} dr = 2dr.$$

Writing the above relation in the form of gradient operator, we have

 $\pounds_{\mathcal{L}} \mathscr{D}r = 2\mathscr{D}r,$

or,

$$\nabla_{\zeta} \mathscr{D}r - \nabla_{\mathscr{D}r} \zeta = 2 \mathscr{D}r.$$

Using the relation (2.4) in the above relation, we lead to

$$\nabla_{\zeta} \mathscr{D}r = \mathscr{D}r - \zeta(r)\zeta. \tag{3.24}$$

Applying ϕ in the above relation (3.24) and using the relations in (2.3) and (3.21), we get

 $\phi \mathscr{D} r = 0.$

This implies

$$\mathscr{D}r = -\zeta(r)\zeta. \tag{3.25}$$

Differentiating (3.25) covariantly w.r.t. \mathscr{X} , it yields

$$\nabla_{\mathscr{X}}\mathscr{D}r = -[g(\nabla_{\mathscr{X}}\mathscr{D}r,\zeta)\zeta - g(\mathscr{D}r,\mathscr{X})\zeta - g(\mathscr{D}r,\zeta)\mathscr{X} - 2g(\mathscr{D}r,\zeta)\eta(\mathscr{X})\zeta],$$
(3.26)

which by contracting over \mathscr{X} gives

$$(div\mathscr{D}r) = (m-3)\zeta(r).$$
(3.27)

Relations (3.24) and (3.25) give

$$\nabla_{\zeta} \mathscr{D} r = -2\zeta(r)\zeta. \tag{3.28}$$

Using relations (3.25), (3.27) and (3.28) in (3.20), we obtain

$$\lambda \eta(U) = \frac{1}{(m-2)} \left[\frac{\zeta(r)}{2} \eta(U) - \frac{1}{2(m-1)} \{ (m-3)\zeta(r)\eta(U) + 2\zeta(r)\eta(U) \} + \frac{(r+1-m)^2}{(m-1)(m-2)} \eta(U) - \frac{1}{(m-2)} \{ |\mathcal{Q}|^2 - (m-1)^2 \} \eta(U) \right].$$
(3.29)

On simplification, relation (3.29) gives

$$\lambda = \frac{1}{(m-2)^2} \left[\frac{(r+1-m)^2}{(m-1)} + (m-1)^2 - |\mathcal{Q}|^2 \right].$$
(3.30)

In the light of the relation (3.30), succeeding theorem holds:

Theorem 3.3. The Bach almost solitons (g, ζ, λ) on an $(LPK)_m$ are expanding, steady and shrinking according as

$$[\frac{(r+1-m)^2}{(m-1)} + (m-1)^2] > |\mathcal{Q}|^2, \ [\frac{(r+1-m)^2}{(m-1)} + (m-1)^2] = |\mathcal{Q}|^2 \ and \ [\frac{(r+1-m)^2}{(m-1)} + (m-1)^2] < |\mathcal{Q}|^2.$$

Consider a Lorentzian para-Kenmotsu space form of *m*-dimension. Then by relation (3.23), we have r = m(m-1). Hence,

$$\lambda = \frac{1}{(m-2)^2} [m(m-1)^2 - |\mathcal{Q}|^2].$$

The above relation leads the following corollary:

Corollary 3.4. The Bach almost solitons (g, ζ, λ) on an LP-Kenmotsu space form of dimension *m* is expanding, steady and shrinking according as $m(m-1)^2 > |\mathcal{Q}|^2$, $m(m-1)^2 = |\mathcal{Q}|^2$ and $m(m-1)^2 < |\mathcal{Q}|^2$.

Definition 3.5. An $(LPK)_m$ is named η -Einstein if its \mathscr{S} satisfies [35]

$$\mathscr{S}(\mathscr{V},U) = ag(\mathscr{V},U) + b\eta(\mathscr{V})\eta(U),$$

 $\forall \mathscr{V}, U$, where, a and b are scalars.

Now, replacing $U = \zeta$ in relation (3.13), we have

$$(\nabla_{\mathscr{W}}C_{0})(\zeta,\mathscr{V})\zeta = -(\nabla_{\mathscr{V}}\mathscr{S})(\zeta,\mathscr{W}) + (\nabla_{\zeta}\mathscr{S})(\mathscr{V},\mathscr{W}) - \frac{1}{2(m-1)}[g(\nabla_{\mathscr{W}}\mathscr{D}r,\zeta)\eta(\mathscr{V}) - g(\nabla_{\mathscr{W}}\mathscr{D}r,\mathscr{V})\eta(\zeta)].$$
(3.31)

Taking the inner product of relation (3.26) with \mathscr{V} and replacing \mathscr{X} by \mathscr{W} , we obtain

$$g(\nabla_{\mathscr{W}}\mathscr{D}r,\mathscr{V}) = -[g(\nabla_{\mathscr{W}}\mathscr{D}r,\zeta)\eta(\mathscr{V}) - g(\mathscr{D}r,\mathscr{W})\eta(\mathscr{V}) - g(\mathscr{D}r,\zeta)g(\mathscr{V},\mathscr{W}) - 2g(\mathscr{D}r,\zeta)\eta(\mathscr{V})\eta(\mathscr{W})].$$
(3.32)

The relations (3.22), (3.31) and (3.32) give

$$(\nabla_{\mathscr{W}}C_0)(\zeta,\mathscr{V})\zeta = g((\nabla_{\zeta}\mathscr{Q})\mathscr{V},\mathscr{W}) - g((\nabla_{\mathscr{V}}\mathscr{Q})\zeta,\mathscr{W}) - \frac{\zeta(r)}{2(m-1)}[g(\mathscr{V},\mathscr{W}) + \eta(\mathscr{V})\eta(\mathscr{W})].$$
(3.33)

In an $(LPK)_m$, the following result holds (for perusal, see [36])

$$(\nabla_{\mathcal{C}}\mathcal{Q})\mathcal{V} = 2\mathcal{Q}\mathcal{V} - 2(m-1)\mathcal{V}.$$
(3.34)

Applying relations (3.22) and (3.34) into (3.33), it yields

$$(\nabla_{\mathscr{W}}C_{0})(\zeta,\mathscr{V})\zeta = g(\mathscr{QV},\mathscr{W}) - (m-1)g(\mathscr{V},\mathscr{W}) - \frac{\zeta(r)}{2(m-1)}[g(\mathscr{V},\mathscr{W}) + \eta(\mathscr{V})\eta(\mathscr{W})].$$
(3.35)

If $(\nabla_{\mathscr{W}} C_0)(\zeta, \mathscr{V})\zeta = 0$ and (3.23), then (3.35) leads to

$$\mathscr{S}(\mathscr{V},\mathscr{W}) = \left(\frac{r}{m-1} - 1\right)g(\mathscr{V},\mathscr{W}) + \left(\frac{r}{m-1} - m\right)\eta(\mathscr{V})\eta(\mathscr{W}).$$
(3.36)

The relation (3.36) leads the following theorem:

Theorem 3.6. An $(LPK)_m$ $(m \ge 4)$ admitting (g, ζ, λ) is an η -Einstein manifold provided $(\nabla_{\mathscr{W}}C_0)(\zeta, \mathscr{V})\zeta = 0, \forall \mathscr{V}, \mathscr{W}$.

4. 3-Dimensional Bach Perfect Fluid Lorentzian Para-Kenmotsu Manifold

We consider an $(LPK)_3$ admitting (g, ζ, λ) . Curvature tensor of Riemannian manifold in dimension 3 states

$$\mathscr{R}(U,\mathscr{V})\mathscr{W} = -\mathscr{S}(U,\mathscr{W})\mathscr{V} + \mathscr{S}(\mathscr{V},\mathscr{W})U - g(U,\mathscr{W})\mathscr{2}\mathscr{V} + g(\mathscr{V},\mathscr{W})\mathscr{2}U - \frac{r}{2}[g(U,\mathscr{W})\mathscr{V} - g(\mathscr{V},\mathscr{W})U], \tag{4.1}$$

 \forall differentiable vector fields U, \mathscr{V} and \mathscr{W} .

Replacing $U = \mathcal{W} = \zeta$ in (4.1) and using (2.1), (2.8), (2.9) and (2.10), we obtain

$$\mathscr{QV} = \left(\frac{r}{2} - 3\right)\eta(\mathscr{V})\zeta + \left(\frac{r}{2} - 1\right)\mathscr{V}.$$
(4.2)

The preceding result gives

 $\mathcal{Q}\phi = \phi \mathcal{Q}.$

The equation (4.2), together with (2.4), gives

$$(\nabla_U \mathscr{Q})\zeta = \mathscr{Q}U - 2U. \tag{4.3}$$

Equation (3.12), together with (4.3) leads to

$$C_0(U,\mathscr{V})\zeta = \frac{1}{4}[\mathscr{V}(r)\eta(U) - U(r)\eta(\mathscr{V})].$$

The covariant differentiation of above result w.r.t. W yields

$$(\nabla_{\mathscr{W}}C_{0})(U,\mathscr{V})\zeta = -(\nabla_{\mathscr{V}}\mathscr{S})(U,\mathscr{W}) - (\nabla_{U}\mathscr{S})(\mathscr{V},\mathscr{W}) - \frac{1}{4}[g(\nabla_{\mathscr{W}}\mathscr{D}r,U)\eta(\mathscr{V}) - g(\nabla_{\mathscr{W}}\mathscr{D}r,\mathscr{V})\eta(U)]$$

Putting $\mathcal{W} = U = \mathcal{E}_i$ and taking sum over i = 1, 2, 3 in above relation, where $\{\mathcal{E}_1, \mathcal{E}_2, \mathcal{E}_3 = \zeta\}$ is orthonormal frame at each point p of $T_p \mathcal{M}$, we have

$$\sum_{e \in \{1,2,3\}} \varepsilon_i(\nabla_{\mathscr{E}_i} C_o)(\mathscr{E}_i,\mathscr{V})\zeta = -\frac{\mathscr{V}(r)}{2} - \frac{1}{4} [(div\mathscr{D}r)\eta(\mathscr{V}) - g(\nabla_{\zeta}\mathscr{D}r,\mathscr{V})].$$

$$\tag{4.4}$$

Taking $\mathscr{V} = \zeta$ in (1.6), we have

$$\mathscr{B}(U,\zeta) = \sum_{i \in \{1,2,3\}} \varepsilon_i(\nabla_{\mathscr{C}_i} C_0)(\mathscr{E}_i, U)\zeta.$$
(4.5)

Equations (3.4), (4.4) and (4.5) taken together give

$$\lambda \eta(U) = -\frac{1}{2}g(\mathscr{D}r, U) - \frac{1}{4}[(div\mathscr{D}r)\eta(U) - g(\nabla_{\zeta}\mathscr{D}r, U)].$$
(4.6)

Replacing ϕU for U in (4.6), we get

 $\phi \nabla_{\zeta} \mathscr{D} r = 2\phi \mathscr{D} r. \tag{4.7}$

We have the relation (3.23) and (3.24), for m = 3, which yields

$$\nabla_{\zeta} \mathscr{D}r = \mathscr{D}r - 2(r-6)\zeta. \tag{4.8}$$

The relations (4.7) and (4.8) provide

$$\mathscr{D}r = -2(r-6)\zeta. \tag{4.9}$$

By the covariant diffentiation of (4.9) w.r.t. $\mathscr X$ yields

 $\nabla_{\mathscr{X}}\mathscr{D}r = -2g(\mathscr{D}r,\mathscr{X})\zeta + 2(r-6)\mathscr{X} + 2(r-6)\eta(\mathscr{X})\zeta.$ (4.10)

By contracting the relation (4.10) over \mathscr{X} , we get

$$(div\mathscr{D}r) = 0. \tag{4.11}$$

Using relations (4.8), (4.9) and (4.11) in (4.6), it yields

$$\lambda = 0. \tag{4.12}$$

With the help of (4.12), the relation (3.4) reduces to

$$\mathscr{B}(U,\mathscr{V}) = \eta(U)\eta(\mathscr{V}) + g(U,\mathscr{V}).$$

The above results imply the succeeding theorem:

Theorem 4.1. Let $(LPK)_3$ admit a (g, ζ, λ) , then the manifold is a Bach perfect fluid and (g, ζ, λ) is always steady.

5. Example

We assume a manifold $\mathcal{M}^3 = \{(u_1, v_1, w_1) \in \mathbb{R}^3 : w_1 > 0\}$, here (u_1, v_1, w_1) are the general coordinates in \mathbb{R}^3 . Consider $\hat{\mathcal{E}}_1, \hat{\mathcal{E}}_2, \hat{\mathcal{E}}_3$, the vector fields on \mathcal{M}^3 given as

$$\hat{\mathscr{E}}_1 = w_1 \frac{\partial}{\partial u_1}, \qquad \hat{\mathscr{E}}_2 = w_1 \frac{\partial}{\partial v_1}, \qquad \hat{\mathscr{E}}_3 = w_1 \frac{\partial}{\partial w_1} = \zeta$$

and are linearly independent at each point of \mathcal{M}^3 . This implies

$$g(\hat{\mathscr{E}}_{i},\hat{\mathscr{E}}_{j}) = \begin{cases} 0, & 1 \le i \ne j \le 3, \\ -1, & i = j = 1, 2, \\ 0, & otherwise. \end{cases}$$

Suppose that η is 1-form on \mathscr{M}^3 given by $\eta(U) = g(U, \hat{\mathscr{E}}_3) = g(U, \zeta), \forall U \in \chi(\mathscr{M}^3)$. Again, assume that ϕ is (1,1) tensor field on \mathscr{M}^3 given below:

$$\phi \hat{\mathscr{E}}_1 = -\hat{\mathscr{E}}_2, \qquad \phi \hat{\mathscr{E}}_2 = -\hat{\mathscr{E}}_1, \qquad \phi \hat{\mathscr{E}}_3 = 0.$$

The linear property of g and ϕ give the following relations

$$\eta(\zeta) = g(\zeta,\zeta) = -1, \phi^2 = U + \eta(U)\zeta, \ g(U,\zeta) = \eta(U), \ \eta(\phi U) = 0, \ g(\phi U, \phi \mathscr{V}) = \eta(U)\eta(\mathscr{V}) + g(U, \mathscr{V}).$$

Assuming ∇ to be Levi-Civita connection w.r.t. Lorentzian metric *g*, then

 $[\hat{\mathscr{E}}_2,\hat{\mathscr{E}}_1]=0,\; [\hat{\mathscr{E}}_3,\hat{\mathscr{E}}_1]=\hat{\mathscr{E}}_1,\; [\hat{\mathscr{E}}_3,\hat{\mathscr{E}}_2]=\hat{\mathscr{E}}_2.$

Applying Koszul's formula, we can comfortably obtain

$$\nabla_{\hat{\mathscr{E}}_{i}}\hat{\mathscr{E}}_{j} = \begin{cases} -\hat{\mathscr{E}}_{3}, & i = j = 1, 2, \\ -\hat{\mathscr{E}}_{i}, & i = 1, 2, j = 3, \\ 0, & otherwise \end{cases}$$
(5.1)

Let $U \in \chi(\mathcal{M}^3)$, then the following relations can also be verified

$$abla_U \zeta + U + \eta(U) \zeta = 0, \ (
abla_U \phi) \mathscr{V} = -g(\phi U, \mathscr{V}) \zeta - \eta(\mathscr{V}) \phi(U).$$

For $U, \mathcal{V}, \mathcal{W} \in \chi(\mathcal{M}^3)$. Equation (5.1) helps to get the following non-vanishing values:

$$\begin{cases} \mathscr{R}(\hat{\mathcal{E}}_{1},\hat{\mathcal{E}}_{2})\hat{\mathcal{E}}_{1} = -\hat{\mathcal{E}}_{2}, \ \mathscr{R}(\hat{\mathcal{E}}_{1},\hat{\mathcal{E}}_{3})\hat{\mathcal{E}}_{1} = -\hat{\mathcal{E}}_{3}, \ \mathscr{R}(\hat{\mathcal{E}}_{1},\hat{\mathcal{E}}_{2})\hat{\mathcal{E}}_{2} = \hat{\mathcal{E}}_{1} \\ \mathscr{R}(\hat{\mathcal{E}}_{2},\hat{\mathcal{E}}_{3})\hat{\mathcal{E}}_{2} = -\hat{\mathcal{E}}_{3}, \ \mathscr{R}(\hat{\mathcal{E}}_{2},\hat{\mathcal{E}}_{3})\hat{\mathcal{E}}_{3} = -\hat{\mathcal{E}}_{2}. \end{cases}$$

The above results help to verify

$$\mathscr{R}(U,\mathscr{V})\mathscr{W} = -g(U,\mathscr{W})\mathscr{V} + g(\mathscr{V},\mathscr{W})U.$$

Hence, \mathcal{M}^3 is a Lorentzian para-Kenmotsu manifold of constant curvature. By contracting (5.2) over W, we obtain

$$\mathscr{S}(U,\mathscr{V}) = 2g(\mathscr{V},\mathscr{W}).$$

This implies

r = 6.

Then, (4.6) provides $\lambda = 0$. Hence, in this manifold, the Bach almost solitons are steady.

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(5.2)

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