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Evaluation of Structural and Dielectric Properties of Eu3+, B3+ co-doped Ba2GdMO⁶ (M=Nb, Ta) Double Perovskite Ceramics

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Abstract: In the study, the structural and dielectric properties of Ba₂GdMO₆ (M=Nb, Ta) double perovskite ceramics produced with solid-state method were examined by co-doping $xEu3+$ and $yB3+$ ($x=10$ mol%, $y=0$, 5, 15, 30, 50, 70 and 100 mol%). XRD (X-ray diffraction) results of the ceramic samples exhibited a singlephase structure with cubic symmetry *Fm*-3*m* space group, while increasing B3+ concentration led to an increase in crystallite sizes and lattice parameters up to 50 mol% in both series. SEM (scanning electron microscopy) examinations revealed the presence of boron-supported grain growth and agglomeration in the grains of both series, and also a slight angularity occurred in grain shape at high B^{3+} concentrations. The dielectric constant (ε') of the ceramic samples in both series increased with increasing boron concentration up to 50 mol%, and it was approximately 33.5 and 35.4 at 20 Hz for the Ba₂Gd_{1-x}NbO₆:xEu³⁺, yB³⁺ and Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+, respectively. The decrease in the dielectric constant after 50 mol% may be attributed to the presence of increased strain in the structure, as shown by the decrease in crystallite size. The increasing B^{3+} concentration caused a decrease in dielectric loss (tan δ) in both series, which was attributed to the suppression of oxygen vacancies due to the increased presence of B^{3+} and hence to a decrease in ionic conductivity and dielectric loss.

Keywords: Ba₂GdNbO₆; Ba₂GdTaO₆; Eu³⁺, B³⁺ co-doping; dielectric properties.

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1. INTRODUCTION

The $A_2BB'O_6$ type double perovskite structure attracts the attention of material scientists and solidstate chemists. This interest stems from their dielectric properties, such as low dielectric loss, and low chemical reactivity, as well as their potential applications, such as their compatibility with some of the commonly used superconductors as a result of having a suitable coefficient of thermal expansion (1- 6). The double perovskites with the general formula $A_2B'B''O_6$ are formed when the B site of ABO₃ is occupied by two different cations ($B' =$ rare earth elements, $B'' =$ transition metal, d-block metals, such as Nb, Ta, Mo, Sb) in the disordered state (7,8). Perovskite oxides with the structure $ABO₃$, which contain two different cations represented as A and B, are structurally stable due to the balanced arrangement of the constituent atoms and valences. Moreover, when the B site is occupied by a transition

metal, $B'O_6$ and $B'O_6$ in $A_2B'B'O_6$ perovskite can reduce the symmetry of the A sites and become suitable hosts for doping rare earth ions to obtain luminescent ferroelectrics (9-16). On the other hand, investigating the electrical properties of ferroelectric materials doped with trivalent europium ion ($Eu³⁺$) is important for the development of versatile optoelectronic devices (17-19). The Eu³⁺ is known for its strong luminescence in the red spectral region. It exhibits interesting spectral properties with $5D_0 \rightarrow 7F_J$ $(J=0, 1, 2, 3, 4, 5, 6)$ transitions, as well as having non-degenerate (J=0) first levels of transitions in both the absorption and luminescence spectrum, and has a great advantage over other RE ions (20-25). Boron is extensively employed as a flux in traditional solid-state reactions due to its relatively low melting point. Additionally, it's utilized to enhance optical and dielectric properties, alongside influencing structural aspects like morphology and crystallinity (26-29). Additionally, there are studies on the effect of boron

on grain morphology and its improvement in dielectric properties in which the doping of boron has the effect of increasing the bulk properties to some extent and can reduce the grain boundaries in the structure (30-32).

In the study, the structural and dielectric properties of Eu³⁺ doped Ba₂GdMO₆ (M=Nb, Ta) ceramics were studied by B^{3+} co-doping. The structural and dielectric characterizations of ceramic samples were carried out by XRD, SEM, and impedance analyses.

2. EXPERIMENTAL

Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70 and 100 mol%) and Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70 and 100 mol%) ceramic samples were fabricated by solidstate reaction. Barium carbonate BaCO₃ (Sigma-Aldrich, 99%), niobium oxide (Nb₂O₅: Alpha Aesar, 99.9%), tantalum oxide $(Ta₂O₅)$: Alpha Aesar, 99.9%) powders and gadolinium oxide (Gd_2O_3 : Alpha Aesar, 99.9%) were used as starting materials in calculated stoichiometric amounts. Europium oxide (Eu₂O₃: Alpha Aesar, 99.9%) and boric acid (H₃BO₃: Kimyalab, %99.9) were used as dopant materials. The stoichiometric amounts of $Ba_2Gd_{0.9}NbO_6$:0.1Eu³⁺ and $Ba_2Gd_{0.9}TaO_6:0.1Eu^{3+}$ starting materials were weighed and mixed in an agate mortar to provide homogeneity. Then, by adding different amounts of boric acid (H_3BO_3) , the final mixture of the powders was thoroughly mixed and ground in an agate mortar for the last time to provide more homogeneity. For sintering process, the sufficient amount of mixture was taken and calcined in an alumina crucible at 1250 °C for 6 h after pelleting.

The phase structure of the samples were examined by XRD (X-ray diffractometer; D2 PHASER, Bruker Corp., Germany) using Cu-K_{α} (1.5406 Å) radiation in between 2θ=15-80 °C with scan speed 2 °C/min. The grain morphology of the samples was investigated by scanning electron microscopy (FEG-SEM; XL 30S, Philips Corp., Netherlands). The crystal structure was visualized via VESTA software. The average crystallite sizes of the samples were determined from the Scherrer Eq. (1) (33):

$$
D = \frac{k \cdot \lambda}{B \cdot \cos \theta} \tag{1}
$$

where *D* stands for particle size in nanometers, with k is a constant (usually taken as 0.9). CuKα represents the wavelength ($\lambda = 0.15406$ Å), and *B* is the full width at half maximum in radians. Frequencydependent changes of real and imaginary permittivity and loss factor were defined using dielectric Eq. (2) and Eq. (3) respectively:

$$
\varepsilon' = \frac{c}{c_0}, \qquad \varepsilon'' = \frac{c}{\omega c_0}, \qquad C_0 = \varepsilon_0 \frac{A}{d} \tag{2}
$$

$$
tan\delta = \frac{\varepsilon^{\prime\prime}}{\varepsilon^{\prime}} \tag{3}
$$

where *C*0 is vacuum capacitance, *C* is capacitance, *w* is angular frequency and *G* is conductance. The dielectric properties of the ceramic samples were carried out using an impedance analyzer (HIOKI, LCR Hitester 3532-50; between frequency 20 Hz-10⁶ Hz, UK) at room temperature.

3. RESULTS AND DISCUSSION

3.1. XRD-SEM Results of Ba2Gd1-*x***NbO6:***x***Eu3+,** *y***B3+ and Ba2Gd1-***x***TaO6:***x***Eu3+,** *y***B3+**

Figure 1a and Figure 2a show the XRD results for Ba2Gd1-*x*MO6:*x*Eu3+, *y*B3+ (M=Nb, Ta), (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70, 100 mol%) samples, respectively. X-ray diffractions of all the sintered samples from 0 to 100 mol% B^{3+} showed the singlephase of $Ba₂GdMO₆$ (M=Nb, Ta). The crystal structure of Ba2GdNbO₆ (JCPDS card no. 47-0378) and Ba2GdTaO₆ (JCPDS card no. 49-1900) was indexed with space group *Fm*3̅*m* (225) and cubic symmetry. Figure 3 shows the schematic representation of Ba₂GdMO₆ (M=Nb, Ta) crystal structure consisting of edge-shared $B'O_6$ (GdO₆) and $B''O_6$ (MO₆) octahedral, where the formation of the single-phase or the absence of any minor phase may be explained to the substitution of $Eu³⁺$ ions with ionic radius 0.947 Å (for 6 CN) by Gd3+ ions (*r*=0.938 Å, for 6 CN), and the dissolution of the boron codoped at different concentrations in double perovskite host. Figure 1b and Figure 2b shows XRD two theta angles (400) for Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ and Ba₂Gd_{1-x}TaO₆:xEu³⁺, yB³⁺, respectively, where the XRD two theta peaks shifted towards smaller angles. Considering the slightly larger ionic radius of Eu³⁺, the substitution of Gd³⁺ with Eu³⁺ would contribute to lattice expansion. In addition, B^{3+} ions, which are characterized by a small ionic radius and are likely to occupy the structure as interstitial atoms, can lead to the expansion of the lattice. Table 1 summarizes the variation of lattice data and crystallite size. The shift is related to the increase in the lattice constant which elevates up to 100 mol% B3+ concentration in both series. The cell data (*a*, *V*) for 0 and 100 mol% B^{3+} concentrations were 8.4729 Å, 608.27 Å³ and 8.5029 Å, 614.75 Å³ respectively, for Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+, while they were 8.4680 Å, 607.22 λ^3 and 8.4942 Å, 612.87 λ^3 , respectively, for Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+. The lattice parameters $Ba₂GdNbO₆$ (34) and $Ba₂GdTaO₆$ (35) compounds are *a*=8.4900 Å, *V*=611.96 Å³ and *a*=8.4780 Å, *V*=609.37 Å³ , respectively. The cell data had values close to the literature data, where the lattice volume increased by about 1% in both series.

Figure 2: (a) X-ray diffractions of Ba₂Gd_{1-x}TaO₆: xEu^{3+} , yB^{3+} co-doped ceramics ($x=10$ mol%, and $y=0$, 5, 15, 30, 50, 70 and 100 mol%), (b) XRD two theta peak (400) shifted to lower angles with $B³⁺$ concentration.

The average crystallite sizes of the samples determined from the Scherrer equation using two theta reflection peaks (220), (400) and (422) are summarized in Table 1. The increasing B^{3+}

concentration led to an increase in crystallite size in both series up to range of 30-50 mol%, and then a decrease occurred at 70 and 100 mol% concentrations. As seen in Table 1, the crystallite

sizes from 0 to 50 mol% B^{3+} varied to 42.36-49.14 nm for Ba₂Gd_{1-x}NbO₆:xEu³⁺, yB³⁺, while the sizes changed in the same range as 47.07-55.61 nm for Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+, respectively. However, increasing boron concentration caused a decrease in crystallinity size at 70 and 100 mol% levels, where the crystallite sizes for 70 and 100 mol% B^{3+} concentrations were 45.79 and 43.38 nm (for Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+) and 49.72 and 46.28 nm (for Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+), respectively. Since the temperature effect will increase with increasing boron, the nucleation rate will slow down, which is related to the increasing crystallite size in the 0-50 mol% range. On the other hand, the decrease in crystallite size at 70 and 100 mol% can be linked to elevated boron concentration, leading to heightened nucleation, lattice deformation, and disruption of the charge balance within the structure, particularly at higher levels of doping (36-38).

Figure 3: Crystal structure visualization of Ba₂GdMO₆ (M=Nb, Ta).

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Figure 4: SEM micrographs of 10 mol% Eu³⁺ and (a) 0, (b) 5, (c) 15, (d) 30, (e) 50, (f) 70 and (g) 100 mol% B^{3+} co-doped Ba₂GdNbO₆ ceramics, at 2000 \times magnification and 2 kV acceleration voltage.

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Figure 5: SEM micrographs of 10 mol% Eu³⁺ and (a) 0, (b) 5, (c) 15, (d) 30, (e) 50, (f) 70 and (g) 100 mol% B^{3+} co-doped Ba₂GdTaO₆ ceramics, at 2000 \times magnification and 2 kV acceleration voltage.

Figure 4(a-g) and Figure 5(a-g) show the SEM micrographs of Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ and Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 30, 50, 70 and 100 mol%) samples, respectively, at 2000× magnifications under 2 kV accelerating voltage. SEM micrographs of both series have an irregular shape and the grains are rounded, while the increasing B^{3+} concentration led to grain growth and agglomeration. As well-known boron's flux effect reduces sintering temperature. Consequently, the increase in grain size may be attributed to the elevated temperatures resulting from boron addition. This higher temperature facilitates the release of stored energy within the grains, leading to an

increase in grain size (39,40). In addition, as seen in Figure 4(f,g) and Figure 5(f,g), the grain shape showed a slight angular trend at high B^{3+} concentrations. Moreover, at a concentration of 100 mol% B^{3+} , the agglomeration and flux effect is evident in the grains (Figure 4g and Figure 5g). The grain size of Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ samples varied mostly in the range of 0.5 -2 μ m at 0 mol% concentration and there are also some grains reaching 3-4 μ m, while the size of large grains ranging from 3 to 6 µm, while the grain size can exceed 10 µm, at 100 mol% concentration (Figure 4(f,g)). The grain sizes of Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+ samples are round-like and/or shapeless grains varied between 0.5-3 μ m at 0 mol%, whilst the sizes of the angular grains at 100 mol% ranged between $4-10 \mu m$ (Figure 5(f,g)).

3.2. Dielectric Properties of Ba2Gd1-*x***MO6:***x***Eu3+,** *y***B3+ (M=Nb, Ta), (***x***=10 mol%,** *y***=0, 5, 15, 30, 50, 70 and 100 mol%) Ceramics**

Figure $6(a,b)$ show the dielectric constants (ε') with frequency for Ba₂Gd_{1-x}NbO₆:xEu³⁺, yB^{3+} and Ba₂Gd₁*^x*TaO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70 and 100 mol%) samples, respectively. In Figure 6a, the dielectric constants of Ba2Gd1-*x*NbO6:*x*Eu3+, yB³⁺ from 0 to 100 mol% B³⁺ changed approximately 21.3 and 33.5 at 20 Hz, respectively. In Figure 6b, at the same range and 20 Hz, the ε' values for the Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+ ceramics were 21.3 and 35.4, respectively. However, the dielectric constant did not change in the high frequency or over $10³$ Hz. As the frequency rises, the electron exchange's capability to align with the applied field diminishes, leading to a reduction in the dielectric constant. At extremely high frequencies, the field reverses before the space charge carriers can respond, thereby preventing them from contributing to polarization (41-47), so the dielectric constant almost stayed unchanged at the high-frequency range. Moreover, as seen in Figure 6(a,b) there is an increase in polarization or dielectric constant from 0 to 50 mol% B^{3+} . The increase can be accounted for by the Maxwell-Wagner theory of external factors. According to this principle, the dielectric constant correlates directly with the grain size of the sample. The larger grain sizes result in greater atom polarizability, leading to an uptick in the dielectric constant (48-52). In the SEM micrographs in Figure 4(a-g) and Figure 5(a-g), it was previously noted that there is an increase in grain size with increasing B^{3+} concentration. Hence, this circumstance may be ascribed to the declining occurrence of grain boundaries in both series, coupled with the rise in polarizability and the value of ε' within the atomic structure. Additionally, although there was a slight increase in grain size at 70 and 100 mol% B^{3+} concentrations in both series, there was a decrease in dielectric constant. Caruntu et al (53) stated that there is a direct correlation between the change of dielectric constant and crystallite size/microstrain, in which the crystallite size increases uniformly from 32

to 94 nm with the increase of sintering temperature from 600 to 1000 °C. On the other hand, Kim et al (54) , in their study on BaTiO₃ and SrTiO₃, found that the presence of an optimal amount of lattice distortion or strain increases the dielectric constant. However, there is no information about crystallite size change in the study. As seen in Table 1, a decrease in crystallite size occurs at 70 and 100 mol% B3+ concentrations. Furthermore, while there is a noticeable rise in both crystallite size and dielectric constant as the concentration increases from 0 to 50 mol%, the augmented cell constant can be attributed to alterations in charge equilibrium and the existence of strain. In this context, the decrease in dielectric constant at 70 and 100 mol% could be linked to heightened strain due to an increased cell constant and deteriorated charge balance, as evidenced by the reduction in crystallite size.

Figure 7(a,b) show the dielectric losses (tan *δ*) with frequency for Ba₂Gd_{1-x}NbO₆:xEu³⁺, yB^{3+} and Ba₂Gd₁. *^x*TaO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70 and 100 mol%) samples, respectively. As seen in Figure 7(a,b), dielectric loss continuously decreases with the increase in frequency. In the low-frequency range where resistance is high, polarization necessitates more energy because of the grain boundary. Conversely, in the high-frequency range where resistance is low, electron transfer across the grain boundary requires minimal energy, resulting in reduced energy loss (54-57). On the other hand, the dielectric loss factor for both sample series decreased with increasing B^{3+} concentration up to 50 mol% and then showed a slight increase at 70 and 100 mol%. Various studies in the literature connect the tan *δ* in ferroelectrics to oxygen vacancies, which contribute to the dielectric loss or leakage current (58-61). Liu et al (62) studied the dielectric loss and oxygen vacancy relation by fabricating $TTB-Ba_4Nd_2Fe_2Nb_8O_{30}$ ceramics at different annealing and sintering temperatures as well as in $O₂$ and $N₂$ atmospheres, where the increase in temperature reduces oxygen vacancies and electrical conductivity. Similarly, in the study reported by Iqbal et al (57) on CuFe₂O₄, an increase in crystallite size and dielectric constant occurred due to the increase in annealing temperature. The results of the flux effect of boron, as seen in the XRD-SEM section, the temperature effect on the structure will increase with increasing boron concentration. Thus, the higher concentration of boron will not only intensify the material's sensitivity to temperature but also impact the presence of oxygen vacancies. Therefore, the decrease in the dielectric loss factor in both series may be attributed to the suppression of oxygen vacancies by the temperature effect that increases with the increase in B^{3+} . In addition, the slightly increased dielectric loss at the 70 and 100 mol% levels is likely due to the high B^{3+} concentration leading to some increase in the mobility of oxygen vacancies in the structure.

Figure 6: Dielectric constants for (a) $Ba_2Gd_{1-x}NbO_6$: xEu^{3+} , yB^{3+} ($x=10$ mol%, $y=0$, 5, 15, 30, 50, 70, 100 mol%), and (b) Ba₂Gd_{1-x}TaO₆:xEu³⁺, yB^{3+} (x=10 mol%, y=0, 5, 15, 30, 50, 70, 100 mol%) ceramics.

Figure 7: Dielectric losses (tan δ) for (a) Ba₂Gd_{1-x}NbO₆: xEu^{3+} , yB^{3+} ($x=10$ mol%, $y=0$, 5, 15, 30, 50, 70, 100 mol%), and (b) Ba2Gd1-*x*TaO6:*x*Eu3+, *y*B3+ (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70, 100 mol%) ceramics.

4. CONCLUSION

The effect of boron on the structural and dielectric properties was studied by using Eu^{3+} , B^{3+} co-doped double perovskite Ba2Gd1-*x*MO6:*x*Eu3+, *y*B3+ (M=Nb, Ta), (*x*=10 mol%, *y*=0, 5, 15, 30, 50, 70 and 100 mol%) ceramics, in which XRD results of the ceramics showed a single-phase structure. $Ba₂Gd₁$ *^x*TaO6:*x*Eu3+, *y*B3+ samples had slightly better crystallite sizes than Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ ones, while the crystallinity of both sample series increased up to about 50 mol% B^{3+} , and then decreased at 70 and 100 mol%. SEM micrographs of Eu^{3+} , B^{3+} codoped samples showed that boron in both series supported aggregation and growth in grains, and a slight angularity in grain shape occurred at high B^{3+} concentrations. The *ε*' values of Ba₂Gd_{1-x}NbO₆: xEu³⁺, *y*B3+, and Ba2Gd1-*x*NbO6:*x*Eu3+, *y*B3+ were measured in the range of 21.3-33.5 and 21.3 and 35.4 at 20 Hz, respectively. For both series, the increasing dielectric constant up to 50 mol% B^{3+} concentration was associated with the developing grain size and crystallinity, while its decrease at 70 and 100 mol% was associated with the reduction in crystallite size. The dielectric loss factor for both series decreased with increasing B^{3+} concentration, whilst Ba2Gd₁. *^x*TaO6:*x*Eu3+, *y*B3+ series exhibited a lower loss factor. The decline in dielectric loss in both series as B³⁺ concentration rises was linked to the reduction of oxygen vacancies. The heightened presence of B^{3+} resulted in diminished ionic conductivity and consequently decreased dielectric loss. The study may be useful in evaluating the dielectric properties of double perovskite ceramics, in terms of controlling the grain morphology and crystallite size.

5. CONFLICT OF INTEREST

There is no conflict of interest

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