# SPACELIKE CURVES OF CONSTANT BREADTH ACCORDING TO BISHOP FRAME IN MINKOWSKI 3-SPACE 

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#### Abstract

In this paper, the spacelike curves of constant breadth according to Bishop frame in Minkowski 3-space are studied. The differential equations characterizing the spacelike curves of constant breadth in $E_{1}^{3}$ are obtained. Furthermore, It is shown that the spacelike curves of constant breadth are connected with slant helix in Minkowski 3-space $E_{1}^{3}$.


## 1. Introduction

Firstly, the curves of constant breadth were introduced by Euler [1]. After him, many mathematicians studied about the curves of constant breadth [2-12]. Reuleaux gave the obtaining method some curves of constant breadth and used in kinematics of machinery [13]. The differential equations characterizing space curves of constant breadth were established and a criterion for these curves were given in [14]. Önder and et al gave the diferential equations characterizing the timelike and spacelike curves of constant breadth in Minkowski 3-space in [15]. Furthermore, they gave a criterion for a timelike or spacelike curve to be curve of constant breadth in $E_{1}^{3}$. Also, Kocayiğit and Önder showed that in $E_{1}^{3}$ spacelike and timelike curves of constant breadth were normal curves, helices and spherical curves in some special cases [16].

In this paper, we study the spacelike curves of constant breadth according to Bishop frame in $E_{1}^{3}$. We obtain the differential equations characterizing the spacelike curves of constant breadth. In addition, we demostrate that spacelike curves of constant breadth are connected with slant helix.

## 2. Preliminaries

The Minkowski 3 -space $E_{1}^{3}$ is the real vector space $\mathbb{R}^{3}$ provided with the standart flat metric given by

$$
\langle,\rangle=-d x_{1}^{2}+d x_{2}^{2}+d x_{3}^{2}
$$

[^0]where $\left(x_{1}, x_{2}, x_{3}\right)$ is a rectangular coordinate system of $E_{1}^{3}$. An arbitrary vector $\vec{a}=\left(a_{1}, a_{2}, a_{3}\right)$ in $E_{1}^{3}$ can have one of three Lorentzian characters; it can be spacelike if $\langle\vec{a}, \vec{a}\rangle>0$ or $\vec{a}=0$, timelike if $\langle\vec{a}, \vec{a}\rangle<0$ and null (lightlike) if $\langle\vec{a}, \vec{a}\rangle=0$ and $\vec{a} \neq 0$. Similarly, an arbitrary curve $\vec{\alpha}=\vec{\alpha}(s)$ can be spacelike, timelike or null (lightlike), if all of its velocity vectors $\vec{\alpha}^{\prime}$ are spacelike, timelike or null (lightlike), respectively. In addition, a timelike vector is future pointing or past pointing if the first compound of the vector is positive or negative, respectively.

For any vectors $\vec{x}=\left(x_{1}, x_{2}, x_{3}\right)$ and $\vec{y}=\left(y_{1}, y_{2}, y_{3}\right)$ in $E_{1}^{3}$, the vector product of $\vec{x}$ and $\vec{y}$ is defined by

$$
\vec{x} \wedge \vec{y}=\left|\begin{array}{ccc}
\overrightarrow{e_{1}} & -\overrightarrow{e_{2}} & -\overrightarrow{e_{3}} \\
x_{1} & x_{2} & x_{3} \\
y_{1} & y_{2} & y_{3}
\end{array}\right|=\left(x_{2} y_{3}-x_{3} y_{2}, x_{1} y_{3}-x_{3} y_{1}, x_{2} y_{1}-x_{1} y_{2}\right)
$$

where $\delta_{i j}=\left\{\begin{array}{ll}1, & i=j \\ 0, & i \neq j\end{array}\right.$, and $\overrightarrow{e_{i}}=\left(\delta_{i 1}, \delta_{i 2}, \delta_{i 3}\right)$ and $\overrightarrow{e_{1}} \wedge \overrightarrow{e_{2}}=-\overrightarrow{e_{3}}, \overrightarrow{e_{2}} \wedge \overrightarrow{e_{3}}=\overrightarrow{e_{1}}$, $\overrightarrow{e_{3}} \wedge \overrightarrow{e_{1}}=-\overrightarrow{e_{2}}$ (See for details [17]).

The parallel transport frame is an allternative approach to defining a moving frame that is well-defined even when the curve has vanishing second derivative. We can parallel transport an orthonormal frame along a curve simply by paralel transporting each component of the frame [18].

Its mathematical properties derive from the observation that, while $\vec{T}(s)$ for a given curve model is unique, we may choose any convenient arbitrary basis $\left(\overrightarrow{N_{1}}(s), \overrightarrow{N_{2}}(s)\right)$ for the remainder of the frame, so long as it is in the normal plane perpendicular to $\vec{T}(s)$ at each point. If the derivatives of $\left(\overrightarrow{N_{1}}(s), \overrightarrow{N_{2}}(s)\right)$ depend only on $\vec{T}(s)$ and not each other, we can make $\overrightarrow{N_{1}}(s)$ and $\overrightarrow{N_{2}}(s)$ vary smoothly throughout the path regardless of the curvature. We may therefore choose the alternative frame equations

$$
\left[\begin{array}{c}
\dot{\vec{T}} \\
\overrightarrow{N_{1}} \\
\overrightarrow{N_{2}}
\end{array}\right]=\left[\begin{array}{ccc}
0 & k_{1} & k_{2} \\
-k_{1} & 0 & 0 \\
-k_{2} & 0 & 0
\end{array}\right]\left[\begin{array}{l}
\vec{T} \\
\overrightarrow{N_{1}} \\
\overrightarrow{N_{2}}
\end{array}\right]
$$

where $\langle\vec{T}, \vec{T}\rangle=\left\langle\overrightarrow{N_{1}}, \overrightarrow{N_{1}}\right\rangle=\left\langle\overrightarrow{N_{2}}, \overrightarrow{N_{2}}\right\rangle=1$ and $\left\langle\vec{T}, \overrightarrow{N_{1}}\right\rangle=\left\langle\overrightarrow{N_{1}}, \overrightarrow{N_{2}}\right\rangle=\left\langle\vec{T}, \overrightarrow{N_{2}}\right\rangle=$ 0 [19,20].

One can show that [19]

$$
\kappa(s)=\sqrt{k_{1}^{2}+k_{2}^{2}}, \quad \theta(s)=\arctan \left(\frac{k_{2}}{k_{1}}\right), \quad \tau(s)=\frac{d \theta(s)}{d s}
$$

so that $k_{1}$ and $k_{2}$ effectively correspond to a Cartesian coordinate system for the polar coordinates $\kappa, \theta$ with $\theta=\tau(s) d s$. A fundamental ambiguity in the parallel transport frame compared to the Frenet frame thus aries from the arbitrary choice of an integration constant for $\theta_{0}$, which disappears from $\tau$ due to the differentation [20].

Denote by $\left\{\vec{T}, \overrightarrow{N_{1}}, \overrightarrow{N_{2}}\right\}, k_{1}$ and $k_{2}$ the moving Bishop frame, the first curvature and the second curvature along the curve $\vec{\alpha}(s)$, respectively in Minkowski 3-space
$E_{1}^{3}$. If $\vec{\alpha}(s)$ is a spacelike curve in $E_{1}^{3}$, then Bishop frame is given by

$$
\left[\begin{array}{c}
\dot{\vec{T}}  \tag{2.1}\\
\dot{\overrightarrow{N_{1}}} \\
\dot{N_{2}}
\end{array}\right]=\left[\begin{array}{ccc}
0 & k_{1} & -k_{2} \\
\epsilon k_{1} & 0 & 0 \\
\epsilon k_{2} & 0 & 0
\end{array}\right]\left[\begin{array}{l}
\vec{T} \\
\overrightarrow{N_{1}} \\
\overrightarrow{N_{2}}
\end{array}\right]
$$

where $\langle\vec{T}, \vec{T}\rangle=1,\left\langle\overrightarrow{N_{1}}, \overrightarrow{N_{1}}\right\rangle=-\epsilon$ and $\left\langle\overrightarrow{N_{2}}, \overrightarrow{N_{2}}\right\rangle=\epsilon$, and [21,22]. Here, $\epsilon$ determines the kind of spacelike curve $\vec{\alpha}(s)$. If $\epsilon=1$, then $\vec{\alpha}(s)$ is a spacelike curve with timelike principal normal and spacelike binormal. If $\epsilon=-1$, then $\vec{\alpha}(s)$ is a spacelike curve with spacelike principal normal and timelike binormal. The relations between $\kappa, \tau, \theta$ and $k_{1}, k_{2}$ are given by as follows.

$$
\kappa(s)=\sqrt{\left|k_{1}^{2}-k_{2}^{2}\right|}, \quad \theta(s)=\operatorname{arctanh}\left(\frac{k_{2}}{k_{1}}\right), \quad \tau(s)=\frac{d \theta(s)}{d s}
$$

Also, Bishop curvatures are defined by

$$
k_{1}=\kappa \cosh (\theta), \quad k_{2}=\sinh (\theta)
$$

and

$$
\vec{T}=\vec{T}, \quad \overrightarrow{N_{1}}=\vec{N} \cosh (\theta)-\vec{B} \sinh (\theta), \quad \overrightarrow{N_{2}}=\vec{N} \sinh (\theta)-\vec{B} \cosh (\theta)
$$

[21,22].
Theorem 2.1. Let $\alpha: I \longrightarrow E_{1}^{3}$ be a unit speed spacelike curve with non-zero natural curvatures. Then $\vec{\alpha}(s)$ is a spacelike slant helix if and only if $\frac{k_{1}}{k_{2}}$ is constant [23].

## 3. The Spacelike Curves of Constant Breadth

In this section, we give differential equations characterizing the spacelike curves of constant breadth according to Bishop frame in Minkowski 3-space. In addition, we show that the spacelike curves of constant breadth are related to slant helix in Minkowski 3-space $E_{1}^{3}$.
Definition 3.1. Let $(C)$ be a spacelike curve. If $(C)$ has parallel tangents in opposite directions at the opposite points $\alpha(s)$ and $\alpha^{*}(s)$, and if the distance between these points is always constant, then $(C)$ is called a spacelike curve of constant breadth. Moreover, a pair of spacelike curves $(C)$ and $\left(C^{*}\right)$ for which the tangents at the corresponding points $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$ are parallel and in opposite directions, and the distance between these points is always constant are called a spacelike curve pair of constant breadth [16].

Let $(C)$ and $\left(C^{*}\right)$ be a pair of unit-speed spacelike curves with non-zero Bishop curvatures in $E_{1}^{3}$ and let those curves have parallel tangents in opposite directions at the corresponding points $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$, respectively. The position vector of the curve $\left(C^{*}\right)$ at the point $\alpha^{*}\left(s^{*}\right)$ can be expressed as

$$
\begin{equation*}
\overrightarrow{\alpha^{*}}\left(s^{*}\right)=\vec{\alpha}(s)+\lambda_{1}(s) \vec{T}(s)+\lambda_{2}(s) \overrightarrow{N_{1}}(s)+\lambda_{3}(s) \overrightarrow{N_{2}}(s) \tag{3.1}
\end{equation*}
$$

where $\lambda_{i}(s) \quad(i=1,2,3)$ are differentiable functions of $s$ which is arc lenght of $(C)$. Denote by $\left\{\vec{T}, \overrightarrow{N_{1}}, \overrightarrow{N_{2}}\right\}, k_{1}$ and $k_{2}$ the moving Bishop frame, Bishop curvatures
along the curve $(C)$ respectively. And denote by $\left\{\overrightarrow{T^{*}}, \overrightarrow{N_{1}^{*}}, \overrightarrow{N_{2}^{*}}\right\}, k_{1}^{*}$ and $k_{2}^{*}$ the moving Bishop frame, Bishop curvatures along the curve $\left(C^{*}\right)$, respectively.

Differentiating Eq.(3.1) with respect to $s$ and using the Bishop formulae given by (2.1), we obtain

$$
\begin{equation*}
\frac{d \overrightarrow{\alpha^{*}}}{d s}=\overrightarrow{T^{*}} \frac{d s^{*}}{d s}=\left(1+\frac{d \lambda_{1}}{d s}+\epsilon k_{1} \lambda_{2}+\epsilon k_{2} \lambda_{3}\right) \vec{T}+\left(k_{1} \lambda_{1}+\frac{d \lambda_{2}}{d s}\right) \overrightarrow{N_{1}}+\left(-k_{2} \lambda_{1}+\frac{d \lambda_{3}}{d s}\right) \overrightarrow{N_{2}} \tag{3.2}
\end{equation*}
$$

Since $\vec{T}=-\overrightarrow{T^{*}}$ at the corresponding points of $(C)$ and $\left(C^{*}\right)$, we gain the following differential equations system

$$
\left\{\begin{array}{c}
\frac{d \lambda_{1}}{d s}=-\frac{d s^{*}}{d s}-1-\epsilon k_{1} \lambda_{2}-\epsilon k_{2} \lambda_{3}  \tag{3.3}\\
\frac{d \lambda_{2}}{d s}=-k_{1} \lambda_{1} \\
\frac{d \lambda_{3}}{d s}=k_{2} \lambda_{1}
\end{array}\right.
$$

It is well known that the curvature $\kappa(s)$ of the curve $(C)$ is

$$
\lim _{\Delta s \rightarrow 0} \frac{\Delta \varphi}{\Delta s}=\frac{d \varphi}{d s}=\kappa
$$

where $\varphi$ is the angle between the tangent of the spacelike curve $(C)$ and a given fixed direction at the point $\alpha(s)$.

Hence, we can rewrite the system (3.3) as follow.

$$
\left\{\begin{array}{c}
\frac{d \lambda_{1}}{d s}=-\epsilon \mu_{1} \lambda_{2}-\epsilon \mu_{2} \lambda_{3}-f  \tag{3.4}\\
\frac{d \lambda_{2}}{d s}=-\mu_{1} \lambda_{1} \\
\frac{d \lambda_{3}}{d s}=\mu_{2} \lambda_{1}
\end{array}\right.
$$

where

$$
\mu_{1}=\rho k_{1}=\frac{k_{1}}{\kappa}=\cosh (\theta), \quad \mu_{2}=\rho k_{2}=\frac{k_{2}}{\kappa}=\sinh (\theta), \quad(\theta=\tau d s)
$$

and

$$
f(\varphi)=\rho+\rho^{*}, \quad \rho=\frac{1}{\kappa}, \quad \rho^{*}=\frac{1}{\kappa^{*}}
$$

Here, $\rho$ and $\rho^{*}$ indicates the radius of curvatures at the points $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$, respectively.

Eliminating $\lambda_{2}, \lambda_{3}$ and their derivatives from the system (3.4), we obtain the following differential equation of third order with respect to $\lambda_{1}$.

$$
\begin{equation*}
a_{1} \lambda_{1}^{\prime \prime \prime}+b_{1} \lambda_{1}^{\prime \prime}+c_{1} \lambda_{1}^{\prime}+d_{1} \lambda_{1}=e_{1} \tag{3.5}
\end{equation*}
$$

where

$$
\begin{aligned}
a_{1}= & -\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{3}-\mu_{1}^{\prime \prime} \mu_{2}^{3}\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right) \\
b_{1}= & -\mu_{2}\left[\left(\mu_{1}^{\prime \prime}\right)^{2} \mu_{2}^{3}+\mu_{1}^{\prime \prime}\left(-\mu_{2}^{2} \mu_{1} \mu_{2}^{\prime \prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\right)+\mu_{2}^{\prime \prime}\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\right] \\
c_{1}= & \mu_{1}^{\prime \prime}\left[-\mu_{2}^{3} \mu_{2}^{\prime \prime} \mu_{1}^{\prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(\epsilon \mu_{2}^{3} \mu_{1}^{2}+\mu_{1}\left(\mu_{2}^{\prime}\right)^{2}-\mu_{2} \mu_{1}^{\prime} \mu_{2}^{\prime}-\epsilon \mu_{2}^{5}\right)\right] \\
& +\left(\mu_{1}^{\prime \prime}\right)^{2} \mu_{2}^{3} \mu_{2}^{\prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\left(\mu_{2}^{\prime \prime} \mu_{2}^{\prime}+\epsilon\left(\mu_{1}^{2}-\mu_{2}^{2}\right)\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\right) \\
d_{1}= & \epsilon \mu_{2} \mu_{1}^{\prime \prime}\left[\left(\mu_{2}^{4} \mu_{1}-\mu_{2}^{2} \mu_{1}^{3}\right) \mu_{2}^{\prime \prime}+\left(\left(\mu_{1}^{3}+\mu_{1}^{2} \mu_{2}-3 \mu_{2}^{3}\right) \mu_{2}^{\prime}-\mu_{1} \mu_{2} \mu_{1}^{\prime}\left(-2 \mu_{2}+\mu_{1}\right)\right)\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\right] \\
& +\epsilon\left(\mu_{2}^{4} \mu_{1}^{2}-\mu_{2}^{6}\right)\left(\mu_{1}^{\prime \prime}\right)^{2}+\epsilon\left[\left(\mu_{1}^{2} \mu_{2}-\mu_{2}^{3}\right) \mu_{2}^{\prime \prime}+3\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(\mu_{1} \mu_{1}^{\prime}-\mu_{2} \mu_{2}^{\prime}\right)\right]\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2} \\
e_{1}= & -\mu_{1}^{\prime \prime}\left[-\mu_{2}^{3}\left(f \mu_{1}^{\prime}-f^{\prime} \mu_{1}\right) \mu_{2}^{\prime \prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(-\mu_{2}^{3} f^{\prime \prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(-\mu_{2} f^{\prime}+\mu_{2}^{\prime} f\right)\right)\right] \\
& -\left(\mu_{1}^{\prime \prime}\right)^{2}\left(\mu_{2}^{3} \mu_{2}^{\prime} f-\mu_{2}^{4} f^{\prime}\right)-\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\left[\left(-\mu_{2} f^{\prime}+\mu_{2}^{\prime} f\right) \mu_{2}^{\prime \prime}-f^{\prime \prime}\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\right] .
\end{aligned}
$$

Here and later ( ${ }^{\prime}$ ) denotes the differentiation with respect to " $\varphi$ ". Similarly, eliminating $\lambda_{1}, \lambda_{3}$ and their derivatives from the system (3.4) we obtain the following differential equation of third order with respect to $\lambda_{2}$.

$$
\begin{equation*}
a_{2} \lambda_{2}^{\prime \prime \prime}+b_{2} \lambda_{2}^{\prime \prime}+c_{2} \lambda_{2}^{\prime}+d_{2} \lambda_{2}=e_{2} \tag{3.6}
\end{equation*}
$$

where

$$
\begin{aligned}
a_{2} & =-\mu_{1}^{2} \mu_{2} \\
b_{2} & =2 \mu_{1}^{\prime} \mu_{1} \mu_{2}+\mu_{2}^{\prime} \mu_{1}^{2} \\
c_{2} & =\mu_{1}^{\prime \prime} \mu_{1} \mu_{2}-2\left(\mu_{1}^{\prime}\right)^{2} \mu_{2}+\epsilon \mu_{1}^{4} \mu_{2}-\mu_{2}^{\prime} \mu_{1}^{\prime} \mu_{1}-\epsilon \mu_{2}^{3} \mu_{1}^{2} \\
d_{2} & =\epsilon \mu_{1}^{3}\left(\mu_{1}^{\prime} \mu_{2}-\mu_{2}^{\prime} \mu_{1}\right) \\
e_{2} & =\mu_{1}^{3}\left(\mu_{2}^{\prime} f-f^{\prime} \mu_{2}\right) .
\end{aligned}
$$

Furthermore, eliminating $\lambda_{1}, \lambda_{2}$ and their derivatives from the system (3.4) we gain the following differential equation of third order with respect to $\lambda_{3}$.

$$
\begin{equation*}
a_{3} \lambda_{3}^{\prime \prime \prime}+b_{3} \lambda_{3}^{\prime \prime}+c_{3} \lambda_{3}^{\prime}+d_{3} \lambda_{3}=e_{3} \tag{3.7}
\end{equation*}
$$

where

$$
\begin{aligned}
a_{3} & =-\mu_{2}^{2} \mu_{1} \\
b_{3} & =2 \mu_{2}^{\prime} \mu_{2} \mu_{1}+\mu_{1}^{\prime} \mu_{2}^{2} \\
c_{3} & =\mu_{2}^{\prime \prime} \mu_{2} \mu_{1}-2\left(\mu_{2}^{\prime}\right)^{2} \mu_{1}-\mu_{1}^{\prime} \mu_{2}^{\prime} \mu_{2}+\epsilon \mu_{1}^{3} \mu_{2}^{2}-\epsilon \mu_{2}^{4} \mu_{1} \\
d_{3} & =\epsilon \mu_{2}^{3}\left(-\mu_{2}^{\prime} \mu_{1}+\mu_{2} \mu_{1}^{\prime}\right) \\
e_{3} & =-\mu_{2}^{3}\left(\mu_{1}^{\prime} f-f^{\prime} \mu_{1}\right) .
\end{aligned}
$$

When the curve $(C)$ and the function $f(\varphi)$ are given, from the solving the systems (3.3), (3.4) or the Eqs.(3.5), (3.6), (3.7), we can find the values of $\lambda_{i} \quad(i=1,2,3)$. Eqs.(3.5), (3.6) and (3.7) express the differential characterizations for the spacelike curves $(C)$ and $\left(C^{*}\right)$ according to coeffcients $\lambda_{i}$.

When the curves $(C)$ and $\left(C^{*}\right)$ are a spacelike curve pair of constant breadth, then the distance $d$ between the corresponding points $\alpha(s)$ and $\alpha^{*}\left(s^{*}\right)$ is constant. Hence,

$$
\begin{equation*}
d^{2}=\|\vec{d}\|^{2}=\left\|\overrightarrow{\alpha^{*}}\left(s^{*}\right)-\vec{\alpha}(s)\right\|^{2}=\left|\lambda_{1}^{2}-\epsilon \lambda_{2}^{2}+\epsilon \lambda_{3}^{2}\right|=k^{2}, \quad k \in \mathbb{R} \tag{3.8}
\end{equation*}
$$

Differentiating Eq.(3.8) with respect to $\varphi$, we gain

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d \varphi}\|\vec{d}\|^{2}=\lambda_{1} \lambda_{1}^{\prime}-\epsilon \lambda_{2} \lambda_{2}^{\prime}+\epsilon \lambda_{3} \lambda_{3}^{\prime}=0 \tag{3.9}
\end{equation*}
$$

Substituting the equalities given by the system (3.4) into the Eq.(3.9) we obtain the following equality.

$$
\begin{equation*}
\lambda_{1} f=0 \tag{3.10}
\end{equation*}
$$

This relation express to be spacelike curve pair of constant breadth of the spacelike curves $(C)$ and $\left(C^{*}\right)$ in $E_{1}^{3}$. Here, there are two main cases.

Case 1. Let $f(\varphi)=0 \quad\left(\frac{d s^{*}}{d s}+1=0\right)$. This means that the spacelike curve $\left(C^{*}\right)$ is a translation of the spacelike curve $(C)$ by the constnat vector

$$
\begin{equation*}
\vec{d}=\lambda_{1} \vec{T}+\lambda_{2} \overrightarrow{N_{1}}+\lambda_{3} \overrightarrow{N_{3}} \tag{3.11}
\end{equation*}
$$

In fact, if $f(\varphi)=0$ then the vector $\vec{d}$ is constant. To verify this fact, differentiate Eq.(3.11) with respect to $\varphi$ and use the equalities given by (3.4) for $f=0$ and Bishop formulae given by (2.1). Hence, we obtain $\frac{d \vec{d}}{d \varphi}=0$. Consequently, if $\frac{d \vec{d}}{d \varphi}=0$ then the vector $\vec{d}$ is constant. In this case we can rewrite the systems (3.3), (3.4) and the Eqs.(3.5), (3.6) and (3.7) as follows:

$$
\left.\begin{array}{c}
\left\{\begin{array}{c}
\frac{d \lambda_{1}}{d s}=-\epsilon k_{1} \lambda_{2}-\epsilon k_{2} \lambda_{3} \\
\frac{d \lambda_{2}}{d s}=-k_{1} \lambda_{1} \\
\frac{d \lambda_{3}}{d s}=k_{2} \lambda_{1}
\end{array}\right. \\
\left\{\begin{array}{c}
\frac{d \lambda_{1}}{d s}=-\epsilon \mu_{1} \lambda_{2}-\epsilon \mu k_{2} \lambda_{3} \\
\frac{d \lambda_{2}}{d s_{3}}=-\mu_{1} \lambda_{1}
\end{array}\right. \\
\frac{d \lambda_{3}}{d s}=\mu_{2} \lambda_{1} \tag{3.14}
\end{array}\right\}
$$

where

$$
\begin{aligned}
a_{1}= & -\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{3}-\mu_{1}^{\prime \prime} \mu_{2}^{3}\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right) \\
b_{1}= & -\mu_{2}\left[\left(\mu_{1}^{\prime \prime}\right)^{2} \mu_{2}^{3}+\mu_{1}^{\prime \prime}\left(-\mu_{2}^{2} \mu_{1} \mu_{2}^{\prime \prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\right)+\mu_{2}^{\prime \prime}\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\right] \\
c_{1}= & \mu_{1}^{\prime \prime}\left[-\mu_{2}^{3} \mu_{2}^{\prime \prime} \mu_{1}^{\prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(\epsilon \mu_{2}^{3} \mu_{1}^{2}+\mu_{1}\left(\mu_{2}^{\prime}\right)^{2}-\mu_{2} \mu_{1}^{\prime} \mu_{2}^{\prime}-\epsilon \mu_{2}^{5}\right)\right] \\
& +\left(\mu_{1}^{\prime \prime}\right)^{2} \mu_{2}^{3} \mu_{2}^{\prime}+\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2}\left(\mu_{2}^{\prime \prime} \mu_{2}^{\prime}+\epsilon\left(\mu_{1}^{2}-\mu_{2}^{2}\right)\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\right) \\
d_{1}= & \epsilon \mu_{2} \mu_{1}^{\prime \prime}\left[\left(\mu_{2}^{4} \mu_{1}-\mu_{2}^{2} \mu_{1}^{3}\right) \mu_{2}^{\prime \prime}+\left(\left(\mu_{1}^{3}+\mu_{1}^{2} \mu_{2}-3 \mu_{2}^{3}\right) \mu_{2}^{\prime}-\mu_{1} \mu_{2} \mu_{1}^{\prime}\left(-2 \mu_{2}+\mu_{1}\right)\right)\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\right] \\
& +\epsilon\left(\mu_{2}^{4} \mu_{1}^{2}-\mu_{2}^{6}\right)\left(\mu_{1}^{\prime \prime}\right)^{2}+\epsilon\left[\left(\mu_{1}^{2} \mu_{2}-\mu_{2}^{3}\right) \mu_{2}^{\prime \prime}+3\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)\left(\mu_{1} \mu_{1}^{\prime}-\mu_{2} \mu_{2}^{\prime}\right)\right]\left(\mu_{1} \mu_{2}^{\prime}-\mu_{2} \mu_{1}^{\prime}\right)^{2} .
\end{aligned}
$$

$$
\begin{equation*}
a_{2} \lambda_{2}^{\prime \prime \prime}+b_{2} \lambda_{2}^{\prime \prime}+c_{2} \lambda_{2}^{\prime}+d_{2} \lambda_{2}=0 \tag{3.15}
\end{equation*}
$$

where

$$
\begin{aligned}
a_{2} & =-\mu_{1}^{2} \mu_{2} \\
b_{2} & =2 \mu_{1}^{\prime} \mu_{1} \mu_{2}+\mu_{2}^{\prime} \mu_{1}^{2} \\
c_{2} & =\mu_{1}^{\prime \prime} \mu_{1} \mu_{2}-2\left(\mu_{1}^{\prime}\right)^{2} \mu_{2}+\epsilon \mu_{1}^{4} \mu_{2}-\mu_{2}^{\prime} \mu_{1}^{\prime} \mu_{1}-\epsilon \mu_{2}^{3} \mu_{1}^{2} \\
d_{2} & =\epsilon \mu_{1}^{3}\left(\mu_{1}^{\prime} \mu_{2}-\mu_{2}^{\prime} \mu_{1}\right) .
\end{aligned}
$$

$$
\begin{equation*}
a_{3} \lambda_{3}^{\prime \prime \prime}+b_{3} \lambda_{3}^{\prime \prime}+c_{3} \lambda_{3}^{\prime}+d_{3} \lambda_{3}=0 \tag{3.16}
\end{equation*}
$$

where

$$
\begin{aligned}
a_{3} & =-\mu_{2}^{2} \mu_{1} \\
b_{3} & =2 \mu_{2}^{\prime} \mu_{2} \mu_{1}+\mu_{1}^{\prime} \mu_{2}^{2} \\
c_{3} & =\mu_{2}^{\prime \prime} \mu_{2} \mu_{1}-2\left(\mu_{2}^{\prime}\right)^{2} \mu_{1}-\mu_{1}^{\prime} \mu_{2}^{\prime} \mu_{2}+\epsilon \mu_{1}^{3} \mu_{2}^{2}-\epsilon \mu_{2}^{4} \mu_{1} \\
d_{3} & =\epsilon \mu_{2}^{3}\left(-\mu_{2}^{\prime} \mu_{1}+\mu_{2} \mu_{1}^{\prime}\right) .
\end{aligned}
$$

Theorem 3.1. The general differential equations and systems characterizing spacelike curve pair of constant breadth according to Bishop frame in $E_{1}^{3}$ are given by (3.12), (3.13) and (3.14), (3.15), (3.16).

Case 2. Let $\lambda_{1}=0$. Then, there are three cases here.
i) We can take $\lambda_{2}=$ const. and $\lambda_{3}=0$ (from (3.4)). Then, $f(\varphi)=-\epsilon \mu_{1} \lambda_{2}$.
ii) We can take $\lambda_{2}=0$ and $\lambda_{3}=$ const. (from (3.4)). Then, $f(\varphi)=-\epsilon \mu_{2} \lambda_{3}$.
iii) Now, we consider the third and important case $\lambda_{2}=$ const. and $\lambda_{3}=$ const. If $\lambda_{2}=$ const., $\lambda_{3}=$ const. and $f(\varphi)=0$ (from (3.4)), then we obtain $\frac{k_{1}}{k_{2}}=-\frac{\lambda_{3}}{\lambda_{2}}=$ const. This means that the curve $(C)$ is a spacelike slant helix according to Bishop frame. Thus we can give following theorem.

Theorem 3.2. Let consider the spacelike curve pair of constant breadth which has the sum of curvature radius at corresponding points is zero according to Bishop frame in $E_{1}^{3}$. If the first normal component $\lambda_{2}=$ const. and the second normal component $\lambda_{3}=$ const. given by Eq.(3.1), then the spacelike curve $(C)$ is a spacelike slant helix in $E_{1}^{3}$.

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[^0]:    Date: Received: June 12, 2014, Accepted: February 20, 2015.
    2010 Mathematics Subject Classification. 53A04, 14H50.
    Key words and phrases. Bishop frame, Curve of constant breadth, Differential characterizations of curve, Minkowski space, Spacelike Curve .

