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Biharmonic maps on kenmotsu manifolds

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Abstract: In this paper we study biharmonic maps on Kenmotsu manifolds. An example for biharmonic map of a three-Kenmotsu manifold is constructed for illustration.

Keywords: Biharmonic map, Kenmotsu manifold, three-Kenmotsu manifold.

1 Introduction

1.1 Kenmotsu manifolds.

The notion of Kenmotsu manifolds is defined by K. Kenmotsu (see [15]). Let $(M^{2n+1}, \phi, \xi, \eta, g)$ be an almost contact Riemannian manifold, where ϕ is an (1,1) tensor field, η is a 1-form and g is the Riemannian metric. We have for any $X, Y \in \Gamma(TM)$ (see [7] and [21])

$$\phi(\xi) = 0, \ \eta(\phi(X)) = 0, \ \eta(\xi) = 1,$$
 (1)

$$\phi^{2}(X) = -X + \eta(X)\xi, \tag{2}$$

$$g(X,\xi) = \eta(X) \tag{3}$$

and

$$g(\phi(X),\phi(Y)) = g(X,Y) - \eta(X)\eta(Y). \tag{4}$$

If, moreover

$$(\nabla_X \phi) Y = -\eta (Y) \phi (X) - g(X, \phi (Y)) \xi \tag{5}$$

and

$$\nabla_X \xi = X - \eta(X) \xi, \tag{6}$$

where ∇ denotes the Riemannian connexion of g, then $(M^{2n+1}, \phi, \xi, \eta, g)$ is called an almost Kenmotsu manifold. In kenmotsu manifolds, we have the following relations (see [15]):

$$(\nabla_X \eta) Y = g(\phi(X), \phi(Y)), \tag{7}$$

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$$R(X,Y)\xi = \eta(X)Y - \eta(Y)X, \tag{8}$$

$$R(\xi, X)Y = \eta(Y)X - g(X, Y)\xi, \tag{9}$$

$$R(\xi, X)\xi = X - \eta(X)\xi \tag{10}$$

and

$$S(X,\xi) = -(n-1)\eta(X),$$
 (11)

where R is the Riemannian curvature tensor and S is the Ricci tensor. Kenmotsu manifolds have been studied by many authors, for example see [8], [14] and [20].

1.2 Harmonic maps.

Let $\phi:(M^m,g)\to (N^n,h)$ be a smooth map between riemannian manifolds. Then ϕ is said to be harmonic if it is a critical point of the energy functional :

$$E(\phi) = \frac{1}{2} \int_{M} |d\phi|^2 dv_g \tag{12}$$

with respect to compactly supported variations. Equivalently, ϕ is harmonic if it satisfies the associated Euler-Lagrange equations:

$$\tau(\phi) = Tr_{g}\nabla d\phi = 0,\tag{13}$$

 $\tau(\phi)$ is called the tension field of ϕ . One can refer to [1], [5], [9], [10], [11] and [12] for background on harmonic maps. Indeed, the Euler-Lagrange equation associated to the energy is the vanishing of the tension field $\tau(\phi) = Tr_g\nabla d\phi$. In the context of harmonic maps, the stress-energy tensor was studied in details by Baird and Eells in [2]. Indeed, the Euler-Lagrange equation associated to the energy is the vanishing of the tension field $\tau(\phi) = Tr_g\nabla d\phi$, and the stress-energy tensor for a map $\phi: (M^m, g) \longrightarrow (N^n, h)$ defined by

$$S(\phi) = e(\phi)g - \phi^*h.$$

The relation between $S(\phi)$ and $\tau(\phi)$ is given by

$$divS(\phi) = -h(\tau(\phi), d\phi).$$

1.3 Biharmonic maps.

The map ϕ is said to be biharmonic if it is a critical point of the bi-energy functional:

$$E_2(\phi) = \frac{1}{2} \int_M |\tau(\phi)|^2 d\nu_g$$
 (14)

Equivalently, ϕ is biharmonic if it satisfies the associated Euler-Lagrange equations:

$$\tau_2(\phi) = -Tr_g \left(\nabla^{\phi}\right)^2 \tau(\phi) - Tr_g R^N(\tau(\phi), d\phi) d\phi = 0, \tag{15}$$

where ∇^{ϕ} is the connection in the pull-back bundle $\phi^{-1}(TN)$ and, if e_i is a local orthonormal frame field on M, then

$$\operatorname{Tr}_{g}\left(\nabla^{\phi}\right)^{2}\tau\left(\phi\right)=\left(\nabla_{e_{i}}^{\phi}\nabla_{e_{i}}^{\phi}-\nabla_{\nabla_{e_{i}}e_{i}}^{\phi}\right)\tau\left(\phi\right),$$

where we sum over repeated indices. We will call the operator $\tau_2(\phi)$, the bi-tension field of the map ϕ . In analogy with harmonic maps, Jiang In [13] has constructed for a map ϕ the stress bi-energy tensor defined by

$$S_2(\phi) = \left(\frac{-1}{2}|\tau(\phi)|^2 + divh(\tau(\phi), d\phi)\right)g - 2symh(\nabla \tau(\phi), d\phi),$$

where

$$\operatorname{symh}\left(\nabla\tau(\phi),d\phi\right)\left(X,Y\right)=\frac{1}{2}\left\{h\left(\nabla_{X}\tau(\phi),d\phi\left(Y\right)\right)+h\left(\nabla_{Y}\tau(\phi),d\phi\left(X\right)\right)\right\},$$

for any $X, Y \in \Gamma(TM)$. The stress bi-energy tensor of ϕ satisfies the following relationship

$$divS_2(\phi) = h(\tau_2(\phi), d\phi).$$

Clearly any harmonic map is biharmonic, therefore it is interesting to construct non-harmonic biharmonic maps. In [4] the authors found new examples of biharmonic maps by conformally deforming the domain metric of harmonic ones. While in [6] the author analyzed the behavior of the biharmonic equation under the conformal change the domain metric, he obtained metrics $\tilde{g} = e^{2\gamma}$ such that the idendity map $Id: (M,g) \longrightarrow (M,\tilde{g})$ is biharmonic non-harmonic. Moreover, in [19] the author gave some extensions of the result in [6] together with some further constructions of biharmonic maps. The author in [18] deform conformally the codomain metric in order to render a semi-conformal harmonic map biharmonic. In [3] the authors studied the case where $\phi: (M^n,g) \longrightarrow (N^n,h)$ is a conformal mapping between equidimensional manifolds where they show that a conformal mapping ϕ is biharmonic if and only if the gradient of its dilation satisfies a second order elliptic partial differential equation. We can refer the reader to [17], for a survey of biharmonic maps. In [16], the author obtain some results concerning harmonic maps and harmonic morphism on Kenmotsu manifolds, he prove that any structure preserving map from a Kenmotsu manifold to a Kahler manifold is harmonic and that there are no nonconstant harmonic holomorphic maps from a Kahler manifold to a Kenmotsu manifold. In this paper, we calculate the bitension field of $\psi: (N^{2n}, J, h) \longrightarrow (M, \phi, \xi, \eta, g)$, where (N^{2n}, J, h) is a Kahlerian manifold and (M, ϕ, ξ, η, g) is a Kenmotsu manifold (Theorem 1), with this setting we obtain new example of biharmonic non-harmonic maps.

2 Statement of results.

As first result, we give another formula for the stress bi-energy tensor for a smooth map $\phi: (M^m, g) \to (N^n, h)$.

Proposition 1. Let $\phi: (M^m, g) \to (N^n, h)$ be a smooth map, then we have

$$S_2(\phi) = \left(\frac{1}{2} |\tau(\phi)|^2 + Tr_g h(\nabla \tau(\phi), d\phi)\right) g - 2symh(\nabla \tau(\phi), d\phi), \tag{16}$$

and the trace of $S_2(\phi)$ is given by

$$Tr_g S_2(\phi) = \frac{m}{2} |\tau(\phi)|^2 + (m-2) Tr_g h(\nabla \tau(\phi), d\phi).$$
 (17)

Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. To give a necessary and sufficient condition for the biharmonicity of the map ψ , we will calculate the bi-tension field $\tau_2(\psi)$. We get the following theorem:

Theorem 1.Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. Then, the bi-tension field of ψ is given by

$$\tau_2(\psi) = -2\left(\Delta\left(e\left(\psi\right)\right)\xi + 2d\psi\left(grad\left(e\left(\psi\right)\right)\right)\right). \tag{18}$$

The biharmonicity of ψ is given by the following result :

Corollary 1. Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. Then ψ is biharmonic if and only if

$$\Delta(e(\psi))\xi + 2d\psi(grad(e(\psi))) = 0.$$

By calculating of the term $\eta \circ \tau_2(\psi)$, we obtain immediately the following Corollary :

Corollary 2.Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. If the map ψ is biharmonic, then the function $e(\psi)$ is harmonic.

In particular, if the function $e(\psi)$ is constant we obtain the following result :

Corollary 3. Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. If the function $e(\psi)$ is constant, then ψ is biharmonic.

As an application of Theorem 1, we construct an example of biharmonic non-harmonic map.

Example 1. Let $(N = \mathbb{R}^2, J, h = dx^2 + dy^2)$ be a complex manifold with complex structure $J\left(\frac{\partial}{\partial x}\right) = \frac{\partial}{\partial y}, J\left(\frac{\partial}{\partial y}\right) = -\frac{\partial}{\partial x}$ and $(M = \mathbb{R}^3 \setminus \{(0,0,0)\}, \phi, \xi, \eta, g)$ be a Kenmotsu manifold with orthonormal basis

$$e_{1} = x_{3} \frac{\partial}{\partial x_{1}},$$

$$e_{2} = x_{3} \frac{\partial}{\partial x_{2}},$$

$$e_{3} = -x_{3} \frac{\partial}{\partial x_{3}},$$

$$\xi = e_{3},$$

$$\eta(X) = g(X, \xi),$$

where

$$g = \frac{1}{x_3^2} \left(dx_1^2 + dx_2^2 + dx_3^2 \right)$$

and

$$\phi = \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Let $\psi: N \longrightarrow M$ be a (J, ϕ) -holomorphic map defined by

$$\psi(x,y) = (\psi_1(x,y), \psi_2(x,y), \psi_3(x,y)).$$

We have

$$\frac{\partial \psi_1}{\partial x} = -\frac{\partial \psi_2}{\partial y},$$

$$\frac{\partial \psi_1}{\partial y} = \frac{\partial \psi_2}{\partial x}$$

and

$$\frac{\partial \psi_3}{\partial x} = \frac{\partial \psi_3}{\partial y} = 0.$$

The last equation give $\psi_3(x,y) = C$, where C is the real constant. A simple calculate give

$$\begin{split} e\left(\psi\right) &= \frac{1}{C^2} \left(\left(\frac{\partial \psi_1}{\partial x}\right)^2 + \left(\frac{\partial \psi_1}{\partial y}\right)^2 \right), \\ grad\left(e\left(\psi\right)\right) &= \frac{2}{C^2} \left(\frac{\partial \psi_1}{\partial x} \left(\frac{\partial^2 \psi_1}{\partial x^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) \frac{\partial}{\partial x} + \frac{\partial \psi_1}{\partial y} \left(\frac{\partial^2 \psi_1}{\partial y^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) \frac{\partial}{\partial y} \right), \\ d\psi\left(grad\left(e\left(\psi\right)\right)\right) &= \frac{2}{C^2} \left(\left(\frac{\partial \psi_1}{\partial x}\right)^2 \left(\frac{\partial^2 \psi_1}{\partial x^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) + \left(\frac{\partial \psi_1}{\partial y}\right)^2 \left(\frac{\partial^2 \psi_1}{\partial y^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) \right) \frac{\partial}{\partial x_1} \\ &+ \frac{2}{C^2} \left(\frac{\partial \psi_1}{\partial x} \frac{\partial \psi_2}{\partial x} \left(\frac{\partial^2 \psi_1}{\partial x^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) + \frac{\partial \psi_1}{\partial y} \frac{\partial \psi_2}{\partial y} \left(\frac{\partial^2 \psi_1}{\partial y^2} + \frac{\partial^2 \psi_1}{\partial x \partial y}\right) \right) \frac{\partial}{\partial x_2} \end{split}$$

and

$$\Delta (e(\psi)) = \frac{2}{C^2} \left(\left(\frac{\partial^2 \psi_1}{\partial x^2} \right)^2 + \left(\frac{\partial^2 \psi_2}{\partial x^2} \right)^2 + \left(\frac{\partial^2 \psi_1}{\partial y^2} \right)^2 + \left(\frac{\partial^2 \psi_2}{\partial y^2} \right)^2 \right) + \frac{2}{C^2} \left(\frac{\partial \psi_1}{\partial x} \left(\frac{\partial^3 \psi_1}{\partial x^3} - \frac{\partial^3 \psi_2}{\partial y^3} \right) + \frac{\partial \psi_2}{\partial x} \left(\frac{\partial^3 \psi_1}{\partial y^3} + \frac{\partial^3 \psi_2}{\partial x^3} \right) \right)$$

In particular, if we consider $\psi_1(x,y) = \psi_1(x)$ and $\psi_2(x,y) = \psi_2(y)$, we obtain

$$\psi(x,y) = (Ax + B, -Ay + D, C),$$

where A,B,C and D are a real constants $(A,C\neq 0)$. In this case $e\left(\psi\right)=\frac{A^{2}}{C^{2}}$ and the map ψ is biharmonic non-harmonic.

Proposition 2. Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. Then, the stress bi-energy tensor of ψ is given by

$$S_2(\psi) = 6\left(e(\psi)\right)^2 h - 4e(\psi)g(d\psi(\cdot), d\psi(\cdot)),\tag{19}$$

where $g(d\psi(\cdot), d\psi(\cdot))(X,Y) = g(d\psi(X), d\psi(Y))$, and the trace of $S_2(\psi)$ is given by

$$Tr_h S_2(\psi) = (12n - 8)(e(\psi))^2$$
.

3 Proof of results.

Proof of Proposition 1. Let $\phi:(M^m,g)\to (N^n,h)$ be a smooth map. By definition, we have

$$S_2(\phi) = \left(\frac{-1}{2}|\tau(\phi)|^2 + divh(\tau(\phi), d\phi)\right)g - 2symh(\nabla \tau(\phi), d\phi). \tag{20}$$



Let $(e_i)_{i=1}^m$ is a local orthonormal frame field on M. For the term $divh(\tau(\phi), d\phi)$, we have

$$\begin{split} \operatorname{divh} \left(\tau \left(\phi \right), \operatorname{d} \phi \right) &= \left(\nabla_{e_i} h \left(\tau \left(\phi \right), \operatorname{d} \phi \right) \right) \left(e_i \right) \\ &= e_i \left(h \left(\tau \left(\phi \right), \operatorname{d} \phi \left(e_i \right) \right) \right) - h \left(\tau \left(\phi \right), \operatorname{d} \phi \left(\nabla_{e_i} e_i \right) \right) \\ &= h \left(\nabla_{e_i}^{\phi} \tau \left(\phi \right), \operatorname{d} \phi \left(e_i \right) \right) + h \left(\tau \left(\phi \right), \nabla_{e_i}^{\phi} \operatorname{d} \phi \left(e_i \right) \right) \\ &- h \left(\tau \left(\phi \right), \operatorname{d} \phi \left(\nabla_{e_i} e_i \right) \right) \\ &= h \left(\nabla_{e_i}^{\phi} \tau \left(\phi \right), \operatorname{d} \phi \left(e_i \right) \right) + h \left(\tau \left(\phi \right), \nabla_{e_i}^{\phi} \operatorname{d} \phi \left(e_i \right) - \operatorname{d} \phi \left(\nabla_{e_i} e_i \right) \right). \end{split}$$

But we know that

$$h\left(\nabla_{e_{i}}^{\phi}\tau\left(\phi\right),d\phi\left(e_{i}\right)\right)=Tr_{g}h\left(\nabla\tau\left(\phi\right),d\phi\right)$$

and

$$h\left(\tau\left(\phi\right),\nabla_{e_{i}}^{\phi}d\phi\left(e_{i}\right)-d\phi\left(\nabla_{e_{i}}e_{i}\right)\right)=h\left(\tau\left(\phi\right),\tau\left(\phi\right)\right)=\left|\tau\left(\phi\right)\right|^{2},$$

it follows that

$$divh(\tau(\phi), d\phi) = Tr_{g}h(\nabla \tau(\phi), d\phi) + |\tau(\phi)|^{2}$$
(21)

If we replace (21) in (20), we obtain

$$S_2(\phi) = \left(\frac{1}{2} |\tau(\phi)|^2 + Tr_g h(\nabla \tau(\phi), d\phi)\right) g - 2 symh(\nabla \tau(\phi), d\phi).$$

By definition the trace of $S_2(\phi)$, we have

$$\begin{split} Tr_g S_2(\phi) &= S_2(\phi) \left(e_i, e_i \right) \\ &= \left(\frac{1}{2} \left| \tau(\phi) \right|^2 + Tr_g h(\nabla \tau(\phi), d\phi) \right) g\left(e_i, e_i \right) \\ &- 2 symh(\nabla \tau(\phi), d\phi) \left(e_i, e_i \right) \\ &= m \left(\frac{1}{2} \left| \tau(\phi) \right|^2 + Tr_g h(\nabla \tau(\phi), d\phi) \right) - 2 h(\nabla_{e_i} \tau(\phi), d\phi \left(e_i \right)) \\ &= \frac{m}{2} \left| \tau(\phi) \right|^2 + m Tr_g h(\nabla \tau(\phi), d\phi) - 2 Tr_g h(\nabla \tau(\phi), d\phi) \end{split}$$

which gives

$$Tr_gS_2(\phi) = \frac{m}{2} |\tau(\phi)|^2 + (m-2)Tr_gh(\nabla \tau(\phi), d\phi).$$

The proof of Proposition 1 is complete. To prove Theorem 1, we need the following Lemma:

Lemma 1. Let (N^{2n}, J, h) be a Kahlerian manifold, (M, ϕ, ξ, η, g) be a Kenmotsu manifold and $\psi : N \longrightarrow M$ be a (J, ϕ) -holomorphic map. Then, we have for any $X \in \Gamma(TM)$

$$(\eta \circ d\psi)(X) = 0.$$

Proof of Lemma 1. Since $\psi: N \longrightarrow M$ is (J, ϕ) -holomorphic map, then for any $X \in \Gamma(TN)$, we have

$$(d\psi \circ J)(X) = (\phi \circ d\psi)(X).$$

Using the fact that (N^{2n}, J, h) is a Kahlerian manifold, then

$$J^{2}(X) = -X.$$

As (M, ϕ, ξ, η, g) is a Kenmotsu manifold and by using the equation 1, we have

$$\eta \circ \phi = 0.$$

Then, we obtain

$$\begin{split} \left(\eta \circ d\psi\right)(X) &= \left(\eta \circ d\psi\right)\left(-J^2\left(X\right)\right) \\ &= -\eta \circ d\psi \circ J(J(X)) \\ &= -\eta \circ \phi \circ d\psi(J(X)) \\ &= 0. \end{split}$$

We are now able to prove Theorem 1.

Proof of Theorem 1. By definition, we have

$$\tau_2(\psi) = -Tr_h(\nabla^{\psi})^2 \tau(\psi) - Tr_h R^M(\tau(\psi), d\psi) d\psi.$$

The tension field of ψ is given by (see [16])

$$\tau(\boldsymbol{\psi}) = 2e(\boldsymbol{\psi})\,\boldsymbol{\xi},$$

then

$$\tau_2(\psi) = -2\left(Tr_h(\nabla^{\psi})^2 e(\psi)\xi + Tr_h R^M(e(\psi)\xi, d\psi)d\psi\right). \tag{22}$$

In the first, we will simplify the term $Tr_h(\nabla^{\psi})^2 e(\psi) \xi$. Let $(e_i)_{i=1}^{2n}$ is a local orthonormal frame field on N, we have

$$Tr_g(\nabla^{\psi})^2 e(\psi) \xi = \nabla^{\psi}_{e_i} \nabla^{\psi}_{e_i} e(\psi) \xi - \nabla^{\psi}_{\nabla^{\omega}_{o}, e_i} e(\psi) \xi.$$
(23)

A direct calculation gives

$$\begin{split} \nabla^{\psi}_{e_{i}} \nabla^{\psi}_{e_{i}} e\left(\psi\right) \xi &= \nabla^{\psi}_{e_{i}} \left(e\left(\psi\right) \nabla^{\psi}_{e_{i}} \xi\right) + \nabla^{\psi}_{e_{i}} \left(e_{i}\left(e\left(\psi\right)\right) \xi\right) \\ &= e\left(\psi\right) \nabla^{\psi}_{e_{i}} \nabla^{\psi}_{e_{i}} \xi + e_{i}\left(e\left(\psi\right)\right) \nabla^{\psi}_{e_{i}} \xi + e_{i}\left(e\left(\psi\right)\right) \nabla^{\psi}_{e_{i}} \xi \\ &+ e_{i}\left(e_{i}\left(e\left(\psi\right)\right)\right) \xi \\ &= e\left(\psi\right) \nabla^{\psi}_{e_{i}} \nabla^{\psi}_{e_{i}} \xi + 2 \nabla^{\psi}_{\mathrm{grad}\left(e\left(\psi\right)\right)} \xi + e_{i}\left(e_{i}\left(e\left(\psi\right)\right)\right) \xi \end{split}$$

and

$$abla_{
abla_{e_{i}}e_{i}}^{\psi}e\left(\psi\right)\xi=e\left(\psi\right)
abla_{
abla_{e_{i}}e_{i}}^{\psi}\xi+
abla_{e_{i}}e_{i}\left(e\left(\psi\right)\right)\xi.$$

Using the fact that

$$\Delta\left(e\left(\psi\right)\right)=e_{i}\left(e_{i}\left(e\left(\psi\right)\right)\right)-\left(\nabla_{e_{i}}^{N}e_{i}\right)\left(e\left(\psi\right)\right)$$

and

$$Tr_{g}\left(\nabla^{\psi}\right)^{2}\xi=\nabla_{e_{i}}^{\psi}\nabla_{e_{i}}^{\psi}\xi-\nabla_{\nabla_{e_{i}}^{N}e_{i}}^{\psi}\xi,$$



we deduce that

$$Tr_{h}(\nabla^{\psi})^{2} e(\psi) \xi = e(\psi) Tr_{h}(\nabla^{\psi})^{2} \xi + \Delta (e(\psi)) \xi + 2\nabla^{\psi}_{\operatorname{grad}(e(\psi))} \xi.$$

$$(24)$$

Calculate the term $Tr_h(\nabla^{\psi})^2 \xi$, we have

$$Tr_h \left(
abla^\psi
ight)^2 \xi =
abla^\psi_{e_i}
abla^\psi_{e_i} \xi -
abla^\psi_{
abla_e,e_i} \xi.$$

By using equation (6), we obtain

$$abla_{e_i}^{\psi} \xi =
abla_{d\psi(e_i)} \xi = d\psi(e_i) - (\eta \circ d\psi)(e_i) \xi,$$

and by Lemma 1, we have

$$(\boldsymbol{\eta} \circ d\boldsymbol{\psi})(e_i) = 0,$$

then

$$\nabla_{e_i}^{\psi} \xi = d\psi(e_i)$$

wish gives us

$$\nabla_{e_i}^{\psi} \nabla_{e_i}^{\psi} \xi = \nabla_{e_i}^{\psi} d\psi(e_i)$$

and

$$\begin{split} \nabla^{\psi}_{\nabla_{e_{i}}e_{i}}\xi &= \nabla_{d\psi\left(\nabla_{e_{i}}e_{i}\right)}\xi \\ &= d\psi\left(\nabla_{e_{i}}e_{i}\right) - \left(\eta \circ d\psi\right)\left(\nabla_{e_{i}}e_{i}\right)\xi \\ &= d\psi\left(\nabla_{e_{i}}e_{i}\right), \end{split}$$

we conclude that

$$Tr_h \left(\nabla^{\psi}\right)^2 \xi = \nabla^{\psi}_{e_i} d\psi(e_i) - d\psi(\nabla_{e_i} e_i) = \tau(\psi).$$

Then

$$Tr_h(\nabla^{\psi})^2 \xi = 2e(\psi) \xi.$$

Finally, we obtain

$$\mathit{Tr}_{\hbar}\left(\nabla^{\psi}\right)^{2}e\left(\psi\right)\xi=2\left(e\left(\psi\right)\right)^{2}\xi+\Delta\left(e\left(\psi\right)\right)\xi+2\nabla_{\mathrm{grad}\left(e\left(\psi\right)\right)}^{\psi}\xi$$

Now simplify the term $abla_{\operatorname{grad}(e(\psi))}^{\psi} \xi$, we have

$$\begin{split} \nabla^{\psi}_{\mathrm{grad}(e(\psi))} \xi &= \nabla^{M}_{d\psi(\mathrm{grad}(e(\psi)))} \xi \\ &= d\psi \left(\mathrm{grad}\left(e\left(\psi \right) \right) \right) - \eta \left(d\psi \left(\mathrm{grad}\left(e\left(\psi \right) \right) \right) \right) \xi \\ &= d\psi \left(\mathrm{grad}\left(e\left(\psi \right) \right) \right) - \left(\eta \circ d\psi \right) \left(\mathrm{grad}\left(e\left(\psi \right) \right) \right) \xi \\ &= d\psi \left(\mathrm{grad}\left(e\left(\psi \right) \right) \right). \end{split}$$

which finally gives us

$$Tr_{h}(\nabla^{\psi})^{2} e(\psi) \xi = 2(e(\psi))^{2} \xi + \Delta(e(\psi)) \xi + 2d\psi(\operatorname{grad}(e(\psi))). \tag{25}$$

To complete the proof, look at the term $Tr_{g}R^{M}\left(e\left(\psi\right) \xi,d\psi\right) d\psi$, we have

$$Tr_{h}R^{M}(e(\psi)\xi,d\psi)d\psi = e(\psi)Tr_{h}R^{M}(\xi,d\psi)d\psi$$
$$= e(\psi)R^{N}(\xi,d\psi(e_{i}))d\psi(e_{i}).$$

Using equation (9), we obtain

$$R^{N}(\xi, d\psi(e_{i})) d\psi(e_{i}) = \eta (d\psi(e_{i})) - g(d\psi(e_{i}), d\psi(e_{i})) \xi$$
$$= (\eta \circ d\psi) (e_{i}) - 2e(\psi) \xi$$
$$= -2e(\psi) \xi.$$

Then

$$Tr_h R^N \left(e\left(\psi \right) \xi, d\psi \right) d\psi = -2 \left(e\left(\psi \right) \right)^2 \xi. \tag{26}$$

If we replace (25) and (26) in (22), wa arrive at

$$\tau_2(\psi) = -2\left(\Delta\left(e\left(\psi\right)\right)\xi + 2d\psi\left(\operatorname{grad}\left(e\left(\psi\right)\right)\right)\right).$$

This completes the proof of Theorem 1. Then $\psi: N \longrightarrow M$ is biharmonic if and only if

$$\Delta (e(\psi)) \xi + 2d\psi(\operatorname{grad}(e(\psi))) = 0.$$

Proof of Corollary 2. If we assume that the map $\psi: N \longrightarrow M$ is biharmonic, by Theorem 1, we have

$$\Delta(e(\psi))\xi + 2d\psi(\operatorname{grad}(e(\psi))) = 0.$$

Then

$$\eta \circ (\Delta(e(\psi))\xi + 2d\psi(\operatorname{grad}(e(\psi)))) = 0,$$

it follows that

$$\Delta\left(e\left(\psi\right)\right)\eta\left(\xi\right)+2\left(\eta\circ d\psi\right)\left(\operatorname{grad}\left(e\left(\psi\right)\right)\right)=0.$$

By using the equation (1), we have

$$\eta(\xi) = 1$$

and by Lemma 1, we have

$$(\eta \circ d\psi)(\operatorname{grad}(e(\psi))) = 0.$$

Which finally gives

$$\Delta (e(\psi)) = 0.$$

We deduce that the function $e(\psi)$ is harmonic.

Proof of Proposition 2. By definition, for any $X,Y \in \Gamma(TN)$ we have

$$S_{2}(\psi)(X,Y) = \left(\frac{1}{2}|\tau(\psi)|^{2} + Tr_{h}g(\nabla\tau(\psi),d\psi)\right)h(X,Y)$$
$$-2symg(\nabla\tau(\psi),d\psi)(X,Y).$$



Since

$$\tau(\psi) = 2e(\psi)\xi,$$

it follows that

$$|\tau(\psi)|^2 = 4\left(e\left(\psi\right)\right)^2.$$

For the term $Tr_{h}g(\nabla \tau(\psi), d\psi)$, we have

$$\begin{split} Tr_{h}g\left(\nabla\tau\left(\psi\right),d\psi\right) &= 2g\left(\nabla_{e_{i}}e\left(\psi\right)\xi,d\psi\left(e_{i}\right)\right) \\ &= 2e\left(\psi\right)g\left(\nabla_{e_{i}}^{\psi}\xi,d\psi\left(e_{i}\right)\right) \\ &+ 2e_{i}\left(e\left(\psi\right)\right)g\left(\xi,d\psi\left(e_{i}\right)\right), \end{split}$$

where $(e_i)_{i=1}^{2n}$ is a local orthonormal frame field on N. By equation (6), we obtain

$$abla_{di}^{\psi}\xi = \nabla_{d\psi(e_i)}^{M}\xi = d\psi(e_i) - \eta\left(d\psi(e_i)\right)\xi = d\psi(e_i)$$

and by (3), we have

$$g(\xi, d\psi(e_i)) = \eta(d\psi(e_i)) = 0,$$

then

$$Tr_{h}g\left(\nabla\tau\left(\psi\right),d\psi\right)=2e\left(\psi\right)g\left(d\psi\left(e_{i}\right),d\psi\left(e_{i}\right)\right)=4\left(e\left(\psi\right)\right)^{2}.$$

it follow that

$$\frac{1}{2}\left|\tau\left(\psi\right)\right|^{2}+Tr_{h}g\left(\nabla\tau\left(\psi\right),d\psi\right)=6\left(e\left(\psi\right)\right)^{2}.$$

Finally, for the term $symg(\nabla \tau(\psi), d\psi)(X, Y)$, we have

$$symg\left(\nabla\tau\left(\psi\right),d\psi\right)\left(X,Y\right) = \frac{1}{2}g\left(\nabla_{X}^{\psi}\tau\left(\psi\right),d\psi\left(Y\right)\right) + \frac{1}{2}g\left(\nabla_{Y}^{\psi}\tau\left(\psi\right),d\psi\left(X\right)\right).$$

A simple calculate gives

$$\begin{split} \frac{1}{2}g\left(\nabla_{X}^{\psi}\tau\left(\psi\right),d\psi\left(Y\right)\right) &= g\left(\nabla_{X}^{\psi}e\left(\psi\right)\xi,d\psi\left(Y\right)\right) \\ &= e\left(\psi\right)g\left(\nabla_{X}^{\psi}\xi,d\psi\left(Y\right)\right) + X\left(e\left(\psi\right)\right)g\left(\xi,d\psi\left(Y\right)\right), \end{split}$$

but by (6), we have

$$\nabla_{\boldsymbol{X}}^{\psi}\boldsymbol{\xi} = \nabla_{d\psi(\boldsymbol{X})}^{M}\boldsymbol{\xi} = d\psi(\boldsymbol{X}) - \eta\left(d\psi(\boldsymbol{X})\right)\boldsymbol{\xi} = d\psi(\boldsymbol{X})$$

and by (3), we obtain

$$g(\xi, d\psi(Y)) = \eta(d\psi(Y)) = 0,$$

then

$$\frac{1}{2}g\left(\nabla_{X}^{\psi}\tau\left(\psi\right),d\psi\left(Y\right)\right)=g\left(d\psi\left(X\right),d\psi\left(Y\right)\right).$$

A similar calculate gives

$$\frac{1}{2}g\left(\nabla_{Y}^{\psi}\tau\left(\psi\right),d\psi\left(X\right)\right)=g\left(d\psi\left(X\right),d\psi\left(Y\right)\right).$$

We deduce that

$$symg\left(\nabla\tau\left(\psi\right),d\psi\right)\left(X,Y\right)=2g\left(d\psi\left(X\right),d\psi\left(Y\right)\right),$$

then

$$S_2(\psi)(X,Y) = 6(e(\psi))^2 g(X,Y) - 4g(d\psi(X), d\psi(Y)).$$

finally, the stress bi-energy tensor is given by

$$S_2(\psi) = 6(e(\psi))^2 g - 4g(d\psi(\cdot), d\psi(\cdot)),$$

where $g(d\psi(\cdot), d\psi(\cdot))(X, Y) = g(d\psi(X), d\psi(Y))$. Now calculate $Tr_hS_2(\psi)$, we have

$$Tr_h S_2(\psi) = S_2(\psi) (e_i, e_i)$$

$$= 6(e(\psi))^2 g(e_i, e_i) - 4g(d\psi(e_i), d\psi(e_i))$$

$$= 12n(e(\psi))^2 - 4|d\psi|^2$$

$$= 12n(e(\psi))^2 - 8(e(\psi))^2$$

$$= (12n - 8)(e(\psi))^2.$$

We conclude that $Tr_gS_2(\psi) = 0$ if and only if the map ψ is constant.

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