

Investigation of spectral analysis of matrix quantum difference equations with spectral singularities

Yelda Aygar*†

Abstract

In this paper, we investigate the Jost solution, the continuous spectrum, the eigenvalues and the spectral singularities of a nonselfadjoint matrix-valued q -difference equation of second order with spectral singularities.

Keywords: Quantum difference equation, Discrete spectrum, Spectral theory, Spectral singularity, Eigenvalue.

2000 AMS Classification: 39A05, 39A70, 47A05, 47A10, 47A55.

Received : 17.02.2015 *Accepted :* 02.07.2015 *Doi :* 10.15672/HJMS.20164513107

1. Introduction

Spectral analysis of nonselfadjoint differential equations including Sturm–Liouville, Schrödinger and Klein–Gordon equations has been treated by various authors since 1960 [23, 9, 11, 22, 12]. Study of spectral theory of nonselfadjoint discrete Schrödinger and Dirac equations were obtained in [1, 20, 8, 10, 7]. Also, spectral analysis of these equations in self-adjoint case is well-known [4, 5]. In addition to differential and discrete equations, spectral theory of q -difference equations has been investigated in recent years [2, 3], and important generalizations and results were given for dynamic equations including q -difference equations as a special case in [14, 13].

Some problems of spectral theory of differential and difference equations with matrix coefficients were studied in [15, 24, 18, 6]. But spectral analysis of the matrix q -difference equations with spectral singularities has not been investigated yet.

In this paper, we let $q > 1$ and use the notation $q^{\mathbb{N}_0} := \{q^n : n \in \mathbb{N}_0\}$, where \mathbb{N}_0 denotes the set of nonnegative integers. Let us introduce the Hilbert space $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$ consisting of all vector sequences $y \in \mathbb{C}^m$, ($y = y(t)$, $t \in q^{\mathbb{N}}$), such that $\sum_{t \in q^{\mathbb{N}}} \mu(t) \|y(t)\|_{\mathbb{C}^m}^2 < \infty$ with the inner product $\langle y, z \rangle_q := \sum_{t \in q^{\mathbb{N}}} \mu(t) (y(t), z(t))_{\mathbb{C}^m}$, where \mathbb{C}^m is m -dimensional ($m < \infty$) Euclidean space, $\mu(t) = (q - 1)t$ for all $t \in q^{\mathbb{N}}$, and $\|\cdot\|_{\mathbb{C}^m}$ and $(\cdot, \cdot)_{\mathbb{C}^m}$ denote

*University of Ankara, Faculty of Science, Department of Mathematics, 06100, Ankara, Turkey, Email: yaygar@science.ankara.edu.tr

†Corresponding Author.

the norm and inner product in \mathbb{C}^m , respectively. We denote by L the operator generated in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$ by the q -difference expression

$$(ly)(t) := qA(t)y(qt) + B(t)y(t) + A\left(\frac{t}{q}\right)y\left(\frac{t}{q}\right), \quad t \in q^{\mathbb{N}},$$

and the boundary condition $y(1) = 0$, where $A(t)$, $t \in q^{\mathbb{N}_0}$ and $B(t)$, $t \in q^{\mathbb{N}}$ are linear operators (matrices) acting in \mathbb{C}^m . Throughout the paper, we will assume that $A(t)$ is invertible and $A(t) \neq A^*(t)$ for all $t \in q^{\mathbb{N}_0}$. Furthermore $B(t) \neq B^*(t)$ for all $t \in q^{\mathbb{N}}$, where $*$ denotes the adjoint operator. It is clear that L is a nonselfadjoint operator in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$. Related to the operator L , we will consider the matrix q -difference equation of second order

$$(1.1) \quad qA(t)y(qt) + B(t)y(t) + A\left(\frac{t}{q}\right)y\left(\frac{t}{q}\right) = \lambda y(t), \quad t \in q^{\mathbb{N}},$$

where λ is a spectral parameter.

The set up of this paper is summarized as follows: Section 2 discusses the Jost solution of (1.1) and contains analytical properties and asymptotic behavior of this solution. In Section 3, we give the continuous spectrum of L , by using the Weyl compact perturbation theorem. In Section 4, we investigate the eigenvalues and the spectral singularities of L . In particular, we prove that L has a finite number of eigenvalues and spectral singularities with a finite multiplicity.

2. Jost solution of L

We assume that the matrix sequences $\{A(t)\}$ and $\{B(t)\}$, $t \in q^{\mathbb{N}}$ satisfy

$$(2.1) \quad \sum_{t \in q^{\mathbb{N}}} (\|I - A(t)\| + \|B(t)\|) < \infty,$$

where $\|\cdot\|$ denotes the matrix norm in \mathbb{C}^m and I is identity matrix. Let $F(\cdot, z)$, denotes the matrix solution of the q -difference equation

$$(2.2) \quad qA(t)y(qt) + B(t)y(t) + A\left(\frac{t}{q}\right)y\left(\frac{t}{q}\right) = 2\sqrt{q} \cos zy(t), \quad t \in q^{\mathbb{N}},$$

satisfying the condition

$$(2.3) \quad \lim_{t \rightarrow \infty} F(t, z) e^{i \frac{\ln t}{\ln q} z} \sqrt{\mu(t)} = I, \quad z \in \overline{\mathbb{C}}_+ := \{z \in \mathbb{C} : \text{Im } z \geq 0\}.$$

The solution $F(\cdot, z)$ is called the Jost solution of (2.2).

2.1. Theorem. *Assume (2.1). Let the solution $F(\cdot, z)$ be the Jost solution of (2.2). Then*

$$(2.4) \quad F(t, z) = \frac{e^{i \frac{\ln t}{\ln q} z}}{\sqrt{\mu(t)}} I + \sum_{s \in [qt, \infty) \cap q^{\mathbb{N}}} \sqrt{\frac{s}{qt}} \frac{\sin\left(\frac{\ln s - \ln t}{\ln q} z\right)}{\sin z} H(s),$$

where

$$H(s) := \left[I - A\left(\frac{s}{q}\right) \right] F\left(\frac{s}{q}, z\right) - B(s)F(s, z) + q[I - A(s)]F(qs, z).$$

Proof. Using (2.2), we obtain

$$(2.5) \quad F\left(\frac{t}{q}\right) + qF(qt) - 2\sqrt{q} \cos zF(t) = H(t).$$

Since $\frac{\exp\left(i\frac{\ln t}{\ln q}z\right)}{\sqrt{\mu(t)}}I$ and $\frac{\exp\left(-i\frac{\ln t}{\ln q}z\right)}{\sqrt{\mu(t)}}I$ are linearly independent solutions of the homogeneous equation

$$F\left(\frac{t}{q}\right) + qF(qt) - 2\sqrt{q}\cos zF(t) = 0,$$

we get the general solution of (2.5) by

$$(2.6) \quad F(t, z) = \frac{e^{i\frac{\ln t}{\ln q}z}}{\sqrt{\mu(t)}}\alpha + \frac{e^{-i\frac{\ln t}{\ln q}z}}{\sqrt{\mu(t)}}\beta + \sum_{s \in [qt, \infty) \cap q^{\mathbb{N}}} \sqrt{\frac{\mu(s)}{q}} \frac{1}{\sqrt{\mu(t)}} \frac{\sin\left(\frac{\ln s - \ln t}{\ln q}z\right)}{\sin z} H(s),$$

where α and β are constants in \mathbb{C}^m . Using (2.1), (2.3), and (2.6), we find $\alpha = I$ and $\beta = 0$. This completes the proof, i.e., $F(t, z)$ satisfies (2.4). \square

2.2. Theorem. *Assume (2.1). Then the Jost solution $F(\cdot, z)$ has a representation*

$$(2.7) \quad F(t, z) = T(t) \frac{e^{i\frac{\ln t}{\ln q}z}}{\sqrt{\mu(t)}} \left(I + \sum_{r \in q^{\mathbb{N}}} K(t, r) e^{i\frac{\ln r}{\ln q}z} \right), \quad t \in q^{\mathbb{N}_0}$$

where $z \in \overline{\mathbb{C}}_+$, $T(t)$ and $K(t, r)$ are expressed in terms of $\{A(t)\}$ and $\{B(t)\}$.

Proof. If we put $F(\cdot, z)$ defined by (2.7) into (2.2), then we have the relations

$$\begin{aligned} A(t)T(t) &= T(qt), \quad K(t, q) - K\left(\frac{t}{q}, q\right) = \frac{1}{\sqrt{q}}T^{-1}(t)B(t)T(t), \\ K\left(\frac{t}{q}, q^2\right) - K(t, q^2) &= T^{-1}(t) \left(T(t) - A^2(t)T(t) - \frac{1}{\sqrt{q}}B(t)T(t)K(t, q) \right), \\ K(t, rq^2) - K\left(\frac{t}{q}, rq^2\right) &= T^{-1}(t) \left(A^2(t)T(t)K(qt, r) + \frac{1}{\sqrt{q}}B(t)T(t)K(t, qr) \right) - K(t, r), \end{aligned}$$

and using these relations, we obtain

$$\begin{aligned} T(t) &= \prod_{p \in [t, \infty) \cap q^{\mathbb{N}}} [A(p)]^{-1}, \quad K(t, q) = -\frac{1}{\sqrt{q}} \sum_{p \in [qt, \infty) \cap q^{\mathbb{N}}} T^{-1}(p)B(p)T(p), \\ K(t, q^2) &= \sum_{p \in [qt, \infty) \cap q^{\mathbb{N}}} T^{-1}(p) \left[-\frac{1}{\sqrt{q}}B(p)T(p)K(p, q) + (I - A^2(p))T(p) \right], \\ K(t, rq^2) &= K(qt, r) + \sum_{p \in [qt, \infty) \cap q^{\mathbb{N}}} T^{-1}(p) [I - A^2(p)] T(p)K(qp, r) \\ &\quad - \frac{1}{\sqrt{q}} \sum_{p \in [qt, \infty) \cap q^{\mathbb{N}}} T^{-1}(p)B(p)T(p)K(p, qr), \end{aligned}$$

for $r \in q^{\mathbb{N}}$ and $t \in q^{\mathbb{N}_0}$. Due to the condition (2.1), the infinite product and the series in the definition of $T(t)$ and $K(t, r)$ are absolutely convergent. \square

Note that, in analogy to the Sturm–Liouville equation the function $F(1, z) := \frac{T(1)}{\sqrt{\mu(1)}} \left(I + \sum_{r \in q^{\mathbb{N}}} K(1, r) e^{i\frac{\ln r}{\ln q}z} \right)$ is called the Jost function.

2.3. Theorem. *Assume*

$$(2.8) \quad \sum_{t \in q^{\mathbb{N}}} \frac{\ln t}{\ln q} (\|I - A(t)\| + \|B(t)\|) < \infty.$$

Then the Jost solution $F(\cdot, z)$ is continuous in $\overline{\mathbb{C}}_+$ and analytic with respect to z in $\mathbb{C}_+ := \{z \in \mathbb{C} : \text{Im } z > 0\}$.

Proof. Using the equalities for $K(t, r)$ given in Theorem 2.2 and mathematical induction, we get

$$(2.9) \quad \|K(t, r)\| \leq C \sum_{p \in \left[tq^{\lfloor \frac{\ln r}{2 \ln q} \rfloor}, \infty \right) \cap q^{\mathbb{N}}} (\|I - A(p)\| + \|B(p)\|),$$

where $C > 0$ is a constant and $\lfloor \frac{\ln r}{2 \ln q} \rfloor$ is the integer part of $\frac{\ln r}{2 \ln q}$. From (2.8) and (2.9), we get that the series

$$\sum_{r \in q^{\mathbb{N}}} K(t, r) e^{i \frac{\ln r}{\ln q} z}, \quad \sum_{r \in q^{\mathbb{N}}} \frac{\ln r}{\ln q} K(t, r) e^{i \frac{\ln r}{\ln q} z}$$

are absolutely convergent in $\overline{\mathbb{C}}_+$ and in \mathbb{C}_+ , respectively. This completes the proof. \square

2.4. Theorem. Under the condition (2.8), the Jost solution satisfies

$$(2.10) \quad F(t, z) = \frac{e^{i \frac{\ln t}{\ln q} z}}{\sqrt{\mu(t)}} (I + o(1)), \quad z \in \overline{\mathbb{C}}_+, \quad t \rightarrow \infty,$$

$$(2.11) \quad F(t, z) = T(t) \frac{e^{i \frac{\ln t}{\ln q} z}}{\sqrt{\mu(t)}} (I + o(1)), \quad t \in q^{\mathbb{N}_0}, \quad \text{Im } z \rightarrow \infty.$$

Proof. It follows from the definition of $T(t)$, (2.8), and (2.9) that

$$(2.12) \quad \lim_{t \rightarrow \infty} T(t) = I,$$

and

$$(2.13) \quad \sum_{r \in q^{\mathbb{N}}} K(t, r) e^{i \frac{\ln r}{\ln q} z} = o(1), \quad z \in \overline{\mathbb{C}}_+, \quad t \rightarrow \infty.$$

From (2.7), (2.12), and (2.13), we get (2.10). Using (2.8) and (2.9), we have

$$(2.14) \quad \sum_{r \in q^{\mathbb{N}}} K(t, r) e^{i \frac{\ln r}{\ln q} z} = o(1), \quad z \in \overline{\mathbb{C}}_+, \quad \text{Im } z \rightarrow \infty.$$

From (2.7) and (2.14), we get (2.11). \square

3. Continuous spectrum of L

Let L_1 and L_2 denote the q -difference operators generated in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$ by the q -difference expressions

$$(l_1 y)(t) = qy(qt) + y\left(\frac{t}{q}\right)$$

and

$$(l_2 y)(t) = q[A(t) - I]y(qt) + B(t)y(t) + \left[A\left(\frac{t}{q}\right) - I\right]y\left(\frac{t}{q}\right)$$

with the boundary condition $y(1) = 0$, respectively. It is clear that $L = L_1 + L_2$.

3.1. Lemma. The operator L_1 is self-adjoint in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$.

Proof. Since

$$\|L_1 y\|_q \leq 2\sqrt{q}\|y\|_q$$

for all $y \in \ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$, L_1 is bounded in the Hilbert space $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$, and since

$$\begin{aligned} \langle l_1 y, z \rangle_q &= \sum_{t \in q^{\mathbb{N}}} \mu(t) (z(t))^* \left(qy(qt) + y\left(\frac{t}{q}\right) \right) \\ &= \sum_{t \in q^{\mathbb{N}}} \mu(t) \left(qz(qt) + z\left(\frac{t}{q}\right) \right)^* y(t) = \langle y, l_1 z \rangle_q, \end{aligned}$$

the operator L_1 is self-adjoint in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$. \square

3.2. Theorem. *Assume (2.8). Then $\sigma_c(L) = [-2\sqrt{q}, 2\sqrt{q}]$, where $\sigma_c(L)$ denotes the continuous spectrum of L .*

Proof. It is easy to see that L_1 has no eigenvalues, so the spectrum of the operator L_1 consists only its continuous spectrum and

$$\sigma(L_1) = \sigma_c(L_1) = [-2\sqrt{q}, 2\sqrt{q}],$$

where $\sigma(L_1)$ denotes the spectrum of the operator L_1 . Using (2.8), we find that L_2 is compact operator in $\ell_2(q^{\mathbb{N}}, \mathbb{C}^m)$ [21]. Since $L = L_1 + L_2$ and $L_1 = (L_1)^*$, we obtain that

$$\sigma_c(L) = \sigma_c(L_1) = [-2\sqrt{q}, 2\sqrt{q}]$$

by using Weyl's theorem of a compact perturbation [19, p.13]. \square

4. Eigenvalues and spectral singularities of L

If we define

$$(4.1) \quad f(z) := \det F(1, z), \quad z \in \overline{\mathbb{C}}_+,$$

then the function f is analytic in \mathbb{C}_+ , $f(z) = f(z + 2\pi)$ and is continuous in $\overline{\mathbb{C}}_+$. Let us define the semi-strips $P_0 = \{z \in \mathbb{C}_+ : 0 \leq \operatorname{Re} z \leq 2\pi\}$ and $P = P_0 \cup [0, 2\pi]$. We will denote the set of all eigenvalues and spectral singularities of L by $\sigma_d(L)$ and $\sigma_{ss}(L)$, respectively. From the definitions of eigenvalues and spectral singularities of nonselfadjoint operators[22, 23], we have

$$(4.2) \quad \sigma_d(L) = \{\lambda \in \mathbb{C} : \lambda = 2\sqrt{q} \cos z, z \in P_0, f(z) = 0\},$$

$$(4.3) \quad \sigma_{ss}(L) = \{\lambda \in \mathbb{C} : \lambda = 2\sqrt{q} \cos z, z \in [0, 2\pi], f(z) = 0\} \setminus \{0\}.$$

4.1. Theorem. *Assume (2.8). Then*

- i) *the set $\sigma_d(L)$ is bounded and countable, and its limit points lie only in the interval $[-2\sqrt{q}, 2\sqrt{q}]$,*
- ii) *$\sigma_{ss}(L) \subset [-2\sqrt{q}, 2\sqrt{q}]$ and the linear Lebesgue measure of the set $\sigma_{ss}(L)$ is zero.*

Proof. In order to investigate the quantitative properties of eigenvalues and spectral singularities of L , it is necessary to discuss the quantitative properties of zeros of f in P from (4.2) and (4.3). Using (2.11) and (4.1), we get

$$(4.4) \quad f(z) = \det T(1) \frac{1}{\mu(1)} [I + o(1)], \quad \operatorname{Im} z > 0, \quad z \in P_0, \quad \operatorname{Im} z \rightarrow \infty,$$

where $\det T(1) \neq 0$. From (4.4), we get the boundedness of zeros of f in P_0 . Since f is a 2π -periodic function and is analytic in \mathbb{C}_+ , we obtain that f has at most a countable number of zeros in P_0 . By the uniqueness of analytic functions, we find that the the limit points of zeros of f in P_0 can lie only in $[0, 2\pi]$. We get $\sigma_{ss}(L) \subset [-2\sqrt{q}, 2\sqrt{q}]$ using

(4.3). Since $f(z) \neq 0$ for all $z \in \mathbb{C}_+$, we get that the linear Lebesgue measure of the set of zeros of f on real axis is not positive, by using the boundary uniqueness theorem of analytic functions [17], i.e., the linear Lebesgue measure of the $\sigma_{ss}(L)$ is zero. \square

4.2. Definition. The multiplicity of a zero of f in P is called the multiplicity of the corresponding eigenvalue or spectral singularity of L .

4.3. Theorem. *If, for some $\varepsilon > 0$,*

$$(4.5) \quad \sup_{t \in q^{\mathbb{N}}} \left\{ e^{\varepsilon \frac{\ln t}{\ln q}} (\|I - A(t)\| + \|B(t)\|) \right\} < \infty$$

then the operator L has a finite number of eigenvalues and spectral singularities, and each of them is of finite multiplicity.

Proof. Since $F(1, z) = \frac{T(1)}{\sqrt{q-1}} \left(I + \sum_{r \in q^{\mathbb{N}}} K(1, r) e^{i \frac{\ln r}{\ln q} z} \right)$, using (2.9) and (4.5), we get that

$$(4.6) \quad \|K(1, r)\| \leq D e^{-\frac{\varepsilon}{4} \frac{\ln r}{\ln q}}, \quad r \in q^{\mathbb{N}},$$

where $D > 0$ is a constant. From (4.1) and (4.6), we obtain that the function f has an analytic continuation to the half-plane $\text{Im } z > -\frac{\varepsilon}{4}$. Because the series

$$\sum_{r \in q^{\mathbb{N}}} iK(1, r) \frac{\ln r}{\ln q} e^{i \frac{\ln r}{\ln q} z}$$

is uniformly convergent in $\text{Im } z > -\frac{\varepsilon}{4}$. Since f is a 2π periodic function, the limit points of its zeros in P cannot lie in $[0, 2\pi]$. Using Theorem 4.1, we find that the bounded sets $\sigma_d(L)$ and $\sigma_{ss}(L)$ have no limit points, i.e., the sets $\sigma_d(L)$ and $\sigma_{ss}(L)$ have a finite number of elements. From the analyticity of f in $\text{Im } z > -\frac{\varepsilon}{4}$, we get that all zeros of f in P have a finite multiplicity. \square

References

- [1] Adivar, M. and Bairamov, E. *Spectral properties of non-selfadjoint difference operators*, J. Math. Anal. Appl. **261** (2), 461-478, 2001.
- [2] Adivar, M. and Bohner, M. *Spectral analysis of q -difference equations with spectral singularities*, Math. Comput. Modelling **43** (7), 695-703, 2006.
- [3] Adivar, M. and Bohner, M. *Spectrum and principal vectors of second order q -difference equations*, Indian J. Math. **48** (1), 17-33, 2006.
- [4] Agarwal, R.P. *Difference equations and inequalities*, Monographs and Textbooks in Pure and Applied Mathematics (Marcel Dekker Inc., New York, 2000).
- [5] Agarwal, R.P. and Wong, P.J.Y. *Advanced topics in difference equations* (Kluwer Academic Publishers Group, Dordrecht, 1997).
- [6] Aygar, Y. and Bairamov, B. *Jost solution and the spectral properties of the matrix-valued difference operators*, Appl. Math. Comput. **218** (19), 9676-9681, 2012.
- [7] Bairamov, E. Aygar, Y. and Koprubasi, T. *The spectrum of eigenparameter-dependent discrete Sturm-Liouville equations*, J. Comput. Appl. Math. **235**, 4519-4523, 2011.
- [8] Bairamov, E. and Coskun, C. *Jost solutions and the spectrum of the system of difference equations*, Appl. Math. Lett. (2) **17** (9), 1039-1045, 2004.
- [9] Bairamov, E. Çakar, Ö. and Çelebi, A.O. *Quadratic pencil of Schrödinger operators with spectral singularities: discrete spectrum and principal functions*, J. Math. Anal. Appl. **216** (1), 303-320, 1997.
- [10] Bairamov, E. Çakar, Ö. and Krall, A.M. *Non-selfadjoint difference operators and Jacobi matrices with spectral singularities*, Math. Nachr. **229**, 5-14, 2001.
- [11] Başcanbaz Tunca, G. *Spectral properties of the Klein-Gordon s -wave equation with spectral parameter-dependent boundary condition*, Int. J. Math. Math. Sci. (25-28), 1437-1445, 2004.

- [12] Başcanbaz Tunca, G. and Bairamov, E. *Discrete spectrum and principal functions of non-selfadjoint differential operator*, Czechoslovak Math. J. **49** (4), 689–700, 1999.
- [13] Bohner, M. Guseinov, G. and Peterson, A. *Introduction to the time scales calculus*. In *Advances in dynamic equations on time scales* (Birkhäuser Boston, MA, 2003), 1–15.
- [14] Bohner, M. and Peterson, A. *Dynamic equations on time scales* (Birkhäuser Boston, MA, 2001).
- [15] Clark, S. Gesztesy, F. and Renger, W. *Trace formulas and Borg-type theorems for matrix-valued Jacobi and Dirac finite difference operators*, J. Differential Equations **219** (1), 144–182, 2005.
- [16] Coskun, C. and Olgun, M. *Principal functions of non-selfadjoint matrix Sturm-Liouville equations*, J. Comput. Appl. Math. **235** (16), 4834–4838, 2011.
- [17] Dolzhenko, E.P., *Boundary-value uniqueness theorems for analytic functions*, Mathematical notes of the Academy of Sciences of the USSR **25**, (6), 437–442, 1979.
- [18] Gesztesy, F. Kiselev, A. and Makarov, K.A. *Uniqueness results for matrix-valued Schrödinger, Jacobi, and Dirac-type operators*, Math. Nachr. **239/240**, 103–145, 2002.
- [19] Glazman, I.M. *Direct methods of qualitative spectral analysis of singular differential operators* (Jerusalem, 1966).
- [20] Krall, A.M. Bairamov, E. and Çakar, Ö. *Spectral analysis of non-selfadjoint discrete Schrödinger operators with spectral singularities*, Math. Nachr. **231**, 89–104, 2001.
- [21] L.A. Lusternik and V.J. Sobolev, *Elements of functional analysis* (Hindustan Publishing Corp., Delhi, 1974).
- [22] Lyance, V.E. *A differential operator with spectral singularities I-II*, Amer. Math. Soc. Transl. **60** (2) 185–225, 227–283, 1967.
- [23] Naimark, M.A. *Linear differential operators. Part II: Linear differential operators in Hilbert space* (Frederick Ungar Publishing Co., New York, 1960).
- [24] Olgun, M. and Coskun, C. *Non-selfadjoint matrix Sturm-Liouville operators with spectral singularities*, Appl. Math. Comput. **216**, (8), 2271–2275, 2011.

