# CONTINUOUS DEPENDENCE ON THE PARAMETERS OF PHASE FIELD EQUATIONS 

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#### Abstract

Phase field equations are analyzed. It is shown that the solution of the problem considered depends continuously on changes in the parameters.


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## 1. Introduction

We consider the problem
(2) $\quad u_{t}+\frac{l}{2} \phi_{t}=K \Delta u+h_{2}(x, t),(x, t) \in Q_{T}$
(3) $\left.\quad \phi\right|_{\Gamma}=\phi_{\partial}(x, t),\left.u\right|_{\Gamma}=u_{\partial}(x, t),(x, t) \in \partial \Omega \times(0, T]$
(4) $\quad \phi(x, 0)=\phi_{0}(x), u(x, 0)=u_{0}(x), x \in \Omega$,
where $Q_{T}=\Omega \times(0, T], T>0, \Omega \subset \mathbb{R}^{n},(n \geq 1)$ is a bounded domain with a sufficiently smooth boundary, $\partial \Omega ; \xi, \tau, l$ and $K$ are positive constants characterizing the length scale, the relaxation time, the latent heat and the thermal diffusivity, respectively. Also, $\phi_{0}(x), u_{0}(x), \phi_{\partial}(x, t), u_{\partial}(x, t), h_{1}(x, t), h_{2}(x, t)$ and $f(x, s)$ are given functions.

In [1], G. Caginalp considered the following system of equations as a model describing the phase transitions with a separation surface of finite thickness:

$$
\begin{align*}
& \tau \phi_{t}=\xi^{2} \Delta \phi+\frac{1}{2}\left(\phi-\phi^{3}\right)+2 u, x \in \Omega, t \in \mathbb{R}^{+}  \tag{5}\\
& u_{t}+\frac{l}{2} \phi_{t}=K \Delta u, x \in \Omega, t \in \mathbb{R}^{+} .
\end{align*}
$$

[^0]Under the assumption $\frac{\xi^{2}}{\tau}<K$, a global existence theorem was proved for the classical solution of the initial boundary value problem for the system (5)-(6) with non-homogeneous Dirichlet boundary conditions of the form

$$
\begin{equation*}
\left.\phi(t, x)\right|_{\partial \Omega}=\phi_{\partial}(x),\left.u(t, x)\right|_{\partial \Omega}=u_{\partial}(x), t \in \mathbb{R}^{+} \tag{7}
\end{equation*}
$$

Several scientists have investigated problems based on Caginalp's model, and made a few modifications. In [2], Caginalp and Hastings investigated the existence of stationary solutions of problem (5)-(7) in $\Omega \subset \mathbb{R}^{1}$. In [3], C. M. Eliott and Song Mu Zheng proved the global unique solvability of initial boundary value problems for the system (5)-(6) in the class $H^{2}(\Omega) \times H^{2}(\Omega), \Omega \subset \mathbb{R}^{n}, n \leq 3$, without the assumption $\frac{\xi^{2}}{\tau}<K$, for boundary conditions of the form (7), as well as for conditions of the form

$$
\begin{aligned}
& \left.\frac{\partial \phi}{\partial n}\right|_{\partial \Omega}=0,\left.\quad \frac{\partial u}{\partial n}\right|_{\partial \Omega}=0, t \in \mathbb{R}^{+} \\
& \left.\phi\right|_{\partial \Omega}=\phi_{\partial}(x),\left.\quad \frac{\partial u}{\partial n}\right|_{\partial \Omega}=0, t \in \mathbb{R}^{+}
\end{aligned}
$$

They also studied the behaviour of the solutions of the system (5),(6) when $t \rightarrow \infty$. In [4], Kalantarov proved that the initial boundary value problem for system (5)-(6), under homogeneous boundary conditions of the form (7), is globally uniquely solvable in $C\left(\mathbb{R}^{+}, X\right), X=\stackrel{0}{H^{1}}(\Omega) \times \stackrel{0}{H^{1}}(\Omega)$, and established the existence of a global attractor. In [5], Brochet, Hilhorst and Chen investigated problem (1)-(4), considering $v=u+\frac{l}{2} \phi$, $f(s)=\sum_{j=0}^{2 p-1} b_{j} s^{j}, b_{2 p-1}>0, p \geq 2, h_{i}(x, t)=0,(i=1,2)$ and homogeneous Neumann boundary conditions, proving this problem to be well posed if $\left(\phi_{0}, u_{0}\right) \in\left(L_{2}(\Omega)\right)^{2}$.

## 2. Continuous Dependence of Solutions

We investigate the continuous dependence on the parameters $\xi, \tau, l$ and $K$ of solutions of problem (1)-(4) in the class $V\left(Q_{T}\right) \times V\left(Q_{T}\right)$, where

$$
\begin{equation*}
V\left(Q_{T}\right)=W_{2}^{1}\left(Q_{T}\right) \cap\left\{v(x, t): \Delta v \in L_{2}\left(Q_{T}\right)\right\} \tag{8}
\end{equation*}
$$

The existence of a solution to this problem can be seen from the general results of [7] and [8], but to the best of our knowledge an investigation of continuous dependence does not occur in the literature. Investigations of this type are of interest in physical problems, and can lead to useful applications.

We assume that $f(x, \phi)$ is the Caratheodory function which satisfies the local Lipschitz condition:

$$
\begin{equation*}
\left|f\left(x, s_{1}\right)-f\left(x, s_{2}\right)\right| \leq c\left(1+\left|s_{1}\right|^{p-1}+\left|s_{2}\right|^{p-1}\right)\left|s_{1}-s_{2}\right|, \forall s_{1}, s_{2} \in \mathbb{R}^{1} \tag{9}
\end{equation*}
$$

where $p \in[1, \infty)$ if $n=1,2$, and $p \in\left[1, \frac{n}{n-2}\right]$ if $n \geq 3$. We have used standard techniques for the calculations (cf. [6], which considers this type of question for a different problem). Let $\left\{\phi_{1}, u_{1}\right\}$ and $\left\{\phi_{2}, u_{2}\right\}$ be the solutions from $V\left(Q_{T}\right) \times V\left(Q_{T}\right)$ of the following initial-boundary value problems for different coefficients $\xi_{1}, \tau_{1}, l_{1}, K_{1}$ and $\xi_{2}, \tau_{2}, l_{2}, K_{2}$
respectively.

$$
\begin{aligned}
& \tau_{1}\left(\phi_{1}\right)_{t}-\xi_{1}^{2} \Delta \phi_{1}+f\left(x, \phi_{1}\right)=2 u_{1}+h_{1}(x, t),(x, t) \in Q_{T} \\
& \left(u_{1}\right)_{t}+\frac{l_{1}}{2}\left(\phi_{1}\right)_{t}=K_{1} \Delta u_{1}+h_{2}(x, t),(x, t) \in Q_{T} \\
& \left.\phi_{1}\right|_{\Gamma}=\phi_{\partial}(x, t),\left.u_{1}\right|_{\Gamma}=u_{\partial}(x, t),(x, t) \in \partial \Omega \times(0, T] \\
& \phi_{1}(x, 0)=\phi_{0}(x), u_{1}(x, 0)=u_{0}(x), x \in \Omega \\
& \tau_{2}\left(\phi_{2}\right)_{t}-\xi_{2}^{2} \Delta \phi_{2}+f\left(x, \phi_{2}\right)=2 u_{2}+h_{1}(x, t),(x, t) \in Q_{T} \\
& \left(u_{2}\right)_{t}+\frac{l_{2}}{2}\left(\phi_{2}\right)_{t}=K_{2} \Delta u_{2}+h_{2}(x, t),(x, t) \in Q_{T} \\
& \left.\phi_{2}\right|_{\Gamma}=\phi_{\partial}(x, t),\left.u_{2}\right|_{\Gamma}=u_{\partial}(x, t),(x, t) \in \partial \Omega \times(0, T] \\
& \phi_{2}(x, 0)=\phi_{0}(x), u_{2}(x, 0)=u_{0}(x), x \in \Omega .
\end{aligned}
$$

We define the difference variables $\phi, u, \xi, \tau, l$ and $K$ by

$$
\begin{aligned}
\phi & =\phi_{1}-\phi_{2}, & \tau & =\tau_{1}-\tau_{2}, \\
u & =u_{1}-u_{2}, & l & =l_{1}-l_{2}, \\
\xi^{2} & =\xi_{1}^{2}-\xi_{2}^{2}, & K & =K_{1}-K_{2}
\end{aligned}
$$

Then $\{\phi, u\}$ satisfies the initial-boundary value problem:

$$
\begin{align*}
& \tau_{1} \phi_{t}+\tau\left(\phi_{2}\right)_{t}-\xi_{1}^{2} \Delta \phi-\xi^{2} \Delta \phi_{2}+f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)=2 u,  \tag{10}\\
& u_{t}+\frac{l_{1}}{2} \phi_{t}+\frac{l}{2}\left(\phi_{2}\right)_{t}=K_{1} \Delta u+K \Delta u_{2} \\
& \left.\phi\right|_{\Gamma}=\left.u\right|_{\Gamma}=0 \\
& \phi(x, 0)=u(x, 0)=0 \tag{13}
\end{align*}
$$

If we take the inner product in $L_{2}(\Omega)$ of (10) by $\phi_{t}+\phi$ and of (11) by $\frac{2 \tau_{1}}{l_{1}^{2}} u_{t}+\frac{4}{l_{1}} u$ and then add the equations obtained, we obtain

$$
\begin{align*}
& \tau_{1}\left\|\phi_{t}\right\|^{2}+\xi_{1}^{2}\|\nabla \phi\|^{2}+\frac{4 K_{1}}{l_{1}}\|\nabla u\|^{2}+\frac{2 \tau_{1}}{l_{1}^{2}}\left\|u_{t}\right\|^{2}+ \\
& +\frac{d}{d t}\left[\frac{\xi_{1}^{2}}{2}\|\nabla \phi\|^{2}+\frac{2}{l_{1}}\|u\|^{2}+\frac{\tau_{1}}{2}\|\phi\|^{2}+\frac{\tau_{1} K_{1}}{l_{1}^{2}}\|\nabla u\|^{2}\right] \leq \\
\leq & \left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi_{t} d x\right|+\left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi d x\right|+  \tag{14}\\
\quad & +2|(u, \phi)|+\frac{4|K|}{l_{1}}\left|\left(\Delta u_{2}, u\right)\right|+\frac{2 \tau_{1}|K|}{l_{1}^{2}}\left|\left(\Delta u_{2}, u_{t}\right)\right|+\frac{\tau_{1}}{l_{1}}\left|\left(\phi_{t}, u_{t}\right)\right|+ \\
& +|\tau|\left|\left(\left(\phi_{2}\right)_{t}, \phi_{t}\right)\right|+|\tau|\left|\left(\left(\phi_{2}\right)_{t}, \phi\right)\right|+\left|\xi^{2}\right|\left|\left(\Delta \phi_{2}, \phi_{t}\right)\right|+\left|\xi^{2}\right|\left|\left(\Delta \phi_{2}, \phi\right)\right|+ \\
& \frac{2|l|}{l_{1}}\left|\left(\left(\phi_{2}\right)_{t}, u\right)\right|+\frac{\tau_{1}|l|}{l_{1}^{2}}\left|\left(\left(\phi_{2}\right)_{t}, u_{t}\right)\right|,
\end{align*}
$$

where $\|$.$\| denotes the norm on L_{2}(\Omega)$. Using (9) and Hölder's inequality, the first term on the right hand side of (14) can be estimated in the following manner:

$$
\begin{aligned}
& \left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi_{t} d x\right| \leq \int_{\Omega}\left|\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right)\right|\left|\phi_{t}\right| d x \leq \\
& \leq \int_{\Omega} c\left(1+\left|\phi_{1}\right|^{p-1}+\left|\phi_{2}\right|^{p-1}\right)|\phi|\left|\phi_{t}\right| d x \leq \\
& \leq c\|\phi\|\left\|\phi_{t}\right\|+c \int_{\Omega}\left|\phi_{1}\right|^{p-1}|\phi|\left|\phi_{t}\right| d x+c \int_{\Omega}\left|\phi_{2}\right|^{p-1}|\phi|\left|\phi_{t}\right| d x \leq \\
& \leq c\|\phi\|\left\|\phi_{t}\right\|+c\|\phi\|_{\frac{2 n}{n-2}}\left\|\phi_{t}\right\|\left(\left\|\phi_{1}\right\|_{L_{(p-1) n}}^{p-1}+\left\|\phi_{2}\right\|_{L_{(p-1) n}}^{p-1}\right) .
\end{aligned}
$$

By the Sobolev imbedding theorem the following inequality holds:

$$
\|\phi\|_{\frac{2 n}{n-2}(\Omega)} \leq c_{2}\|\nabla \phi\| .
$$

Therefore,

$$
\begin{aligned}
& \left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi_{t} d x\right| \leq c\|\phi\|\left\|\phi_{t}\right\|+ \\
& +c c_{2}\left\|\phi_{t}\right\|\|\nabla \phi\|\left(\left\|\phi_{1}\right\|_{L_{(p-1) n}}^{p-1}+\left\|\phi_{2}\right\|_{L_{(p-1) n}}^{p-1}\right) .
\end{aligned}
$$

Since $\left\{\phi_{i}, u_{i}\right\} \in V\left(Q_{T}\right) \times V\left(Q_{T}\right)$ are fixed,

$$
\left\|\phi_{1}\right\|_{L_{(p-1) n}}^{p-1}+\left\|\phi_{2}\right\|_{L_{(p-1) n}}^{p-1} \leq c_{1}(t) .
$$

Hence

$$
\begin{equation*}
\left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi_{t} d x\right| \leq c\|\phi\|\left\|\phi_{t}\right\|+c c_{1}(t) c_{2}\left\|\phi_{t}\right\|\|\nabla \phi\|, \tag{15}
\end{equation*}
$$

and similarly,

$$
\begin{equation*}
\left|\int_{\Omega}\left(f\left(x, \phi_{1}\right)-f\left(x, \phi_{2}\right)\right) \phi d x\right| \leq c\|\phi\|^{2}+c c_{1}(t) c_{2}\|\phi\|\|\nabla \phi\| . \tag{16}
\end{equation*}
$$

Taking into account (15) and (16), we obtain from (14)

$$
\begin{align*}
& \quad \tau_{1}\left\|\phi_{t}\right\|^{2}+\xi_{1}^{2}\|\nabla \phi\|^{2}+\frac{4 K_{1}}{l_{1}}\|\nabla u\|^{2}+\frac{2 \tau_{1}}{l_{1}^{2}}\left\|u_{t}\right\|^{2}+ \\
& \quad+\frac{d}{d t}\left[\frac{\xi_{1}^{2}}{2}\|\nabla \phi\|^{2}+\frac{2}{l_{1}}\|u\|^{2}+\frac{\tau_{1}}{2}\|\phi\|^{2}+\frac{\tau_{1} K_{1}}{l_{1}^{2}}\|\nabla u\|^{2}\right] \leq \\
& \leq \quad 2|(u, \phi)|+\frac{4|K|}{l_{1}}\left|\left(\Delta u_{2}, u\right)\right|+\frac{2 \tau_{1}|K|}{l_{1}^{2}}\left|\left(\Delta u_{2}, u_{t}\right)\right|+\frac{\tau_{1}}{l_{1}}\left|\left(\phi_{t}, u_{t}\right)\right|+  \tag{17}\\
& \quad+|\tau|\left|\left(\left(\phi_{2}\right)_{t}, \phi_{t}\right)\right|+|\tau|\left|\left(\left(\phi_{2}\right)_{t}, \phi\right)\right|+\left|\xi^{2}\right|\left|\left(\Delta \phi_{2}, \phi_{t}\right)\right|+\left|\xi^{2}\right|\left|\left(\Delta \phi_{2}, \phi\right)\right|+ \\
& \quad+\frac{2|l|}{l_{1}}\left|\left(\left(\phi_{2}\right)_{t}, u\right)\right|+\frac{\tau_{1}|l|}{l_{1}^{2}}\left|\left(\left(\phi_{2}\right)_{t}, u_{t}\right)\right|+c\|\phi\|\left\|\phi_{t}\right\|+ \\
& \quad+c c_{1}(t) c_{2}\left\|\phi_{t}\right\|\|\nabla \phi\|+c\|\phi\|^{2}+c c_{1}(t) c_{2}\|\phi\|\|\nabla \phi\| .
\end{align*}
$$

Making use of Cauchy's inequality with $\varepsilon$, the right hand side of (17) can be estimated. If we select the number $\varepsilon>0$ sufficiently small then we obtain

$$
\begin{array}{ll} 
& \frac{\tau_{1}}{4}\left\|\phi_{t}\right\|^{2}+\frac{4 K_{1}}{l_{1}}\|\nabla u\|^{2}+\frac{\tau_{1}}{2 l_{1}^{2}}\left\|u_{t}\right\|^{2}+ \\
& \frac{d}{d t}\left[\frac{\xi_{1}^{2}}{2}\|\nabla \phi\|^{2}+\frac{2}{l_{1}}\|u\|^{2}+\frac{\tau_{1}}{2}\|\phi\|^{2}+\frac{\tau_{1} K_{1}}{l_{1}^{2}}\|\nabla u\|^{2}\right] \leq \\
\leq & a_{1}(t)\|\nabla \phi\|^{2}+\left(\frac{3+l_{1}}{l_{1}}\right)\|u\|^{2}+a_{2}(t)\|\phi\|^{2}+  \tag{18}\\
& +\left[\left(\frac{\tau_{1}+2 l_{1}}{2 l_{1}^{2}}\right) l^{2}+\left(\frac{8+\tau_{1}}{2 \tau_{1}}\right) \tau^{2}\right]\left\|\left(\phi_{2}\right)_{t}\right\|^{2}+ \\
& +\left(\frac{2\left(\tau_{1}+l_{1}\right)}{l_{1}^{2}}\right) K^{2}\left\|\Delta u_{2}\right\|^{2}+\left(\frac{8+\tau_{1}}{2 \tau_{1}}\right) \xi^{4}\left\|\Delta \phi_{2}\right\|^{2},
\end{array}
$$

where,

$$
a_{1}(t)=\frac{4 c^{2} c_{1}^{2}(t) c_{2}^{2}}{\tau_{1}} \text { and } a_{2}(t)=\frac{4 c^{2}}{\tau_{1}}+\frac{c^{2} c_{1}^{2}(t) c_{2}^{2}}{4 \xi_{1}^{2}}+c+2 .
$$

If we set

$$
c_{2}(t)=\max \left\{\frac{2 a_{1}(t)}{\xi_{1}^{2}}, \frac{3+l_{1}}{2}, \frac{2 a_{2}(t)}{\tau_{1}}, 1\right\},
$$

and

$$
Y(t)=\frac{\xi_{1}^{2}}{2}\|\nabla \phi\|^{2}+\frac{2}{l_{1}}\|u\|^{2}+\frac{\tau_{1}}{2}\|\phi\|^{2}+\frac{\tau_{1} K_{1}}{l_{1}^{2}}\|\nabla u\|^{2}
$$

then from (18) we obtain

$$
\left\{\begin{aligned}
\frac{d Y(t)}{d t} & \leq c_{2}(t) Y(t)+\left(c_{3} l^{2}+c_{4} \tau^{2}\right)\left\|\left(\phi_{2}\right)_{t}\right\|^{2}+c_{4} \xi^{4}\left\|\Delta \phi_{2}\right\|^{2}+c_{5} K^{2}\left\|\Delta u_{2}\right\|^{2} \\
Y(0) & =0
\end{aligned}\right.
$$

where, $c_{3}=\frac{\tau_{1}+2 l_{1}}{2 l_{1}^{2}}, c_{4}=\frac{8+\tau_{1}}{2 \tau_{1}}$ and $c_{5}=\frac{2\left(\tau_{1}+l_{1}\right)}{l_{1}^{2}}$. According to Gronwall's lemma, we have

$$
\begin{align*}
& Y(t) \leq \exp \left\{\int_{0}^{T} c_{2}(s) d s\right\}\left\{\left(c_{3} l^{2}+c_{4} \tau^{2}\right)\left\|\left(\phi_{2}\right)_{t}\right\|_{L_{2}\left(Q_{T}\right)}^{2}+\right.  \tag{19}\\
&\left.+c_{4} \xi^{4}\left\|\Delta \phi_{2}\right\|_{L_{2}\left(Q_{T}\right)}^{2}+c_{5} K^{2}\left\|\Delta u_{2}\right\|_{L_{2}\left(Q_{T}\right)}^{2}\right\}
\end{align*}
$$

Since $\left\{\phi_{i}, u_{i}\right\} \in V\left(Q_{T}\right) \times V\left(Q_{T}\right)$, we have

$$
\begin{aligned}
\left\|\left(\phi_{2}\right)_{t}\right\|_{L_{2}\left(Q_{T}\right)}^{2} & \leq C \\
\left\|\Delta \phi_{2}\right\|_{L_{2}\left(Q_{T}\right)}^{2} & \leq C, \text { and } \\
\left\|\Delta u_{2}\right\|_{L_{2}\left(Q_{T}\right)}^{2} & \leq C .
\end{aligned}
$$

If we set

$$
\max \left\{c_{3} C, c_{4} C, c_{4} C, c_{5} C\right\}=C_{6}
$$

and

$$
C_{6} \exp \left\{\int_{0}^{T} c_{2}(s) d s\right\}=C_{7}
$$

then from (19) we have

$$
Y(t) \leq C_{7}\left[K^{2}+\xi^{4}+\tau^{2}+l^{2}\right]
$$

Hence we have proved the following theorem.
2.1. Theorem. Assume that (9) is satisfied. Then the solution of problem (1)-(4) from $V\left(Q_{T}\right) \times V\left(Q_{T}\right)$ depends continuously on the parameters $\xi, \tau, l$ and $K$. Moreover,

$$
\left\|\phi_{1}-\phi_{2}\right\|_{C\left(0, T ; W_{2}^{1}(\Omega)\right)}^{2} \leq C_{7}\left[\left(K_{1}-K_{2}\right)^{2}+\left(\xi_{1}-\xi_{2}\right)^{4}+\left(\tau_{1}-\tau_{2}\right)^{2}+\left(l_{1}-l_{2}\right)^{2}\right]
$$

and

$$
\left\|u_{1}-u_{2}\right\|_{C\left(0, T ; W_{2}^{1}(\Omega)\right)}^{2} \leq C_{7}\left[\left(K_{1}-K_{2}\right)^{2}+\left(\xi_{1}-\xi_{2}\right)^{4}+\left(\tau_{1}-\tau_{2}\right)^{2}+\left(l_{1}-l_{2}\right)^{2}\right]
$$

## References

[1] Caginalp, G. An analysis of a phase field model of a free boundary, Arch. Rat. Mech. Anal. 92, 205-245, 1986.
[2] Caginalp, G. and Hastings, S. Properties of some ordinary differential equations related to free boundary problems, Proc. Roy. Soc. Edinburgh 104 A, 217-234, 1986.
[3] Elliot, C. M. and Zheng, S. M. Global existence and stability of solutions to the phase field equations, in: Free boundary value problems, Internat. Ser. Number. Math. 95, Birkhause, Basel, 46-58, 1990.
[4] Kalantarov, V. K. On the minimal global attractor of a system of phase field equations, Zap. Nauchn. Semin. LOMI 188, 70-86, 1991.
[5] Brochet, D., Hilhorst, D. and Chen, X. Finite dimensional exponential attractor for the phase field model, Appl. Analysis 49, 197-212, 1993.
[6] Payne, L. E. and Straughan, B. Convergence and continuous dependence for the BirkmanForchmeier equations, Studies in applied mathematics 102, 419-439, 1999.
[7] Soltanov, K. N. On nonlinear equations of the form $F(x, u, D u, \Delta u)=0$, Russian Acad. Sic. Sb. Math. 80 no:2, 367-3923, 1995.
[8] Soltanov, K. N. Some imbedding theorems and nonlinear differential equations, Trans. Acad. Sci. Azerb. Ser. Phys.-Tech. Math. Sci. 19, 125-146, 1999.


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