

PRINCIPAL FUNCTIONS OF NON-SELFADJOINT MATRIX STURM – LIOUVILLE OPERATORS WITH BOUNDARY CONDITIONS DEPENDENT ON THE SPECTRAL PARAMETER

CAFER COSKUN, DENIZ KATAR AND MURAT OLGUN

ABSTRACT. Let L denote operator generated in $L_2(\mathbb{R}_+, E)$ by the differential expression

$$l(y) = -y'' + Q(x)y, \quad x \in \mathbb{R}_+ := [0, \infty),$$

and the boundary condition $Y'(0, \lambda) - (\beta_0 + \beta_1\lambda + \beta_2\lambda^2)Y(0, \lambda) = 0$, where Q is a non-selfadjoint matrix-valued function and $\beta_0, \beta_1, \beta_2$ are non-selfadjoint matrices, also β_2 is invertible. In this paper, we investigate the principal functions corresponding to the eigenvalues and the spectral singularities of L .

1. INTRODUCTION

Let us consider the boundary value problem (BVP)

$$-y'' + q(x)y = \lambda^2 y, \quad x \in \mathbb{R}_+, \quad (1.1)$$

$$y(0) = 0, \quad (1.2)$$

in $L^2(\mathbb{R}_+)$, where q is a complex-valued function. The spectral theory of the BVP (1.1)–(1.2) with continuous and point spectrum was investigated by Naimark [1]. He showed the existence of the spectral singularities in the continuous spectrum of the BVP (1.1)–(1.2). Note that the eigenfunctions and the associated functions (principal functions) corresponding to the spectral singularities are not the elements of $L^2(\mathbb{R}_+)$. Also, the spectral singularities belong to the continuous spectrum and are the poles of the resolvent's kernel, but are not the eigenvalues of the BVP (1.1)–(1.2). The spectral singularities in the spectral expansion of the BVP (1.1)–(1.2) in terms of the principal functions have been investigated in [2]. The spectral analysis of the quadratic pencil of Schrödinger, Dirac and Klein-Gordon operators with spectral singularities were studied in [3, 4, 5, 6, 7, 8, 9]. The spectral analysis

Received by the editors Feb. 12, 2014; Accepted: May 12, 2014.

2000 *Mathematics Subject Classification.* 53C15, 53B05(53B20), 53C40.

Key words and phrases. Eigenvalues, spectral singularities, spectral analysis, Sturm - Liouville Operator, non-selfadjoint matrix operator.

of the non-selfadjoint operator, generated in $L^2(\mathbb{R}_+)$ by (1.1) and the boundary condition

$$\frac{y'(0)}{y(0)} = \alpha_0 + \alpha_1\lambda + \alpha_2\lambda^2,$$

where $\alpha_i \in \mathbb{C}$, $i = 0, 1, 2$ with $\alpha_2 \neq 0$ was investigated by Bairamov et al. [20]. The properties of the principal functions corresponding to the eigenvalues and the spectral singularities were studied in [14, 20, 21, 22]. Spectral analysis of the selfadjoint differential and difference equations with matrix coefficients are studied in [10, 11, 12, 13].

Let E be an n -dimensional ($n < \infty$) Euclidian space with the norm $\|\cdot\|$ and let us introduce the Hilbert space $L^2(\mathbb{R}_+, E)$ consisting of vector-valued functions with the values in E . We will consider the BVP

$$-y'' + Q(x)y = \lambda^2 y, \quad x \in \mathbb{R}_+, \quad (1.3)$$

$$y(0) = 0, \quad (1.4)$$

in $L^2(\mathbb{R}_+, E)$, where Q is a non-selfadjoint matrix-valued function (i. e., $Q \neq Q^*$). It is clear that, the BVP (1.3)–(1.4) is non-selfadjoint. In [15, 16] discrete spectrum of the non-selfadjoint matrix Sturm–Liouville operator and properties of the principal functions corresponding to the eigenvalues and the spectral singularities was investigated.

Let us consider the BVP in $L_2(\mathbb{R}_+, E)$

$$-y'' + Q(x)y = \lambda^2 y, \quad x \in \mathbb{R}_+, \quad (1.5)$$

$$y'(0, \lambda) - (\beta_0 + \beta_1\lambda + \beta_2\lambda^2)y(0, \lambda) = 0, \quad (1.6)$$

where Q is a non-selfadjoint matrix-valued function and $\beta_0, \beta_1, \beta_2$ are non-selfadjoint matrices also β_2 is invertible. In this paper, which is an extension of [23], we aim to investigate the properties of the principal functions corresponding to the eigenvalues and the spectral singularities of the BVP (1.5)–(1.6).

2. Jost Solution of (1.5)

We will denote the solution of (1.5) satisfying the condition

$$\lim_{x \rightarrow \infty} y(x, \lambda)e^{-i\lambda x} = I, \quad \lambda \in \overline{\mathbb{C}}_+ := \{\lambda : \lambda \in \mathbb{C}, \text{Im } \lambda \geq 0\}, \quad (2.1)$$

by $E(x, \lambda)$. The solution $E(x, \lambda)$ is called the Jost solution of (1.5).

Under the condition

$$\int_0^\infty x \|Q(x)\| dx < \infty, \quad (2.2)$$

the Jost solution has a representation

$$E(x, \lambda) = e^{i\lambda x} I + \int_x^\infty K(x, t)e^{i\lambda t} dt \quad (2.3)$$

for $\lambda \in \bar{\mathbb{C}}_+$, where the kernel matrix function $K(x, t)$ satisfies

$$\begin{aligned} K(x, t) &= \frac{1}{2} \int_{\frac{x+t}{2}}^{\infty} Q(s) ds + \frac{1}{2} \int_x^{\frac{x+t}{2}} \int_{t+x-s}^{t+s-x} Q(s) K(s, v) dv ds \\ &\quad + \frac{1}{2} \int_{\frac{x+t}{2}}^{\infty} \int_s^{t+s-x} Q(s) K(s, v) dv ds \end{aligned} \quad (2.4)$$

Moreover, $K(x, t)$ is continuously differentiable with respect to its arguments and

$$\|K(x, t)\| \leq c\sigma\left(\frac{x+t}{2}\right), \quad (2.5)$$

$$\|K_x(x, t)\| \leq \frac{1}{4} \left\| Q\left(\frac{x+t}{2}\right) \right\| + c\sigma\left(\frac{x+t}{2}\right), \quad (2.6)$$

$$\|K_t(x, t)\| \leq \frac{1}{4} \left\| Q\left(\frac{x+t}{2}\right) \right\| + c\sigma\left(\frac{x+t}{2}\right), \quad (2.7)$$

where $\sigma(x) = \int_x^{\infty} \|Q(s)\| ds$ and $c > 0$ is a constant. Therefore, $E(x, \lambda)$ is analytic with respect to λ in $\mathbb{C}_+ := \{\lambda : \lambda \in \mathbb{C}, \operatorname{Im} \lambda > 0\}$ and continuous on the real axis ([17; chp.1]).

Let $\hat{E}^{\pm}(x, \lambda)$ denote the solutions of (1.5) subject to the conditions

$$\lim_{x \rightarrow \infty} \hat{E}^{\pm}(x, \lambda) e^{\pm i\lambda x} = I, \quad \lim_{x \rightarrow \infty} \hat{E}_x^{\pm}(x, \lambda) e^{\pm i\lambda x} = \pm i\lambda I, \quad \lambda \in \bar{\mathbb{C}}_{\pm}. \quad (2.8)$$

Then

$$W \left[E(x, \lambda), \hat{E}^{\pm}(x, \lambda) \right] = \mp 2i\lambda I, \quad \lambda \in \mathbb{C}_{\pm}, \quad (2.9)$$

$$W \left[E(x, \lambda), E(x, -\lambda) \right] = -2i\lambda I, \quad \lambda \in \mathbb{R}, \quad (2.10)$$

where $W[f_1, f_2]$ is the Wronskian of f_1 and f_2 .

Let $\varphi(x, \lambda)$ denote the solution of (1.5) subject to the initial conditions $\varphi(0, \lambda) = I$, $\varphi'(0, \lambda) = \beta_0 + \beta_1 \lambda + \beta_2 \lambda^2$. Therefore $\varphi(x, \lambda)$ is an entire function of λ .

Let us define the following functions:

$$A_{\pm}(\lambda) = \varphi(0, \lambda) E_x(0, \pm\lambda) - \varphi'(0, \lambda) E(0, \pm\lambda) \quad \lambda \in \bar{\mathbb{C}}_{\pm}, \quad (2.11)$$

where $\bar{\mathbb{C}}_{\pm} = \{\lambda : \lambda \in \mathbb{C}, \pm \operatorname{Im} \lambda \geq 0\}$. It is obvious that the functions $A_+(\lambda)$ and $A_-(\lambda)$ are analytic in \mathbb{C}_+ and \mathbb{C}_- respectively and continuous on the real axis. The functions A_+ and A_- are called Jost functions of L .

3. Eigenvalues and Spectral Singularities of L

The resolvent of L defined by

$$R_\lambda(L)f = \int_0^\infty G(x, t; \lambda)g(t)dt, \quad g \in L_2(\mathbb{R}_+, E), \quad (3.1)$$

where

$$G(x, t; \lambda) = \begin{cases} G_+(x, t; \lambda), & \lambda \in \mathbb{C}_+ \\ G_-(x, t; \lambda), & \lambda \in \mathbb{C}_- \end{cases}. \quad (3.2)$$

and

$$G_\pm(x, t; \lambda) = \begin{cases} -E(x, \pm\lambda)A_\pm^{-1}(\lambda)\varphi^T(t, \lambda), & 0 \leq t \leq x \\ -\varphi(x, \lambda)[A_\pm^T(\lambda)]^{-1}E^T(t, \pm\lambda), & x \leq t < \infty \end{cases} \quad (3.3)$$

We will show the set of eigenvalues and the set of spectral singularities of the operator L by σ_d and σ_{ss} respectively.

Let us suppose that

$$H_\pm(\lambda) = \det A_\pm(\lambda). \quad (3.4)$$

From (2.3) and (3.1)–(3.4)

$$\begin{aligned} \sigma_d &= \{\lambda : \lambda \in \mathbb{C}_+, H_+(\lambda) = 0\} \cup \{\lambda : \lambda \in \mathbb{C}_-, H_-(\lambda) = 0\} \\ \sigma_{ss} &= \{\lambda : \lambda \in \mathbb{R}^*, H_+(\lambda) = 0\} \cup \{\lambda : \lambda \in \mathbb{R}^*, H_-(\lambda) = 0\}, \end{aligned} \quad (3.5)$$

where $\mathbb{R}^* = \mathbb{R} \setminus \{0\}$.

We see from that, the functions

$$K^+(\lambda) = \frac{\hat{A}_+(\lambda)}{2i\lambda}E(x, \lambda) - \frac{A_+(\lambda)}{2i\lambda}\hat{E}^+(x, \lambda), \quad \lambda \in \mathbb{C}_+, \quad (3.6)$$

$$K^-(\lambda) = \frac{\hat{A}_-(\lambda)}{2i\lambda}E(x, -\lambda) - \frac{A_-(\lambda)}{2i\lambda}\hat{E}^-(x, \lambda), \quad \lambda \in \mathbb{C}_-, \quad (3.7)$$

$$K(\lambda) = \frac{A_+(\lambda)}{2i\lambda}E(x, -\lambda) - \frac{A_-(\lambda)}{2i\lambda}E(x, \lambda), \quad \lambda \in \mathbb{R}^*, \quad (3.8)$$

are the solutions of the boundary problem (1.5)–(1.6) where

$$\hat{A}_\pm(\lambda) = \hat{E}_x^\pm(0, \lambda) - (\beta_0 + \beta_1\lambda + \beta_2\lambda^2)\hat{E}^\pm(0, \lambda). \quad (3.9)$$

Now let us assume that

$$Q \in AC(\mathbb{R}_+), \quad \lim_{x \rightarrow \infty} Q(x) = 0, \quad \sup_{x \in \mathbb{R}_+} \left[e^{\varepsilon\sqrt{x}} \|Q'(x)\| \right] < \infty, \quad \varepsilon > 0. \quad (3.10)$$

Theorem 3.1. *Under the condition (3.10), the operator L has a finite number of eigenvalues and spectral singularities, and each of them is of finite multiplicity.*

4. Principal Functions of L

Under the condition (3.10), let $\lambda_1, \dots, \lambda_j$ and $\lambda_{j+1}, \dots, \lambda_k$ denote the zeros H^+ in \mathbb{C}_+ and H^- in \mathbb{C}_- (which are the eigenvalues of L) with multiplicities m_1, \dots, m_j and m_{j+1}, \dots, m_k , respectively. It is obvious that from the definition of the Wronskian

$$\left\{ \frac{d^n}{d\lambda^n} W [K^+(x, \lambda), E(x, \lambda)] \right\}_{\lambda=\lambda_p} = \left\{ \frac{d^n}{d\lambda^n} A_+(\lambda) \right\}_{\lambda=\lambda_p} = 0 \quad (4.1)$$

for $n = 0, 1, \dots, m_p - 1$, $p = 1, 2, \dots, j$, and

$$\left\{ \frac{d^n}{d\lambda^n} W [K^-(x, \lambda), E(x, -\lambda)] \right\}_{\lambda=\lambda_p} = \left\{ \frac{d^n}{d\lambda^n} A_-(\lambda) \right\}_{\lambda=\lambda_p} = 0 \quad (4.2)$$

for $n = 0, 1, \dots, m_p - 1$, $p = j + 1, \dots, k$.

Theorem 4.1. *The following formulae:*

$$\left\{ \frac{\partial^n}{\partial \lambda^n} K^+(x, \lambda) \right\}_{\lambda=\lambda_p} = \sum_{m=0}^n F_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\lambda_p}, \quad (4.3)$$

$n = 0, 1, \dots, m_p - 1$, $p = 1, 2, \dots, j$, where

$$F_m(\lambda_p) = \binom{n}{m} \left\{ \frac{\partial^{n-m}}{\partial \lambda^{n-m}} \hat{A}_+(\lambda) \right\}_{\lambda=\lambda_p}, \quad (4.4)$$

$$\left\{ \frac{\partial^n}{\partial \lambda^n} K^-(x, \lambda) \right\}_{\lambda=\lambda_p} = \sum_{m=0}^n N_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, -\lambda) \right\}_{\lambda=\lambda_p}, \quad (4.5)$$

$n = 0, 1, \dots, m_p - 1$, $p = j + 1, \dots, k$, where

$$N_m(\lambda_p) = \binom{n}{m} \left\{ \frac{\partial^{n-m}}{\partial \lambda^{n-m}} \hat{A}_-(\lambda) \right\}_{\lambda=\lambda_p} \quad (4.6)$$

hold.

Proof. We will prove only (4.3) using the method induction, because the case of (4.5) is similar. Let be $n = 0$. Since $K^+(x, \lambda)$ and $E(x, \lambda)$ are linearly dependent from (4.1), we get

$$K^+(x, \lambda_p) = f_0(\lambda_p) E(x, \lambda_p) \quad (4.7)$$

where $f_0(\lambda_p) \neq 0$. Let us assume that $1 \leq n_0 \leq m_p - 2$, (4.3) holds; that is,

$$\left\{ \frac{\partial^{n_0}}{\partial \lambda^{n_0}} K^+(x, \lambda) \right\}_{\lambda=\lambda_p} = \sum_{m=0}^{n_0} F_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\lambda_p}. \quad (4.8)$$

We will prove that (4.3) holds for $n_0 + 1$. If $Y(x, \lambda)$ is a solution of (1.5), then $\frac{\partial^{n_0}}{\partial \lambda^{n_0}} Y(x, \lambda)$ satisfies

$$\left[-\frac{d^2}{dx^2} + Q(x) - \lambda^2 \right] \frac{\partial^{n_0}}{\partial \lambda^{n_0}} Y(x, \lambda) = 2\lambda n_0 \frac{\partial^{n_0-1}}{\partial \lambda^{n_0-1}} Y(x, \lambda) + n_0(n_0-1) \frac{\partial^{n_0-2}}{\partial \lambda^{n_0-2}} Y(x, \lambda). \quad (4.9)$$

Writing for (4.9) $K^+(x, \lambda)$ and $E(x, \lambda)$, and using (4.8), we find

$$\left[-\frac{d^2}{dx^2} + Q(x) - \lambda^2 \right] g_{n_0+1}(x, \lambda_p) = 0, \quad (4.10)$$

where

$$g_{n_0+1}(x, \lambda_p) = \left\{ \frac{\partial^{n_0+1}}{\partial \lambda^{n_0+1}} K^+(x, \lambda) \right\}_{\lambda=\lambda_p} - \sum_{m=0}^{n_0+1} F_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\lambda_p}. \quad (4.11)$$

From (4.1), we have

$$W [g_{n_0+1}(x, \lambda_p), E(x, \lambda_p)] = \left\{ \frac{d^{n_0+1}}{d\lambda^{n_0+1}} W [K^+(x, \lambda), E(x, \lambda)] \right\}_{\lambda=\lambda_p} = 0. \quad (4.12)$$

Hence there exists a constant $f_{n_0+1}(\lambda_p)$ such that

$$g_{n_0+1}(x, \lambda_p) = f_{n_0+1}(\lambda_p) E(x, \lambda_p). \quad (4.13)$$

This shows that (4.3) holds for $n = n_0 + 1$. \square

Using (4.3) and (4.5), define the functions

$$U_{n,p}(x) = \left\{ \frac{\partial^n}{\partial \lambda^n} K^+(x, \lambda) \right\}_{\lambda=\lambda_p} = \sum_{m=0}^n F_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\lambda_p}, \quad (4.14)$$

$n = 0, 1, \dots, m_p - 1$, $p = 1, 2, \dots, j$ and

$$U_{n,p}(x) = \left\{ \frac{\partial^n}{\partial \lambda^n} K^-(x, \lambda) \right\}_{\lambda=\lambda_p} = \sum_{m=0}^n N_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, -\lambda) \right\}_{\lambda=\lambda_p}, \quad (4.15)$$

$n = 0, 1, \dots, m_p - 1$, $p = j + 1, \dots, k$.

Then for $\lambda = \lambda_p$, $p = 1, 2, \dots, j, j + 1, \dots, k$,

$$\begin{aligned} l(U_{0,p}) &= 0, \\ l(U_{1,p}) + \frac{1}{1!} \frac{\partial}{\partial \lambda} l(U_{0,p}) &= 0, \\ l(U_{n,p}) + \frac{1}{1!} \frac{\partial}{\partial \lambda} l(U_{n-1,p}) + \frac{1}{2!} \frac{\partial^2}{\partial \lambda^2} l(U_{n-2,p}) &= 0, \end{aligned} \quad (4.16)$$

$n = 2, 3, \dots, m_p - 1$,

hold, where $l(u) = -u'' + Q(x)u - \lambda^2 u$ and $\frac{\partial^m}{\partial \lambda^m} l(u)$ denote the differential expressions whose coefficients are the m -th derivatives with respect to λ of the corresponding coefficients of the differential expression $l(u)$. (4.16) shows that $U_{0,p}$ is the eigenfunction corresponding to the eigenvalue $\lambda = \lambda_p$; $U_{1,p}, U_{2,p}, \dots, U_{m_p-1,p}$ are the associated functions of $U_{0,p}$ [18, 19].

$U_{0,p}, U_{1,p}, \dots, U_{m_p-1,p}$, $p = 1, 2, \dots, j, j + 1, \dots, k$ are called the principal functions corresponding to the eigenvalue $\lambda = \lambda_p$, $p = 1, 2, \dots, j, j + 1, \dots, k$ of L .

Theorem 4.2.

$$U_{n,p} \in L_2(\mathbb{R}_+, E), \quad n = 0, 1, \dots, m_p - 1, \quad p = 1, 2, \dots, j, j+1, \dots, k. \quad (4.17)$$

Proof. Let be $0 \leq n \leq m_p - 1$ and $1 \leq p \leq j$. Using (2.2), (3.10) and (4.14) we obtain that

$$\|K(x, t)\| \leq ce^{-\epsilon\sqrt{\frac{x+t}{2}}}. \quad (4.18)$$

From (2.3) we get

$$\left\| \left\{ \frac{\partial^n}{\partial \lambda^n} E(x, \lambda) \right\}_{\lambda=\lambda_p} \right\| \leq x^n e^{-x \operatorname{Im} \lambda_p} + c \int_x^\infty t^n e^{-\epsilon\sqrt{\frac{x+t}{2}}} e^{-t \operatorname{Im} \lambda_p} dt, \quad (4.19)$$

where $c > 0$ is a constant. Since $\operatorname{Im} \lambda_p > 0$ for the eigenvalues $\lambda_p, p = 1, 2, \dots, j$, of L , (4.19) implies that

$$\left\{ \frac{\partial^n}{\partial \lambda^n} E(x, \lambda) \right\}_{\lambda=\lambda_p} \in L_2(\mathbb{R}_+, E), \quad n = 0, 1, \dots, m_p - 1, \quad p = 1, 2, \dots, j. \quad (4.20)$$

So we get $U_{n,p} \in L_2(\mathbb{R}_+, E)$. Similarly we prove the results for $0 \leq n \leq m_p - 1, j+1 \leq p \leq k$. This completes the proof. \square

Let μ_1, \dots, μ_v and μ_{v+1}, \dots, μ_l be the zeros of A_+ and A_- in \mathbb{R}^* with multiplicities n_1, \dots, n_v and n_{v+1}, \dots, n_l , respectively.

We can show

$$\left\{ \frac{\partial^n}{\partial \lambda^n} K(x, \lambda) \right\}_{\lambda=\mu_p} = \sum_{m=0}^n C_m(\lambda_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\mu_p} \quad (4.21)$$

$n = 0, 1, \dots, n_p - 1, p = 1, 2, \dots, v,$

where

$$C_m(\mu_p) = -\binom{n}{m} \left\{ \frac{\partial^{n-m}}{\partial \lambda^{n-m}} A_-(\lambda) \right\}_{\lambda=\mu_p}, \quad (4.22)$$

$$\left\{ \frac{\partial^n}{\partial \lambda^n} K(x, \lambda) \right\}_{\lambda=\mu_p} = \sum_{m=0}^n R_m(\mu_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, -\lambda) \right\}_{\lambda=\mu_p},$$

$n = 0, 1, \dots, n_p - 1, p = v+1, \dots, l,$

where

$$R_m(\mu_p) = \binom{n}{m} \left\{ \frac{\partial^{n-m}}{\partial \lambda^{n-m}} A_+(\lambda) \right\}_{\lambda=\mu_p}. \quad (4.23)$$

Now define the generalized eigenfunctions and generalized associated functions corresponding to the spectral singularities of L by the following :

$$V_{n,p}(x) = \left\{ \frac{\partial^n}{\partial \lambda^n} K(x, \lambda) \right\}_{\lambda=\mu_p} = \sum_{m=0}^n C_m(\mu_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, \lambda) \right\}_{\lambda=\mu_p} \quad (4.24)$$

$n = 0, 1, \dots, n_p - 1, p = 1, 2, \dots, v,$

$$V_{n,p}(x) = \left\{ \frac{\partial^n}{\partial \lambda^n} K(x, \lambda) \right\}_{\lambda=\mu_p} = \sum_{m=0}^n R_m(\mu_p) \left\{ \frac{\partial^m}{\partial \lambda^m} E(x, -\lambda) \right\}_{\lambda=\mu_p},$$

$n = 0, 1, \dots, n_p - 1, p = v + 1, \dots, l.$

Then $V_{n,p}, n = 0, 1, \dots, n_p - 1, p = 1, 2, \dots, v, v + 1, \dots, l,$ also satisfy the equations analogous to (4.16).

$V_{0,p}, V_{1,p}, \dots, V_{n_p-1,p}, p = 1, 2, \dots, v, v + 1, \dots, l$ are called the principal functions corresponding to the spectral singularities $\lambda = \mu_p, p = 1, 2, \dots, v, v + 1, \dots, l$ of L .

Theorem 4.3.

$$V_{n,p} \notin L_2(\mathbb{R}_+, E), \quad n = 0, 1, \dots, n_p - 1, p = 1, 2, \dots, v, v + 1, \dots, l.$$

Proof. For $0 \leq n \leq n_p - 1$ and $1 \leq p \leq v$ using (2.3), we obtain

$$\left\| \left\{ \frac{\partial^n}{\partial \lambda^n} E(x, \lambda) \right\}_{\lambda=\mu_p} \right\| \leq \left\| (ix)^n e^{i\mu_p x} I + \int_x^\infty (it)^n K(x, t) e^{i\mu_p t} dt \right\|,$$

since $\text{Im } \mu_p = 0, p = 1, 2, \dots, v,$ we find that

$$\int_0^\infty \|(ix)^n e^{i\mu_p x} I\|^2 dx = \int_0^\infty x^{2n} dx = \infty.$$

So we obtain $V_{n,p} \notin L_2(\mathbb{R}_+, E), n = 0, 1, \dots, n_p - 1, p = 1, 2, \dots, v.$ Using the similar way, we may also prove the results for $0 \leq n \leq n_p - 1, v + 1 \leq p \leq l.$ \square

Now define the Hilbert spaces of vector-valued functions with values in E by

$$H_n : = \left\{ f : \int_0^\infty (1 + |x|)^{2n} \|f(x)\|^2 dx < \infty \right\}, \quad n = 1, 2, \dots, \quad (4.25)$$

$$H_{-n} : = \left\{ g : \int_0^\infty (1 + |x|)^{-2n} \|g(x)\|^2 dx < \infty \right\}, \quad n = 1, 2, \dots, \quad (4.26)$$

with the norms

$$\|f\|_n^2 := \int_0^\infty (1 + |x|)^{2n} \|f(x)\|^2 dx$$

and

$$\|g\|_n^2 := \int_0^\infty (1 + |x|)^{-2n} \|g(x)\|^2 dx$$

respectively. Then

$$H_{n+1} \subsetneq H_n \subsetneq L_2(\mathbb{R}_+, E) \subsetneq H_{-n} \subsetneq H_{-(n+1)}, \quad n = 1, 2, \dots,$$

and H_{-n} is isomorphic to the dual of H_n .

Theorem 4.4.

$$V_{n,p} \in H_{-(n+1)}, \quad n = 0, 1, \dots, n_p - 1, \quad p = 1, 2, \dots, v, v + 1, \dots, l.$$

Proof. For $0 \leq n \leq n_p - 1$ and $1 \leq p \leq v$ using (2.3) and (4.24), we get

$$\int_0^\infty (1 + |x|)^{-2(n+1)} \|V_{n,p}\|^2 dx \leq M \int_0^\infty (1 + |x|)^{-2(n+1)} \left\{ \begin{array}{l} \{ \|E(x, \lambda)\|^2 \}_{\lambda=\mu_p} + \dots \\ + \left\{ \left\| \frac{\partial^n}{\partial \lambda^n} E(x, \lambda) \right\|^2 \right\}_{\lambda=\mu_p} \end{array} \right\} dx,$$

where $M > 0$ is a constant. Using (2.3), we have

$$\int_0^\infty (1 + |x|)^{-2(n+1)} \|(ix)^n e^{i\mu_p x} I\|^2 dx < \infty$$

and

$$\int_0^\infty (1 + |x|)^{-2(n+1)} \left\| \int_x^\infty (it)^n K(x, t) e^{i\mu_p t} dt \right\|^2 dx < \infty.$$

Consequently $V_{n,p} \in H_{-(n+1)}$ for $0 \leq n \leq n_p - 1$ and $1 \leq p \leq v$. Similarly, we obtain $V_{n,p} \in H_{-(n+1)}$ for $0 \leq n \leq n_p - 1$ and $v + 1 \leq p \leq l$. \square

Theorem 4.5.

$$V_{n,p} \in H_{-n_0}, \quad n = 0, 1, \dots, n_p - 1, \quad p = 1, 2, \dots, v, v + 1, \dots, l.$$

REFERENCES

- [1] M. A. Naimark, Investigation of the spectrum and the expansion in eigenfunctions of a non-selfadjoint operators of second order on a semi-axis, *AMS Translations* 2(16) 1960, 103-193.
- [2] V. E. Lyance, A differential operator with spectral singularities, I, II, *AMS Translations* 2 (60) 1967, 185-225, 227-283.
- [3] E. Bairamov, O. Cakar and A. O. Çelebi, Quadratic pencil of Schrödinger operators with spectral singularities: Discrete spectrum and principal functions, *J. Math. Anal. Appl.* 216 (1997), 303-320.
- [4] E. Bairamov, O. Cakar and A. M. Krall, Spectrum and spectral singularities of a quadratic pencil of a Schrödinger operator with a general boundary condition, *J. Differential Equations* 151 (1999), no. 2, 252-267.
- [5] E. Bairamov, O. Cakar and A. M. Krall, An eigenfunction expansion for a quadratic pencil of a Schrödinger operator with spectral singularities, *J. Differential Equations* 151 (1999), 268-289.
- [6] E. Bairamov and A. O. Celebi, Spectrum and spectral expansion for the non-selfadjoint discrete Dirac operators, *Quart. J. Math. Oxford Ser. (2)* 50 (1999), no. 200, 371-384.
- [7] E. Bairamov and A. O. Çelebi, Spectral properties of the Klein-Gordon s-wave equation with complex potential, *Indian J. Pure Appl. Math.* 28 (1997), 813-824.
- [8] E. Bairamov and G. B. Tunca, Discrete spectrum and principal functions of non-selfadjoint differential operators, *Czechoslovak Math. J.* 49 (1999), 689-700.

- [9] A. M. Krall, E. Bairamov and O. Cakar, Spectral analysis of a non-selfadjoint discrete Schrödinger operators with spectral singularities, *Math. Nachr.* 231 (2001), 89-104.
- [10] R. Carlson, An inverse problem for the matrix Schrödinger equation, *J. Math. Anal. Appl.* 267(2002), 564-575.
- [11] S. Clark and F. Gesztesy, Weyl-Titchmarsh M-function asymptotics, local uniqueness results, trace formulas and Borg-type theorems for Dirac operators, *Trans Amer. Math. Soc.* 354(2002), 3475-3534.
- [12] S. Clark, F. Gesztesy and W. Renger Trace formulas and Borg-type theorems for matrix-valued Jacobi and Dirac finite difference operators, *J. Differential Equations* 219 (2005), 144-182.
- [13] F. Gesztesy, A. Kiselev and K. A. Makarov, Uniqueness results for matrix-valued Schrödinger, Jacobi and Dirac-type operators, *Math. Nachr.* 239 (2002), 103-145.
- [14] N. Yokus, Principal functions of non-selfadjoint Sturm-Liouville problems with eigenvalue-dependent boundary conditions, *Abstract and Applied Analysis*, (2011) 1-12.
- [15] M. Olgun, C. Coskun, Non-selfadjoint matrix Sturm-Liouville operators with spectral singularities *Applied Mathematics and Computations*, 216 (8) (2010)
- [16] C. Coskun, M. Olgun, Principal functions of non-selfadjoint matrix Sturm-Liouville equations, *Journal of Computational and Applied Mathematics*, 235(2011).
- [17] Z. S. Agranovich, V. A. Marchenko, *The Inverse Problem of Scattering Theory*, Gordon and Breach, 1965.
- [18] M. V. Keldysh, " On the completeness of the eigenfunctions of some classes of non-selfadjoint linear operators," *Soviet Mathematics-Doklady*, vol. 77, no. 4, pp. 11-14, 1951.
- [19] M. V. Keldysh, " On the completeness of the eigenfunctions of some classes of non-selfadjoint linear operators," *Russian Mathematical Surveys*, vol. 26, no. 4, pp. 15-41, 1971.
- [20] E. Bairamov, and N. Yokus, Spectral Singularities of Sturm-Liouville Problems with Eigenvalue-Dependent Boundary Conditions, *Abstract and Applied Analysis*, 2009, 1-8.
- [21] M. Olgun, T. Koprubasi and Y. Aygar, Principal Functions of Non-Selfadjoint Difference Operator with Spectral Parameter in Boundary Conditions, *Abstract and Applied Analysis*, pp.1-11, 2011.
- [22] Y. Aygar, ;M. Olgun and T. Koprubasi, Principal Functions of Nonselfadjoint Discrete Dirac Equations with Spectral Parameter in Boundary Conditions, *Abstract and Applied Analysis*, pp.1-15,2012.
- [23] M. Olgun, Non-Selfadjoint Matrix Sturm-Liouville Operators with Eigenvalue-Dependent Boundary Conditions, *Hacettepe J Math & Stat.*, (Accepted).

Current address: Ankara University, Faculty of Sciences

Department of Mathematics Ankara, TURKEY

E-mail address: ccoskun@ankara.edu.tr, deniz.ktr@hotmail.com, olgun@ankara.edu.tr

URL: <http://communications.science.ankara.edu.tr/index.php?series=A1>