# $\left(G^{\prime} / G, 1 / G\right)$-expansion method for analytical solutions of Jimbo-Miwa equation 

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#### Abstract

The main goal of this study is obtaining analytical solutions for (3+1)-dimensional JimboMiwa Equation which the second equation in the well-known KP hierarchy of integrable systems. For the (3+1DJM) equation, hyperbolic, trigonometric, complex trigonometric and rational traveling wave solutions have been constructed by applying the ( $\left.G^{\prime} / G, 1 / G\right)$ expansion method. Then, real and imaginary graphics are presented by giving special values to the constants in the solutions obtained. These graphics are a special solution of the (3+1DJM) equation and represent a stationary wave of the equation. The $\left(G^{\prime} / G, 1 / G\right)$ expansion method is an effective and powerful method for solving nonlinear evolution equations (NLEEs). Ready computer package program is used to obtain the solutions and graphics presented in this study.


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## 1. Introduction

Non-linear science phenomena take an important place in applied mathematics and mathematical physics. The appearance of the solitary wave in nature is quite common particularly in plasmas, condensed matter physics, optical fibers, solid-state physics, chemical kinematics, electrical circuits, etc. Solutions of nonlinear evolutionary equations are the cornerstone of many physical phenomena. For example; such as signaling, ocean waves, shallow water waves, heat dissipation. In wave theory, these physical events are mathematized with traveling wave solutions. Therefore, the traveling wave solution is an indispensable instrument of both mathematics and physics. Last year, NLEEs are applied in many areas of science. Various methods have been used to obtain solutions for different types of NLEEs. Some of these are the extended trial equation method [1], the new extended direct algebraic method [2], $\left(G^{\prime} / G\right)$-expansion method [3], ( $G^{\prime} / G, 1 /$ $G)$-expansion method [4], Clarkson-Kruskal (CK) direct method [5], improved $\tan \phi(\xi) / 2$-expansion method [6], Sumudu transform method [7], new expansion method [8], extended sinh-Gordon equation expansion method [9], the modified Kudryashov method [10], (1/G')-expansion method [11-15], collocation method [16,17], first integral method [18],

F-expansion method [19], the modified exponential function method [20], Difference scheme method [21] and so on [22-29].

In this article, the authors attained the exact solutions of the (3+1DJM) equation. Consider the form of the (3+1DJM) equation [30],
$u_{x x x y}+3 u_{y} u_{x x}+3 u_{x} u_{x y}+2 u_{y t}-3 u_{x z}=0$.
Jimbo-Miwa equation was examined first by Jimbo and Miwa [31] and later by several authors. Some of these are as follows: Wazwaz has been attained multiple soliton solutions of two extended (3+1DJM) equation using the simplified Hirota's method [32], Zhang has been obtained different type solutions of (3+1DJM) equation using Hirota bilinear form [33], Ma and Lee have been attained (3+1DJM) equation using Bäcklund transformation [34], Sun and Chen have been attained lump solutions, the lump-kink of (3+1DJM) equation via bilinear forms [35], Yang and Ma have been obtained Lump-type solutions of the (3+1DJM) equation by applying Hirota bilinear form [36], Liu and Jiang have been attained new exact solutions of the ( $3+1$ DJM) equation with the help of extended homogeneous balance method [37], Öziș and Aslan have been attained analytical and explicit generalized solitary solutions of the (3+1DJM) equation using the Exp-function method [38], Tang

[^0]and Liang have been obtained variable separation solutions of the (3+1DJM) equation [39], Ma has been
attained four classes of lump-type solutions for the (3+1DJM) equation using Hirota bilinear form [40].

## 2. $\left(G^{\prime} / G, 1 / G\right)$-Expansion Method

Consider the general form of NLEEs containing two or more arguments to be analyzed is written as follows [23]:
$T\left(u, \frac{\partial u}{\partial t}, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial^{2} u}{\partial x^{2}}, \ldots\right)=0$.
Let $u=u(x, y, z, t)=U(\xi)=U, \quad \xi=\alpha x+\beta y+\gamma z-c t$ and transmutation Eq. (1) can be converted into following nODE for $U(\xi)$ :
$\mathrm{L}\left(U, U^{\prime}, U^{\prime \prime}, U U^{\prime}, \ldots\right)=0$,
where prime refers to derivatives related to $\xi$. Complexity can be reduced by integrating Eq. (3). By the nature of this method, $G(\xi)=G$ is solution function of the second-order ODE as
$G^{\prime \prime}(\xi)+\lambda G(\xi)=\mu$.
It as $\phi=\phi(\xi)=G^{\prime} / G$ and $\psi=\psi(\xi)=1 / G$ supply operational esthetic. We may write derivatives of the functions described
$\phi^{\prime}=-\phi^{2}+\mu \psi-\lambda, \psi^{\prime}=-\phi \psi$,
We can offer the behavior of solution function Eq. (4) according to the state of $\lambda$ taking into account the equations given by the Eq. (5).
i) If $\lambda<0$,
$G(\xi)=c_{1} \sinh \left(\sqrt{-\lambda} \xi+c_{2} \cosh \sqrt{-\lambda} \xi\right)+\frac{\mu}{\lambda}$,
whereas $c_{1}$ and $c_{2}$ are real numbers. Considering Eq. (6);
$\psi^{2}=\frac{-\lambda}{\lambda^{2} \sigma+\mu^{2}}\left(\phi^{2}-2 \mu \psi+\lambda\right), \sigma=c_{1}^{2}-c_{2}^{2}$,
Eq. (7) is obtained.
ii) If $\lambda>0$,
$G(\xi)=c_{1} \sin \left(\sqrt{\lambda} \xi+c_{2} \cos \sqrt{\lambda} \xi\right)+\frac{\mu}{\lambda}$,
whereas $c_{1}$ and $c_{2}$ are real numbers. Considering Eq. (8),
$\psi^{2}=\frac{\lambda}{\lambda^{2} \sigma-\mu^{2}}\left(\phi^{2}-2 \mu \psi+\lambda\right), \sigma=c_{1}^{2}+c_{2}^{2}$,
Eq. (9) is written.
iii) If $\lambda=0$,
$G(\xi)=\frac{\mu}{2} \xi^{2}+c_{1} \xi+c_{2}$,
here $c_{1}$ and $c_{2}$ are real numbers. Considering Eq. (10),
$\psi^{2}=\frac{1}{c_{1}^{2}-2 \mu c_{2}}\left(\phi^{2}-2 \mu \psi\right)$.
Eq. (11) is written.
For $\psi$ and $\phi$ polynomials, solution of Eq. (3) is
$u(\xi)=\sum_{i=0}^{n} a_{i} \phi^{i}+\sum_{i=1}^{n} b_{i} \phi^{i-1} \psi$.

Here, $a_{i}$ and $b_{i}$ are real numbers be identified. $n$ is a positive equilibrium term which may be attained by comparing the maximum order derivative with the maximum order nonlinear term in Eq. (8). Whether Eq. (12)
is written in Eq. (3) along with equations (5), (7), (9), or (11) is written. Each coefficient of $\phi^{i} \psi^{i}(i=$ $0,1, \ldots, n)(j=1, \ldots, n)$ terms of the attained polynomial functions are equals zero and there is a system of algebraic equations for $a_{i}, b_{i}, c_{1}, c_{2}, \alpha, \beta, c$, and $\lambda,(i=0,1, \ldots, n)$. The necessary coefficients are obtained by solving an algebraic equation with the help of computer package programs. Obtained coefficients are put into Eq. (12) and $U(\xi)$ solution function of the ODE given as (3) is attained and if $\xi=\alpha x+\beta y+\gamma z-c t$ transmutation is employed, we will attain exact solution $u(x, y, z, t)$ of Eq. (2).

## 3. Solutions of the (3+1DJM) Equation

We consider Eq. (1) and using transformation $\quad u=u(x, y, z, t)=U(\xi)=U, \quad \xi=\alpha x+\beta y+\gamma z-c t$, where $\beta, \gamma, \alpha$ and $c$ are constants, once obtained (ODE) after integration, we get
$\alpha^{3} \beta U^{\prime \prime \prime}+3 \beta \alpha^{2}\left(U^{\prime}\right)^{2}-(2 \beta c+3 \alpha \gamma) U^{\prime}=0$,
where the integration constant is zero. In Eq. (13), we get term $n=1$ from the definition of balancing term and the following situation is obtained in Eq. (5),

$$
\begin{equation*}
U(\xi)=a_{0}+a_{1} \phi[\xi]+b_{1} \psi[\xi] . \tag{14}
\end{equation*}
$$

Replacing Eq. (14) into Eq. (13) and the coefficients of Eq. (1) are equal to zero, we may establish the following algebraic equation systems
Const: $2 c \beta \lambda a_{1}+3 \alpha \gamma \lambda a_{1}-2 \alpha^{3} \beta \lambda^{2} a_{1}+\frac{3 \alpha^{3} \beta \lambda^{2} \mu^{2} a_{1}}{\mu^{2}+\lambda^{2} \sigma}+3 \alpha^{2} \beta \lambda^{2} a_{1}^{2}-\frac{3 \alpha^{2} \beta \lambda^{2} \mu^{2} a_{1}^{2}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\phi[\xi]:-\frac{6 \alpha^{3} \beta \lambda^{2} \mu b_{1}}{\mu^{2}+\lambda^{2} \sigma}+\frac{6 \alpha^{2} \beta \lambda^{2} \mu a_{1} b_{1}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\phi[\xi]^{2}: 2 c \beta a_{1}+3 \alpha \gamma a_{1}-8 \alpha^{3} \beta \lambda a_{1}+\frac{3 \alpha^{3} \beta \lambda \mu^{2} a_{1}}{\mu^{2}+\lambda^{2} \sigma}+6 \alpha^{2} \beta \lambda a_{1}^{2}-\frac{3 \alpha^{2} \beta \lambda \mu^{2} a_{1}^{2}}{\mu^{2}+\lambda^{2} \sigma}-\frac{3 \alpha^{2} \beta \lambda^{2} b_{1}^{2}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\phi[\xi]^{3}:-\frac{6 \alpha^{3} \beta \lambda \mu b_{1}}{\mu^{2}+\lambda^{2} \sigma}+\frac{6 \alpha^{2} \beta \lambda \mu a_{1} b_{1}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\psi[\xi]:-2 c \beta \mu a_{1}-3 \alpha \gamma \mu a_{1}+5 \alpha^{3} \beta \lambda \mu a_{1}-\frac{6 \alpha^{3} \beta \lambda \mu^{3} a_{1}}{\mu^{2}+\lambda^{2} \sigma}-6 \alpha^{2} \beta \lambda \mu a_{1}^{2}+\frac{6 \alpha^{2} \beta \lambda \mu^{3} a_{1}^{2}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\phi[\xi] \psi[\xi]: 2 c \beta b_{1}+3 \alpha \gamma b_{1}-5 \alpha^{3} \beta \lambda b_{1}+\frac{12 \alpha^{3} \beta \lambda \mu^{2} b_{1}}{\mu^{2}+\lambda^{2} \sigma}+6 \alpha^{2} \beta \lambda a_{1} b_{1}-\frac{12 \alpha^{2} \beta \lambda \mu^{2} a_{1} b_{1}}{\mu^{2}+\lambda^{2} \sigma}=0$,
$\phi[\xi]^{2} \psi[\xi]: 12 \alpha^{3} \beta \mu a_{1}-6 \alpha^{2} \beta \mu a_{1}^{2}+\frac{6 \alpha^{2} \beta \lambda \mu b_{1}^{2}}{\mu^{2}+\lambda^{2} \sigma}=0$,
the aim with ready computer package program, reaching the solutions of system (15) and we obtained the following stations.
If $\lambda<0$,
Case 1.
$a_{1}=\alpha, \quad b_{1}=i \alpha \sqrt{\lambda} \sqrt{\sigma}, c=\frac{-3 \alpha \gamma-\alpha^{3} \beta \lambda}{2 \beta}, \mu=0, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (16) into Eq. (14) and attain the following hyperbolic exact solutions for Eq. (1):
$u_{1}(x, y, z, t)=\frac{2 \alpha\left(c_{2} \sqrt{-\lambda} \operatorname{Cosh}[\sqrt{-\lambda} \xi]+c_{1} \sqrt{-\lambda} \operatorname{Sinh}[\sqrt{-\lambda} \xi]\right)}{c_{1} \operatorname{Cosh}[\sqrt{-\lambda} \xi]+c_{2} \operatorname{Sinh}[\sqrt{-\lambda} \xi]}+a_{0}$.


Figure 1. Graphs of $u_{1}(x, y, z, t)=u_{1}$ for $\lambda=-4, \beta=1, \gamma=5, \delta_{1}=4, \delta_{2}=-1, a_{0}=1, \alpha=2, c_{1}=1, c_{2}=2$, $y=1, z=1$.

## Case 2.

$a_{1}=2 \alpha, b_{1}=0, c=\frac{-3 \alpha \gamma-4 \alpha^{3} \beta \lambda}{2 \beta}, \mu=0, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (18) into Eq. (14) and attain the following hyperbolic exact solution for Eq. (1):
$u_{2}(x, y, z, t)=\frac{2 \alpha\left(c_{2} \sqrt{-\lambda} \operatorname{Cosh}[\sqrt{-\lambda} \xi]+c_{1} \sqrt{-\lambda} \operatorname{Sinh}[\sqrt{-\lambda} \xi]\right)}{c_{1} \operatorname{Cosh}[\sqrt{-\lambda \xi}]+c_{2} \operatorname{Sinh}[\sqrt{-\lambda} \xi]}+a_{0}$.


Figure 2. Graphs of $u_{2}(x, y, z, t)=u_{2}$ for $\lambda=-1, \beta=1, a_{0}=1, \gamma=5, \alpha=2, c_{1}=1, c_{2}=2, \mathrm{y}=1, \mathrm{z}=1$.
ii) If $\lambda>0$,

## Case 3.

$a_{1}=2 \alpha, b_{1}=0, c=\frac{-3 \alpha \gamma-4 \alpha^{3} \beta \lambda}{2 \beta}, \mu=0, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (20) into Eq. (14) and attain the following trigonometric exact solution for Eq. (1):

$$
\begin{equation*}
u_{3}(x, y, z, t)=\frac{2 \alpha\left(c_{2} \sqrt{\lambda} \operatorname{Cos}[\sqrt{\lambda} \xi]-c_{1} \sqrt{\lambda} \operatorname{Sin}[\sqrt{\lambda} \xi]\right)}{c_{1} \operatorname{Cos}[\sqrt{\lambda} \xi]+c_{2} \operatorname{Sin}[\sqrt{\lambda} \xi]}+a_{0} . \tag{21}
\end{equation*}
$$



Figure 3. Graphs of $u_{3}(x, y, z, t)=u_{3}$ for $\lambda=1, \beta=1, a_{0}=3, \gamma=1, \alpha=0.8, c_{1}=-1, c_{2}=1, \mathrm{y}=0, \mathrm{z}=0$.

Case 4.
$a_{1}=\alpha, b_{1}=-\frac{\alpha \sqrt{-\mu^{2}+\lambda^{2} \sigma}}{\sqrt{\lambda}}, \beta=-\frac{3 \alpha \gamma}{2 c+\alpha^{3} \lambda}, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (22) into Eq. (14) and attain the following trigonometric exact solution for Eq. (1):
$u_{4}(x, y, z, t)=-\frac{\alpha \sqrt{\left(c_{1}^{2}+c_{2}{ }^{2}\right) \lambda^{2}-\mu^{2}}}{\sqrt{\lambda}\left(\frac{\mu}{\lambda}+c_{1} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]+c_{2} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]\right)}+$
$\frac{\alpha\left(c_{2} \sqrt{\lambda} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]-c_{1} \sqrt{\lambda} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]\right)}{\frac{\mu}{\lambda}+c_{1} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]+c_{2} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]}+a_{0}$.


Figure 4. Graphs of $u_{4}(x, y, z, t)=u_{4}$ for $\lambda=16, c=1, a_{0}=5, \gamma=1, \alpha=1, \mu=1, c_{1}=-1, c_{2}=1, \mathrm{y}=0, \mathrm{z}=$ 0.

## Case 5.

$a_{1}=\alpha, b_{1}=\alpha \sqrt{\lambda} \sqrt{\sigma}, c=\frac{-3 \alpha \gamma-\alpha^{3} \beta \lambda}{2 \beta}, \mu=0, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (24) into Eq. (14) and attain the following trigonometric exact solution for Eq. (1):
$u_{5}(x, y, z, t)=\frac{\sqrt{c_{1}^{2}+c_{2}^{2}} \alpha \sqrt{\lambda}}{c_{1} \operatorname{Cos}[\sqrt{\lambda} \xi]+c_{2} \operatorname{Sin}[\sqrt{\lambda} \xi]}+\frac{\alpha\left(c_{2} \sqrt{\lambda} \operatorname{Cos}[\sqrt{\lambda} \xi]-c_{1} \sqrt{\lambda} \operatorname{Sin}[\sqrt{\lambda} \xi]\right)}{c_{1} \operatorname{Cos}[\sqrt{\lambda} \xi]+c_{2} \operatorname{Sin}[\sqrt{\lambda} \xi]}+a_{0}$.


Figure 5. Graphs of $u_{5}(x, y, z, t)=u_{5}$ for $\lambda=4, \beta=10, \gamma=1, \alpha=1, a_{0}=10, c_{1}=1, c_{2}=-2, \mathrm{y}=1, \mathrm{z}=1$,
$c=1$.
Case 6.
$a_{1}=\alpha, b_{1}=\alpha \sqrt{\lambda} \sqrt{\sigma}, \mu=0, \beta=-\frac{3 \alpha \gamma}{2 c+\alpha^{3} \lambda}, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (26) into Eq. (14) and attain the following complex trigonometric exact solution for Eq. (1):
$u_{6}(x, y, z, t)=\frac{\sqrt{c_{1}^{2}+c_{2}^{2}} \alpha \sqrt{\lambda}}{c_{1} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]+c_{2} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]}+$
$\frac{\alpha\left(c_{2} \sqrt{\lambda} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]-c_{1} \sqrt{\lambda} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]\right)}{c_{1} \operatorname{Cos}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]+c_{2} \operatorname{Sin}\left[\sqrt{\lambda}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c+\alpha^{3} \lambda}\right)\right]}+a_{0}$.



Figure 6. Real and imaginary graphs of $u_{6}(x, y, z, t)=u_{6}$ for $\lambda=4, c=1, a_{0}=10, \gamma=8, c_{1}=2, c_{2}=10, \mathrm{y}=1$, $\mathrm{z}=1$.
iii) If $\lambda=0$,

## Case 7.

$a_{1}=\alpha, b_{1}=-\sqrt{c_{2}^{2} \alpha^{2}-2 c_{1} \alpha^{2} \mu}, c=-\frac{3 \alpha \gamma}{2 \beta}, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (28) into Eq. (14) and attain the following rational exact solution for Eq. (1):
$u_{7}(x, y, z, t)=-\frac{\sqrt{c_{2}^{2} \alpha^{2}-2 c_{1} \alpha^{2} \mu}}{c_{1}+c_{2}\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right)+\frac{1}{2}\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right)^{2} \mu}+\frac{\alpha\left(c_{2}+\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right) \mu\right)}{c_{1}+c_{2}\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right)+\frac{1}{2}\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right)^{2} \mu}+$ $a_{0}$.


Figure 7. Graphs of $u_{7}(x, y, z, t)=u_{7}$ for $\mu=2, \beta=1, a_{0}=10, \gamma=5, \alpha=2, c_{1}=3, c_{2}=2, \mathrm{y}=1, \mathrm{z}=1$.

## Case 8.

$a_{1}=2 \alpha, b_{1}=0, c=-\frac{3 \alpha \gamma}{2 \beta}, \mu=0, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (30) into Eq. (14) and attain the following rational exact solution for Eq. (1):
$u_{8}(x, y, z, t)=\frac{2 c_{2} \alpha}{c_{1}+c_{2}\left(x \alpha+y \beta+z \gamma+\frac{3 t \alpha \gamma}{2 \beta}\right)}+a_{0}$.


Figure 8. Graphs of $u_{8}(x, y, z, t)=u_{8}$ for $\beta=4, a_{0}=15, \gamma=3, a_{0}=5, \alpha=5, c_{1}=1, c_{2}=1, \mathrm{y}=1, \mathrm{z}=1$.

## Case 9.

$a_{1}=2 \alpha, b_{1}=0, \mu=0, \beta=-\frac{3 \alpha \gamma}{2 c}, \xi=\alpha x+\beta y+\gamma z-c t$,
replacing values Eq. (32) into Eq. (14) and attain the following rational exact solution for Eq. (1):
$u_{9}(x, y, z, t)=\frac{2 c_{2} \alpha}{c_{1}+c_{2}\left(-c t+x \alpha+z \gamma-\frac{3 y \alpha \gamma}{2 c}\right)}+a_{0}$.


Figure 9. Graphs of $u_{9}(x, y, z, t)=u_{9}$ for $c=1, \beta=1, a_{0}=5, \gamma=4, \alpha=7, c_{1}=5, c_{2}=3, \mathrm{z}=1.5, \mathrm{y}=2$.

## 4. Conclusion

The $\left(G^{\prime} / G, 1 / G\right)$-expansion method was used to establish hyperbolic, trigonometric, complex trigonometric, and rational traveling wave solutions for the (3+1DJM) equation. For the solutions found, real and imaginary graphics are presented for different values given to the constants. This equation has been presented by many authors by applying different methods, traveling wave solutions. The $\left(G^{\prime} / G, 1 / G\right)$ expansion method was applied to the (3+1DJM) equation for the first time and the solutions produced by this method are of different types. The traveling wave solutions produced by the $\left(G^{\prime} / G, 1 / G\right)$ expansion method are trigonometric, complex trigonometric, hyperbolic and rational type solutions. There is a singular point within these solutions. Therefore, constants in solutions presented in shock wave theory and asymptotic behavior analysis will be much more valuable when the physical meaning is loaded. In this study, the applied method is effective, powerful, and will be used in future works to establish new exact solutions of many other NLEEs. Also, the ready computer package program is used for graphics and computations in this letter.

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## Conflict of Interest Statement

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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