

RESEARCH ARTICLE

Spectral properties of non-selfadjoint Sturm-Liouville operator equation on the real axis

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Abstract

In this paper, we analyze the non-selfadjoint Sturm-Liouville operator L defined in the Hilbert space $L_2(\mathbb{R}, H)$ of vector-valued functions which are strongly-measurable and square-integrable in \mathbb{R} . L is defined

$$L(y) = -y'' + Q(x)y, \quad x \in \mathbb{R},$$

for every $y \in L_2(\mathbb{R}, H)$ where the potential Q(x) is a non-selfadjoint, completely continuous operator in a separable Hilbert space H for each $x \in \mathbb{R}$. We obtain the Jost solutions of this operator and examine the analytic and asymptotic properties. Moreover, we find the point spectrum and the spectral singularities of L and also obtain the sufficient condition which assures the finiteness of the eigenvalues and spectral singularities of L.

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1. Introduction

Non-selfadjoint operators are seen in physical systems which do not involve the conservation of energy law. Some selfadjoint problems also give us non-selfadjoint operators after separation of variables. The theory of non-selfadjoint operators has initially begun to analyze ordinary differential equations. M.V. Keldysh played a significant role to develop a general theory for non-selfadjoint operators by inventing a new method for establishing the resolvent of an arbitrary completely continuous, non-selfadjoint operator of finite order [16, 17].

Spectral analysis of non-selfadjoint differential operators has been studied by M.A. Naimark [24, 25]. In particular, he analyzed the non-selfadjoint Sturm-Liouville operator defined by

$$l(y) = -y'' + p(x)y, \quad 0 < x < \infty,$$
(1.1)

$$y'(0) - hy(0) = 0, (1.2)$$

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where p(x) is a complex-valued function satisfying

$$\int_0^\infty (1+t^2)|p(t)|dt < \infty, \tag{1.3}$$

and $h \in \mathbb{C}$. Several authors investigated the non-selfadjoint Sturm-Liouville operator defined by Equations (1.1) and (1.2) in detail [22–25, 28]. The results of Naimark [24, 25] have been generalized in [7,8] to the operator l_0 generated in $L_2(\mathbb{R})$ which is defined by

$$l_o(y) = -y'' + q(x)y, \quad x \in \mathbb{R},$$

where the potential q is a complex-valued function. The authors generalized the results of [24] and applied to the non-selfadjoint Schrödinger operator in the three-dimensional space [13].

Non-selfadjoint Hamiltonians and complex extensions of Quantum Mechanics have been studied by many mathematicans, recently. Moreover, spectral properties of the selfadjoint matrix differential and difference equations have been examined [9, 10, 15]. For the non-selfadjoint case, discrete spectrum and the spectral singularities of the matrix Sturm-Liouville operator were investigated [4, 11, 26, 27]. Further, the authors examined a system of non-selfadjoint Sturm-Liouville equations [2, 5, 6].

B. M. Levitan et al. have studied the point spectrum of the following Sturm-Liouville operator equation in detail [14,19–21]. Let H be a separable Hilbert space and $L_2(\mathbb{R}_+, H)$ denote the space of vector-valued functions f(x) defined on $(0, \infty)$ which are strongly-integrable and also square-integrable on $(0, \infty)$ i.e.,

$$\int_0^\infty \|f(x)\|^2 \, dx < \infty.$$

Consider the operator l_1 defined on $L_2(\mathbb{R}_+, H)$ by

$$l_1(Y) = -Y'' + Q(x)Y, \quad 0 < x < \infty,$$
(1.4)

and the boundary condition Y(0) = 0 where Q(x) is a completely continuous, selfadjoint operator defined on H for every $x \in (0, \infty)$. Equation (1.4) is called Sturm-Liouville operator equation.

In our previous paper [3], we considered the non-selfadjoint analogue of the above problem and investigated the spectral properties of the non-selfadjoint Sturm-Liouville operator equation on the half line on the contrary to [14, 19-21]. We also generalized the results in [2, 4, 11, 26, 27] by considering the coefficients as operators not only finite dimensional matrices. In this study, we extend these results to the whole real axis. Explicitly, we focus on the following non-selfadjoint operator.

Assume H is a separable Hilbert space and $H_1 := L_2(\mathbb{R}, H)$ denotes the space of vectorvalued functions f(x) defined on \mathbb{R} which are strongly-integrable and square-integrable. Note that H_1 is a Hilbert space with the inner product (see [29]);

$$(f,g)_1 = \int_{-\infty}^{\infty} (f(x),g(x))_H dx.$$

Let us consider the non-selfadjoint operator L defined in H_1 ;

$$L(y) = -y'' + Q(x)y, \quad x \in \mathbb{R},$$
(1.5)

where the potential Q(x) is a non-selfadjoint, completely continuous operator in H for each $x \in \mathbb{R}$. In this paper, we specify the domain of L and express the Jost solutions. Then, we find the discrete spectrum and the set of spectral singularities of L by using the properties of the Jost solutions. Finally, we prove that L has a finite number of eigenvalues and spectral singularities. The domain D(L) of L is the subspace consisting of all $y \in H_1$ which satisfies the following conditions;

(i) y is twice strongly-differentiable,

(ii) $L(y) \in H_1$.

Let us consider the eigenvalue equation;

$$-y'' + Q(x)y = \lambda^2 y, \quad x \in \mathbb{R}.$$
(1.6)

2. The Jost solutions of (1.6)

We shall also focus on the equation

$$-Y'' + Q(x)Y = \lambda^2 Y, \quad x \in \mathbb{R},$$
(2.1)

where Y(x) is an operator-valued function i.e, Y(x) is an operator in H for each $x \in \mathbb{R}$.

Lemma 2.1. Every sequence of solutions of Equation (1.6) can be represented as an operator-valued function which satisfies Equation (2.1). Conversely, one can construct a sequence of vector-valued functions which satisfy Equation (1.6) for a given operator-valued solution of Equation (2.1).

Proof. Since H is a separable Hilbert space, there exists an orthonormal basis $(u_n)_{n \in \mathbb{N}}$. Suppose vector-valued functions $(y_n(x))_{n \in \mathbb{N}}$ satisfy Equation (1.6). We can construct an operator-valued function Y(x) such that $Y(x)u_n = y_n(x)$ for every $n \in \mathbb{N}$. It is clear that Y(x) satisfies Equation (2.1).

Conversely, suppose operator-valued function Y(x) satisfies Equation (2.1). Let $y_n(x) = Y(x)u_n$ for every $n \in \mathbb{N}$. Then, it is clear that $(y_n(x))$ satisfies Equation (1.6) for every $n \in \mathbb{N}$.

As a result of this one to one correspondence, we can focus on the solutions of only one of the Equations (1.6)-(2.1).

We shall use the notations;

$$\sigma^{+}(x) = \int_{x}^{\infty} \|Q(t)\| dt, \quad \sigma_{1}^{+}(x) = \int_{x}^{\infty} \sigma^{+}(t) dt,$$
$$\sigma^{-}(x) = \int_{-\infty}^{x} \|Q(t)\| dt, \quad \sigma_{1}^{-}(x) = \int_{-\infty}^{x} \sigma^{-}(t) dt.$$

Suppose that the condition

$$\int_{-\infty}^{\infty} (1+|t|) \, \|Q(t)\| \, dt < \infty, \tag{2.2}$$

holds. Then, Equation (2.1) has operator solutions $E^+(x,\lambda)$ and $F^-(x,\lambda)$ satisfying the initial conditions;

$$\lim_{x \to \infty} e^{-i\lambda x} E^+(x,\lambda) = I, \quad \text{Im}\lambda \ge 0,$$
(2.3)

and

$$\lim_{x \to -\infty} e^{i\lambda x} F^{-}(x,\lambda) = I, \quad \text{Im}\lambda \ge 0,$$
(2.4)

respectively. Indeed, consider the integral equation

$$E^{+}(x,\lambda) = e^{i\lambda x}I + \frac{1}{\lambda} \int_{x}^{\infty} \sin\left(\lambda\left(t-x\right)\right) Q(t)E^{+}(t,\lambda)dt, \quad \text{Im}\lambda \ge 0,$$

which is easily seen to be a solution of Equation (2.1) satisfying (2.3). Similarly, if we define

$$F^{-}(x,\lambda) = E^{+}(-x,\lambda), \quad \text{Im}\lambda \ge 0,$$

it easily follows that $F^{-}(x, \lambda)$ satisfies (2.4). Under the condition (2.2), the solution $E^{+}(x, \lambda)$ can be represented (see [1]);

$$E^{+}(x,\lambda) = e^{i\lambda x}I + \int_{x}^{\infty} e^{i\lambda t}K^{+}(x,t)dt, \quad \text{Im}\lambda \ge 0.$$
(2.5)

Let us consider the equation;

$$-Z'' + ZQ(x) = \lambda^2 Z, \quad x \in \mathbb{R},$$
(2.6)

where Z(x) is an operator-valued function. Similarly, Equation (2.6) has an operator solution $E^{-}(x, \lambda)$ which satisfies the initial condition;

$$\lim_{x \to -\infty} e^{i\lambda x} E^{-}(x,\lambda) = I, \quad \text{Im}\lambda \ge 0,$$

and has the representation

$$E^{-}(x,\lambda) = e^{-i\lambda x}I + \int_{-\infty}^{x} e^{-i\lambda t}K^{-}(x,t)dt, \quad \text{Im}\lambda \ge 0.$$

Further, the operator-valued kernels $K^{+}(x,t)$ are differentiable with respect to x and t and satisfy

$$\left\| K^{+}(x,t) \right\| \leq \frac{1}{2} \sigma^{+}(\frac{x+t}{2}) \exp\left[\sigma^{+}_{1}(x) - \sigma^{+}_{1}(\frac{x+t}{2}) \right], \qquad (2.7)$$

$$\left|K_{x}^{+}(x,t) - \frac{1}{4}Q(\frac{x+t}{2})\right| \leq \frac{1}{2}\sigma_{1}^{+}(x)\sigma_{-}^{+}(\frac{x+t}{2})\exp\sigma_{1}^{+}(x),$$
(2.8)

$$\left\|K_{t}^{+}(x,t) - \frac{1}{4}Q(\frac{x+t}{2})\right\| \leq \frac{1}{2}\sigma_{1}^{+}(t)\sigma_{-}^{+}(\frac{x+t}{2})\exp\sigma_{1}^{+}(t),$$
(2.9)

As a result, the solutions $E^+(x,\lambda)$ and $E^-(x,\lambda)$ are analytic for $\text{Im}\lambda > 0$ and continuous for $\text{Im}\lambda \ge 0$. $E^+(x,\lambda)$ and $E^-(x,\lambda)$ are called the Jost solutions of Equation (1.6). The proofs of above results are very similar to the matrix coefficient case which have been obtained in [1,4]. In addition, we obtained analogous properties in our previous paper [3]. Hence, we omitted the proofs.

Lemma 2.2. Let Y(x) be a solution of Equation (2.1) and Z(x) be a solution of Equation (2.6). Then, the Wronskian W[Y, Z](x) := Z'(x)Y(x) - Z(x)Y'(x) is independent of x.

Proof. We have

$$-Y'' + Q(x)Y = \lambda^2 Y,$$

$$-Z'' + ZQ(x) = \lambda^2 Z.$$

If we multiply the first equality from the left with Z and the second equality from the right with Y and subtract them, we have

$$Z''(x)Y(x) - Z(x)Y''(x) = 0$$

which implies W[Y, Z](x) is constant and hence independent of the variable x.

Let us define the function

$$D(\lambda) := W\left[E^{-}(x,\lambda), E^{+}(x,\lambda)\right], \quad \text{Im}\lambda \ge 0.$$

Since the Wronskian of $E^+(x,\lambda)$ and $E^-(x,\lambda)$ is independent of x, $D(\lambda)$ is a function of λ which is also analytic for $\text{Im}\lambda > 0$ and continuous for $\text{Im}\lambda \ge 0$. The function $D(\lambda)$ is called the Jost function of Equation (1.6).

Theorem 2.3. The function $D(\lambda)$ satisfies

$$D(\lambda) = 2i\lambda I - 2K^{+}(0,0) - 2K^{-}(0,0) + \int_{0}^{\infty} e^{i\lambda t} F(t)dt, \qquad (2.10)$$

where

$$F(t) = K_x^+(0,t) - K_x^-(0,-t) - K^-(0,0)K^+(0,t) - K^+(0,0)K^-(0,-t) + K^-(0,-t) * K_x^+(0,t) - K_x^-(0,-t) * K^+(0,t) + K_t^-(0,-t) - K_t^+(0,t),$$
(2.11)

and $F \in L_1(\mathbb{R}, H)$ where "*" is the convolution operation.

Proof. Since the Wronskian of $E^+(x, \lambda)$ and $E^-(x, \lambda)$ is independent of x, we put x = 0 and obtain

$$D(\lambda) = W\left[E^{-}(x,\lambda), E^{+}(x,\lambda)\right] = E_{x}^{+}(\lambda)E^{-}(\lambda) - E^{+}(\lambda)E_{x}^{-}(\lambda)$$

By using the integral representations of $E^+(x, \lambda)$ and $E^-(x, \lambda)$ we get (2.10) and (2.11). From (2.7)-(2.9) we have $F \in L_1(\mathbb{R}, H)$.

Theorem 2.4. The following asymptotic relations hold;

$$D(\lambda) = 2i\lambda I - 2K^{+}(0,0) - 2K^{-}(0,0) + o(1), \quad \text{Im}\lambda \ge 0, \quad |\lambda| \to \infty, \quad (2.12)$$

$$D(\lambda) = 2i\lambda I + O(1), \quad \text{Im}\lambda \ge 0, \quad |\lambda| \to \infty.$$
 (2.13)

Proof. Let $\lambda \in \mathbb{R}$. By Riemann-Lebesgue Lemma for Fourier transforms [18] we have

$$\int_0^\infty e^{i\lambda t} F(t)dt = o(1), \quad \lambda \in \mathbb{R}, \quad |\lambda| \to \infty.$$
(2.14)

Now, let $\text{Im}\lambda > 0$. Lebesgue Theorem [18] implies

$$\int_0^\infty e^{i\lambda t} F(t) dt = o(1), \quad \text{Im}\lambda > 0, \quad |\lambda| \to \infty.$$
(2.15)

If we use (2.14), (2.15) we get (2.12). The proof is similar for (2.13).

3. Point spectrum and spectral singularities of L

Now, we introduce the point spectrum and the set of spectral singularities of L according to the definitions given in [22–24]

$$\sigma_d(L) = \left\{ \lambda^2 : \text{Im}\lambda > 0, \ D(\lambda) \text{ is not invertible} \right\},$$

$$\sigma_{ss}(L) = \left\{ \lambda^2 : \ \lambda \in \mathbb{R} \smallsetminus \{0\}, \ D(\lambda) \text{ is not invertible} \right\}.$$

Now, we try to examine the eigenvalues of L by employing the results in [17]. Let us recall;

$$D(\lambda) = 2i\lambda I - 2K^{+}(0,0) - 2K^{-}(0,0) + \int_{0}^{\infty} e^{i\lambda t} F(t)dt, \quad \text{Im}\lambda \ge 0.$$
(3.1)

Let

$$\begin{aligned} A(\lambda) &:= \frac{1}{2i\lambda} \left[-2K^{+}(0,0) - 2K^{-}(0,0) + \int_{0}^{\infty} e^{i\lambda t} F(t) dt \right], \\ G(\lambda) &:= \frac{1}{2i\lambda} D(\lambda). \end{aligned}$$

Then,

$$G(\lambda) = I + A(\lambda), \quad \text{Im}\lambda \ge 0,$$

and for $\lambda \neq 0$, it follows $D(\lambda)$ is invertible iff $G(\lambda)$ is invertible. Hence

$$\sigma_d(L) = \left\{ \lambda^2 : \operatorname{Im} \lambda > 0, \ G(\lambda) \text{ is not invertible} \right\},$$

$$\sigma_{ss}(L) = \left\{ \lambda^2 : \ \lambda \in \mathbb{R} \smallsetminus \{0\}, \ G(\lambda) \text{ is not invertible} \right\}.$$

1690

Let us define $M_1 := \{\lambda : \text{Im}\lambda > 0, \ G(\lambda) \text{ is not invertible}\}$. It follows $\sigma_d(L) = \{\lambda^2 : \lambda \in M_1\}$.

Since $\int_{-\infty}^{\infty} (1 + |t|) ||Q(t)|| dt < \infty$ and Q(x) is completely continuous operator for each $x \in \mathbb{R}$, it follows F(t) is completely continuous operator for each $0 < t < \infty$ and as a result $A(\lambda)$ is completely continuous for $\mathrm{Im}\lambda > 0$. Also, since $D(\lambda)$ is analytic for $\mathrm{Im}\lambda > 0$, $A(\lambda)$ is also analytic in the same domain. Now, we can use the results in [17].

Definition 3.1. If

$$(I-R)\left(I+A\right) = I,$$

holds, then the operator R is called the resolvent of the operator A, [17].

Let us denote the resolvent of $A(\lambda)$ by $R(\lambda)$. It follows

$$I - R(\lambda) = (I + A(\lambda))^{-1} = (G(\lambda))^{-1}$$

If $I - R(\lambda)$ exists at least for one λ which means $G(\lambda)$ is invertible, this implies $I - R(\lambda)$ exists on the domain $\mathbb{C}_+ := \{z \in \mathbb{C} : \operatorname{Im} \lambda > 0\}$ except for a set of isolated points, and also $I - R(\lambda)$ is a meromorphic operator function in the same domain [17]. It is obvious that $M_1 \neq \mathbb{C}_+$. This implies there exists at least one λ such that $I - R(\lambda)$ is defined. [17] implies that $I - R(\lambda)$ should exist on the domain \mathbb{C}_+ except for a set of isolated points. These isolated points are obviously the eigenvalues of L. Moreover, $I - R(\lambda)$ is a meromorphic operator function on \mathbb{C}_+ . Therefore, we can express $I - R(\lambda)$ as a ratio of two analytical functions in the domain \mathbb{C}_+ as;

$$(G(\lambda))^{-1} = I - R(\lambda) = \frac{S(\lambda)}{d(\lambda)},$$
(3.2)

where $S(\lambda)$ is an operator function and $d(\lambda)$ is a scalar function on \mathbb{C}_+ . Moreover, the above isolated points are poles of the function $I - R(\lambda)$ and they are the zeros of the function $d(\lambda)$. As a result, it follows

$$M_1 = \{\lambda : \text{Im}\lambda > 0, \ d(\lambda) = 0\}.$$
 (3.3)

Theorem 3.2. If the condition (2.2) holds, then $\sigma_d(L)$ is a bounded and countable set. Further, the limit points of $\sigma_d(L)$ should lie in a bounded interval of the real axis.

Proof. The relation (2.12) implies

$$G(\lambda) = I + o(1), \quad \text{Im}\lambda \ge 0, \quad |\lambda| \to \infty,$$

which means for sufficiently large $\lambda \in \overline{\mathbb{C}_+}$, $G(\lambda) \to I$ and thus $G(\lambda)$ is invertible. Therefore, M_1 is bounded. Since the function $d(\lambda)$ is analytic, its zeros are isolated. This implies M_1 is countable. Further, the limit points of the zeros of $d(\lambda)$ should lie in an interval of the real line [12]. The proof is complete since

$$\sigma_d(L) = \left\{ \lambda^2 : \lambda \in M_1 \right\}.$$

Now, let us assume that the condition

$$\int_{-\infty}^{\infty} e^{\epsilon |t|} \left\| Q(t) \right\| dt < \infty, \ \epsilon > 0, \tag{3.4}$$

holds.

Theorem 3.3. L has a finite number of eigenvalues.

Proof. From the equalities (2.7)-(2.9) and (3.4) we have

$$\left\| K^{+}(x,t) \right\|, \ \left\| K^{+}_{x}(x,t) \right\|, \ \left\| K^{+}_{t}(x,t) \right\| \le c \exp\left(-\epsilon(\frac{x+t}{2})\right),$$

and hence

$$\|F(t)\| \le c \exp\left(-\epsilon(\frac{t}{2})\right), \ \forall t \in [0,\infty),$$

where c is a positive constant. Further,

$$\|F(t)\| \left| e^{i\lambda t} \right| \le c \exp^{-t\left(\frac{\epsilon}{2} + \operatorname{Im}\lambda\right)}, \quad \forall t \in [0,\infty),$$

and thus

$$\begin{split} \left\| \int_0^\infty e^{i\lambda t} F(t) dt \right\| &\leq \int_0^\infty \left| e^{i\lambda t} \right| \|F(t)\| \, dt \\ &\leq \int_0^\infty c \exp^{-t\left(\frac{\epsilon}{2} + \operatorname{Im}\lambda\right)} dt, \\ \int_0^\infty c \exp^{-t\left(\frac{\epsilon}{2} + \operatorname{Im}\lambda\right)} dt < \infty \Leftrightarrow \operatorname{Im}\lambda + \frac{\epsilon}{2} > 0. \end{split}$$

The Uniform Convergence Test implies that the integral $\int_0^\infty e^{i\lambda t} F(t) dt$ is uniformly convergent in the domain $\text{Im}\lambda > -\frac{\epsilon}{2}$. This implies $D(\lambda)$ and also $G(\lambda)$ have analytic continuations to the domain $\text{Im}\lambda > -\frac{\epsilon}{2}$. Since the analytic continuation is unique, it follows

$$(G(\lambda))^{-1} = I - R(\lambda) = \frac{S(\lambda)}{d(\lambda)}, \ \operatorname{Im}\lambda > -\frac{\epsilon}{2}.$$

Let us recall $M_1 = \{\lambda : \text{Im}\lambda > 0, \ d(\lambda) = 0\}, \ \sigma_d(L) = \{\lambda^2 : \lambda \in M_1\}$ and M_1 is bounded. Suppose that M_1 is not finite. Let us recall Bolzano-Weierstrass Theorem which states that each bounded sequence in \mathbb{R} has a convergent subsequence. Bolzano-Weierstrass Theorem implies that M_1 must have one limit point. Also, Theorem 3.2 states that the limit points of M_1 can only lie on the real axis. However, since $d(\lambda)$ is analytic in the domain $\text{Im}\lambda > -\frac{\epsilon}{2}$, the limit points of M_1 should be on the boundary of this domain [12]. This contradicts with the assumption $\epsilon > 0$. Thus, M_1 and $\sigma_d(L)$ are finite.

Let $M_2 := \{\lambda : \lambda \in \mathbb{R}, G(\lambda) \text{ is not invertible}\}$. It is obvious that

$$\sigma_{ss}(L) = \left\{ \lambda^2 : \ \lambda \in M_2 \right\} \smallsetminus \{0\}.$$

From the representation (3.2), we have

$$\frac{G(\lambda)}{d(\lambda)}S(\lambda) = I, \quad \lambda \in \mathbb{C}_+,$$

and this implies $S(\lambda)$ is invertible iff $G(\lambda)$ is invertible or equivalently $d(\lambda) \neq 0$ for $\lambda \in \mathbb{C}_+$. If $d(\lambda) \neq 0$ it follows

$$(S(\lambda))^{-1} = \frac{G(\lambda)}{d(\lambda)}, \ \lambda \in \mathbb{C}_+,$$

and also

$$G(\lambda) = d(\lambda) (S(\lambda))^{-1}, \quad \lambda \in \mathbb{C}_+.$$
(3.5)

Since $G(\lambda)$ is continuous on the real line, Equation (3.5) suggests that the functions $S(\lambda)$ and $d(\lambda)$ are continuous on the real line. Hence, we can extend the representation (3.2) continuously to the real line and obtain

$$(G(\lambda))^{-1} = \frac{S(\lambda)}{d(\lambda)}, \ \lambda \in \overline{\mathbb{C}_+}.$$
(3.6)

(3.6) implies that $G(\lambda)$ is invertible iff $d(\lambda) \neq 0$ for $\lambda \in \mathbb{R}$. Thus, we have

$$M_2 = \{\lambda \in \mathbb{R} : d(\lambda) = 0\}$$

Theorem 3.4. M_2 is compact and has zero Lebesgue measure under the condition (2.2).

Proof. Theorem 3.2 implies M_2 is bounded. We only have to show that M_2 is closed. Let $\{\lambda_n\} \subset M_2$ such that $\lambda_n \to \lambda_0$. $\{\lambda_n\} \subset M_2$ implies $\lambda_n \in \mathbb{R}$ and $G(\lambda_n)^{-1}$ doesn't exist for $n \in \mathbb{N}$. Further, $\lambda_n \to \lambda_0$ implies $\lambda_0 \in \mathbb{R}$. We have $G(\lambda)$ is a continuous operator function on the real line. Now, $\lambda_n \to \lambda_0$ suggests that $G(\lambda_n) \to G(\lambda_0)$ where the latter convergence is strong.

Let $GL(H) := \{A : A \text{ is invertible, bounded, linear operator on } H\}$. It is well known that GL(H) is an open subset of the space B(H) of bounded, linear operators on H. It follows $B(H) \setminus GL(H)$ is a closed set. This implies $G(\lambda_0) \in B(H) \setminus GL(H)$ and $\lambda_0 \in M_2$. Finally, Privalov's Theorem states that M_2 has zero Lebesgue measure [12].

Corollary 3.5. $\sigma_{ss}(L)$ is bounded and has zero Lebesgue measure, under the condition (2.2).

Theorem 3.6. L has a finite number of spectral singularities, under the condition (3.4).

Proof. It can be shown similarly (see the proof of Theorem 3.3) that $G(\lambda)$ has an analytic continuation to the domain $\text{Im}\lambda > -\frac{\epsilon}{2}$ for arbitrary $\epsilon > 0$. Since this analytic continuation is unique, it follows

$$(G(\lambda))^{-1} = I - R(\lambda) = \frac{S(\lambda)}{d(\lambda)}, \text{ Im}\lambda > -\frac{\epsilon}{2}.$$

Suppose that M_2 is not finite. Since M_2 is bounded (see Theorem 3.2), Bolzano-Weierstrass Theorem implies that M_2 has a limit point. This limit point as a zero of the analytic function $d(\lambda)$ should lie on the boundary of the domain $\text{Im}\lambda > -\frac{\epsilon}{2}$ [12]. It contradicts with $\epsilon > 0$.

References

- Z.S. Agranovic, V.A. Marchenko, The Inverse Problem of Scattering Theory, Gordon and Breach, 1965.
- [2] E.K. Arpat and G. Mutlu, Spectral properties of Sturm-Liouville system with eigenvalue-dependent boundary conditions, Internat. J. Math. 26 (10), 1550080-1550088, 2015.
- [3] E. Bairamov, E.K. Arpat and G. Mutlu, Spectral properties of non-selfadjoint Sturm-Liouville operator with operator coefficient, J. Math. Anal. Appl. 456 (1), 293-306, 2017.
- [4] E. Bairamov and Ş. Cebesoy, Spectral singularities of the matrix Schrödinger equations, Hacet. J. Math. Stat. 45 (4), 1007-1014, 2016.
- [5] E. Bairamov and E. Kir, Principal functions of the non-selfadjoint operator generated by system of differential equations, Math. Balkanica (N.S.) 13 (1-2), 85–98, 1999.
- [6] E. Bairamov and E. Kir, Spectral properties of a finite system of Sturm-Liouville differential operators, Indian J. Pure Appl. Math. 35 (2), 249–256, 2004.
- [7] E. Bairamov and G.B. Tunca, Discrete spectrum and principial functions of nonselfadjoint differential operator, Czechoslavak Math. J. 49 (124), 689-700, 1999.
- [8] B.B. Blashak, On the second-order differential operator on the whole axis with spectral singularities (In Russian), Dokl. Akad. Nauk Ukr. SSR I, 38-41, 1966.
- [9] R. Carlson, An inverse problem for the matrix Schrödinger equation, J. Math. Anal. Appl. 267, 564-575, 2002.
- [10] S. Clark, F. Gesztesy and W. Renger, Trace formulas and Borg-type theorems for matrix-valued Jacobi and Dirac finite difference operators, J. Differ. Equations 219, 144-182, 2005.
- [11] C. Coskun and M. Olgun, Principal functions of non-selfadjoint matrix Sturm-Liouville equations, J. Comput. Appl. Math. 235, 4834-4838, 2011.

- [12] E.P Dolzhenko, Boundary value uniqueness theorems for analytic functions, Math. Notes 26 (6), 437-442, 1979.
- [13] M.G. Gasymov, Expansion in solutions of the scattering problem for a nonselfadjoint Schrödinger equation (In Russian), Dokl. Akad. Nauk AzSSR 22 (10), 9-12, 1966.
- [14] M.G. Gasymov, V.V. Zikov and B.M. Levitan, Conditions for discreteness and finiteness of the negative spectrum of Schrödinger's operator equation (in Russian), Mat. Zametki 2, 531–538, 1967.
- [15] F. Gesztesy, A. Kiselev and K.A. Makarov, Uniqueness results for matrix-valued Schrodinger, Jacobi and Dirac-type operators, Math. Nachr. 239, 103-145, 2002.
- [16] M.V. Keldysh, On eigenvalues and eigenfunctions of some classes of nonselfadjoint equations (in Russian), Dokl. Akad. Nauk. SSSR 77 (1), 11-14, 1951.
- [17] M.V. Keldysh, On the completeness of the eigenfunctions of some classes of nonselfadjoint linear operators, Russ. Math. Surv. 26 (4), 1544, 1971.
- [18] A.N. Kolmogorov and S.V. Fomin, *Introductory Real Analysis*, Dover Publ., New York, 1975.
- [19] A.G. Kostjucenko and B. M. Levitan, Asymptotic behavior of eigenvalues of the operator Sturm-Liouville problem (in Russian), Funkcional. Anal. i Prilozen 1, 86–96, 1967.
- [20] B.M. Levitan, Investigation of the Green's function of a Sturm-Liouville equation with an operator coefficient (in Russian), Mat. Sb. (N.S.) 76 (118), 239–270, 1968.
- [21] B.M. Levitan and G. A. Suvorcenkova, Sufficient conditions for discreteness of the spectrum of a Sturm-Liouville equation with operator coefficient (in Russian), Funkcional. Anal. i Prilozen 2 (2), 56–62, 1968.
- [22] V.E. Lyance, A differential operator with spectral singularities, II, Amer. Math. Soc. Transl. Ser. 2 60, 227283, 1967.
- [23] B. Nagy, Operators with spectral singularities, J. Operat. Theor. 15 (2), 307-325, 1986.
- [24] M.A. Naimark, Investigation of the spectrum and the expansion in eigenfunctions of a non-selfadjoint operator of second order on a semi-axis, Tr. Mosk. Mat. Obs. 3, 181270, 1954.
- [25] M.A. Naimark, Linear differential operators, II, Ungar, New York, 1968.
- [26] M. Olgun, Non-selfadjoint matrix Sturm-Liouville operators with eigenvaluedependent boundary conditions, Hacet. J. Math. Stat. 44 (3), 607614, 2015.
- [27] M. Olgun and C. Coskun, Non-selfadjoint matrix Sturm-Liouville operators with spectral singularities, Appl. Math. Comput. 216, 2271-2275, 2010.
- [28] B.S. Pavlov, The Nonself-Adjoint Schrödinger Operator, in: Birman M.S. (eds) Spectral Theory and Wave Processes, Topics in Mathematical Physics, 1, 87-114, Springer, Boston, MA, 1967.
- [29] K. Yosida, Functional Analysis, Springer-Verlag, Berlin-Göttingen-Heidelberg, 1980.