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# **BL-Algebras with Permuting Tri-Derivations**

## **Damla Yılmaz**<sup>1</sup> **[ID](https://orcid.org/0000-0002-6741-8669)**

**Article Info** *Received*: 29 Apr 2023 *Accepted*: 11 Jul 2023 *Published*: 30 Sep 2023 doi[:10.53570/jnt.1289799](https://doi.org/10.53570/jnt.1289799) Research Article **Abstract** − Basic logic algebras (BL-algebras) were introduced by Hajek. Multi-value algebras (MV-algebras), Gödel algebras, and product algebras are particular cases of BLalgebras. Moreover, BL-algebras are algebraic structures, and their principal examples are the real interval [0*,* 1] with the structure given by a continuous *t*-norm and abelian *l*-groups. In this article, we consider a type of derivation structure on BL-algebras. We study (⊙*,* ∨) permuting tri-derivations of BL-algebras and their examples and basic properties. We obtain results regarding the trace of  $(⊙, ∨)$ -permuting derivations on Gödel BL-algebras. Finally, the article presents that the results herein can be generalized in future research.

**Keywords** *BL-algebra, permuting tri-derivation, Boolean algebra, isotone, trace*

**Mathematics Subject Classification (2020)** 08A05, 03G25

#### **1. Introduction**

The interest in multi-value algebras (MV-algebras) introduced by Chang [1] continues to increase. One of the most important examples of MV-algebra is the interval [0*,* 1] of the commutative *l*-group:  $(M, \text{max}, \text{min}, +, ', 0)$  equipped with a continuous t-conorm and continuous t-norm defined by  $m \oplus n =$  $\min \{m + n, 1\}$  and  $m \odot n = \max \{m + n - 1, 0\}$  and with the negation defined by  $m' = 1 - m$ . After MV-algebras, Gödel algebras and product algebras have been investigated [2–5]. These three structures form the most important algebraic model structure for fuzzy logic, which are Lukasiewicz logic, Gödel logic and product logic, respectively. These logics studying these algebras are of logical interest, as well as their connection to some mathematical structures, as they correspond to the most important continuous *t*-norms on [0*,* 1] and their associated residues. Basic logic algebras (BL-algebras) [6] has been introduced by Hajek. MV-algebras, Gödel algebras, and product algebras are particular cases of BL-algebras. Various derivation studies [7–9] have been done on BL-algebras. In this article, we investigate a derivation type defined by some authors [10–13] in rings, lattices, and MV-algebras. More precisely, we define a type of permuting tri-derivations on BL-algebras and study some of its properties.

This article is organized as follows: The next section reminds some results and basic properties of BL-algebras. Section 3 defines (⊙*,* ∨)-permuting tri-derivation structure in BL-algebras and obtains their some results. Moreover, it explores many properties by the trace of the  $(\odot, \vee)$ -permuting triderivation on BL-algebras. Finally, the conclusion briefly overviews this type of algebra and discusses future studies.

<sup>1</sup>damla.yilmaz@erzurum.edu.tr (Corresponding Author)

<sup>&</sup>lt;sup>1</sup>Department of Mathematics, Faculty of Science, Erzurum Technical University, Erzurum, Türkiye

## **2. Preliminaries**

This section provides some basic notions to be needed for the next section. A *t*-norm ⊗ on the real interval [0*,* 1] is a commutative and associative operation on [0*,* 1] such that

*i.* If  $\beta \leq \delta$ , then  $\beta \otimes \eta \leq \delta \otimes \eta$ , for all  $\beta, \delta, \eta \in [0, 1]$ *ii.*  $1 \otimes \beta = \beta$ , for all  $\beta \in [0, 1]$ 

From *i* and  $ii, 0 \otimes \beta = 0$ , for all  $\beta \in [0, 1]$ . A *t*-norm  $\otimes$  on [0, 1] is continuous if it is continuous in the usual sense as a function  $\otimes$  :  $[0,1] \times [0,1] \rightarrow [0,1]$ . On [0,1], Lukasiewicz *t*-norm, Gödel *t*-norm, and Product *t*-norm, the most important continuous *t*-norms, are defined as follows, respectively:

$$
\beta \odot \delta = \max \{ \beta + \delta - 1, 0 \}
$$

$$
\beta \odot \delta = \min \{ \beta, \delta \}
$$

and

*β* ⊙ *δ* = *βδ*

Note that every continuous *t*-norm  $\otimes$  on [0, 1] induces a binary operation  $\rightarrow$  defined by

$$
\beta \to \delta = \max \{ \eta \in [0,1] | \beta \otimes \eta \le \delta \}
$$

called the associated residuum. The binary operation satisfies the statement  $\beta \otimes \delta \leq \eta$  if and only if (iff)  $\beta \leq \delta \to \eta$ , for all  $\beta, \delta, \eta \in [0,1]$ . These three norms above are related to fuzzy logic along with the residues Lukasiewicz implication, Gödel implication, and Product implication, defined as follows, respectively:

$$
\beta \to \delta = \min \{ 1, 1 - \beta - \delta \}
$$

$$
\beta \to \delta = \begin{cases} 1, & \beta \le \delta \\ \delta, & \text{otherwise} \end{cases}
$$

and

$$
\beta \to \delta = \begin{cases} 1, & \beta \le \delta \\ \frac{\delta}{\beta}, & \text{otherwise} \end{cases}
$$

If  $\otimes$  is a continuous *t*-norm on [0, 1] and  $\rightarrow$  is the associated residuum, then the structure

 $([0, 1], \max, \min, \otimes, \rightarrow, 0, 1)$ 

is the starting point in describing and investigating Basic Logic and corresponds to the basic logic system: BL-algebras.

**Definition 2.1.** [6] A basic logic algebra (BL-algebra) is a structure  $(\Lambda, \wedge, \vee, \odot, \rightarrow, 0, 1)$  such that

- *i.*  $(\Lambda, \wedge, \vee, 0, 1)$  is a bounded lattice
- *ii.*  $(\Lambda, \odot, 1)$  is a commutative monoid
- *iii.*  $\beta \leq \delta \to \eta$  iff  $\delta \odot \beta \leq \eta$  (residuation)
- *iv.*  $\beta \wedge \eta = \beta \odot (\beta \rightarrow \eta)$  (divisibility)
- *v.*  $(\beta \to \eta) \lor (\eta \to \beta) = 1$  (prelinearity)

For any  $\beta, \eta \in \Lambda$ , we define

$$
\beta^* = \beta \to 0, \quad \beta \oplus \eta = (\beta^* \odot \eta^*)^*, \quad \text{and} \quad \beta \ominus \eta = \beta \odot \eta^*
$$

We recall that a BL-algebra  $\Lambda$  is an MV-algebra iff  $(\beta^*)^* = \beta$ , for all  $\beta \in \Lambda$  [4]. A Gödel algebra is a

BL-algebra  $\Lambda$  satisfying the condition  $\beta \odot \beta = \beta$ , for all  $\beta \in \Lambda$  [6]. Moreover, a product algebra is a BL-algebra  $\Lambda$  satisfying the conditions  $i$ .  $\beta \wedge \beta^* = 0$  and  $ii$ .  $(\delta^*)^* \odot (\beta \odot \delta \rightarrow \eta \odot \delta) \leq \beta \rightarrow \eta$ , for all  $β, η, δ \in Λ$  [6].

**Proposition 2.2.** [2] Suppose that 
$$
\Lambda
$$
 is a BL-algebra and  $\beta, \eta, \delta \in \Lambda$ . Thus, the following should:  
\ni.  $\beta \leq \eta$  iff  $\beta \to \eta = 1$   
\nii.  $\beta \to (\eta \to \delta) = (\beta \odot \eta) \to \delta = \eta \to (\beta \to \delta)$   
\niii. If  $\beta \leq \eta$ , then  $\eta \to \delta \leq \beta \to \delta$ ,  $\delta \to \beta \leq \delta \to \eta$ ,  $\beta \odot \delta \leq \eta \odot \delta$ , and  $\eta^* \leq \beta^*$   
\niv.  $\beta \leq (\eta \to \beta) \to \beta$ ,  $\eta \leq (\eta \to \beta) \to \beta$ , and  $\beta \vee \eta = ((\beta \to \eta) \to \eta) \vee ((\eta \to \beta) \to \beta)$   
\n*v.*  $\beta \odot \eta \leq \beta$ ,  $\beta \odot \eta \leq \eta$ ,  $\beta \odot \eta \leq \beta \land \eta$ ,  $\beta \odot 0 = 0$ , and  $\beta \odot \beta^* = 0$   
\n*vi.*  $1 \to \beta = \beta$ ,  $\beta \to \beta = 1$ ,  $\beta \leq \eta \to \beta$ ,  $\beta \to 1 = 1$ , and  $0 \to \beta = 1$   
\n*viii.*  $\beta \odot \eta = 0$  iff  $\beta \leq \eta^*$   
\n*viii.*  $\beta \odot (\eta \wedge \sigma) = (\beta \odot \eta) \wedge (\beta \odot \sigma)$   
\n*ix.*  $\beta \odot (\eta \vee \sigma) = (\beta \odot \eta) \vee (\beta \odot \sigma)$ 

A BL-algebra satisfying  $\beta \vee \beta^* = 1$  is called a Boolean algebra. For a BL-algebra  $\Lambda$ , if

$$
B(\Lambda) = \{ \beta \in \Lambda : \beta \oplus \beta = \beta \} = \{ \beta \in \Lambda : \beta \odot \beta = \beta \}
$$

then  $(B(\Lambda), \oplus, *, 0)$  is the largest subalgebra of  $\Lambda$  and a Boolean algebra. Hence,  $B(\Lambda)$  is called Boolean center of Λ.

**Theorem 2.3.** [2] Let  $\Lambda$  be a BL-algebra and  $\beta, \eta \in \Lambda$ . Then, the following conditions are equivalent:

*i.*  $\beta \in B(\Lambda)$ *ii.*  $\beta \odot \beta = \beta$  and  $\beta^{**} = \beta$ *iii.*  $\beta \odot \beta = \beta$  and  $\beta^* \rightarrow \beta = \beta$ *iv.*  $\beta^* \vee \beta = 1$ *v.*  $(\beta \to \eta) \to \beta = \beta$ *vi. β* ∧ *η* = *β* ⊙ *η*

#### **3.**(⊙*,* ∨)**-Permuting Tri-derivations on BL-algebras**

This section defines (⊙*,* ∨)-permuting tri-derivation structures in BL-algebras and explores their some results. It investigates the properties provided by the trace of the (⊙*,* ∨)-permuting tri-derivation on BL-algebras. Throughout this paper, Λ denotes a BL-algebra unless otherwise specified.

**Definition 3.1.** A map  $\Phi : \Lambda \times \Lambda \times \Lambda \to \Lambda$  is called a permuting mapping if  $\Phi(\beta, \delta, \theta) = \Phi(\beta, \theta, \delta) =$  $\Phi(\delta, \beta, \theta) = \Phi(\delta, \theta, \beta) = \Phi(\theta, \beta, \delta) = \Phi(\theta, \delta, \beta)$ , for all  $\beta, \delta, \theta \in \Lambda$ .

Furthermore, a map  $\varphi : \Lambda \to \Lambda$  defined by  $\varphi(\beta) = \Phi(\beta, \beta, \beta)$  is referred to as the trace of  $\Phi$  such that  $\Phi$  is a permuting mapping. Hereinafter, for brevity, we use the notation  $\varphi\beta$  instead of  $\varphi(\beta)$ .

**Definition 3.2.** Let  $\Phi : \Lambda \times \Lambda \times \Lambda \to \Lambda$  be a permuting mapping. If  $\Phi$  satisfies the following condition

$$
\Phi(\beta \odot \eta, \delta, \theta) = (\Phi(\beta, \delta, \theta) \odot \eta) \oplus (\beta \odot \Phi(\eta, \delta, \theta))
$$

for all  $\beta, \eta, \delta, \theta \in \Lambda$ , then  $\Phi$  is called a permuting tri-derivation. Clearly, if  $\Phi$  is a permuting tri-

derivation on  $\Lambda$ , then the following relations hold:

$$
\Phi(\beta,\delta\odot\eta,\theta)=(\Phi(\beta,\delta,\theta)\odot\eta)\oplus(\delta\odot\Phi(\beta,\eta,\theta))
$$

and

$$
\Phi(\beta,\delta,\theta\odot\eta)=(\Phi(\beta,\delta,\theta)\odot\eta)\oplus(\theta\odot\Phi(\beta,\delta,\eta))
$$

for all  $\beta, \eta, \delta, \theta \in \Lambda$ .

**Example 3.3.** Let  $\Lambda = \{0, \beta, \theta, 1\}$  such that  $0 < \beta < \theta < 1$  and the binary operations  $\odot$  and  $\rightarrow$  be defined as follows:



Then,  $(\Lambda, \wedge, \vee, \odot, \rightarrow, 0, 1)$  is a BL-algebra. Define a map  $\Phi : \Lambda \times \Lambda \times \Lambda \rightarrow \Lambda$  by

$$
\Phi(x_1, x_2, x_3) = \begin{cases} \beta, & (x_1, x_2, x_3) \in \left\{ \begin{array}{l} (\beta, \beta, \beta), (\beta, \beta, \theta), (\beta, \theta, \beta), (\theta, \beta, \beta), (\beta, \beta, 1), \\ (\beta, 1, \beta), (1, \beta, \beta), (\theta, \theta, 1), (\theta, 1, \theta), (1, \theta, \theta) \end{array} \right\} \\ \theta, & (x_1, x_2, x_3) = (\theta, \theta, \theta) \\ 0, & \text{otherwise} \end{cases}
$$

It is easy to observe that  $\Phi$  is a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ .

**Proposition 3.4.** Assume that  $\Phi$  is a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ ,  $\varphi$  is the trace of  $\Phi$ , and  $\beta, \delta, \theta \in \Lambda$ . Then, the following conditions are valid:

*i.* 
$$
\varphi 0 = 0
$$
  
\n*ii.*  $\varphi \beta \odot \beta^* = \beta \odot \varphi \beta^* = 0$   
\n*iii.*  $\varphi \beta = \varphi \beta \vee (\beta \odot \Phi(\beta, \beta, 1))$   
\n*iv.* If  $\beta \in B(\Lambda)$ , then  $\beta \le (\Phi(\beta, \beta, \beta^*))^*$   
\n*v.* If  $\beta \in B(\Lambda)$ , then  $\Phi(\beta, \delta, \theta) \le \beta$  and  $\Phi(\beta^*, \delta, \theta) \le \beta^*$   
\nPROOF.  
\n*i.* For all  $\beta \in \Lambda$ ,  
\n
$$
\Phi(\beta, \beta, 0) = \Phi(\beta, \beta, 0 \odot 0)
$$
\n
$$
= (\Phi(\beta, \beta, 0) \odot 0) \vee (0 \odot \Phi(\beta, \beta, 0)))
$$
\n
$$
= 0 \vee 0 = 0
$$
\nSince  $\varphi$  is the trace of  $\Phi$ ,  
\n
$$
\varphi 0 = \Phi(0, 0, 0)
$$
\n
$$
= \Phi(0 \odot 0, 0, 0)
$$
\n
$$
= (\Phi(0, 0, 0) \odot 0) \vee (0 \odot \Phi(0, 0, 0))
$$
\n
$$
= 0 \vee 0 = 0
$$
\n*ii.* For all  $\beta \in \Lambda$ ,  
\n
$$
0 = \Phi(\beta, \beta, \beta)
$$
\n
$$
= \Phi(\beta, \beta, \beta) \odot \beta^*)
$$
\n
$$
= (\Phi(\beta, \beta, \beta) \odot \beta^*) \vee (\beta \odot \Phi(\beta, \beta, \beta^*))
$$

Hence,  $\varphi\beta \odot \beta^* = 0$  and  $\beta \odot \Phi(\beta, \beta, \beta^*) = 0$ . Similarly,  $\beta \odot \varphi\beta^* = 0$ , for all  $\beta \in \Lambda$ . *iii.* For all  $\beta \in \Lambda$ ,

$$
\varphi \beta = \Phi(\beta, \beta, \beta)
$$
  
=  $\Phi(\beta, \beta, \beta \odot 1)$   
=  $(\Phi(\beta, \beta, \beta) \odot 1) \vee (\beta \odot \Phi(\beta, \beta, 1))$   
=  $\varphi \beta \vee (\beta \odot \Phi(\beta, \beta, 1))$ 

*iv.* Let  $\beta \in B(\Lambda)$ . Since  $\Phi(\beta, \beta, \beta^*) \odot \beta = 0$ , then  $\Phi(\beta, \beta, \beta^*) \leq \beta^*$ . From Theorem 2.3,  $\beta^{**} = \beta$ because  $\beta \in B(\Lambda)$ . From Proposition 2.2 *iii*,  $\beta \leq (\Phi(\beta, \beta, \beta^*))^*$ .

*v.* Let *β* ∈ *B*( $Λ$ ). For all  $δ, θ ∈ Λ$ ,

$$
\Phi(\beta \odot \beta^*, \delta, \theta) = (\Phi(\beta, \delta, \theta) \odot \beta^*) \vee (\beta \odot \Phi(\beta^*, \delta, \theta))
$$

which implies that  $\Phi(\beta, \delta, \theta) \odot \beta^* = 0$  and  $\beta \odot \Phi(\beta^*, \delta, \theta) = 0$ . Thus,  $\Phi(\beta, \delta, \theta) \leq \beta$  and  $\Phi(\beta^*, \delta, \theta) \leq \beta^*$ .

 $\Box$ 

**Proposition 3.5.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ ,  $\varphi$  be the trace of  $\Phi$ , and  $\beta$ ,  $\theta \in \Lambda$ such that  $\beta \leq \theta$ . Then, the followings hold:

- *i.*  $\varphi(\beta \odot \theta^*) = 0$
- *ii.*  $\varphi \theta^*$  ≤  $\beta^*$
- *iii.* If  $\beta \in B(\Lambda)$ , then  $\varphi \beta \odot \varphi \theta^* = 0$

Proof*.*

*i.* Since  $\beta \leq \theta$ , then  $\beta \odot \theta^* \leq \theta \odot \theta^* = 0$ . Hence,  $\beta \odot \theta^* = 0$ . Thus,  $\varphi(\beta \odot \theta^*) = 0$ .

*ii.* Since  $\beta \leq \theta$ , then  $\beta \odot \varphi \theta^* \leq \theta \odot \varphi \theta^* = 0$ . Hence,  $\varphi \theta^* \odot \beta = 0$ . From Proposition 2.2 *vii*,  $\beta \in B(\Lambda)$ . Then,  $\varphi \theta^* \leq \beta^*$ .

*iii.* Let  $\beta \in B(\Lambda)$ . For all  $\theta \in \Lambda$ ,

$$
0 = \Phi(\beta \odot \beta^*, \theta, \theta) = (\Phi(\beta, \theta, \theta) \odot \beta^*) \vee (\beta \odot \Phi(\beta^*, \theta, \theta))
$$

Hence,  $\Phi(\beta, \theta, \theta) \odot \beta^* = 0$ . Therefore,  $\Phi(\beta, \theta, \theta) \leq \beta$ . Thus, we replace  $\theta$  by  $\beta$  in the last relation,  $\varphi\beta \leq \beta$ . Therefore,  $\varphi\beta \odot \varphi\theta^* \leq \beta \odot \varphi\theta^* = 0$ .

 $\Box$ 

**Proposition 3.6.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ ,  $\varphi$  be the trace of  $\Phi$ , and  $\beta \in B(\Lambda)$ . Then,

*i.*  $\varphi \beta \odot \varphi \beta^* = 0$ 

*ii.*  $\varphi \beta^* = (\varphi \beta)^*$  iff  $\varphi$  is the identity on  $\Lambda$ 

Proof*.*

*i*. By Proposition 3.5 *iii*, as  $\beta \leq \beta$  and  $\beta \in B(\Lambda)$ , then  $\varphi \beta \odot \varphi \beta^* = 0$ .

*ii.* ( $\Rightarrow$ ) : Since  $\beta \odot \varphi \beta^* = 0$ , then  $\beta \odot (\varphi \beta)^* = 0$ . Thus,  $\varphi \beta \leq \beta$  and  $\beta \leq \varphi \beta$  which implies that  $\varphi\beta = \beta$ . Hence,  $\varphi$  is identity on  $\Lambda$ .

 $(\Leftarrow)$ : If  $\varphi$  is an identity on  $\Lambda$ , then  $\varphi \beta^* = (\varphi \beta)^*$ , for all  $\beta \in \Lambda$ .

**Definition 3.7.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ . If  $\beta \leq \delta$  implies  $\Phi(\beta, \theta, \eta) \leq$  $\Phi(\delta, \theta, \eta)$ , for all  $\beta, \theta, \eta \in \Lambda$ , then  $\Phi$  is called an isotone. If  $\varphi$  is the trace of  $\Phi$ , and  $\Phi$  is an isotone, then  $\beta \leq \delta$  implies  $\varphi \beta \leq \varphi \delta$ , for all  $\beta, \delta \in \Lambda$ .

**Example 3.8.** Let  $\Lambda$  be a BL-algebra and  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation in Example 3.3. Then,  $\Phi$  is not isotone, because  $\Phi(\beta, \beta, 1) \nleq \Phi(1, \beta, 1)$ .

**Example 3.9.** Consider  $\Lambda = \{0, \beta, \theta, 1\}$  in Example 3.3. Define a map  $\Phi : \Lambda \times \Lambda \times \Lambda \to \Lambda$  by

$$
\Phi(x_1, x_2, x_3) = \begin{cases}\n\beta, (x_1, x_2, x_3) \in \begin{cases}\n(\beta, \beta, \beta), (\beta, \beta, \theta), (\beta, \theta, \beta), (\theta, \beta, \beta), (\beta, \beta, 1), (\beta, 1, \beta), (1, \beta, \beta), \\
(\beta, \theta, \theta), (\theta, \beta, \theta), (\theta, \theta, \beta), (\beta, \theta, 1), (\beta, 1, \theta), (\theta, \beta, 1), (\theta, 1, \beta), \\
(1, \beta, \theta), (1, \theta, \beta), (\beta, 1, 1), (1, \beta, 1), (1, \beta, 1), (1, 1, \beta)\n\end{cases} \\
\theta, \quad (x_1, x_2, x_3) \in \left\{ (\theta, \theta, \theta), (\theta, \theta, 1), (\theta, 1, \theta), (1, \theta, \theta), (\theta, 1, 1), (1, \theta, 1), (1, 1, \theta)\right\} \\
1, \quad (x_1, x_2, x_3) = (1, 1, 1) \\
0, \quad \text{otherwise}\n\end{cases}
$$

Then,  $\Phi$  is an isotone  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ .

**Proposition 3.10.** Suppose that  $\Phi$  is a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$  and  $\varphi$  is the trace of  $\Phi$ . If  $\varphi \beta^* = \varphi \beta$ , for all  $\beta \in \Lambda$ , then the followings hold:

- *i.*  $\varphi 1 = 0$
- *ii.*  $\varphi \beta \odot \varphi \beta = 0$
- *iii.* If  $\Phi$  is an isotone on  $\Lambda$ , then  $\varphi = 0$

Proof*.*

*i*. Replacing  $\beta$  by 0 in the relation  $\varphi \beta^* = \varphi \beta$ , since  $0^* = 1$  and  $\varphi 0 = 0$ , then  $\varphi 1 = 0$ .

*ii.* Using hypothesis,  $\varphi \beta \odot \varphi \beta = \varphi \beta \odot \varphi \beta^* = 0$ , for all  $\beta \in \Lambda$ .

*iii.* Let  $\Phi$  be an isotone on  $\Lambda$ . For  $\beta \in \Lambda$ ,  $\varphi\beta = 0$ , since  $\varphi\beta < \varphi 1 = 0$ . Thus,  $\varphi = 0$ .

 $\Box$ 

**Definition 3.11.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on Λ. If  $\Phi(\beta \odot \theta, \delta, \eta) = \Phi(\beta, \delta, \eta) \odot$  $\Phi(\theta, \delta, \eta)$ , for all  $\beta, \theta, \delta, \eta \in \Lambda$ , then  $\Phi$  is called a tri-multiplicative  $(\odot, \vee)$ -permuting tri-derivation on Λ.

**Theorem 3.12.** Suppose that  $\Phi$  is a tri-multiplicative  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$  and  $\varphi$  is the trace of  $\Phi$ . Hence,  $\varphi(B(\Lambda)) \subseteq B(\Lambda)$ .

Proof*.*

Let  $\beta \in \varphi(B(\Lambda))$ . Then,  $\beta = \varphi(\theta)$ , for some  $\theta \in B(\Lambda)$ . Thus,  $\beta \odot \beta = \varphi \theta \odot \varphi \theta = \Phi(\theta \odot \theta, \theta, \theta) = \varphi \theta = \beta$ . Therefore,  $\beta \in B(\Lambda)$ . Hence,  $\varphi(B(\Lambda)) \subseteq B(\Lambda)$ .  $\square$ 

**Theorem 3.13.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on  $\Lambda$ . If there exists a  $\beta \in \Lambda$  such that  $\beta \odot \Phi(\theta, \delta, \eta) = 1$ , for all  $\theta, \delta, \eta \in \Lambda$ , then  $\beta = 1$ .

Proof*.*

Assume that there exists a  $\beta \in \Lambda$  such that  $\Phi(\theta, \delta, \eta) \odot \beta = 1$ , for all  $\theta, \delta, \eta \in \Lambda$ . Since  $\Phi$  is a (⊙*,* ∨)-permuting tri-derivation on Λ,

$$
1 = \Phi(\theta \odot \beta, \delta, \eta) \odot \beta
$$
  
= (( $\Phi(\theta, \delta, \eta) \odot \beta$ )  $\vee$  ( $\theta \odot \Phi(\beta, \delta, \eta)$ )  $\odot \beta$   
= (1  $\vee$  ( $\theta \odot \Phi(\beta, \delta, \eta)$ )  $\odot \beta$   
= 1  $\odot \beta$   
=  $\beta$ 

 $\Box$ 

**Theorem 3.14.** Let  $\Phi$  be a  $(\odot, \vee)$ -permuting tri-derivation on  $B(\Lambda)$ . Thus,

*i.* Φ is a permuting tri-derivation on a lattice, i.e.,

$$
\Phi(\beta \wedge \eta, \theta, \delta) = (\Phi(\beta, \theta, \delta) \wedge \eta) \vee (\beta \wedge \Phi(\eta, \theta, \delta)), \text{ for all } \beta, \theta, \delta, \eta \in B(\Lambda)
$$

*ii.*  $\Phi(\beta, \theta, \delta) = \Phi(\beta, \theta, \delta) \wedge \beta$ , for all  $\beta, \theta, \delta \in B(\Lambda)$ 

Proof*.*

*i.* For all  $\beta, \theta, \delta, \eta \in B(\Lambda)$ ,

$$
\Phi(\beta \wedge \eta, \theta, \delta) = \Phi(\beta \odot \eta, \theta, \delta)
$$
  
=  $(\Phi(\beta, \theta, \delta) \odot \eta) \vee (\beta \odot \Phi(\eta, \theta, \delta))$   
=  $(\Phi(\beta, \theta, \delta) \wedge \eta) \vee (\beta \wedge \Phi(\eta, \theta, \delta))$ 

*ii.* For all  $\beta, \theta, \delta \in B(\Lambda)$ ,

$$
\Phi(\beta, \theta, \delta) = \Phi(\beta \odot \beta, \theta, \delta)
$$
  
=  $(\Phi(\beta, \theta, \delta) \odot \beta) \vee (\beta \odot \Phi(\beta, \theta, \delta))$   
=  $\Phi(\beta, \theta, \delta) \odot \beta$   
=  $\Phi(\beta, \theta, \delta) \wedge \beta$ 

 $\Box$ 

**Theorem 3.15.** Suppose that  $\Phi$  is a  $(\odot, \vee)$ -permuting tri-derivation on Gödel BL-algebra  $\Lambda$ ,  $\varphi$  is the trace of  $\Phi$ , and  $\beta$ ,  $\theta \in \Lambda$ . Then, the followin conditions are valid:

$$
i. \varphi \beta \le \beta
$$

\n- ii. If 
$$
\beta \leq \Phi(\beta, \beta, 1)
$$
, then  $\varphi\beta = \beta$
\n- iii. If  $\beta \geq \Phi(\beta, \beta, 1)$ , then  $\Phi(\beta, \beta, 1) \leq \varphi\beta$
\n- iv. If  $\beta \leq \theta$ , then  $\varphi\beta = \beta$  or  $\varphi\beta \geq \Phi(\beta, \beta, \theta)$
\n
\nPROOF.

*i.*

$$
\varphi \beta = \Phi(\beta, \beta, \beta) = \Phi(\beta \odot \beta, \beta, \beta)
$$
  
=  $(\varphi \beta \odot \beta) \vee (\beta \odot \varphi \beta) = \varphi \beta \odot \beta$   
=  $\min{\{\varphi \beta, \beta\}}$ 

Thus,  $\varphi \beta \leq \beta$ .

*ii.* Let  $\beta \leq \Phi(\beta, \beta, 1)$ , for  $\beta \in \Lambda$ . Thus,

$$
\varphi \beta = \Phi(\beta, \beta, \beta)
$$
  
=  $\Phi(\beta \odot 1, \beta, \beta)$   
=  $(\varphi \beta \odot 1) \vee (\beta \odot \Phi(1, \beta, \beta))$   
=  $\varphi \beta \vee \min{\{\beta, \Phi(1, \beta, \beta)\}}$   
=  $\varphi \beta \vee \beta$   
=  $\beta$ , by *i*.

*iii.* Let  $\beta > \Phi(\beta, \beta, 1)$ , for  $\beta \in \Lambda$ . Hence,

$$
\varphi \beta = \Phi(\beta, \beta, \beta)
$$
  
=  $\Phi(\beta \odot 1, \beta, \beta)$   
=  $\varphi \beta \vee \min{\{\beta, \Phi(1, \beta, \beta)\}}$   
=  $\varphi \beta \vee \Phi(1, \beta, \beta)$ 

Thus,  $\Phi(\beta, \beta, 1) \leq \varphi \beta$ .

*iv.* Let  $\beta \leq \theta$ . From *i*,  $\varphi\beta \leq \beta \leq \theta$ . Thus,  $\varphi\beta \leq \theta$ . Thereby,

$$
\varphi \beta = \Phi(\beta, \beta, \beta)
$$
  
=  $\Phi(\beta \odot \theta, \beta, \beta)$   
=  $(\varphi \beta \odot \theta) \vee (\beta \odot \Phi(\theta, \beta, \beta))$   
=  $\varphi \beta \vee (\beta \odot \Phi(\theta, \beta, \beta))$ 

If  $\beta \leq \Phi(\theta, \beta, \beta)$ , then  $\varphi\beta = \beta$  by *i*. If  $\beta \geq \Phi(\theta, \beta, \beta)$ , then  $\varphi\beta = \varphi\beta \vee \Phi(\theta, \beta, \beta)$ . Hence,  $\varphi\beta \geq$ Φ(*β, β, θ*).

 $\Box$ 

#### **4. Conclusion**

BL-algebras were introduced by Hajek [6] to investigate many-valued logic. One of the reasons for his motivations for introducing BL-algebras was providing an algebraic counterpart of propositional logic, called Basic Logic (BL-logic), which embodies a fragment common to some of the most important many-valued logics, namely Lukasiewicz Logic, Gödel Logic, and Product Logic. This BL-logic is proposed as "the most general" many-valued logic with truth values in [0*,* 1]. Another reason to work was to provide an algebraic mean for the study of continuous t-norms on [0*,* 1]. In this work, we introduce one type of permuting tri-derivation on BL-algebras and investigate its some properties. Moreover, we propose many of the basic properties that the trace of (⊙*,* ∨)-permuting tri-derivation provides. In the future, different permuting tri-derivations can be defined in these algebras, and generalized permuting tri-derivations can be studied in BL-algebras.

#### **Author Contributions**

The author read and approved the final version of the paper.

#### **Conflicts of Interest**

The author declares no conflict of interest.

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# **New Exact Solutions of the Drinfeld-Sokolov System by the Generalized Unified Method**

## Tuğba Aydemir<sup>1</sup> <sup>D</sup>

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doi[:10.53570/jnt.1294322](https://doi.org/10.53570/jnt.1294322) in more general forms. In addition, the GUM has great potential to apply to a wide range **Abstract** − In this study, we apply the generalized unified method (GUM), an enhanced version of the unified method, to find novel exact solutions of the Drinfeld-Sokolov System (DSS) that models the dispersive water waves in fluid dynamics. Moreover, 3D and 2D graphs of some of the obtained exact solutions are plotted to present how various characteristic forms they have. The results show that the presented method simplifies the computation process on the computer in a highly reliable and straightforward manner while providing the solutions of problems, including nonlinear partial differential equations (NPDEs) and fractional partial differential equations (FPDEs) for finding exact solutions.

**Keywords** *Drinfeld-Sokolov system, unified method, generalized unified method, exact solution, traveling wave solutions* **Mathematics Subject Classification (2020)** 35B08, 76B15

## **1. Introduction**

Nonlinear partial differential equations (NPDEs), usually represented by an equation that describes a relationship between an unknown function and its partial derivatives, are great tools for modeling these real-world phenomena in many scientific fields. Finding solutions for NPDEs is essential because it improves our understanding of the physical phenomena these equations model, accelerates technological advances, and optimizes the production process of these technological tools. Therefore, finding numerical and exact solutions to NPDEs corresponding to the physical problems in science and engineering becomes increasingly important.

In this study, the exact solutions of the nonlinear Drinfeld-Sokolov System (DSS), considered to model for dispersive water in fluid dynamics, are investigated by using the generalized unified method (GUM). The explicit mathematical form of the system is given as follows:

$$
\begin{cases}\n u_t(x,t) + (v^2(x,t))_x = 0 \\
 v_t(x,t) - \alpha v_{xxx}(x,t) + 3\beta u_x(x,t)v(x,t) + 3\gamma u(x,t)v_x(x,t) = 0\n\end{cases}
$$

where  $\alpha$ ,  $\beta$ , and  $\gamma$  are real constants. This system was first proposed by Drinfeld and Sokolov [1] as an extension of the Korteweg-de Vries (KdV), which possesses Lax pairs of a special form for affine Lie algebras. The solutions of the DSS exhibit different characteristics, such as static solitons

<sup>1</sup> tgb.aydemir@gmail.com (Corresponding Author)

<sup>&</sup>lt;sup>1</sup>Department of International Trade and Finance, Faculty of Economics and Administrative Sciences, Yalova University, Yalova, Türkiye

that interact with moving solitons without deformations [2]. Morris and Kara [3] showed that the interaction between the underlying symmetries and conservation laws of a partial differential equation system results in double reductions of a class of Drinfeld–Sokolov–Wilson equations.

Due to its significant applications in fluid dynamics and optical fiber, many researchers have developed several methods to obtain exact and numerical solutions for the Drinfeld-Sokolov family. Wazwaz [4] used the sine–cosine and tanh methods to have exact solutions for the Drinfeld-Sokolov System. Arora and Kumar [5] applied the homotopy analysis method to obtain an approximation of the analytic solution for the coupled Drinfeld's–Sokolov–Wilson. He et al. [6] obtained exact solutions for the classic Drinfeld-Sokolov-Wilson equation using the F-expansion method combined with the exp-function method. G´omez [7] studied the generalized Drinfeld-Sokolov-Wilson equation to obtain exact solutions by applying the improved tanh-coth method. Düşünceli [8] employed the improved Bernoulli sub-equation function method (IBSEFM) to have exact solutions of the Drinfeld-Sokolov equation. Zhang [9] solved the Drinfeld-Sokolov-Wilson equation using a variational approach. Günay et al. [10] derived solitary wave solutions to the DSS using the generalized exponential rational function method. Salim et al. [11] solved the Drinfeld-Sokolov-Wilson system numerically using the modified Adomian decomposition method. Alam et al. [12] established some exact solutions of the Drinfeld-Sokolov-Wilson equation with  $S(\xi)$ -expansion method.

Recently, fractional partial differential equations (FPDEs) have also been studied by many authors because fractional order derivatives provide a powerful and enhanced model for expressing some realworld phenomena. In particular, these fractional equations are encountered in modeling considerable complex problems in networks, optics, and fluid dynamics. Jaradat et al. [13] investigated the analytical solution of the time-fractional Drinfeld–Sokolov–Wilson system through the residual power series method. Bhatter et al. [14] considered the fractionalized homotopy analysis transform method to solve the fractional Drinfeld–Sokolov–Wilson model numerically. Gao et al. [15] found the solutions of fractional Drinfeld–Sokolov–Wilson equation using amalgamations of Laplace transform technique with q-homotopy analysis scheme, called q-homotopy analysis transform method (q-HATM) with Atangana-Baleanu derivative. Taşbozan et al. [16] applied the Sine-Gordon expansion method to obtain exact solutions and the perturbation-iteration algorithm to obtain approximate solutions for the fractional Drinfeld-Sokolov-Wilson system. Wang and Wang [17] found the numerical solution for the time-space fractional nonlinear Drinfeld–Sokolov–Wilson system using He's variational method. Noor et al. [18] used the homotopy perturbation transform method and Sumudu transform decomposition method to solve approximately the time-fractional Drinfeld–Sokolov–Wilson system.

In this paper, we have applied the GUM [19] to derive exact solutions for the DSS. The GUM is an enhanced version of the unified method [20,21] that many authors [22–36] have successfully applied to solve different types of NPDEs. Compared with other methods, the GUM requires less computational work with high reliability to solve NPDEs.

This paper is structured as follows: After describing briefly the GUM in section 2, the exact solutions of the DSS are obtained by the GUM, and graphical illustrations of some selected solutions are displayed in section 3. In section 4, the derivation of the exact solution sets obtained by some methods from the solutions obtained by the GUM is discussed. Lastly, conclusive remarks are given in section 5.

## **2. The Generalized Unified Method**

In this section, a brief explanation of the GUM for constructing wave solutions of NPDEs is provided. In the first step, when applying the GUM, the NPDE is converted to a nonlinear ordinary differential equation (NODE) by using wave transformation  $u(x, t) = U(\eta)$  such that  $\eta = x - \nu t + \eta_0$ . The solution form of the obtained NODE is expressed by an ansatz statement as follows:

$$
U(\eta) = a_0 + \sum_{m=1}^{M} \left[ a_m \phi^m + b_m \phi^{-m} \right]
$$
 (1)

After finding the balance value *M* in the ansatz statement, this statement and its derivatives are substituted into NODE to obtain an algebraic polynomial system. Here,  $\phi(\eta)$  is considered Riccati differential equation in the form  $\phi'(\eta) = \phi^2(\eta) - \mu^2$  with  $\phi' = \frac{d\phi}{d\eta}$  and  $\mu = (c + id)$  where *c* and *d* are parameters. The polynomial system with powers  $\phi$  provides to get the values of the coefficients in Equality 1. Finally, wave solutions of the NPDE are obtained in closed form with free parameters *A*, *B*, and *C*, using these coefficients and the solutions of the Riccati differential equation as given by:

$$
\phi(\eta) = \begin{cases}\n\phi_1 = \frac{\mp (c+id)\sqrt{A^2 + (B+iC)^2} - A(c+id)\cosh(2(c+id)(\eta + \eta_0))}{(B+iC) + A\sinh(2(c+id)(\eta + \eta_0))} \\
\phi_2 = \frac{\mp (c+id)(-A + e^{\mp 2(c+id)(\eta + \eta_0)})}{(A + e^{\mp 2(c+id)(\eta + \eta_0)})} \\
\phi_3 = -\frac{1}{\eta + \eta_0}\n\end{cases}
$$

where  $A \neq 0$ , *B*, and *C* are real arbitrary parameters.

## **3. Application of The Generalized Unified Method to Drinfeld-Sokolov System**

In this section, the GUM is applied to the DSS to obtain exact wave solutions by following the steps as explained in Section 2. The DSS is given by

$$
\begin{cases}\n u_t(x,t) + (v^2(x,t))_x = 0 \\
 v_t(x,t) - \alpha v_{xxx}(x,t) + 3\beta u_x(x,t)v(x,t) + 3\gamma u(x,t)v_x(x,t) = 0\n\end{cases}
$$
\n(2)

where  $\alpha$ ,  $\beta$ , and  $\gamma$  are real constants. Converting NPDE System 2 to ordinary differential equations (ODEs) by using the wave transformation  $\eta = x - \nu t + \eta_0$  in  $u(x, t)$  and  $v(x, t)$  gives:

$$
\begin{cases}\n-\nu U'(\eta) + 2V(\eta)V'(\eta) = 0 \\
-\nu V'(\eta) - \alpha V'''(\eta) + 3\beta U'(\eta)V(\eta) + 3\gamma U(\eta)V'(\eta) = 0\n\end{cases}
$$
\n(3)

Here,  $U(\eta)$  and  $V(\eta)$  denote the shape of the nonlinear wave with the wave variable  $\eta = x - \nu t + \eta_0$ . Integrating the first equation with respect to *η* yields the equality  $U(\eta) = \frac{V^2(\eta)}{V^2(\eta)}$  $\frac{\gamma}{\nu}$  in ODE System 3. Substituting this equality into the second equation in ODE System 3, this second one can be rewritten in the form as follows:

$$
\nu V(\eta) + \alpha V''(\eta) - \frac{6\beta + 3\gamma}{\nu} V^3(\eta) = 0
$$
\n(4)

The solutions of the DSS in Equation 2 can be written in the form with balance value  $M = 1$ , found from the highest order  $V''$  and the nonlinear term  $V^3$ . Therefore, the solutions in closed form for Equation 2 are described as follows:

$$
\begin{cases}\nV(\eta) = a_0 + a_1 \phi + \frac{b_1}{\phi} \\
U(\eta) = \frac{V^2(\eta)}{\nu}\n\end{cases}
$$
\n(5)

where  $a_0, a_1$ , and  $b_1$  are coefficients of  $\phi$ , determined later. Substituting the ansatz statement in Solution 5 and its derivatives into Equation 4 gives an algebraic polynomial equation system with  $a_0, a_1, b_1$ , and  $\mu$ . Solving this polynomial equation system by using any symbolic computation program yields the following sets for the unknowns  $a_0, a_1, b_1$ , and  $\mu$ . Substituting these values into Solution 5

with  $\phi_1$  and  $\phi_2$  gives rise to the following exact solutions to the DSS, respectively.

**Case 1.** For  $a_0 = 0$ ,  $a_1 = \pm \sqrt{\frac{2\alpha}{K}}$  $\frac{2\alpha}{K}$ ,  $b_1 = 0$ , and  $\mu = \pm \sqrt{\frac{\nu}{2\alpha}}$  such that  $K = \frac{6\beta + 3\gamma}{\nu}$  $\frac{+3\gamma}{\nu}$ , the exact solutions are as follows:  $\sqrt{ }$  $\int u_1(x,t) = \frac{v_1^2(x,t)}{v}$ *ν*

$$
v_1(x,t) = \pm \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}} \left(\sqrt{A^2 + (B + iC)^2} \mp A \cosh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)\right)}{(B + iC) \mp A \sinh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)}
$$

and

$$
\begin{cases}\nu_2(x,t) = \frac{v_2^2(x,t)}{\nu} \\
v_2(x,y,t) = \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}} \left(-A + e^{\mp\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)}\right)}{\left(A + e^{\mp\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)}\right)}\n\end{cases}
$$



**Figure 1.** 3D and 2D graphs of the solution  $v_1$  for the parameter  $A = 1$ ,  $B = 0$ , and  $C = 0$ 

**Case 2.** For  $a_0 = 0$ ,  $a_1 = 0$ ,  $b_1 = \pm \nu \sqrt{\frac{1}{2K\alpha}}$ , and  $\mu = \pm \sqrt{\frac{\nu}{2\alpha}}$  such that  $K = \frac{6\beta + 3\gamma}{\nu}$  $\frac{+3\gamma}{\nu}$ , the exact solutions are as follows:  $\int u_3(x,t) = \frac{v_3^2(x,t)}{v}$ 

$$
u_3(x,t) = \frac{1}{\nu}
$$

$$
v_3(x,t) = \mp \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}}(B+iC)\mp A\sinh\left(\sqrt{\frac{2\nu}{\alpha}}(x-\nu t+\eta_0)\right)}{\left(\sqrt{A^2 + (B+iC)^2}\mp A\cosh\left(\sqrt{\frac{2\nu}{\alpha}}(x-\nu t+\eta_0)\right)\right)}
$$

and

$$
\begin{cases}\nu_4(x,t) = \frac{v_4^2(x,t)}{\nu} \\
v_4(x,y,t) = \frac{\sqrt{\frac{\nu^2}{6\beta+3\gamma}} \left(A + e^{\mp\sqrt{\frac{2c}{\alpha}}(x-\nu t + \eta_0)}\right)}{\left(-A + e^{\mp\sqrt{\frac{2\nu}{\alpha}}(x-\nu t + \eta_0)}\right)}\n\end{cases}
$$



**Figure 2.** 3D and 2D graphs of real part for the solution  $v_3$  for the parameter  $A = 1$ ,  $B = 1$ , and  $C=1$ 

**Case 3.** For  $a_0 = 0$ ,  $a_1 = \frac{\nu}{2K}$  $\frac{\nu}{2Kb_1}$ ,  $b_1 = \pm \nu \sqrt{\frac{1}{8K\alpha}}$ , and  $\mu = \pm \frac{b_1}{\nu}$ √  $\sqrt{-2\nu K}$  such that  $K = \frac{6\beta + 3\gamma}{\nu}$  $\frac{+3\gamma}{\nu}$ , the exact solutions are as follows:

$$
\begin{cases}\nu_5(x,t) = \frac{v_5^2(x,t)}{\nu} \\
v_5(x,t) = \pm \sqrt{\frac{\nu^2}{-12\beta - 6\gamma}} \left( \frac{(B+iC)\mp A \sinh\left(\sqrt{\frac{-\nu}{\alpha}}(x-\nu t + \eta_0)\right)}{\left(\sqrt{A^2 + (B+iC)^2} \mp A \cosh\left(\sqrt{\frac{-\nu}{\alpha}}(x-\nu t + \eta_0)\right)\right)} + \frac{\left(\sqrt{A^2 + (B+iC)^2} \mp A \cosh\left(\sqrt{\frac{-\nu}{\alpha}}(x-\nu t + \eta_0)\right)\right)}{(B+iC)\mp A \sinh\left(\sqrt{\frac{-\nu}{\alpha}}(x-\nu t + \eta_0)\right)} \right)\n\end{cases}
$$

and

$$
\begin{cases}\nu_6(x,t) = \frac{v_6^2(x,t)}{\nu} \\
v_6(x,t) = \pm \sqrt{\frac{\nu^2}{-12\beta - 6\gamma}} \left( \frac{A + e^{\mp \sqrt{\frac{-\nu}{\alpha}}(x - \nu t + \eta_0)}}{-A + e^{\mp \sqrt{\frac{-\nu}{\alpha}}(x - \nu t + \eta_0)}} + \frac{-A + e^{\mp \sqrt{\frac{-\nu}{\alpha}}(x - \nu t + \eta_0)}}{A + e^{\mp \sqrt{\frac{-\nu}{\alpha}}(x - \nu t + \eta_0)}} \right)\n\end{cases}
$$



**Figure 3.** 3D and 2D graphs of the real part for the solution  $v_5$  for the parameter choices  $A = 1$ ,  $B = 1$ , and  $C = 0$ 



**Figure 4.** 3D and 2D graphs of the imaginary part for the solution  $v<sub>5</sub>$  for the parameter choices  $A = 1$ ,  $B = 1$ , and  $C = 0$ 

**Case 4.** For  $a_0 = 0$ ,  $a_1 = \frac{\nu}{4K}$  $\frac{\nu}{4Kb_1}$ ,  $b_1 = \pm \nu \sqrt{\frac{1}{32K\alpha}}$ , and  $\mu = \pm \frac{2b_1}{\nu}$ √  $\overline{\nu K}$  such that  $K = \frac{6\beta + 3\gamma}{\nu}$  $\frac{+3\gamma}{\nu}$ , the exact solutions are as follows:

$$
\begin{cases}\nu_{7}(x,t) = \frac{v_{7}^{2}(x,t)}{\nu} \\
v_{7}(x,t) = \mp \frac{\nu}{2\sqrt{6\beta+3\gamma}} \left( \frac{(B+iC)\mp A\sinh(\sqrt{\frac{\nu}{2\alpha}}(x-\nu t+\eta_{0}))}{(\sqrt{A^{2}+(B+iC)^{2}}\mp A\cosh(\sqrt{\frac{\nu}{2\alpha}}(x-\nu t+\eta_{0}))})} + \frac{(\sqrt{A^{2}+(B+iC)^{2}}\mp A\cosh(\sqrt{\frac{\nu}{2\alpha}}(x-\nu t+\eta_{0}))}{(B+iC)\mp A\sinh(\sqrt{\frac{\nu}{2\alpha}}(x-\nu t+\eta_{0}))} \right)\n\end{cases}
$$

and

$$
\begin{cases} u_8(x,t) = \frac{v_8^2(x,t)}{\nu} \\ v_8(x,y,t) = \mp \frac{\nu}{2\sqrt{6\beta+3\gamma}} \left( \frac{A + e^{\mp\sqrt{\frac{\nu}{2\alpha}}(x-\nu t + \eta_0)}}{-A + e^{\mp\sqrt{\frac{\nu}{2\alpha}}(x-\nu t + \eta_0)}} + \frac{-A + e^{\mp\sqrt{\frac{\nu}{2\alpha}}(x-\nu t + \eta_0)}}{A + e^{\mp\sqrt{\frac{\nu}{2\alpha}}(x-\nu t + \eta_0)}} \right) \end{cases}
$$



**Figure 5.** 3D and 2D graphs of imaginary part for the solution  $v_7$  for the parameter choices  $A = 1$ ,  $B = 1$ , and  $C = 0$ 

The 3D and 2D graphs of the solutions  $v_1$ ,  $v_3$ ,  $v_5$ , and  $v_7$  are shown as in Figures 1-5 for the parameters  $\nu = 2, \ \alpha = 1, \ \beta = 1, \text{ and } \ \gamma = 1 \text{ on the intervals } -10 < x < 10 \text{ and } -10 < t < 10. \text{ From these figures, }$ it is observed that the solutions obtained by the GUM display a wide variety of characteristics, represented by hyperbolic and trigonometric functions. 2D graphs of the solutions are also plotted for  $x = 0$  to follow properly the characteristics of the solutions. The solutions in different shapes can be obtained by using free parameters for the corresponding physical problems.

#### **4. Results and Discussion**

 $\lambda$ 

We have obtained more general forms of exact solutions for the Drinfeld-Sokolov System by using the GUM in Section 3. In this section, we show more concretely how to derive solutions from some other methods from these solutions, considering some of the above solution sets.

The first group of hyperbolic and trigonometric solutions of the unified method can be obtained by taking  $B = 0$  and  $C = 0$ , respectively, depending on the wave velocity in  $\{u_1, v_1\}$  as follows:

$$
\begin{cases}\nu_{1,1}(x,t) = \frac{v_1^2(x,t)}{\nu} \\
v_{1,1}(x,t) = \mp \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}} \left(\sqrt{A^2 + B^2} \mp A \cosh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)\right)}{B \mp A \sinh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)}\n\end{cases}
$$

 $22$ 

and

$$
\begin{cases}\nu_{1,2}(x,t) = \frac{v_1^2(x,t)}{\nu} \\
v_{1,2}(x,t) = \pm \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}} \left(\sqrt{A^2 - C^2} \mp A \cosh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)\right)}{iC \mp A \sinh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)}\n\end{cases}
$$

In addition, the solution sets obtained by the tanh method can be derived by taking  $B = 0$  and  $C = 0$ with the identities  $\cosh(2x) = 2\cosh^2(x) - 1 = 2\sinh^2(x) + 1$  and  $\sinh(2x) = 2\sinh(x)\cosh(x)$  as follows:

$$
\begin{cases}\nu_{1,3}(x,t) = \frac{v_1^2(x,t)}{\nu} \\
v_{1,3}(x,t) = \pm \frac{\sqrt{\frac{\nu^2}{6\beta + 3\gamma}} \left(A \mp A \cosh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)\right)}{A \sinh\left(\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)\right)}\n\end{cases}
$$

Similarly, we can also obtain the second group of hyperbolic and trigonometric solutions of the unified method depending on the wave velocity after some simplification in  $\{u_2, v_2\}$ .

$$
\begin{cases}\nu_{21}(x,t) = \frac{v_2^2(x,t)}{\nu} \\
v_{21}(x,y,t) = \sqrt{\frac{\nu^2}{6\beta + 3\gamma}} - \frac{2A\sqrt{\frac{\nu^2}{6\beta + 3\gamma}}}{\left(A + e^{\mp\sqrt{\frac{2\nu}{\alpha}}(x - \nu t + \eta_0)}\right)}\n\end{cases}
$$

In the studies [20, 21], it is explained that the unified method gives many more solutions than the family of the tanh method, the  $\left(\frac{G'}{G}\right)$  $\left(\frac{G'}{G}\right)$ -expansion method, and the extended homogeneous balance method. Therefore, it can be concluded that the GUM gives more general solution sets than these methods.

The GUM, an enhanced version of the unified method, provides a more general solution structure for the models in mathematical physics. Considering hyperbolic-trigonometric identities and the value of  $\mu$  in the case of a real number or pure imaginary number, the solutions obtained by the unified method can be derived by GUM as above while setting  $B = 0$  and  $C = 0$ , respectively. Further, when  $\mu$  is a complex number, solutions combining trigonometric-hyperbolic solution sets can also be obtained. Moreover, performing the GUM as a new expansion method on a computer contributes to a very simple, straightforward, effective, and accurate way to solve a wide range of NPDEs. Since a more general closed solution form in a compact way is obtained by the GUM with free parameters, we applied this method to solve the Drinfeld-Sokolov System.

## **5. Conclusion**

In this study, the GUM has been successfully applied to obtain solutions for the DSS. The main contributions of this study are:

*i.* More general forms of exact solutions are obtained for DSS using the GUM. Some of the obtained solutions have been visualized by plotting graphs to show how diverse characteristics the solutions have.

*ii.* These obtained solutions can be more functional in explaining the physical characteristics of various models arising in science and engineering.

*iii.* We have shown that the solution sets obtained by some other exact solution methods, including the unified method, can be derived from the GUM.

*iv.* It is obvious that the GUM is quite simple to perform on any symbolic computation program and gives reliable and straightforward results to find exact solutions for NPDEs.

*v.* Moreover, the results show that the GUM can be applied in future studies for NPDEs and fractional NPDEs.

## **Author Contributions**

The author read and approved the final version of the paper.

## **Conflicts of Interest**

The author declares no conflict of interest.

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# New Theory

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# **An Extension of the UEHL Distribution Based on the DUS Transformation**

**Murat Genç<sup>1</sup> , Ömer Özbilen<sup>2</sup>**

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doi[:10.53570/jnt.1317652](https://doi.org/10.53570/jnt.1317652) investigate the behavior of the maximum likelihood estimates in various conditions. Furthermore, we **Abstract** − In this study, we propose a new distribution based on the Dinesh, Umesh, and Sanjay (DUS) transformation by using the Unit Exponentiated Half-Logistic (UEHL) distribution as the baseline distribution, a member of the family of proportional hazard rate models. Moreover, we study several properties, such as moments, skewness, kurtosis, stress-strength reliability, and likelihood ratio ordering. Further, we discuss the statistical inference on the parameters of the proposed distribution by the maximum likelihood estimation (MLE) method. Besides, we conduct a simulation based on the new distribution to present a numerical example to show the performance of the distribution on a real-life data set. Finally, we discuss the need for further research.

**Keywords** *DUS transformation, UEHL distribution, maximum likelihood estimator, data analysis*

**Mathematics Subject Classification (2020)** 62E10, 60E05

## **1. Introduction**

A variety of models exist in the literature for examining lifetime data. The Weibull distribution originated as an extension of the exponential distribution. The distribution is more flexible in shape than the exponential distribution. This property makes the Weibull distribution useful for lifetime modeling and modeling data from economics and business administration [1-3]. In reliability theory, the Weibull distribution is convenient for modeling monotonic hazard rates. However, it does not provide a reasonable parametric fit for non-monotonic failure rates [4, 5]. To improve the ability to model non-monotonic failure rates, several generalizations and modifications of the Weibull distribution have been introduced in the literature [6, 7].

As an alternative to the Weibull distribution, the omega distribution was introduced as a new probability distribution, and its applications in reliability theory were discussed [8]. Dombi et al. [8] showed that the asymptotic omega hazard rate function is the Weibull hazard rate function, which implies that the asymptotic omega distribution is just the Weibull distribution. Moreover, the omega distribution belongs to the class of proportional hazard rate models. Özbilen and Genç [9] introduced the unit exponentiated-half logistic (UEHL) distribution corresponding to the omega distribution with a special case of the parameter value of the distribution. The distribution corresponds to the exponentiated half-logistic distribution with a simple transformation, which has many applications in reliability theory [10, 11]. **2DEVALUATION 2DEATY** Amount of Section and Mersine in the DUS Transfer of Geomatic Engineering, Faculty of Engineering, Faculty (Section 2013) and ENGINEERING CONTENT Mersing Control and Control and American Control an

In statistical literature, there are various ways to generate new distributions using some baseline distributions [12, 13]. In this context, Kumar et al. [14] proposed the Dinesh, Umesh, and Sanjay (DUS) transformation to

<sup>1</sup>muratgenc@tarsus.edu.tr (Corresponding Author); <sup>2</sup>ozbilen@mersin.edu.tr

<sup>1</sup>Department of Management Information Systems, Faculty of Economics and Administrative Sciences, Tarsus University, Mersin, Türkiye

obtain a new distribution family. Let  $f(x)$  and  $F(x)$  denote the probability density function (PDF) and cumulative distribution function (CDF) of the baseline distribution, respectively. Then, the PDF and CDF of the DUS family are given by

$$
f_{DUS}(x) = \frac{f(x)e^{F(x)}}{e-1}, \quad x \in D
$$
 (1)

and

$$
F_{DUS}(x) = \frac{e^{F(x)} - 1}{e - 1}
$$
 (2)

respectively. In Equation [\(1\),](#page--1-0)  $\hat{D}$  is the domain of the baseline distribution. As observed from Equation[s \(1\)](#page--1-0) and [\(2\)](#page--1-1) that the DUS family does not contain any new parameters in addition to those used in the baseline distribution. Hence, the DUS family results in a parsimonious distribution in terms of computation and interpretation.

There exist several studies in the literature where the DUS transformation is used. Kumar et al. [14] applied the DUS transformation to an exponential CDF. Deepthi and Chacko [15] proposed the DUS-Lomax distribution, considering the Lomax distribution as the baseline distribution for the DUS transformation. Kavya and Manoharan [16] obtained the DUS-Weibull distribution with the same approach. Maurya et al. [17] proposed a generalization of the DUS transformation and studied the exponential baseline in the proposed method. Karakaya et al. [18] considered the DUS transformation with the Kumaraswamy distribution and investigated the properties of the distribution.

The performance of the UEHL distribution can be improved by applying the DUS transformation to the CDF of the UEHL distribution. In this paper, we use the DUS transformation to the UEHL distribution to enhance distribution in the scope of modeling proportional data.

The rest of the paper is organized as follows: Section 2 introduces the UEHL distribution with DUS transformation and obtains the survival and hazard rate functions. Section 3 discusses some analytical characteristics, such as the moments, quantile function, stress-strength reliability, likelihood ratio ordering, and the maximum likelihood estimation (MLE) of the method. Section 4 performs a simulation study to validate the performance of the maximum likelihood estimates of the distribution. Section 5 provides a numerical example to illustrate the performance of the proposed distribution in modeling the real data. The last section concludes the paper.

#### **2. DUS-UEHL Distribution**

Recently, the UEHL distribution [9] has been proposed as a special case of omega distribution [8], a member of the class of proportional hazard rate models. The distribution corresponds to the transformation  $Z = e^{-X}$ where  $X$  has the exponentiated half-logistic distribution. The UEHL distribution has many applications in reliability theory through exponentiated half-logistic distribution [10, 19, 20]. This section applies the DUS transformation to the UEHL distribution to introduce a new class of distributions. The PDF and CDF of the UEHL distribution are given as follows, respectively:

$$
f_{UEHL}(x) = 2\lambda \theta x^{\theta - 1} \frac{\left(1 - x^{\theta}\right)^{\lambda - 1}}{\left(1 + x^{\theta}\right)^{\lambda + 1}}, \quad 0 < x < 1 \tag{3}
$$

and

$$
F_{UEHL}(x) = 1 - \left(\frac{1 - x^{\theta}}{1 + x^{\theta}}\right)^{\lambda}, \quad 0 < x < 1 \tag{4}
$$

where  $\theta > 0$  and  $\lambda > 0$  are the scale and shape parameters of the distribution, respectively. By applying the transformations given in Equations [\(1\)](#page--1-0) and [\(2\)](#page--1-1) to Equations [\(3\)](#page--1-2) and [\(4\),](#page--1-3) the PDF and CDF of the new distribution are obtained as follows, respectively:

$$
f_{D-UEHL}(x) = \frac{1}{e-1} 2\lambda \theta x^{\theta-1} \frac{\left(1-x^{\theta}\right)^{\lambda-1}}{(1+x^{\theta})^{\lambda+1}} e^{1-\left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)^{\lambda}}
$$
(5)

and

$$
F_{D-UEHL}(x) = \frac{1}{e-1} \left( e^{1 - \left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)^{\lambda}} - 1 \right)
$$
(6)

The random variable  $X$  with the CDF in Equation [\(6\)](#page--1-4) is said to follow the two-parameter DUS-UEHL distribution and is denoted by  $D$ -UEHL $(\theta, \lambda)$ . Considering the derivative of PDF in Equation 5, we obtain

$$
f'_{D-UEHL}(x) = \frac{2\lambda\theta x^{\theta - 2}(-2\lambda\theta x^{\theta} + (\theta + 1)x^{2\theta} + \theta - 1)\left(\frac{2}{x^{\theta} + 1} - 1\right)^{\lambda}}{(x^{2\theta} - 1)^2}
$$

which has the critical points  $x_{1,2} = \left(\frac{(\lambda \theta \pm \sqrt{\lambda^2 \theta^2 - \theta^2 + 1})}{(\theta + 1)}\right)$  $\frac{\lambda^{6} \theta^{6} + 1}{(\theta+1)}$  $1/\theta$ and the extreme points 0 and 1. Given  $\theta > 1$ , this indicates that the D-UEHL( $\theta$ ,  $\lambda$ ) distribution is unimodal with the unique mode at  $x_1$  or  $x_2$  depending on the scale and shape parameters of the distribution. Moreover, the PDF can be an increasing, decreasing, or Ushaped function of  $x$  depending on the values of the PDF at the extreme points.

Figure 1 shows the PDF of the D-UEHL $(\theta, \lambda)$  distribution for various values of the parameters  $\theta$  and  $\lambda$ . According to Figure 1, the  $D$ -UEHL $(\theta, \lambda)$  distribution can take different shapes, including unimodal, increasing, decreasing, and U-shaped functions.



Figure 1. PDF plots of the D-UEHL( $\theta$ ,  $\lambda$ ) distribution for various values of the parameters  $\theta$  and  $\lambda$ 

The survival function and the hazard rate functions of the  $D$ -UEHL distribution are provided as follows, respectively:

$$
S_{D-UEHL}(x) = \frac{e - e^{-1 - \left(\frac{1 - x^{\theta}}{1 + x^{\theta}}\right)^{\lambda}}}{e - 1}
$$
 (7)

and

$$
h_{D-UEHL}(x) = \frac{2\lambda\theta x^{\theta-1} \frac{\left(1-x^{\theta}\right)^{\lambda-1}}{\left(1+x^{\theta}\right)^{\lambda+1}} e^{1-\left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)^{\lambda}}}{e-e^{1-\left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)^{\lambda}}}
$$
(8)

Figure 2 manifests the hazard rate function for selected values of the parameters  $\theta$  and  $\lambda$ . According to Figure 2, the hazard function of  $D$ -UEHL( $\theta$ , $\lambda$ ) is an increasing function.



Figure 2. Hazard rate function plots of the  $D$ - $UEHL(\theta, \lambda)$  distribution for various values of the parameters  $\theta$ and  $\lambda$ 

## **3.** Some Analytical Characteristics of the  $D$ - $UEHL(\theta, \lambda)$  Distribution

This section discusses some statistical characteristics, including the moments, quantile function, stress-strength reliability, likelihood ratio ordering, and MLE of the  $D$ -UEHL $(\theta, \lambda)$  distribution.

#### **3.1. Moments**

The moments are useful for understanding the various characteristics of a statistical distribution. In this section, we consider the moments of a  $D$ -UEHL $(\theta, \lambda)$  random variable. Let X be a  $D$ -UEHL $(\theta, \lambda)$  random variable with PDF given by Equation [\(5\).](#page--1-5) Then, for  $r \in \{1,2,3,\dots\}$ , the r-th raw moment of X is

$$
E(X^r) = \frac{2\lambda e}{e-1} \int_0^1 x^r \left( \theta x^{\theta-1} \frac{\left(1-x^{\theta}\right)^{\lambda-1}}{(1+x^{\theta})^{\lambda+1}} \right) e^{1-\left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)^{\lambda}} dx
$$

By substituting  $v = x^{\theta}$ , the expectation is written as

$$
E(X^{r}) = \frac{2\lambda e}{e-1} \int_{0}^{1} v^{r/\theta} (1-v)^{\lambda-1} (1+v)^{-\lambda-1} e^{-\left(\frac{1-v}{1+v}\right)^{\lambda}} dv
$$

By using the Taylor expansion  $e^x = \sum_{i=0}^{\infty} \frac{x^i}{i!}$ j!  $\int_{0}^{\infty}$   $\frac{x^{j}}{i!}$ 

$$
E(X^{r}) = \frac{2\lambda e}{e-1} \sum_{i=0}^{\infty} \left[ \frac{(-1)^{i}}{i!} \int_{0}^{1} (v^{r/\theta}(1-v)^{\lambda(1+i)-1}(1+v)^{-\lambda(1+i)-1}) dv \right]
$$

Using Equation 3.197.3 provided in [21], i.e.,  $\int_0^1 x^{\lambda-1}(1-x)^{\mu-1}(1-\beta x)^{-\nu}dx = B(\lambda, \mu) \int_2 F_1(\nu, \lambda; \lambda + \mu; \beta)$ such that Re  $\lambda > 0$ , Re  $\mu > 0$ , and  $|\beta| < 1$ , the r-th raw moment is obtained as

$$
E(X^{r}) = \frac{2\lambda e}{e-1} \sum_{i=0}^{\infty} \frac{(-1)^{i}}{i!} B\left(1 + \frac{r}{\theta}, \lambda(1+i)\right) {}_{2}F_{1}\left(\lambda(1+i) + 1, 1 + \frac{r}{\theta}; 1 + \frac{r}{\theta} + \lambda(1+i); -1\right)
$$

where

$$
B(a,b) = \int_0^1 t^{a-1} (1-t)^{b-1} dt \text{ and } {}_2F_1(a,b;c;z) = \sum_{n=0}^\infty \frac{(a)_n (b)_n}{(c)_n} \frac{z^n}{n!}
$$

are the beta and Gauss hypergeometric functions, respectively, and  $(a)_n = a(a + 1) \cdots (a + n + 1)$ . Table 1 shows the mean, variance, skewness, and kurtosis of  $D$ - $UEHL(\theta, \lambda)$  distribution for selected pairs of  $(\theta, \lambda)$ . The results indicate that the D-UEHL $(\theta, \lambda)$  has various shapes depending on the values of the distribution parameters.

		Mean	Variance	<b>Skewness</b>	Kurtosis	
0.7	0.7	0.4893	0.1057	0.0478	1.6247	
	1.5	0.2600	0.0557	0.9766	3.0931	
		0.1179	0.0168	1.8289	6.8839	
1.5	0.7	0.6574	0.0725	$-0.5416$	2.1708	
	1.5	0.4688	0.0573	0.1228	2.1052	
	3	0.3153	0.0322	0.5336	2.7771	
3	0.7	0.7874	0.0374	$-1.0363$	3.4153	
	1.5	0.6571	0.0369	$-0.4670$	2.5985	
		0.5355	0.0285	$-0.1805$	2.5863	

Table 1. Mean, variance, skewness, and kurtosis of the  $D$ -UEHL( $\theta$ ,  $\lambda$ ) distribution

#### **3.2. Quantile Function**

The quantile function of the  $D$ -UEHL( $\theta$ ,  $\lambda$ ) distribution is given by

$$
Q_{D-UEHL}(u; \theta, \lambda) = F^{-1}(u; \theta, \lambda) = \left[2\big(1 + \big(1 - \log\big(1 + u(e-1)\big)\big)^{1/\lambda}\big)^{-1} - 1\right]^{1/\theta}
$$
(9)

Based on Equation [\(9\),](#page--1-6) the median of the  $D$ -UEHL( $\theta$ ,  $\lambda$ ) distribution is given by the following function, a function of the parameters  $\theta$  and  $\lambda$ :

$$
Q_{D-UEHL}(0.5; \theta, \lambda) = \left[2\left(1 + (1 - \log 2 - \log (1 + e))\right)^{1/\lambda}\right)^{-1} - 1\right]^{1/\theta}
$$
(10)

The median values of the  $D$ -UEHL $(\theta, \lambda)$  distribution for selected parameter values are provided in Table 2.

Table 2. Median of the D-UEHL( $\theta$ , $\lambda$ ) distribution									
	θ	0.1	0.5		2.5		10		
	0.1	0.99874892	0.05466397	0.00033596	0.00000007	$< 10^{-8}$	$< 10^{-8}$		
	0.5	0.99974966	0.55916532	0.20195643	0.03655548	0.00930982	0.00233835		
λ		0.99987482	0.74777357	0.44939563	0.19119487	0.09648743	0.04835653		
	2.5	0.99994993	0.89024193	0.72619213	0.51592966	0.39245358	0.29770219		
	5	0.99997496	0.94352633	0.85216907	0.71828244	0.62646116	0.54562093		
	10	0.99998748	0.97135283	0.92313004	0.84751545	0.79149299	0.73866158		

 $T_{\rm eff}$  2.  $\Lambda$  and  $T_{\rm eff}$  of the  $D$  distribution  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  and  $T_{\rm eff}$  an

#### **3.3. Stress-Strength Reliability**

Given the stress and strength random variables, Y and X, the stress-strength reliability is defined as  $R =$  $P(Y < X)$ . In this section, we compute the stress-strength reliability for the model D-UEHL( $\theta$ ,  $\lambda$ ).

**Proposition 1.** Let  $Y$  and  $X$  be independent stress and strength random variables following  $D$ -UEHL distribution with parameters  $(\theta, \lambda_1)$  and  $(\theta, \lambda_2)$ , respectively. Then, the stress-strength reliability is

$$
R = \frac{e^2}{(e-1)^2} \sum_{j=0}^{\infty} \frac{(-1)^j}{j!} \gamma \left(\frac{\lambda_1}{\lambda_2} j + 1, 1\right) - \frac{1}{e-1}
$$
 (11)

where  $\lambda_1 > 0$ ,  $\lambda_2 > 0$ , and  $\gamma(a, x) = \int_a^x t^{a-1}$  $\int_a^x t^{a-1} e^{-t} dt$  is the incomplete gamma function.

PROOF.

By the definition, the stress-strength reliability can be written as

$$
R = \int_{0}^{1} P(Y < X \mid X = x) f_X(x) dx
$$
\n
$$
= \int_{0}^{1} \frac{1}{e - 1} \left( e^{1 - \left(\frac{1 - x^{\theta}}{1 + x^{\theta}}\right)^{\lambda_1}} - 1 \right) \left( \frac{1}{e - 1} 2\lambda_2 \theta x^{\theta - 1} \frac{\left(1 - x^{\theta}\right)^{\lambda_2 - 1}}{\left(1 + x^{\theta}\right)^{\lambda_2 + 1}} e^{1 - \left(\frac{1 - x^{\theta}}{1 + x^{\theta}}\right)^{\lambda_2}} \right) dx
$$

By substituting  $t = \left(\frac{1-x^{\theta}}{1+x^{\theta}}\right)$  $\frac{1-x}{1+x^{\theta}}$  $\lambda_1$ and applying the Taylor expansion of  $e^x$ ,

$$
R = \frac{e}{(e-1)^2} \int_0^1 \left( e^{-t^{\lambda_1/\lambda_2} + 1} - 1 \right) e^{-t} dt
$$
  
\n
$$
= \frac{e^2}{(e-1)^2} \int_0^1 \left( \sum_{j=0}^\infty \frac{(-1)^j t^{\lambda_1 j/\lambda_2}}{j!} - \frac{1}{e} \right) e^{-t} dt
$$
  
\n
$$
= \frac{e^2}{(e-1)^2} \sum_{j=0}^\infty \left\{ \frac{(-1)^j}{j!} \left( \int_0^1 t^{\lambda_1 j/\lambda_2} e^{-t} dt \right) \right\} - \frac{1}{e-1}
$$
  
\n
$$
= \frac{e^2}{(e-1)^2} \sum_{j=0}^\infty \left\{ \frac{(-1)^j}{j!} \gamma \left( \frac{\lambda_1}{\lambda_2} j + 1, 1 \right) \right\} - \frac{1}{e-1}
$$

◻

#### **3.4. Likelihood Ratio Ordering**

Let X and Y be random variables with PDFs  $f_X(\cdot)$  and  $f_Y(\cdot)$ , respectively. If  $\frac{f_X(x)}{f_Y(x)}$  is non-decreasing in x, then it is said to be the random variable X is smaller than Y in likelihood ratio ordering and denoted by  $X \leq_{lr} Y$ . Proposition 2 gives the property of likelihood ratio ordering for the  $D$ -UEHL distribution.

**Proposition 2.** Let  $X \sim D$ -UEHL( $\theta_1$ ,  $\lambda_1$ ) and  $Y \sim D$ -UEHL( $\theta_2$ ,  $\lambda_2$ ). If  $\theta_1 = \theta_2$ , then  $X \leq_{lr} Y$ .

PROOF*.*

Let  $\theta_1 = \theta_2 = \theta$ . The ratio of the PDFs of X and Y is given by

$$
g(x) = \frac{2\lambda_1 \theta x^{\theta-1} \frac{\left(1-x^{\theta}\right)^{\lambda_1-1}}{\left(1+x^{\theta}\right)^{\lambda_1+1}}}{2\lambda_2 \theta x^{\theta-1} \frac{\left(1-x^{\theta}\right)^{\lambda_2-1}}{\left(1+x^{\theta}\right)^{\lambda_2+1}}}
$$

After simplifications, we obtain  $g'(x) < 0$ , for  $0 < x < 1$  and  $\theta > 0$ . Hence,  $g(\cdot)$  is a non-decreasing function of  $x$ , which completes the proof.

#### **3.5. Maximum Likelihood Estimation**

Let  $X_1, X_2, \dots, X_n$  be an identically independent distributed sample from  $D$ -UEHL( $\theta$ , $\lambda$ ). Then, the likelihood and log-likelihood functions are written as

$$
L(\theta,\lambda) = \prod_{i=1}^{n} \left( \frac{1}{e-1} 2\lambda \theta x_i^{\theta-1} \frac{\left(1-x_i^{\theta}\right)^{\lambda-1}}{\left(1+x_i^{\theta}\right)^{\lambda+1}} e^{1-\left(\frac{1-x_i^{\theta}}{1+x_i^{\theta}}\right)^{\lambda}} \right)
$$
(12)

and

$$
\ell(\theta, \lambda) = -n \log(e - 1) + n(1 + \log 2) + n \log \lambda + n \log \theta + (\theta - 1) \sum_{i=1}^{n} \log(x_i)
$$
  
+  $(\lambda - 1) \sum_{i=1}^{n} \log(1 - x_i^{\theta}) - (\lambda + 1) \sum_{i=1}^{n} \log(1 + x_i^{\theta}) - \sum_{i=1}^{n} \left(\frac{1 - x_i^{\theta}}{1 + x_i^{\theta}}\right)^{\lambda}$  (13)

respectively. Differentiating the log-likelihood function with respect to the parameters, we get the loglikelihood equations as

$$
\frac{\partial \ell(\theta, \lambda)}{\partial \theta} = \frac{n}{\theta} + \sum_{i=1}^{n} \log(x_i) - (\lambda - 1) \sum_{i=1}^{n} \frac{x_i^{\theta} \log x_i}{1 - x_i^{\theta}} + (\lambda + 1) \sum_{i=1}^{n} \frac{x_i^{\theta} \log x_i}{1 + x_i^{\theta}} + 2\theta \sum_{i=1}^{n} \frac{x_i^{\theta - 1}}{(1 + x_i^{\theta})^2} = 0 \quad (14)
$$

and

$$
\frac{\partial \ell(\theta, \lambda)}{\partial \lambda} = \frac{n}{\lambda} + \sum_{i=1}^{n} \log \left( \frac{1 - x_i^{\theta}}{1 + x_i^{\theta}} \right) - \sum_{i=1}^{n} \left( \frac{1 - x_i^{\theta}}{1 + x_i^{\theta}} \right)^{\lambda} \log \left( \frac{1 - x_i^{\theta}}{1 + x_i^{\theta}} \right) = 0
$$
\n(15)

Equations [\(14\)](#page--1-7) and [\(15\)](#page--1-8) do not have a closed-form solution. To solve these equations, some iterative methods can be employed. We apply the optim procedure in  $R$  to obtain the solution of the equation system.

#### **4. Simulation Study**

In this section, a simulation study is carried out to examine the properties of the MLE, which we discussed in detail in Section 3. Table 3 shows the biases and mean squared errors (MSEs) of the parameter estimates with 5000 replications for various values of the parameters  $(n, \theta, \lambda)$ . According to the results in Table 3, the MLEs are asymptotically unbiased. Moreover, the bias and MSE values of the MLEs decrease to zero as the sample size increases as desired.
Table 3. Bias and MSEs of MLE estimators for selected parameter values											
			<b>Bias</b>	<b>MSE</b>							
$\theta$	λ	п	$\widehat{\theta}$	Â	Ô	Â					
$\overline{2}$	$\overline{2}$	50	0.06108	0.10841	0.09290	0.19997					
		100	0.02521	0.04550	0.04196	0.08155					
		200	0.01127	0.02281	0.01972	0.03706					
		300	0.00690	0.01386	0.01311	0.02379					
		500	0.00304	0.00752	0.00779	0.01394					
3	1.5	50	0.09989	0.06955	0.24153	0.09187					
		100	0.04119	0.02920	0.10856	0.03866					
		200	0.01851	0.01493	0.05089	0.01787					
		300	0.01147	0.00909	0.03392	0.01148					
		500	0.00511	0.00496	0.02015	0.00676					
1.5	0.8	50	0.06884	0.02793	0.09443	0.01843					
		100	0.02873	0.01179	0.04116	0.00817					
		200	0.01308	0.00629	0.01895	0.00389					
		300	0.00834	0.00381	0.01267	0.00250					
		500	0.00396	0.00216	0.00753	0.00148					

**5. Real Data Application**

In this section, we investigate the flexibility performance of the  $D$ -UEHL( $\theta$ , $\lambda$ ) distribution with a real data application. For this purpose, we consider the reservoir data set obtained from the monthly water capacity of the Shasta Reservoir in California [22]. We compare the  $D$ -UEHL( $\theta$ ,  $\lambda$ ) distribution with well-known Weibull, Beta, Kumaraswamy, and UEHL distributions. The PDFs of the distributions used for comparison are given as follows:

Weibull distribution

$$
f_W(x; \theta, \lambda) = \frac{\theta}{\lambda} \left(\frac{x}{\lambda}\right)^{\theta - 1} e^{-(x/\lambda)^{\theta}}, \quad \theta, \lambda > 0
$$

Beta distribution

$$
f_B(x; \theta, \lambda) = \frac{1}{B(\theta, \lambda)} x^{\theta - 1} (1 - x)^{\lambda - 1}, \quad \theta, \lambda > 0
$$

Kumaraswamy distribution

$$
f_{Kw}(x;\theta,\lambda) = \theta\lambda x^{\theta-1} (1-x^{\theta})^{\lambda-1}, \quad \theta,\lambda > 0
$$

We use the maximum likelihood method to estimate the unknown parameters of the considered distributions. We test the goodness of fits of the models with Kolmogorov-Smirnov test statistic (K-S (stat)) and associate  $p$ -value (K-S (p-value)). We report the Akaike information criterion (AIC) and Bayesian information criterion (BIC) to compare the proposed model with the other models. AIC and BIC are computed as

$$
AIC = 2k - 2\ell(\theta, \lambda) \quad \text{and} \quad BIC = k \log n - 2\ell(\theta, \lambda)
$$

where k is the number of parameters,  $n$  is the number of observations, and  $\ell$  is the maximum value of the likelihood function for the underlying distribution.

Table 4. Dias and MSES Of MEE estimators for selected barameter values									
			AIC	BIC-	$-2\ell$	$K-S$ (stat)	$K-S$ (p-value)		
Weibull	7.2987	0.7748	$-20.5347$	$-18.5432$	$-24.5347$	0.2220	0.2396		
Beta	7.3157	2.9099	$-21.1238$	$-19.1324$	$-25.1238$	0.2359	0.1834		
Kumaraswamy	6.3476	4.4894	-22.9494	$-20.9580$	$-26.9494$	0.2209	0.2447		
<b>UEHL</b>	6.9010	2.8883	$-21.2699$	$-19.2784$	$-25.2699$	0.2254	0.2248		
D-UEHL	6.4316	3.2682	$-23.0543$	$-21.0629$	$-27.0543$	0.2046	0.3267		

Table 4. Bias and MSEs of MLE estimators for selected parameter values

Table 4 shows the parameter estimates, AIC values, BIC values, and  $-2\ell$  values together with K-S (stat) and K-S ( $p$ -value) for all models fitted to the reservoir data set. The K-S test results indicate that the compared distributions are appropriate to fit the underlying data. From Table 4, it can be observed that the  $D$ -UEHL( $\theta$ ,  $\lambda$ ) distribution yields the lowest AIC and BIC values, followed by the Kumaraswamy distribution. Hence, we conclude that the  $D$ -UEHL $(\theta, \lambda)$  model has the best fit for the reservoir data set among the compared distributions in the study. Moreover, Figure 3 manifests the empirical and fitted curves based on the reservoir data set for illustrative purposes.



Figure 3: Empirical and fitted distribution functions based on the reservoir data set

#### **6. Conclusion**

In this study, we introduce  $D$ -UEHL distribution for record data based on DUS transformation on unit exponentiated half-logistic distribution CDF. We obtain various analytical characteristics, including moments, quantile function, stress-strength reliability, and likelihood ratio ordering of the proposed distribution. We perform a simulation study that illustrates the properties of the maximum likelihood estimates of the  $D$ -UEHL distribution. The real data analysis on a record data shows that the D-UEHL model performs better on the data than the other well-known models in terms of AIC and BIC criteria.

The findings of this article can be extended by applying the DUS transform to the omega distribution. In this context, the effect of the extra support parameter of the omega distribution with the DUS transformation is a topic for further research.

#### **Author Contributions**

All the authors equally contributed to this work. They all read and approved the final version of the paper.

#### **Conflict of Interest**

All the authors declare no conflict of interest.

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### **On Dual Quaternions with** *k*−**Generalized Leonardo Components**

 $\operatorname{Cijden}$  Zeynep Yılmaz<sup>ı ()</sup> , Gülsüm Yeliz Saçlı<sup>2 ()</sup>

*Received*: 17 Jul 2023 *Accepted*: 21 Sep 2023 *Published*: 30 Sep 2023 doi[:10.53570/jnt.1328605](https://doi.org/10.53570/jnt.1328605) the present paper is that one can reduce the calculations of Leonardo-like dual quaternions Research Article **Abstract** − In this paper, we define a one-parameter generalization of Leonardo dual quaternions, namely *k*−generalized Leonardo-like dual quaternions. We introduce the properties of *k*-generalized Leonardo-like dual quaternions, including relations with Leonardo, Fibonacci, and Lucas dual quaternions. We investigate their characteristic relations, involving the Binet-like formula, the generating function, the summation formula, Catalan-like, Cassini-like, d'Ocagne-like, Tagiuri-like, and Hornsberger-like identities. The crucial part of by considering  $k$ . For  $k = 1$ , these results are generalizations of the ones for ordered Leonardo quadruple numbers. Finally, we discuss the need for further research.

**Keywords** *Leonardo sequence, recurrence relations, dual quaternions*

**Mathematics Subject Classification (2020)** 11B37, 11B39

#### **1. Introduction**

**Article Info**

The well-known Fibonacci sequence  ${F_n}_{n>2}$  and the Lucas sequence  ${L_n}_{n>2}$  are defined recursively by  $F_n = F_{n-1} + F_{n-2}$ , and  $L_n = L_{n-1} + L_{n-2}$  with initial conditions  $F_0 = 0$ ,  $F_1 = 1$ , and  $L_0 = 2$ ,  $L_1 = 1$ , respectively [1]. The Binet's formulas of the Fibonacci and Lucas sequences are as follows, respectively:

$$
F_n = \frac{\alpha^n - \beta^n}{\alpha - \beta} \tag{1}
$$

and

 $L_n = \alpha^n + \beta^n$ 

where  $\alpha = \frac{1+\sqrt{5}}{2}$  $\frac{-\sqrt{5}}{2}$  and  $\beta = \frac{1-\sqrt{5}}{2}$  $\frac{x^2}{2}$  are roots of characteristic equation  $x^2 - x - 1 = 0$  [1]. Many generalizations of the Fibonacci and Lucas sequences have been studied by several researchers. In this study, we consider the Leonardo sequence. The Leonardo sequence  ${L_{e_n}}_{n>2}$  is defined nonhomogeneous recursively by

$$
Le_n = Le_{n-1} + Le_{n-2} + 1
$$

or

$$
Le_{n+1} = 2Le_n - Le_{n-2}
$$

<sup>&</sup>lt;sup>1</sup>cigdem.yilmaz@std.yildiz.edu.tr; <sup>2</sup>yeliz.sacli@yildiz.edu.tr (Corresponding Author)

<sup>&</sup>lt;sup>1,2</sup>Department of Mathematics, Faculty of Arts and Sciences, Yildiz Technical University, İstanbul, Türkiye

with initial conditions  $Le_0 = Le_1 = 1$  and  $Le_2 = 3$  [2]. The Binet-like formula of the Leonardo sequence is

$$
Le_n = 2\left(\frac{\alpha^{n+1} - \beta^{n+1}}{\alpha - \beta}\right) - 1
$$

There exist many identities between Fibonacci, Lucas, and Leonardo numbers. For  $n \geq 0$ , the fundamental relationships between the Fibonacci, Lucas, and Leonardo sequences are [2]:

$$
Le_n = 2F_{n+1} - 1
$$

$$
Le_n = 2\left(\frac{L_n + L_{n+2}}{5}\right) - 1
$$

$$
Le_{n+3} = \frac{L_{n+1} + L_{n+7}}{5} - 1
$$

and

$$
Le_n = L_{n+2} - F_{n+2} - 1
$$

Although the Fibonacci, Lucas, and Leonardo sequences are closely related, they exhibit distinct characteristic properties. Several different properties and generalizations of the Leonardo sequence were previously studied by various researchers [3–15]. Recently, a one-parameter generalized Leonardo sequence has been defined as non-homogeneous recursively by

$$
\mathsf{Le}_n^{(k)} = \mathsf{Le}_{n-1}^{(k)} + \mathsf{Le}_{n-2}^{(k)} + k, \quad n \ge 2
$$
 (2)

with the initial conditions  $\mathsf{Le}_0^{(k)} = \mathsf{Le}_1^{(k)} = 1$ . Here, k is a fixed positive integer [4]. The Binet-like formula of the *k*−generalized Leonardo sequence is

$$
\mathsf{Le}_n^{(k)} = (k+1) \left( \frac{\alpha^{n+1} - \beta^{n+1}}{\alpha - \beta} \right) - k
$$

The *k*−generalized Leonardo sequence is related to the Fibonacci and Lucas sequences. For *n* ≥ 0, the fundamental relationships between Fibonacci, Lucas, and the *k*−generalized Leonardo sequences are [4]:

$$
\mathsf{Le}_n^{(k)} = (k+1)F_{n+1} - k \tag{3}
$$

and

$$
\mathsf{Le}_n^{(k)} = (k+1) \left( L_n - F_{n-1} \right) - k
$$

The summation formulas of the *k*−generalized Leonardo numbers are [4]:

$$
\sum_{s=0}^{n} \mathsf{Le}_{s}^{(k)} = \mathsf{Le}_{n+2}^{(k)} - k(n+1) - 1
$$
\n(4)

$$
\sum_{s=0}^{n} \mathsf{Le}_{2s}^{(k)} = \mathsf{Le}_{2n+1}^{(k)} - kn \tag{5}
$$

and

$$
\sum_{s=0}^{n} \mathsf{Le}_{2s+1}^{(k)} = \mathsf{Le}_{2n+2}^{(k)} - k(n+2)
$$
 (6)

The *k*−generalized Leonardo sequence is the key concept of the present paper. For *k* = 1, this sequence is the classical Leonardo sequence, i.e.,  $\mathsf{Le}_n^{(1)} = Le_n$ . In this case, we may omit the superscript (1) in the notation.

There are several ways to define new special sequences but the most popular method is to define a

sequence with different hypercomplex number components. Horadam [16] defined the real quaternions with the classic Fibonacci sequence  $\{QF_n\}_{n\geq 2}$  and the classic Lucas sequence  $\{QL_n\}_{n\geq 2}$  recursively by

$$
QF_n = F_n + F_{n+1}e_1 + F_{n+2}e_2 + F_{n+3}e_3
$$

and

$$
QL_n = L_n + L_{n+1}e_1 + L_{n+2}e_2 + L_{n+3}e_3
$$

where  $F_n$  and  $L_n$  are the *n*-th classic Fibonacci and Lucas numbers, respectively. Here, the quaternionic units  $\{e_1, e_2, e_3\}$  satisfy the following multiplication rules:

$$
e_1^2 = e_2^2 = e_3^2 = -1
$$
,  $e_1e_2 = -e_2e_1 = e_3$ ,  $e_2e_3 = -e_3e_2 = e_1$ , and  $e_3e_1 = -e_1e_3 = e_2$ 

In general, a real quaternion *q* is of the form  $q = a + e_1b + e_2c + e_3d$  where  $a, b, c, d \in \mathbb{R}$ . The quaternions form a four-dimensional associative and non-commutative algebra over the real numbers. For a deeper discussion of the quaternions, see [17–19]. Changing conditions in the multiplication rules produces different type of quaternions. In this study, we consider the dual quaternions. The dual quaternionic units obey the following multiplication rules:

$$
e_1^2 = e_2^2 = e_3^2 = 0
$$
 and  $e_1e_2 = -e_2e_1 = e_2e_3 = -e_3e_2 = e_3e_1 = -e_1e_3 = 0$  (7)

Yüce et al. [20] defined the dual quaternions with the Fibonacci sequence  $\{\hat{\mathcal{F}}_n\}_{n\geq 2}$  and the Lucas sequence  $\{\hat{\mathcal{L}}_n\}_{n\geq 2}$  recursively by

$$
\hat{\mathcal{F}}_n = F_n + F_{n+1}e_1 + F_{n+2}e_2 + F_{n+3}e_3
$$

and

$$
\hat{\mathcal{L}}_n = L_n + L_{n+1}e_1 + L_{n+2}e_2 + L_{n+3}e_3
$$

where  $F_n$  and  $L_n$  are the *n*-th classic Fibonacci and Lucas numbers, respectively. Here, the non-real dual quaternionic units  $\{e_1, e_2, e_3\}$  satisfy Equation 7. For a deeper discussion of the dual quaternions, see [21–27]. Nurkan et al. [13] defined the dual quaternions with the classic Leonardo sequence  $\{\mathcal{\hat{L}}e_n\}_{n\geq 2}$  recursively by

$$
\hat{\mathcal{L}}e_n = Le_n + Le_{n+1}e_1 + Le_{n+2}e_2 + Le_{n+3}e_3
$$

where  $Le_n$  is the *n*-th classic Leonardo number.

Considering all these details, a natural question is whether the paper [13] can be generalized. In this study, we aim to determine the dual quaternions with the *k*−generalized Leonardo sequence. We consider the coefficients of the dual quaternions as the *k*−generalized Leonardo sequence.

#### **2. The** *k*−**Generalized Leonardo-like Dual Quaternion Sequence**

This section introduces the dual quaternions with one parameter generalized Leonardo sequence and investigates their characteristic properties.

**Definition 2.1.** The *n*-th *k*−generalized Leonardo-like dual quaternion is defined by

$$
\hat{\mathsf{Le}}_{n}^{(k)} = \mathsf{Le}_{n}^{(k)} + \mathsf{Le}_{n+1}^{(k)} e_1 + \mathsf{Le}_{n+2}^{(k)} e_2 + \mathsf{Le}_{n+3}^{(k)} e_3 \tag{8}
$$

 $\mathsf{where} \ \mathsf{Le}_n^{(k)} \ \text{is the $n$-th $k$--generalized Leonardo number, $k$ is a fixed positive integer, and \{e_1, e_2, e_3\}$ are$ the set of all the dual quaternionic units. The *k*−generalized Leonardo-like dual quaternion sequence is denoted by  $\left\{\mathcal{L}\mathbf{e}_n^{(k)}\right\}$ *n*  $\mathcal{L}$ *n*≥2 .

Note that, if  $k = 1$ , the generalized Leonardo-like dual quaternion sequence  $\left\{\hat{\mathsf{Le}}_n^{(k)}\right\}$ *n*  $\lambda$ *n*≥2 is the Leonardo quadruple sequence. In this case, we may omit the superscript (1) in the notation.

Let  $\hat{\sf Le}_n^{(k)}$  and  $\hat{\sf Le}_m^{(k)}$  be two  $k-$ generalized Leonardo-like dual quaternions. The  $k-$ generalized Leonardo number  $\mathsf{Le}_n^{(k)}$  is called scalar (real) part of  $\hat{\mathsf{Le}}_n^{(k)}$  and denoted by  $\mathcal{S}_{\hat{\mathsf{Le}}_n^{(k)}}$  and the vector

$$
\operatorname{Le}_{n+1}^{(k)} e_1 + \operatorname{Le}_{n+2}^{(k)} e_2 + \operatorname{Le}_{n+3}^{(k)} e_3
$$

is called the pure part of  $\hat{\mathsf{Le}}_n^{(k)}$  and denoted by  $\mathcal{V}_{\hat{\mathsf{Le}}_n^{(k)}}$ . The addition is defined component-wise as

$$
\hat{\mathsf{Le}}_n^{(k)} \pm \hat{\mathsf{Le}}_m^{(k)} = \left( \mathsf{Le}_n^{(k)} \pm \mathsf{Le}_m^{(k)} \right) + \left( \mathsf{Le}_{n+1}^{(k)} \pm \mathsf{Le}_{m+1}^{(k)} \right) e_1 + \left( \mathsf{Le}_{n+2}^{(k)} \pm \mathsf{Le}_{m+2}^{(k)} \right) e_2 + \left( \mathsf{Le}_{n+3}^{(k)} \pm \mathsf{Le}_{m+3}^{(k)} \right) e_3
$$

whereas multiplication is defined by

$$
\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m}^{(k)} = \left(\mathsf{Le}_{n}^{(k)}\mathsf{Le}_{m}^{(k)}\right) + \left(\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m+1}^{(k)} + \hat{\mathsf{Le}}_{m}^{(k)}\hat{\mathsf{Le}}_{n+1}^{(k)}\right)e_{1} + \left(\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m+2}^{(k)} + \hat{\mathsf{Le}}_{m}^{(k)}\hat{\mathsf{Le}}_{n+2}^{(k)}\right)e_{2} + \left(\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m+3}^{(k)} + \hat{\mathsf{Le}}_{m}^{(k)}\hat{\mathsf{Le}}_{n+3}^{(k)}\right)e_{3}
$$

or

$$
\hat{\mathsf{Le}}_n^{(k)}\hat{\mathsf{Le}}_m^{(k)} = \mathcal{S}_{\hat{\mathsf{Le}}_n^{(k)}}\mathcal{S}_{\hat{\mathsf{Le}}_m^{(k)}} + \mathcal{S}_{\hat{\mathsf{Le}}_n^{(k)}}\mathcal{V}_{\hat{\mathsf{Le}}_m^{(k)}} + \mathcal{S}_{\hat{\mathsf{Le}}_m^{(k)}}\mathcal{V}_{\hat{\mathsf{Le}}_n^{(k)}}
$$

The conjugate and norm of any *k*−generalized Leonardo-like dual quaternion  $\hat{\mathsf{Le}}_n^{(k)}$  $\int_{n}^{\infty}$  is given by

$$
\widehat{\mathsf{Le}_n^{(k)}} = \mathcal{S}_{\widehat{\mathsf{Le}}_n^{(k)}} - \mathcal{V}_{\widehat{\mathsf{Le}}_n^{(k)}} = \widehat{\mathsf{Le}}_n^{(k)} - \widehat{\mathsf{Le}}_{n+1}^{(k)} e_1 - \widehat{\mathsf{Le}}_{n+2}^{(k)} \, e_2 - \widehat{\mathsf{Le}}_{n+3}^{(k)} e_3
$$

and

$$
||\hat{\mathsf{Le}}_{n}^{(k)}|| = \hat{\mathsf{Le}}_{n}^{(k)} \overline{\hat{\mathsf{Le}}_{n}^{(k)}} = \left(\hat{\mathsf{Le}}_{n}^{(k)}\right)^2 \in \mathbb{R}
$$

respectively.

**Theorem 2.2.** The recurrence relation of the *k*−generalized Leonardo-like dual quaternion sequence is

$$
\hat{\mathsf{Le}}_n^{(k)} = \hat{\mathsf{Le}}_{n-1}^{(k)} + \hat{\mathsf{Le}}_{n-2}^{(k)} + \hat{\mathcal{K}}, \quad n \ge 2
$$

where  $\hat{\mathcal{K}} = k(1 + e_1 + e_2 + e_3)$  with initial conditions  $\hat{\mathsf{Le}}_0^{(k)} = 1 + e_1 + (2 + k) e_2 + (3 + 2k) e_3$  and  $\hat{\text{Le}}_1^{(k)} = 1 + (2 + k) e_1 + (3 + 2k) e_2 + (5 + 4k) e_3.$ 

Proof*.*

From Definition 2.1, it follows that

$$
\begin{array}{ll}\hat{{\rm Le}}_{n-1}^{(k)}+\hat{{\rm Le}}_{n-2}^{(k)}+\hat{{\cal K}}=&\left({\rm Le}_{n-1}^{(k)}+{\rm Le}_{n-2}^{(k)}+k\right)+\left({\rm Le}_{n}^{(k)}+{\rm Le}_{n-1}^{(k)}+k\right)e_1+\left({\rm Le}_{n+1}^{(k)}+{\rm Le}_{n}^{(k)}+k\right)e_2\\&+\left({\rm Le}_{n+2}^{(k)}+{\rm Le}_{n+1}^{(k)}+k\right)e_3\end{array}
$$

By applying Equation 2, we complete the proof.  $\Box$ 

Throughout this paper, let  $\hat{\mathcal{K}} = k(1 + e_1 + e_2 + e_3)$ .

**Theorem 2.3.** The other recurrence relation of  $\left\{\hat{\mathsf{Le}}_n^{(k)}\right\}$ *n*  $\mathcal{L}$ *n*≥2 is  $\hat{\mathsf{Le}}_{n+1}^{(k)} = 2 \hat{\mathsf{Le}}_n^{(k)} - \hat{\mathsf{Le}}_{n-1}^{(k)}$ *n*−2

Proof*.*

By Theorem 2.2, the proof is straightforward.  $\Box$ 

Afterward, we state the Binet-like formula for the *k*−generalized Leonardo-like dual quaternion  $\hat{\mathsf{Le}}_n^{(k)}$  $\binom{n}{n}$ . Thus, we derive some well-known mathematical properties.

**Theorem 2.4.** The Binet-like formula of the *k*−generalized Leonardo-like dual quaternion  $\hat{\mathsf{Le}}_n^{(k)}$  $\int_n^{\infty}$  is

$$
\hat{\mathsf{Le}}_n^{(k)} = (k+1) \left( \frac{\alpha^* \alpha^{n+1} - \beta^* \beta^{n+1}}{\alpha - \beta} \right) - \hat{\mathcal{K}} \tag{9}
$$

where  $\alpha^* = 1 + \alpha e_1 + \alpha^2 e_2 + \alpha^3 e_3$  and  $\beta^* = 1 + \beta e_1 + \beta^2 e_2 + \beta^3 e_3$ .

Proof*.*

By using Definition 2.1 and Equation 3,

$$
\hat{\mathsf{Le}}_{n}^{(k)} = \mathsf{Le}_{n}^{(k)} + \mathsf{Le}_{n+1}^{(k)} e_1 + \mathsf{Le}_{n+2}^{(k)} e_2 + \mathsf{Le}_{n+3}^{(k)} e_3
$$
\n
$$
= (k+1) \left( F_{n+1} + F_{n+2} e_1 + F_{n+3} e_2 + F_{n+4} e_3 \right) - k \left( 1 + e_1 + e_2 + e_3 \right)
$$

Applying the Binet's formula of the Fibonacci sequence in Equation 1 and then taking  $\alpha^* = 1 + \alpha e_1 + \alpha e_2$  $\alpha^2 e_2 + \alpha^3 e_3$ ,  $\beta^* = 1 + \beta e_1 + \beta^2 e_2 + \beta^3 e_3$ , and  $\hat{\mathcal{K}} = k(1 + e_1 + e_2 + e_3)$ ,

$$
\hat{\mathsf{Le}}_{n}^{(k)} = (k+1) \left( \frac{\alpha^{n+1} \left( 1 + \alpha e_1 + \alpha^2 e_2 + \alpha^3 e_3 \right) - \beta^{n+1} \left( 1 + \beta e_1 + \beta^2 e_2 + \beta^3 e_3 \right)}{\alpha - \beta} \right) - \hat{\mathcal{K}}
$$
\n
$$
= (k+1) \left( \frac{\alpha^* \alpha^{n+1} - \beta^* \beta^{n+1}}{\alpha - \beta} \right) - \hat{\mathcal{K}}
$$

is obtained.  $\square$ 

Here, we state some relations between *k*−generalized Leonardo-like dual quaternions, Fibonacci dual quaternions, Lucas dual quaternions, and Fibonacci and Lucas numbers.

**Theorem 2.5.** Let  $\hat{\mathsf{L}}_{e_n}^{(k)}$  be the *n*-th *k*-generalized Leonardo-like dual quaternion,  $\hat{\mathcal{F}}_n$  be the *n*-th Fibonacci dual quaternion, and  $\mathcal{L}_n$  be the *n*-th Lucas dual quaternion. For positive integers *n*, *m*, *r*, and *t* with  $n \geq r$  and  $n \geq m$ , the following relations hold:

*i.* 
$$
\hat{L}e_n^{(k)} = (k+1)\hat{\mathcal{F}}_{n+1} - \hat{\mathcal{K}}
$$
  
\n*ii.*  $\hat{L}e_n^{(k)} = (k+1)(\hat{\mathcal{L}}_n - \hat{\mathcal{F}}_{n-1}) - \hat{\mathcal{K}}$   
\n*iii.*  $\hat{L}e_{n+r}^{(k)} + \hat{L}e_{n-r}^{(k)} = (k+1)\begin{cases} L_r\hat{\mathcal{F}}_{n+1} - 2\hat{\mathcal{K}}, \ r = 2t \\ F_r\hat{\mathcal{L}}_{n+1} - 2\hat{\mathcal{K}}, \ r = 2t + 1 \end{cases}$   
\n*iv.*  $\hat{L}e_{n+r}^{(k)} - \hat{L}e_{n-r}^{(k)} = (k+1)\begin{cases} F_r\hat{\mathcal{L}}_{n+1}, \ r = 2t \\ L_r\hat{\mathcal{F}}_{n+1}, \ r = 2t + 1 \end{cases}$   
\n*v.*  $\hat{L}e_{n+m}^{(k)} + (-1)^m \hat{L}e_{n-m}^{(k)} = L_m \hat{L}e_n^{(k)} + \hat{\mathcal{K}}(L_m - (-1)^m - 1)$   
\n*vi.*  $\hat{L}e_{n+m}^{(k)} - (-1)^m \hat{L}e_{n-m}^{(k)} = (k+1) F_m\hat{\mathcal{L}}_{n+1} - \hat{\mathcal{K}} (1 - (-1)^m)$ 

Proof*.*

Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion,  $\hat{\mathcal{F}}_n$  be the *n*-th Fibonacci dual quaternion, and  $\mathcal{\hat{L}}_n$  be the *n*-th Lucas dual quaternion.

*i.* According to Equation 3,

$$
\hat{\mathsf{Le}}_{n}^{(k)} = ((k+1) F_{n+1} - k) + ((k+1) F_{n+2} - k) e_1 + ((k+1) F_{n+3} - k) e_2 + ((k+1) F_{n+4} - k) e_3
$$
  
= (k+1)  $\hat{\mathcal{F}}_{n+1} - \hat{\mathcal{K}}$ 

*iii.*

$$
\begin{aligned} \hat{\mathsf{Le}}_{n+r}^{(k)}+\hat{\mathsf{Le}}_{n-r}^{(k)} &=\left(\mathsf{Le}_{n+r}^{(k)}+\mathsf{Le}_{n-r}^{(k)}\right)+\left(\mathsf{Le}_{n+1+r}^{(k)}+\mathsf{Le}_{n+1-r}^{(k)}\right)e_1+\left(\mathsf{Le}_{n+2+r}^{(k)}+\mathsf{Le}_{n+2-r}^{(k)}\right)e_2\\ &+\left(\mathsf{Le}_{n+3+r}^{(k)}+\mathsf{Le}_{n+3-r}^{(k)}\right)e_3 \end{aligned}
$$

Considering Equation 3,

$$
\hat{\mathsf{Le}}_{n+r}^{(k)} + \hat{\mathsf{Le}}_{n-r}^{(k)} = (k+1) \left( (F_{n+1+r} + F_{n+1-r}) + (F_{n+2+r} + F_{n+2-r}) e_1 + (F_{n+3+r} + F_{n+3-r}) e_2 + (F_{n+4+r} + F_{n+4-r}) e_3 \right) - 2\hat{\mathcal{K}}
$$

By using the definition of Fibonacci dual quaternion and the following relation of Fibonacci numbers (see [1])

$$
F_{n+r} + F_{n-r} = \begin{cases} F_n L_r, & r = 2t \\ F_r L_n, & r = 2t + 1 \end{cases}
$$

we complete the proof.

*v.* From Theorem  $2.5(i)$  and the Binet's formulas of Fibonacci and Lucas numbers,

$$
\hat{\mathsf{Le}}_{n+m}^{(k)} + (-1)^m \hat{\mathsf{Le}}_{n-m}^{(k)} = ((k+1)\hat{\mathcal{F}}_{n+m+1} - \hat{\mathcal{K}}) + (-1)^m ((k+1)\hat{\mathcal{F}}_{n-m+1} - \hat{\mathcal{K}})
$$
\n
$$
= (k+1)\hat{\mathcal{F}}_{n+1}L_m + \hat{\mathcal{K}}(-1 - (-1)^m)
$$
\n
$$
= ((k+1)\hat{\mathcal{F}}_{n+1} - \hat{\mathcal{K}}) L_m + \hat{\mathcal{K}}(L_m - (-1)^m - 1)
$$
\n
$$
= L_m \hat{\mathsf{Le}}_n^{(k)} + \hat{\mathcal{K}}(L_m - (-1)^m - 1)
$$

 $\Box$ 

**Corollary 2.6.** Using the identities *iii* and *iv* presented in Theorem 2.5, the following basic identities are obtained:

*i*.  $\hat{\mathsf{Le}}_{n+1}^{(k)} + \hat{\mathsf{Le}}_{n-1}^{(k)} = (k+1)\hat{\mathcal{L}}_{n+1} - 2\hat{\mathcal{K}}$  $ii. \ \hat{\text{Le}}_{n+1}^{(k)} - \hat{\text{Le}}_{n-1}^{(k)} = (k+1)\hat{\mathcal{F}}_{n+1}$  $iii$ .  $\hat{\text{Le}}_{n+2}^{(k)} + \hat{\text{Le}}_{n-2}^{(k)} = 3(k+1)\hat{\mathcal{F}}_{n+1} - 2\hat{\mathcal{K}}$  $iv. \ \hat{\text{Le}}_{n+2}^{(k)} - \hat{\text{Le}}_{n-2}^{(k)} = (k+1)\hat{\mathcal{L}}_{n+1}$ 

**Theorem 2.7.** Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion. Then, the following relations hold:

$$
\begin{aligned}\ni. \ \hat{\mathsf{Le}}_{n}^{(k)} - \hat{\mathsf{Le}}_{n+1}^{(k)} e_1 - \hat{\mathsf{Le}}_{n+2}^{(k)} e_2 - \hat{\mathsf{Le}}_{n+3}^{(k)} e_3 &= \mathsf{Le}_{n}^{(k)} \\
ii. \ \hat{\mathsf{Le}}_{n}^{(k)} + \overline{\hat{\mathsf{Le}}_{n}^{(k)}} &= 2\mathsf{Le}_{n}^{(k)} \\
iii. \ \left(\hat{\mathsf{Le}}_{n}^{(k)}\right)^2 &= 2\mathsf{Le}_{n}^{(k)} \hat{\mathsf{Le}}_{n}^{(k)} - \left(\mathsf{Le}_{n}^{(k)}\right)^2 \\
iv. \ \hat{\mathsf{Le}}_{n}^{(k)} \overline{\hat{\mathsf{Le}}_{n}^{(k)}} + \hat{\mathsf{Le}}_{n+1}^{(k)} \overline{\hat{\mathsf{Le}}_{n+1}^{(k)}} &= (k+1) \mathsf{Le}_{2n+2}^{(k)} - 2k \mathsf{Le}_{n+2}^{(k)} + k (k+1) \\
v. \ \hat{\mathsf{Le}}_{n+1}^{(k)} \overline{\hat{\mathsf{Le}}_{n+1}^{(k)}} - \hat{\mathsf{Le}}_{n-1}^{(k)} \overline{\hat{\mathsf{Le}}_{n-1}^{(k)}} &= (k+1) \mathsf{Le}_{2n+1}^{(k)} - 2k \mathsf{Le}_{n}^{(k)} - k (k-1) \\
\end{aligned}
$$

Proof*.*

Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion.

*iv.* Using Equation 3 and the Binet's formula of the Fibonacci sequence in Equation 1,

$$
\begin{split}\n\hat{\mathsf{Le}}_{n}^{(k)} \overline{\hat{\mathsf{Le}}_{n}^{(k)}} + \hat{\mathsf{Le}}_{n+1}^{(k)} \overline{\hat{\mathsf{Le}}_{n+1}^{(k)}} &= \left( \mathsf{Le}_{n}^{(k)} \right)^{2} + \left( \mathsf{Le}_{n+1}^{(k)} \right)^{2} \\
&= \left( (k+1) \, F_{n+1} - k \right)^{2} + \left( (k+1) \, F_{n+2} - k \right)^{2} \\
&= (k+1)^{2} \left( F_{n+1}^{2} + F_{n+2}^{2} \right) - 2k \left( k+1 \right) \left( F_{n+1} + F_{n+2} \right) + 2k^{2} \\
&= (k+1)^{2} \, F_{2n+3} - 2k \left( k+1 \right) F_{n+3} + 2k^{2} \\
&= (k+1) \left( (k+1) \, F_{2n+3} - k \right) + k \left( k+1 \right) - 2k \left( (k+1) \, F_{n+3} - k \right) \\
&= (k+1) \, \mathsf{Le}_{2n+2}^{(k)} - 2k \, \mathsf{Le}_{n+2}^{(k)} + k \left( k+1 \right)\n\end{split}
$$

is obtained.

*v.* Applying Equation 3 and the Binet's formula of the Fibonacci sequence in Equation 1,

$$
\hat{\mathsf{Le}}_{n+1}^{(k)} \overline{\hat{\mathsf{Le}}_{n+1}^{(k)}} - \hat{\mathsf{Le}}_{n-1}^{(k)} \overline{\hat{\mathsf{Le}}_{n-1}^{(k)}} = (\mathsf{Le}_{n+1}^{(k)})^2 - (\mathsf{Le}_{n-1}^{(k)})^2
$$
\n
$$
= ((k+1) F_{n+2} - k)^2 - ((k+1) F_n - k)^2
$$
\n
$$
= (k+1)^2 (F_{n+2}^2 - F_n^2) + 2k (k+1) (F_n - F_{n+2})
$$
\n
$$
= (k+1) ((k+1) F_{2n+2} - k) + k (k+1) - 2k ((k+1) F_{n+1} - k) - 2k^2
$$
\n
$$
= (k+1) \mathsf{Le}_{2n+1}^{(k)} - 2k \mathsf{Le}_n^{(k)} - k (k-1)
$$

is obtained.

 $\Box$ 

**Theorem 2.8.** Let  $\hat{\mathsf{L}}_{n}^{(k)}$  be the *n*-th *k*-generalized Leonardo-like dual quaternion. For  $n \geq 2$ , the generating function  $\overline{\mathsf{G}(x)} = \sum_{n=0}^{\infty}$ *n*=0  $\hat{\mathsf{Le}}_n^{(k)}$  $n^{(n)}$  is as follows:  $G(x) =$  $\hat{\mathsf{Le}}_0^{(k)} + \hat{\mathsf{Le}}_1^{(k)} - 2\hat{\mathsf{Le}}_0^{(k)}$ 0  $x + (\hat{L}_{2}^{(k)} - 2\hat{L}_{1}^{(k)})$ 1  $\Big) x^2$  $1 - 2x + x^3$ =  $(1-x+kx^2)+(1+ kx-x^2) e_1+(2+k-x+x^2) e_2+(3+2k-x+(-2-k)x^2) e_3$  $1 - 2x + x^3$ 

where  $1 - 2x + x^3 \neq 0$ .

#### Proof*.*

The proof is similar to the proof of the generating function of the Leonardo sequence in [2].  $\Box$ 

**Theorem 2.9.** For  $n \geq 0$ , the following summation formulas are satisfied:

$$
i. \sum_{s=0}^{n} \hat{\mathsf{Le}}_{s}^{(k)} = \hat{\mathsf{Le}}_{n+2}^{(k)} - \hat{\mathcal{K}}(n+2) + (k-1) - 2e_1 + (-k-3)e_2 + (-5-3k)e_3
$$
  

$$
ii. \sum_{s=0}^{n} \hat{\mathsf{Le}}_{2s}^{(k)} = \hat{\mathsf{Le}}_{2n+1}^{(k)} - \hat{\mathcal{K}}n - (2k)e_1 + (-k-1)e_2 + (-3k-1)e_3
$$
  

$$
iii. \sum_{s=0}^{n} \hat{\mathsf{Le}}_{2s+1}^{(k)} = \hat{\mathsf{Le}}_{2n+2}^{(k)} - \hat{\mathcal{K}}n - (2k) + (-k-1)e_1 + (-3k-1)e_2 + (-3k-3)e_3
$$

Proof*.*

Let  $n \geq 0$ .

*i.* Using Equations 4 and 8,

$$
\sum_{s=0}^{n} \hat{\mathsf{Le}}_{s}^{(k)} = \sum_{s=0}^{n} \mathsf{Le}_{s}^{(k)} + \left(\sum_{s=0}^{n} \mathsf{Le}_{s+1}^{(k)}\right) e_{1} + \left(\sum_{s=0}^{n} \mathsf{Le}_{s+2}^{(k)}\right) e_{2} + \left(\sum_{s=0}^{n} \mathsf{Le}_{s+3}^{(k)}\right) e_{3}
$$
\n
$$
= \left(\mathsf{Le}_{n+2}^{(k)} - k(n+1) - 1\right) + \left(\sum_{s=0}^{n+1} \mathsf{Le}_{s}^{(k)} - \mathsf{Le}_{0}^{(k)}\right) e_{1} + \left(\sum_{s=0}^{n+2} \mathsf{Le}_{s}^{(k)} - \mathsf{Le}_{1}^{(k)}\right) e_{2}
$$
\n
$$
+ \left(\sum_{s=0}^{n+3} \mathsf{Le}_{s}^{(k)} - \mathsf{Le}_{0}^{(k)} - \mathsf{Le}_{1}^{(k)} - \mathsf{Le}_{2}^{(k)}\right) e_{3}
$$
\n
$$
= \left(\mathsf{Le}_{n+2}^{(k)} - k(n+1) - 1\right) + \left(\mathsf{Le}_{n+3}^{(k)} - k(n+2) - 1 - 1\right) e_{1}
$$
\n
$$
+ \left(\mathsf{Le}_{n+4}^{(k)} - k(n+3) - 1 - 1 - 1\right) e_{2} + \left(\mathsf{Le}_{n+5}^{(k)} - k(n+4) - 1 - 1 - 1 - (2 + k)\right) e_{3}
$$
\n
$$
= \hat{\mathsf{Le}}_{n+2}^{(k)} - \hat{K}(n+2) + (k-1) - 2e_{1} + (-k-3) e_{2} + (-5 - 3k) e_{3}
$$

*ii.* Using Equations 5, 6, and 8,

$$
\sum_{s=0}^{n} \hat{L}e_{2s}^{(k)} = \sum_{s=0}^{n} L e_{2s}^{(k)} + \left(\sum_{s=0}^{n} L e_{2s+1}^{(k)}\right) e_1 + \left(\sum_{s=0}^{n} L e_{2s+2}^{(k)}\right) e_2 + \left(\sum_{s=0}^{n} L e_{2s+3}^{(k)}\right) e_3
$$
\n
$$
= \left(L e_{2n+1}^{(k)} - kn\right) + \left(L e_{2n+2}^{(k)} - k (n+2)\right) e_1 + \left(\sum_{s=0}^{n+1} L e_{2s}^{(k)} - L e_0^{(k)}\right) e_2 + \left(\sum_{s=0}^{n+1} L e_{2s+1}^{(k)} - L e_1^{(k)}\right) e_3
$$
\n
$$
= \left(L e_{2n+1}^{(k)} - kn\right) + \left(L e_{2n+2}^{(k)} - k (n+2)\right) e_1 + \left(L e_{2n+3}^{(k)} - k (n+1) - 1\right) e_2
$$
\n
$$
+ \left(L e_{2n+4}^{(k)} - k (n+3) - 1\right) e_3
$$
\n
$$
= \hat{L}e_{2n+1}^{(k)} - \hat{K}n + (-2k) e_1 + (-k-1) e_2 + (-3k-1) e_3
$$

 $\Box$ 

**Theorem 2.10** (The Honsberger-like identity). Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*-generalized Leonardo-like dual quaternion. For positive integers *n* and *m*,

$$
\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m}^{(k)} + \hat{\mathsf{Le}}_{n+1}^{(k)}\hat{\mathsf{Le}}_{m+1}^{(k)} = (k+1)\left(2\hat{\mathsf{Le}}_{n+m+2}^{(k)} + k\right) - k\,\hat{\mathsf{Le}}_{n+2}^{(k)} - k\,\hat{\mathsf{Le}}_{m+2}^{(k)} - (k+1)\,\mathsf{Le}_{n+m+2}^{(k)}
$$
\n
$$
- (e_1 + e_2 + e_3)\left(k\,\mathsf{Le}_{n+2}^{(k)} + k\,\mathsf{Le}_{m+2}^{(k)} - 2k\,(k+1)\right)
$$

where  $\mathsf{Le}_n^{(k)}$  is the *n*-th *k*−generalized Leonardo number.

Proof*.*

Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion and *n* and *m* be positive integers. Then,

$$
\begin{aligned} \hat{\mathsf{Le}}_{n}^{(k)} \hat{\mathsf{Le}}_{m}^{(k)} + \hat{\mathsf{Le}}_{n+1}^{(k)} \hat{\mathsf{Le}}_{m+1}^{(k)} &= \left( \mathsf{Le}_{n}^{(k)} \, \mathsf{Le}_{m}^{(k)} + \mathsf{Le}_{n+1}^{(k)} \, \mathsf{Le}_{m+1}^{(k)} \right) \\ &+ \left( \left( \mathsf{Le}_{n}^{(k)} \, \mathsf{Le}_{m+1}^{(k)} + \mathsf{Le}_{n+1}^{(k)} \, \mathsf{Le}_{m+2}^{(k)} \right) + \left( \mathsf{Le}_{n+1}^{(k)} \, \mathsf{Le}_{m}^{(k)} + \mathsf{Le}_{n+2}^{(k)} \, \mathsf{Le}_{m+1}^{(k)} \right) \right) e_{1} \\ &+ \left( \left( \mathsf{Le}_{n}^{(k)} \, \mathsf{Le}_{m+2}^{(k)} + \mathsf{Le}_{n+1}^{(k)} \, \mathsf{Le}_{m+3}^{(k)} \right) + \left( \mathsf{Le}_{n+2}^{(k)} \, \mathsf{Le}_{m}^{(k)} + \mathsf{Le}_{n+3}^{(k)} \, \mathsf{Le}_{m+1}^{(k)} \right) \right) e_{2} \\ &+ \left( \left( \mathsf{Le}_{n}^{(k)} \, \mathsf{Le}_{m+3}^{(k)} + \mathsf{Le}_{n+1}^{(k)} \, \mathsf{Le}_{m+4}^{(k)} \right) + \left( \mathsf{Le}_{n+3}^{(k)} \, \mathsf{Le}_{m}^{(k)} + \mathsf{Le}_{n+4}^{(k)} \, \mathsf{Le}_{m+1}^{(k)} \right) \right) e_{3} \end{aligned}
$$

We conclude from Equation 3 and  $F_nF_m + F_{n+1}F_{m+1} = F_{n+m+1}$  [1] that

$$
\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m}^{(k)} + \hat{\mathsf{Le}}_{n+1}^{(k)}\hat{\mathsf{Le}}_{m+1}^{(k)} = (k+1)^{2} \left(2\hat{\mathcal{F}}_{n+m+3} - F_{n+m+3}\right) - k(k+1) \left(F_{n+3} + F_{m+3}\right) \left(e_{1} + e_{2} + e_{3}\right)
$$

$$
-k(k+1) \left(\hat{\mathcal{F}}_{m+3} + \hat{\mathcal{F}}_{n+3}\right) + 4k^{2} \left(1 + e_{1} + e_{2} + e_{3}\right) - 2k^{2}
$$

Then, applying Theorem 2.5(*i*),

$$
\hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{m}^{(k)} + \hat{\mathsf{Le}}_{n+1}^{(k)}\hat{\mathsf{Le}}_{m+1}^{(k)} = (k+1)\left(2\hat{\mathsf{Le}}_{n+m+2}^{(k)} + k\right) - k\hat{\mathsf{Le}}_{n+2}^{(k)} - k\hat{\mathsf{Le}}_{m+2}^{(k)} - (k+1)\mathsf{Le}_{n+m+2}^{(k)} - (e_{1} + e_{2} + e_{3})\left(k\mathsf{Le}_{n+2}^{(k)} + k\mathsf{Le}_{m+2}^{(k)} - 2k(k+1)\right)
$$

is obtained.  $\quad \Box$ 

Across this study, let  $\alpha^* \beta^* = 1 + e_1 + 3e_2 + 4e_3$ .

**Theorem 2.11** (The Catalan-like identity). Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion. For positive integers *n* and *r* with  $n \geq r$ ,

$$
\left(\hat{\mathsf{Le}}_{n}^{(k)}\right)^{2} - \hat{\mathsf{Le}}_{n+r}^{(k)}\hat{\mathsf{Le}}_{n-r}^{(k)} = (k+1)^{2} \alpha^{*} \beta^{*} (-1)^{n-r+1} (F_{r})^{2} - \hat{\mathcal{K}} \left(2\hat{\mathsf{Le}}_{n}^{(k)} - \hat{\mathsf{Le}}_{n+r}^{(k)} - \hat{\mathsf{Le}}_{n-r}^{(k)}\right)
$$

where  $F_n$  is the *n*-th Fibonacci number.

Proof*.*

By using the Binet-like formula in Equation 9,

$$
\begin{split}\n\left(\hat{\mathsf{Le}}_{n}^{(k)}\right)^{2} - \hat{\mathsf{Le}}_{n+r}^{(k)}\hat{\mathsf{Le}}_{n-r}^{(k)} &= \left((k+1)\left(\frac{\alpha^{*}\alpha^{n+1}-\beta^{*}\beta^{n+1}}{\alpha-\beta}\right)-\hat{\mathcal{K}}\right)^{2} \\
&- \left((k+1)\left(\frac{\alpha^{*}\alpha^{n+r+1}-\beta^{*}\beta^{n+r+1}}{\alpha-\beta}\right)-\hat{\mathcal{K}}\right)\left((k+1)\left(\frac{\alpha^{*}\alpha^{n-r+1}-\beta^{*}\beta^{n-r+1}}{\alpha-\beta}\right)-\hat{\mathcal{K}}\right) \\
&= \frac{(k+1)^{2}}{5}\alpha^{*}\beta^{*}\left(\alpha^{n-r+1}\beta^{n-r+1}\right)\left(\alpha^{2r}+\beta^{2r}-2\alpha^{r}\beta^{r}\right) \\
&- \hat{\mathcal{K}}\left(2\left(\hat{\mathsf{Le}}_{n}^{(k)}+\hat{\mathcal{K}}\right)-\left(\hat{\mathsf{Le}}_{n+r}^{(k)}+\hat{\mathcal{K}}\right)-\left(\hat{\mathsf{Le}}_{n-r}^{(k)}+\hat{\mathcal{K}}\right)\right) \\
&= (k+1)^{2}\alpha^{*}\beta^{*}\left(-1\right)^{n-r+1}\left(F_{r}\right)^{2}-\hat{\mathcal{K}}\left(2\hat{\mathsf{Le}}_{n}^{(k)}-\hat{\mathsf{Le}}_{n+r}^{(k)}-\hat{\mathsf{Le}}_{n-r}^{(k)}\right)\n\end{split}
$$

 $\Box$ 

**Theorem 2.12** (The Cassini-like identity). Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion. For positive integer *n* with  $n \geq 3$ ,

$$
\hat{\mathsf{Le}}_{n-1}^{(k)} \hat{\mathsf{Le}}_{n+1}^{(k)} - \left(\hat{\mathsf{Le}}_{n}^{(k)}\right)^2 = (k+1)^2 \left(-1\right)^{n+1} \alpha^* \beta^* - \hat{\mathcal{K}} \hat{\mathsf{Le}}_{n-3}^{(k)} - \hat{\mathcal{K}}^2
$$

Proof*.*

From the Cassini identity and the recurrence relation of the Fibonacci dual quaternion sequence (see [20]),

$$
\hat{\mathsf{Le}}_{n-1}^{(k)} \hat{\mathsf{Le}}_{n+1}^{(k)} - (\hat{\mathsf{Le}}_{n}^{(k)})^2 = ((k+1)\hat{\mathcal{F}}_{n} - \hat{\mathcal{K}}) ((k+1)\hat{\mathcal{F}}_{n+2} - \hat{\mathcal{K}}) - ((k+1)\hat{\mathcal{F}}_{n+1} - \hat{\mathcal{K}})^2
$$
\n
$$
= ((k+1))^{2} \left( \hat{\mathcal{F}}_{n} \hat{\mathcal{F}}_{n+2} - (\hat{\mathcal{F}}_{n+1})^2 \right) - \hat{\mathcal{K}} (k+1) (\hat{\mathcal{F}}_{n} + \hat{\mathcal{F}}_{n+2} - 2\hat{\mathcal{F}}_{n+1})
$$
\n
$$
= (k+1)^2 (-1)^{n+1} (1 + e_1 + 3e_2 + 4e_3) - \hat{\mathcal{K}} (k+1) ((\hat{\mathcal{F}}_{n+2} - \hat{\mathcal{F}}_{n+1}) - (\hat{\mathcal{F}}_{n+1} - \hat{\mathcal{F}}_{n}))
$$
\n
$$
= (k+1)^2 (-1)^{n+1} (1 + e_1 + 3e_2 + 4e_3) - \hat{\mathcal{K}} ((k+1)\hat{\mathcal{F}}_{n-2} - \hat{\mathcal{K}}) - \hat{\mathcal{K}}^2
$$
\n
$$
= (k+1)^2 (-1)^{n+1} (1 + e_1 + 3e_2 + 4e_3) - \hat{\mathcal{K}} \hat{\mathsf{Le}}_{n-3}^{(k)} - \hat{\mathcal{K}}^2
$$

**Theorem 2.13** (The d'Ocagne-like identity). Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*−generalized Leonardo-like dual quaternion. For positive integers *m* and *n*,

$$
\hat{\mathsf{Le}}_{m}^{(k)} \hat{\mathsf{Le}}_{n+1}^{(k)} - \hat{\mathsf{Le}}_{m+1}^{(k)} \hat{\mathsf{Le}}_{n}^{(k)} = (k+1)^2 \alpha^* \beta^* (-1)^{n+1} F_{m-n} + \hat{\mathcal{K}} \left( \hat{\mathsf{Le}}_{m-1}^{(k)} - \hat{\mathsf{Le}}_{n-1}^{(k)} \right)
$$

where  $F_n$  is the *n*-th Fibonacci number.

Proof*.*

Applying the Binet-like formula in Equation 9,

$$
\hat{\mathsf{Le}}_{m}^{(k)} \hat{\mathsf{Le}}_{n+1}^{(k)} - \hat{\mathsf{Le}}_{m+1}^{(k)} \hat{\mathsf{Le}}_{n}^{(k)} = \left( (k+1) \left( \frac{\alpha^{*} \alpha^{m+1} - \beta^{*} \beta^{m+1}}{\alpha - \beta} \right) - \hat{\mathcal{K}} \right) \left( (k+1) \left( \frac{\alpha^{*} \alpha^{n+2} - \beta^{*} \beta^{n+2}}{\alpha - \beta} \right) - \hat{\mathcal{K}} \right)
$$
\n
$$
- \left( (k+1) \left( \frac{\alpha^{*} \alpha^{m+2} - \beta^{*} \beta^{m+2}}{\alpha - \beta} \right) - \hat{\mathcal{K}} \right) \left( (k+1) \left( \frac{\alpha^{*} \alpha^{n+1} - \beta^{*} \beta^{n+1}}{\alpha - \beta} \right) - \hat{\mathcal{K}} \right)
$$
\n
$$
= \frac{(k+1)^{2}}{\alpha - \beta} \alpha^{*} \beta^{*} \left( \alpha^{m+1} \beta^{n+1} - \alpha^{n+1} \beta^{m+1} \right)
$$
\n
$$
- \hat{\mathcal{K}} \left( \left( \hat{\mathsf{Le}}_{m}^{(k)} + \hat{\mathcal{K}} \right) + \left( \hat{\mathsf{Le}}_{n+1}^{(k)} + \hat{\mathcal{K}} \right) - \left( \hat{\mathsf{Le}}_{m+1}^{(k)} + \hat{\mathcal{K}} \right) - \left( \hat{\mathsf{Le}}_{n}^{(k)} + \hat{\mathcal{K}} \right) \right)
$$
\n
$$
= (k+1)^{2} \alpha^{*} \beta^{*} \alpha^{n+1} \beta^{n+1} \frac{(\alpha^{m-n} - \beta^{m-n})}{\alpha - \beta} + \hat{\mathcal{K}} \left( \left( \hat{\mathsf{Le}}_{m+1}^{(k)} - \hat{\mathsf{Le}}_{m}^{(k)} \right) - \left( \hat{\mathsf{Le}}_{n+1}^{(k)} - \hat{\mathsf{Le}}_{n}^{(k)} \right) \right)
$$
\n
$$
= (k+1)^{2} \alpha^{*} \beta^{*} (-1)^{n+1} F_{m-n} + \hat{\mathcal{K}} \left( \hat{\mathsf{Le}}_{m-1}^{
$$

 $\Box$ 

**Theorem 2.14** (The Tagiuri-like identity). Let  $\hat{\mathsf{Le}}_n^{(k)}$  be the *n*-th *k*-generalized Leonardo-like dual quaternion. For positive integers  $n, n + r$  and  $n + s$ ,

$$
\hat{\mathsf{Le}}_{n+r}^{(k)}\hat{\mathsf{Le}}_{n+s}^{(k)} - \hat{\mathsf{Le}}_{n}^{(k)}\hat{\mathsf{Le}}_{n+r+s}^{(k)} = \frac{(k+1)^2}{5}\alpha^*\beta^*(-1)^{n+1}(L_{r+s}-(-1)^sL_{r-s}) + \hat{\mathcal{K}}(\hat{\mathsf{Le}}_{n}^{(k)} + \hat{\mathsf{Le}}_{n+r+s}^{(k)} - \hat{\mathsf{Le}}_{n+r}^{(k)} - \hat{\mathsf{Le}}_{n+s}^{(k)})
$$

where  $L_n$  is the *n*-th Lucas number.

Proof*.*

The proof is straightforward from applying the Binet-like formula in Equation 9.  $\Box$ 

Note that the d'Ocagne-like, Catalan-like, and Cassini-like identities are the special cases of the Tagiuri-like identity.

#### **3. Conclusion**

Taking *k* = 1 gives the analogous relations for the Leonardo sequence with the dual-quaternions coefficients. Hence, we can say that our main results presented here generalize the paper [13]. These results can trigger further research on the subjects of the Leonardo sequence and the dual quaternions. Additionally, this study opens the door for future research on sequences; for instance, one may define non-commutative quaternions (real, split, semi-split, etc.) with the *k*−generalized Leonardo sequence.

#### **Author Contributions**

All the authors equally contributed to this work. They all read and approved the final version of the paper.

#### **Conflicts of Interest**

All the authors declare no conflict of interest.

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# **Inverse Problems for a Conformable Fractional Diffusion Operator**

#### Yaşar Çakmak<sup>1</sup> <sup>D</sup>

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Research Article

doi[:10.53570/jnt.1335702](https://doi.org/10.53570/jnt.1335702) mined by the specification of its zeros for each fixed *α*. The obtained results in this paper can **Abstract** − In this paper, we consider a diffusion operator with discrete boundary conditions, which include the conformable fractional derivatives of order  $\alpha$  such that  $0 < \alpha < 1$ instead of the ordinary derivatives in the classical diffusion operator. We prove that the coefficients of the given operator are uniquely determined by the Weyl function and spectral data, which consist of a spectrum and normalizing numbers. Moreover, using the well-known Hadamard's factorization theorem, we prove that the characteristic function  $\Delta_{\alpha}(\rho)$  is deterbe regarded as partial *α*-generalizations of similar findings obtained for the classical diffusion operator.

**Keywords** *Inverse problem, diffusion operator, conformable fractional derivative, Weyl function, spectral data* **Mathematics Subject Classification (2020)** 34A55, 26A33

#### **1. Introduction**

Inverse spectral problems aim to reconstruct the coefficients of an operator from given data such as the Weyl function, nodal points, and spectral data (two spectra or a spectrum and normalizing numbers). For the last century, these kinds of problems for various classical Sturm-Liouville, diffusion, and Dirac operators have been extensively investigated; for more details, see [1–7].

The beginning of the fractional derivative dates back to 1695, and many fractional derivative concepts have been proposed until today, such as the Riemann-Liouville fractional derivative, the Caputo fractional derivative, and the Atangana fractional derivative. In 2014, Khalil et al. [8] introduced the conformable fractional derivative. Then, many researchers identified important and fundamental properties of this derivative in [9–14]. In 2017, Jarad et al. [15] showed that this derivative is necessary and useful for generating new types of fractional operators. In recent years, numerous significant studies [16–20] have been conducted on inverse problems related to various conformable fractional operators, including the diffusion operator.

We consider a conformable fractional diffusion operator with discrete boundary conditions, denoted as  $L_{\alpha} = L_{\alpha}(p(x), q(x), h, H)$ . The form of this operator is as follows:

$$
\ell_{\alpha}y := -T_x^{\alpha}T_x^{\alpha}y + [2\rho p(x) + q(x)]y = \rho^2 y, \quad 0 < x < \pi
$$
\n
$$
U_{\alpha}(y) := T_x^{\alpha}y(0) - hy(0) = 0
$$
\n
$$
(1)
$$

1 ycakmak@cumhuriyet.edu.tr(Corresponding Author)

 $1$ Department of Mathematics, Faculty of Science, Sivas Cumhuriyet University, Sivas, Türkiye

and

$$
V_{\alpha}(y) := T_x^{\alpha} y(\pi) + Hy(\pi) = 0
$$

where  $\rho$  is the spectral parameter,  $h, H \in \mathbb{R}$ ,  $q(x) \in W_{2,\alpha}^1[0,\pi]$  and  $p(x) \in W_{2,\alpha}^2[0,\pi]$  are real-valued functions such that  $p(x) \ne const$ ,  $T_x^{\alpha}y$  is a conformable fractional derivative of order  $\alpha \in (0,1]$  of *y* at *x*,

$$
T_x^{\alpha}y(x) = \lim_{h \to 0} \frac{y(x + hx^{1-\alpha}) - y(x)}{h}, \quad \text{for all } x > 0
$$

 $W_{2,\alpha}^1[0,\pi] = \{f(x) \mid f(x) \text{ is absolutely continuous on } [0,\pi] \text{ and } T_x^{\alpha}f(x) \in L_{2,\alpha}(0,\pi)\}\$ 

 $W_{2,\alpha}^2[0,\pi] = \{f(x) \mid f(x) \text{ and } T_x^{\alpha}f(x) \text{ are absolutely continuous on } [0,\pi] \text{ and } T_x^{\alpha}T_x^{\alpha}f(x) \in L_{2,\alpha}(0,\pi)\}\$ 

and the space  $L_{2,\alpha}(0,\pi)$  consists of all the functions  $f:[0,\pi]\to\mathbb{R}$  satisfying the condition

$$
\left(\int_0^\pi |f(x)|^2 \, d\alpha x\right)^{1/2} = \left(\int_0^\pi |f(x)|^2 \, x^{\alpha - 1} \, dx\right)^{1/2} < \infty
$$

This operator is referred to as the Conformable Fractional Diffusion Operator (CFDO).

In this paper, we have proved that the coefficients of the given operator can be uniquely determined by the Weyl function and spectral data, which consist of a spectrum and normalizing numbers.

#### **2. Preliminaries**

This section provides some basic notions to be needed in the following sections. Let the functions  $\varphi = \varphi(x, \rho; \alpha)$ ,  $\psi = \psi(x, \rho; \alpha)$ , and  $S = S(x, \rho; \alpha)$  be the solutions of Equation 1 satisfying the following initial conditions

$$
\varphi(0, \rho; \alpha) = 1
$$
 and  $T_x^{\alpha} \varphi(0, \rho; \alpha) = h$  (2)

$$
\psi(\pi, \rho; \alpha) = 1 \quad \text{and} \quad T_x^{\alpha} \psi(\pi, \rho; \alpha) = -H \tag{3}
$$

$$
S(0, \rho; \alpha) = 0 \quad \text{and} \quad T_x^{\alpha} S(0, \rho; \alpha) = 1 \tag{4}
$$

respectively. From [19, 21, 22], these solutions satisfy the following asymptotic formulas, for  $|\rho| \to \infty$ and each fixed  $\alpha$ ,

$$
\varphi = \cos\left(\frac{\rho x^{\alpha}}{\alpha} - \theta(x)\right) + O\left(\frac{1}{|\rho|}\exp\left(\frac{|\text{Im}\rho| x^{\alpha}}{\alpha}\right)\right)
$$
(5)

$$
T_x^{\alpha} \varphi = -\rho \sin \left(\frac{\rho x^{\alpha}}{\alpha} - \theta(x)\right) + O\left(\exp\left(\frac{|\text{Im}\rho| \, x^{\alpha}}{\alpha}\right)\right) \tag{6}
$$

$$
\psi = \cos\left(\frac{\rho\left(\pi^{\alpha} - x^{\alpha}\right)}{\alpha} - \theta(\pi) + \theta(x)\right) + O\left(\frac{1}{|\rho|}\exp\left(\frac{|\text{Im}\rho|\left(\pi^{\alpha} - x^{\alpha}\right)}{\alpha}\right)\right) \tag{7}
$$

$$
T_x^{\alpha} \psi = \rho \sin \left( \frac{\rho \left( \pi^{\alpha} - x^{\alpha} \right)}{\alpha} - \theta(\pi) + \theta(x) \right) + O \left( \exp \left( \frac{|\text{Im}\rho| \left( \pi^{\alpha} - x^{\alpha} \right)}{\alpha} \right) \right)
$$
(8)

$$
S = \frac{1}{\rho} \sin\left(\frac{\rho x^{\alpha}}{\alpha} - \theta(x)\right) + O\left(\frac{1}{\rho^{2}} \exp\left(\frac{|\text{Im}\rho| x^{\alpha}}{\alpha}\right)\right)
$$
(9)

$$
T_x^{\alpha} S = \cos\left(\frac{\rho x^{\alpha}}{\alpha} - \theta(x)\right) + O\left(\frac{1}{\rho}\exp\left(\frac{|\text{Im}\rho| x^{\alpha}}{\alpha}\right)\right)
$$
(10)

where

*T*

$$
\theta(x) = \int\limits_0^x p(t) \, d_{\alpha}t
$$

We denote

$$
\Delta_{\alpha}(\rho) = W_{\alpha}[\psi, \varphi] = \begin{vmatrix} \psi & \varphi \\ T_x^{\alpha} \psi & T_x^{\alpha} \varphi \end{vmatrix} = \psi T_x^{\alpha} \varphi - \varphi T_x^{\alpha} \psi \tag{11}
$$

where  $W_\alpha[\psi, \varphi]$  is the fractional Wronskian of the functions  $\psi$  and  $\varphi$ . Furthermore, the  $\Delta_\alpha(\rho)$  is called as the characteristic function of the operator  $L_{\alpha}$  and is entire function in  $\rho$  for each fixed  $\alpha$ .

**Lemma 2.1.** [23] For each fixed  $\alpha$ ,  $\Delta_{\alpha}(\rho)$  does not depend on *x* and can be written as

$$
\Delta_{\alpha}(\rho) = V_{\alpha}(\varphi) = -U_{\alpha}(\psi)
$$
\n(12)

**Lemma 2.2.** [23] The zeros  $\{\rho_n\}$  of the function  $\Delta_{\alpha}(\rho)$  are coincide with the eigenvalues of the operator  $L_{\alpha}$ , and for eigenfunctions  $\psi(x, \rho_n; \alpha)$  and  $\varphi(x, \rho_n; \alpha)$ , there exists a sequence  $\{\beta_n\}$  such that

$$
\psi(x, \rho_n; \alpha) = \beta_n \varphi(x, \rho_n; \alpha) \quad \text{and} \quad \beta_n \neq 0 \tag{13}
$$

are satisfied for each fixed *α*.

It is clear from Equations 2, 3, and 13 that  $\beta_n = \psi(0, \rho_n; \alpha) = \frac{1}{\varphi(\pi, \rho_n; \alpha)}$ .

**Lemma 2.3.** [23] The equality  $\dot{\Delta}_{\alpha}(\rho_n) = -2\rho_n \beta_n \alpha_n$  is valid where  $\dot{\Delta}_{\alpha}(\rho) = \frac{d\Delta_{\alpha}(\rho)}{d\rho}$  and the normalizing numbers are

$$
\alpha_n = \int_0^{\pi} \varphi^2(x, \rho_n; \alpha) d\alpha x - \frac{1}{\rho_n} \int_0^{\pi} p(x) \varphi^2(x, \rho_n; \alpha) d\alpha x
$$

**Definition 2.4.** The data  $\{\rho_n, \alpha_n\}_{n\geq 1}$  are called the spectral data of the operator  $L_\alpha$ .

Let  $\{\rho_n\}$  be the eigenvalues set of the operator  $L_\alpha$ . From [23], the numbers  $\rho_n$  hold the following estimate:

$$
\rho_n = \frac{n\alpha}{\pi^{\alpha-1}} + c_{\alpha,0} + \frac{c_{\alpha,1}}{n} + o\left(\frac{1}{n}\right), \quad n \to \infty
$$

where

$$
c_{\alpha,0} = \frac{\alpha}{\pi^{\alpha}} \int_{0}^{\pi} p(x) d_{\alpha} x
$$

and

$$
c_{\alpha,1} = \frac{1}{\pi} \left[ h + H + \frac{1}{2} \int\limits_0^{\pi} \left( q(x) + p^2(x) \right) d_{\alpha} x \right]
$$

Let  $G_{\delta} = \left\{ \rho \mid \left| \rho - \frac{n\alpha}{\pi^{\alpha-1}} \right| \right\}$  $\leq \delta, n \in \{1, 2, ...\}$  where  $\delta$  is a sufficiently small positive number. It is obvious from Equations 5, 6, and 12 that the function  $\Delta_{\alpha}(\rho)$  satisfies the inequality

$$
|\Delta_{\alpha}(\rho)| \ge c_{\delta} |\rho| \exp\left(\frac{|\text{Im}\rho|}{\alpha}\pi^{\alpha}\right), \quad \rho \in G_{\delta}
$$
\n(14)

#### **3. Main Results**

This section proves uniqueness theorems for the solution of inverse problems according to the Weyl function and spectral data, which consist of a spectrum and normalizing numbers. Together with  $L_{\alpha}$ , we consider a second operator  $L_{\alpha} = L_{\alpha}(p(x), \tilde{q}(x), h, H)$  of the following form

$$
\widetilde{\ell}_{\alpha}y := -T_x^{\alpha}T_x^{\alpha}y + [2\rho p(x) + \widetilde{q}(x)]y = \rho^2 y, \quad 0 < x < \pi
$$
\n
$$
\widetilde{U}_{\alpha}(y) := T_x^{\alpha}y(0) - \widetilde{h}y(0) = 0
$$

and

$$
\widetilde{V}_{\alpha}(y) := T_x^{\alpha} y(\pi) + \widetilde{H} y(\pi) = 0
$$

We note that if a certain symbol  $\sigma$  denotes an object related to  $L_{\alpha}$ , then  $\tilde{\sigma}$  will denote an analogous object related to  $L_{\alpha}$ .

It can be observed that  $W_\alpha[\varphi, S]|_{x=0} = 1 \neq 0$ . Thus, the functions  $\varphi$  and *S* are linearly independent, and the function  $\psi$  can be written as

$$
\psi(x,\rho;\alpha) = c_1(\rho;\alpha)\varphi(x,\rho;\alpha) + c_2(\rho;\alpha) S(x,\rho;\alpha)
$$
\n(15)

where  $c_1(\rho; \alpha)$  and  $c_2(\rho; \alpha)$  are arbitrary constant for each fixed  $\alpha$ . It is clear from Equation 15 that

$$
\psi(0,\rho;\alpha) = c_1(\rho;\alpha)\,\varphi(0,\rho;\alpha) + c_2(\rho;\alpha)\,S(0,\rho;\alpha)
$$

and

$$
T_x^{\alpha} \psi (0, \rho; \alpha) = c_1 (\rho; \alpha) T_x^{\alpha} \varphi (0, \rho; \alpha) + c_2 (\rho; \alpha) T_x^{\alpha} S (0, \rho; \alpha)
$$

From Equations 2, 4, and 12,

$$
c_1(\rho;\alpha) = \psi(0,\rho;\alpha)
$$

and

$$
c_2(\rho;\alpha) = T_x^{\alpha} \psi(0,\rho;\alpha) - h\psi(0,\rho;\alpha) = -\Delta_{\alpha}(\rho)
$$

Consequently, Equation 15 is rewritten as

$$
\psi(x, \rho; \alpha) = \psi(0, \rho; \alpha) \varphi(x, \rho; \alpha) - \Delta_{\alpha}(\rho) S(x, \rho; \alpha)
$$

or

$$
-\frac{\psi(x,\rho;\alpha)}{\Delta_{\alpha}(\rho)} = -\frac{\psi(0,\rho;\alpha)}{\Delta_{\alpha}(\rho)}\varphi(x,\rho;\alpha) + S(x,\rho;\alpha)
$$
\n(16)

If we denote

$$
\Phi(x,\rho;\alpha) := -\frac{\psi(x,\rho;\alpha)}{\Delta_{\alpha}(\rho)} \quad \text{and} \quad M_{\alpha}(\rho) := \Phi(0,\rho;\alpha) = -\frac{\psi(0,\rho;\alpha)}{\Delta_{\alpha}(\rho)} \tag{17}
$$

then, from Equation 16,

$$
\Phi(x,\rho;\alpha) = S(x,\rho;\alpha) + M_{\alpha}(\rho) \varphi(x,\rho;\alpha)
$$
\n(18)

The functions  $\Phi(x, \rho; \alpha)$  and  $M_\alpha(\rho)$  are called as the Weyl solution and the Weyl function of the operator  $L_{\alpha}$ , respectively. It is obvious that  $\Phi(x, \rho; \alpha)$  is the solution of Equation 1 under the conditions  $U_{\alpha}(\Phi) = 1$ ,  $V_{\alpha}(\Phi) = 0$ , and  $M_{\alpha}(\rho)$  is a meromorphic function with poles in  $\{\rho_n\}$ .

**Theorem 3.1.** If  $M_{\alpha}(\rho) = \widetilde{M}_{\alpha}(\rho)$  for each fixed  $\alpha$ , then  $q(x) = \widetilde{q}(x)$ , almost everywhere in [0,  $\pi$ ],  $h = \tilde{h}$ , and  $H = \tilde{H}$ . Thus, the Weyl function uniquely determines the operator  $L_{\alpha}$ .

#### Proof*.*

Consider the functions  $P_1(x, \rho; \alpha)$  and  $P_2(x, \rho; \alpha)$  defined by

$$
P_1(x,\rho;\alpha) = \varphi(x,\rho;\alpha) T_x^{\alpha} \tilde{\Phi}(x,\rho;\alpha) - \Phi(x,\rho;\alpha) T_x^{\alpha} \tilde{\varphi}(x,\rho;\alpha)
$$
(19)

and

$$
P_2(x, \rho; \alpha) = \Phi(x, \rho; \alpha) \tilde{\varphi}(x, \rho; \alpha) - \varphi(x, \rho; \alpha) \tilde{\Phi}(x, \rho; \alpha)
$$
\n(20)

From Equation 18,

$$
P_1(x, \rho; \alpha) = \varphi(x, \rho; \alpha) T_x^{\alpha} \tilde{S}(x, \rho; \alpha) - S(x, \rho; \alpha) T_x^{\alpha} \tilde{\varphi}(x, \rho; \alpha) + \left[ \widetilde{M}_{\alpha}(\rho) - M_{\alpha}(\rho) \right] \varphi(x, \rho; \alpha) T_x^{\alpha} \tilde{\varphi}(x, \rho; \alpha)
$$

and

$$
P_2(x, \rho; \alpha) = S(x, \rho; \alpha) \tilde{\varphi}(x, \rho; \alpha) - \varphi(x, \rho; \alpha) \tilde{S}(x, \rho; \alpha) + \left[ M_\alpha(\rho) - \widetilde{M}_\alpha(\rho) \right] \varphi(x, \rho; \alpha) \tilde{\varphi}(x, \rho; \alpha)
$$

Since  $M_{\alpha}(\rho) = \widetilde{M}_{\alpha}(\rho)$ , the functions  $P_1(x, \rho; \alpha)$  and  $P_2(x, \rho; \alpha)$  are entire in  $\rho$ , for each fixed  $\alpha$ . Moreover, from Equations 11 and 17,

$$
W_{\alpha}\left[\varphi\left(x,\rho;\alpha\right),\Phi\left(x,\rho;\alpha\right)\right]=-\frac{W_{\alpha}\left[\varphi\left(x,\rho;\alpha\right),\psi\left(x,\rho;\alpha\right)\right]}{\Delta_{\alpha}\left(\rho\right)}=1
$$

and similarly,

$$
W_{\alpha}\left[\widetilde{\varphi}\left(x,\rho;\alpha\right),\widetilde{\Phi}\left(x,\rho;\alpha\right)\right]=1
$$

Thus, Equation 19 can be rewritten as

$$
P_1(x, \rho; \alpha) = 1 + \varphi(x, \rho; \alpha) \left[ T_x^{\alpha} \widetilde{\Phi}(x, \rho; \alpha) - T_x^{\alpha} \Phi(x, \rho; \alpha) \right] + \Phi(x, \rho; \alpha) \left[ T_x^{\alpha} \varphi(x, \rho; \alpha) - T_x^{\alpha} \widetilde{\varphi}(x, \rho; \alpha) \right]
$$

It follows from the asymptotic formulas of Equations 5-10 and Equality 14 that

$$
|P_1(x, \rho; \alpha) - 1| \leq \frac{C_{\delta}}{|\rho|}
$$
 and  $|P_2(x, \rho; \alpha)| \leq \frac{C_{\delta}}{|\rho|}$ ,  $x \in [0, \pi]$ ,  $|\rho| \in G_{\delta}$ 

Therefore, since  $\lim_{|\rho|\to\infty} |P_1(x,\rho;\alpha)-1| = \lim_{|\rho|\to\infty} |P_2(x,\rho;\alpha)| = 0$  by the well-known Liouville's theorem, we obtain for  $x \in [0, \pi]$  and each fixed  $\alpha$  that

$$
P_1(x, \rho; \alpha) = 1
$$
 and  $P_2(x, \rho; \alpha) = 0$  (21)

Hence, by using Equations 19-21, we get the following system

$$
\begin{cases}\n\varphi(x,\rho;\alpha) T_x^{\alpha} \tilde{\Phi}(x,\rho;\alpha) - \Phi(x,\rho;\alpha) T_x^{\alpha} \tilde{\varphi}(x,\rho;\alpha) = 1 \\
\Phi(x,\rho;\alpha) \tilde{\varphi}(x,\rho;\alpha) - \varphi(x,\rho;\alpha) \tilde{\Phi}(x,\rho;\alpha) = 0\n\end{cases}
$$
\n(22)

If System 22 is solved according to functions  $\varphi(x, \rho; \alpha)$  and  $\Phi(x, \rho; \alpha)$ , then

$$
\varphi(x,\rho;\alpha) = \widetilde{\varphi}(x,\rho;\alpha)
$$

and

$$
\Phi(x,\rho;\alpha) = \widetilde{\Phi}(x,\rho;\alpha)
$$

is obtained, for all *x* and *ρ* and each fixed  $\alpha$ . Thus,  $q(x) = \tilde{q}(x)$ , almost everywhere in [0,  $\pi$ ],  $h = h$ , and  $H = H$ . Consequently,  $L_{\alpha} = L_{\alpha}$ .

**Lemma 3.2.** For each fixed *α*, the characteristic function  $\Delta_{\alpha}(\rho)$  is determined by the specification of its zeros as:

$$
\Delta_{\alpha}(\rho) = C\rho \exp(C_1\rho) \prod_{n=1}^{\infty} \left(1 - \frac{\rho}{\rho_n}\right) \exp\left(\frac{\rho}{\rho_n}\right)
$$

where

$$
C = \sin \theta \left(\pi\right) \prod_{n=1}^{\infty} \frac{\rho}{\rho_n^0}, \quad C_1 = -\frac{\pi^{\alpha}}{\alpha} \cot \theta \left(\pi\right) + \sum_{n=1}^{\infty} \left(\frac{1}{\rho_n^0} - \frac{1}{\rho_n}\right), \quad \rho_n^0 = \left(n + \frac{\theta\left(\pi\right)}{\pi}\right) \frac{\alpha}{\pi^{\alpha-1}}, \quad n \in \{1, 2, \dots\}
$$

Proof*.*

It is clear from Equations 5, 6, and 12 that the characteristic function  $\Delta_{\alpha}(\rho)$  holds the following asymptotic representation:

$$
\Delta_{\alpha}(\rho) = -\rho \sin\left(\frac{\rho \pi^{\alpha}}{\alpha} - \theta(\pi)\right) + O\left(\exp\left(\frac{|\text{Im}\rho|}{\alpha}\pi^{\alpha}\right)\right)
$$
\n(23)

Consider the function

$$
\Delta_{\alpha}^{0}(\rho) = -\rho \sin\left(\frac{\rho \pi^{\alpha}}{\alpha} - \theta(\pi)\right)
$$
\n(24)

The zeros of the function  $\Delta_{\alpha}^{0}(\rho)$  are  $\rho = 0$  and  $\rho_{n}^{0} = \left(n + \frac{\theta(\pi)}{\pi}\right)$  $\left(\frac{\pi}{\pi}\right)$   $\frac{\alpha}{\pi^{\alpha-1}}$  such that  $n \in \{1, 2, 3, \cdots\}$ *.* Since  $\Delta^0_\alpha(\rho)$  is an entire function, according to the Hadamard's factorization theorem,

$$
\Delta_{\alpha}^{0}(\rho) = -\rho^{m} \exp\left(g(\rho)\right) \prod_{n=1}^{\infty} E_{p}\left(\frac{\rho}{\rho_{n}^{0}}\right)
$$
\n(25)

where  $m \geq 0$  is the multiplicity of the zero eigenvalue,  $g(\rho)$  is a polynomial with der  $(g(\rho)) = p$ , and

$$
E_p(\xi) = \begin{cases} (1-\xi), & n = 0\\ (1-\xi)\exp\left(\frac{\xi}{1} + \frac{\xi^2}{2} + \ldots + \frac{\xi^n}{n}\right), \text{ otherwise} \end{cases}
$$

Since the multiplicity of the zero is 1,  $m = 1$ . Besides, for every  $r > 0$  and for  $p = 1$ , the series  $\sum_{i=1}^{\infty}$ *n*=1 *r* 1+*p*  $\frac{r^{1+p}}{|\rho_n^0|^{1+p}}$  converges. Therefore, Equation 25 can rewrite as

$$
\Delta_{\alpha}^{0}(\rho) = -\rho \exp (a\rho + b) \prod_{n=1}^{\infty} \left( 1 - \frac{\rho}{\rho_n^{0}} \right) \exp \left( \frac{\rho}{\rho_n^{0}} \right)
$$

If we consider the following equalities to find the constants *a* and *b*,

$$
\lim_{\rho \to 0} \sin \left( \frac{\rho \pi^{\alpha}}{\alpha} - \theta(\pi) \right) = \lim_{\rho \to 0} \exp (a\rho + b) \prod_{n=1}^{\infty} \left( 1 - \frac{\rho}{\rho_n^0} \right) \exp \left( \frac{\rho}{\rho_n^0} \right)
$$

and

$$
\lim_{\rho \to 0} \frac{d}{d\rho} \ln \left[ \sin \left( \frac{\rho \pi^{\alpha}}{\alpha} - \theta(\pi) \right) \right] = \lim_{\rho \to 0} \frac{d}{d\rho} \ln \left[ C^{0} \exp (a\rho) \right] \prod_{n=1}^{\infty} \left( 1 - \frac{\rho}{\rho_n^0} \right) \exp \left( \frac{\rho}{\rho_n^0} \right) \right]
$$

then

$$
C^0 = \exp{(b)} = -\sin{\theta(\pi)}
$$

and

$$
C_1^0 = a = -\frac{\pi^{\alpha}}{\alpha} \cot \theta(\pi)
$$

respectively. Thus,

$$
\Delta_{\alpha}^{0}(\rho) = -\rho C^{0} \exp\left(C_{1}^{0}\rho\right) \prod_{n=1}^{\infty} \left(1 - \frac{\rho}{\rho_{n}^{0}}\right) \exp\left(\frac{\rho}{\rho_{n}^{0}}\right)
$$
\n(26)

Moreover,

$$
\Delta_{\alpha}(\rho) = C \exp(C_1 \rho) \rho^m \prod_{n=1}^{\infty} \left(1 - \frac{\rho}{\rho_n}\right) \exp\left(\frac{\rho}{\rho_n}\right)
$$
\n(27)

where *C* and  $C_1$  are constants and  $m \geq 0$ . According to Equations 23 and 24,

$$
\frac{\Delta_{\alpha}(\rho)}{\Delta_{\alpha}^{0}(\rho)} = 1 + O\left(\frac{1}{\rho}\right), \quad |\rho| \to \infty
$$

Then, together with Equations 26 and 27,

$$
\frac{\Delta_{\alpha}(\rho)}{\Delta_{\alpha}^{0}(\rho)} = -\frac{C}{C^{0}}\rho^{m-1}\prod_{n=1}^{\infty}\frac{\rho_{n}^{0}}{\rho_{n}}\prod_{n=1}^{\infty}\left(1 + \frac{\rho_{n}-\rho_{n}^{0}-\rho}{\rho_{n}^{0}-\rho}\right)\exp\left(\sum_{n=1}^{\infty}\frac{\rho_{n}^{0}-\rho_{n}}{\rho_{n}\rho_{n}^{0}} + C_{1} - C_{1}^{0}\right)\rho
$$

Consequently,

$$
m = 1
$$
,  $C = -C^0 \prod_{n=1}^{\infty} \frac{\rho_n}{\rho_n^0}$ , and  $C_1 = C_1^0 + \sum_{n=1}^{\infty} \left( \frac{1}{\rho_n^0} - \frac{1}{\rho_n} \right)$ 

 $\Box$ 

**Theorem 3.3.** If  $\{\rho_n, \alpha_n\}_{n\geq 1} = \{\tilde{\rho}_n, \tilde{\alpha}_n\}_{n\geq 1}$  for each fixed  $\alpha$ , then  $q(x) = \tilde{q}(x)$ , almost everywhere in  $[0, \pi]$ ,  $h = h$ , and  $H = H$ . Thus, the spectral data  $\{\rho_n, \alpha_n\}_{n \geq 1}$  uniquely determines the operator  $L_{\alpha}$ .

#### Proof*.*

Since  $\rho_n = \tilde{\rho}_n$ , according to Lemma 3.2,  $\Delta_{\alpha}(\rho) = \Delta_{\alpha}(\rho)$ . Using Lemma 2.3 and  $\alpha_n = \tilde{\alpha}_n$ ,  $\beta_n = \beta_n$ and thus  $\psi(0, \rho_n; \alpha) = \tilde{\psi}(0, \rho_n; \alpha)$ . For each fixed  $\alpha$ , let

$$
H_{\alpha}(\rho) := \frac{\psi(0,\rho;\alpha) - \widetilde{\psi}(0,\rho;\alpha)}{\Delta_{\alpha}(\rho)}
$$

It is clear that  $H_{\alpha}(\rho)$  is entire on  $\rho$ . Moreover, by using Equations 7 and 14,

$$
H_{\alpha}(\rho) := O\left(\frac{1}{\rho^2}\right), \quad |\rho| \to \infty
$$

Hence,  $H_{\alpha}(\rho) \equiv 0$  and  $\psi(0, \rho; \alpha) = \tilde{\psi}(0, \rho; \alpha)$ . Consequently, from Equation 17,  $M_{\alpha}(\rho) \equiv \widetilde{M}_{\alpha}(\rho)$ . Thus, the proof is completed by Theorem 3.1.  $\Box$ 

#### **4. Conclusion**

The Weyl function and spectral data are very natural and useful spectral characteristics in inverse problem theory. Until today, by using these concepts, many inverse problems have been studied for various classes of operators, such as regular or singular Sturm-Liouville operators, diffusion operators, and Dirac operators, including the classical derivatives. In [18], some inverse problems for the Sturm-Liouville operator, including conformable fractional derivatives, are investigated.

In this study, the diffusion operator, which includes conformable fractional derivatives, is considered, and the inverse problems are investigated for this operator for the first time according to both the Weyl function and spectral data. This study can be considered as a partial *α*-generalization of similar findings for the classical diffusion operator.

Considering this study's results, some inverse problems can be investigated in the future for various conformable operators with jump conditions, parameter-dependent boundary conditions, or non-local boundary conditions.

#### **Author Contributions**

The author read and approved the final version of the paper.

#### **Conflicts of Interest**

The author declares no conflict of interest.

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# New Theory

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# **Results of Paired-Domination of Some Special Graph Families on Transformation Graphs:**  $G^{xy+}$  **and**  $G^{xy-}$

**Hande Tunçel Gölpek<sup>1</sup>**

**Article Info** *Received*: 4 Aug 2023 *Accepted*: 22 Sep 2023 *Published*: 30 Sep 2023 doi[:10.53570/jnt.1337633](https://doi.org/10.53570/jnt.1337633) Research Article **Abstract** − In this study, transformation graphs obtained from the concept of the total graph and the result of its paired domination number for some special graph families are discussed. If a subset  $S$  of the vertex set of the graph  $G$  dominates and the induced subgraph  $\langle S \rangle$  has a perfect matching that covers every vertex of the graph, then  $S$  is called a paired-dominating set of  $G$ . A paired dominating set with the smallest cardinality is denoted by  $\gamma_{pr}$ -set. Haynes and Slater introduced paired domination parameters. The present study commences with assessing outcomes stemming from eight permutations within the realm of path graphs. Subsequently, building upon this foundational structure, the results are extrapolated from the realm of cycle transformation graph structures based on findings from path transformation graphs.

**Keywords** *Graph theory, graph vulnerability, paired domination number, transformation graphs*

**Mathematics Subject Classification (2020)** 05C69, 05C76

#### **1. Introduction**

In graph theory, the investigation of domination numbers and their various forms has garnered significant attention, leading to a continual expansion of knowledge in this domain. Let  $G$  be a graph with the vertex set  $V(G)$  (or simply V) and the edge set  $E(G)$  (or simply E). Then,  $S \subseteq V$  is called the dominating set of G if every vertex in  $V - S$  is adjacent to any vertex in S, and the domination number of G is denoted by  $\gamma(G)$ . The domination number is the minimum cardinality of a dominating set.

Furthermore, the discourse extends to exploring the total domination number, a parameter that has drawn extensive discussion. A total dominating set S of G, denoted by  $TD$ -set, is a dominating set with no isolated vertex and total domination number, which is the minimum cardinality of a total dominating set abbreviated by  $\gamma_t(G)$ . Total domination was introduced by Cockayne et al. [1].

This paper delves into another alteration of the domination parameter, termed paired domination. Let  $S \subseteq V$ . Then,  $S$  is called a paired-dominating set of  $G$  and is denoted  $PD$ -set, if it dominates, and the induced subgraph ⟨⟩ has a perfect matching that covers every vertex of the graph. A paired dominating set with the smallest cardinality is denoted by  $\gamma_{pr}$ -set. This concept was pioneered by Haynes and Slater [2]. To explain the meaning of these three domination parameters in terms of application, consider each  $s \in S$  as a protector that can guard each vertex where s dominates then every protector guards itself. Regarding total domination, every protector must be guarded by another protector. As for paired domination, the guards' positions must be chosen as pairs of adjacent vertices so that each guard is assigned to each other and exists as substitutes for each other. There **1DEW Theory** (Theory ( $\mathbf{F}$  Marine Control of Nome Special Graph Faculty, *Down Theory Theory Theory Theory Theory Theory Theory Theory Chemical Section (* $\Phi$  *Marine Chemical Section 2014) (* $\Phi$  *Marine Chemical Sectio* 

<sup>1</sup>hande.tuncel@deu.edu.tr (Corresponding Author)

are many domination parameters in different categories, and the paired domination parameters discussed in this paper are considered under the category with restricted forms based on  $\langle S \rangle$ . We can present the reader with references to some significant studies, particularly in this work, that address the behavior of the paired domination number under various graph operations, as [3-11] in the literature. In addition, regarding the domination parameter, the three new books, which provide even more comprehensive coverage of current studies on domination parameters, can be offered to the reader as references [12-14].

In this paper, we consider simple graphs and refer to the book [15] for notation and terminology. The open neighborhood of any vertex in  $V(G)$ , denoted by  $N(v) = \{u \in V(G) | (uv) \in E(G)\}\)$ . The degree of a vertex v is the number of edges that are incident to it and is denoted by  $deg(v)$ . Moreover, a vertex whose degree is one is called a pendant vertex. Furthermore, the maximum degree of G, denoted by  $\Delta(G)$ , is the degree of the vertex with the maximum number of edges incident to it. Similarly, the minimum degree of  $G$ , abbreviated by  $\delta(G)$ , is the degree of a vertex with the minimum number of edges incident to it. The distance between two vertices u and v is the shortest path between them and is denoted by the notation  $d(u, v)$ .

In this study, transformation graphs  $G^{xy-}$  and  $G^{xy+}$  are considered, and values of their paired-domination number for some special graph families are obtained. Transformation graph definitions are bottomed on the total graph [16]. Let  $V(G) \cup E(G)$  be the vertex set of a total graph where two vertices are adjacent if and only if the vertices and edges are adjacent or incident in  $G$ . The transformation graphs introduced by Wu and Meng can be provided, denoted by  $G^{xyz}$ . Their paper considers its eight structure cases and generalizations of the total graph concept [17]. Before defining transformation graphs, we must explain some fundamental graph constructions such as complement and line graphs. The complement graph of G, denoted by the notation  $\overline{G}$ , is based on the same vertex set  $V(G)$  but a different edge set completes the graph G to complement the graph. The line graph  $L(G)$  of G is the graph that takes the edge set  $E(G)$  as its vertex set. Two vertices are adjacent in the line graph  $L(G)$  if and only if corresponding edges are adjacent in graph G [18]. Thus, the concept of the transformation graph represented by  $G^{xyz}$  can now be addressed, where x, y, and z can take values + or -. The transformation graph  $G^{xyz}$  of G is a simple graph having as the vertex set  $V(G) \cup E(G)$ , and for  $\alpha, \beta \in V(G)$   $\cup$  $E(G)$  and their relationships in  $G^{xyz}$  can be explained as follows:

*i*. Let  $\alpha, \beta \in V(G)$ . Then,  $\alpha$  and  $\beta$  are adjacent in  $G^{xyz}$  if  $x = +$ ; otherwise  $x = -$ .

*ii.* Let  $\alpha, \beta \in E(G)$ . Then,  $\alpha$  and  $\beta$  are adjacent in  $G^{xyz}$  if  $y = +$ ; otherwise  $y = -$ .

*iii.* Let  $\alpha \in V(G)$  and  $\beta \in E(G)$ . Then,  $\alpha$  and  $\beta$  are incidents in  $G^{xyz}$  if  $z = +$ ; otherwise  $z = -$ .

There are eight distinct constructions from three signature permutations {+, −}. Then, it may be obtained transformation graphs where  $G^{+++}$  is the total graph of G, and  $G^{---}$  is the complement of it. Moreover,  $G^{-+-}$ ,  $G^{--+}$ , and  $G^{-++}$  are the complements of  $G^{+-+}$ ,  $G^{++-}$ , and  $G^{+--}$ , respectively. Moreover, the transformation graph can be expressed as a relationship between G (or  $\overline{G}$ ) and its line graph  $L(G)$  (or  $\overline{L(G)}$ ) depending on the sign of x, y, and z. If  $x = +$ , then the transformation graph includes G as a subgraph, otherwise includes  $\overline{G}$  as a subgraph, and if  $y = +$ , then the transformation graph includes  $L(G)$  as a subgraph, otherwise includes  $L(G)$  as a subgraph alike. Thus, depending on the sign of z, the relationship between G (or  $\overline{G}$ ) and  $L(G)$  (or  $\overline{L(G)}$ ) has been structured based on the aforesaid situation *(iii)*. Let  $V(G) \cup V(L(G))$  be the vertex set of the transformation graph where  $V(G) = \{1, 2, \dots, n\}$  and  $V(L(G)) = \{(ij) : (ij) \in E(G)\}$ . The vertex set will be expressed in this form unless otherwise stated. We direct readers' attention to select papers that offer comprehensive insights into the outcomes concerning specific domination parameters in transformation graphs. These referenced works provide detailed and in-depth information that can facilitate a deeper understanding of the subject matter of transformation graphs [19- 26].

#### **2. Preliminaries**

This section presents pivotal outcomes from scholarly literature concerning the paired domination number.

**Proposition 2.1.** [2] If  $u$  is adjacent to an end vertex of  $G$ , then  $u$  is in every paired dominating set. Furthermore, for certain specific families of graphs,  $\gamma_{pr}(K_n) = \gamma_{pr}(K_{r,s}) = \gamma_{pr}(W_n) = 2$  and  $\gamma_{pr}(P_n) =$  $\gamma_{pr}(C_n) = 2\left[\frac{n}{4}\right]$  $\frac{n}{4}$ .

Here, the notation  $[$ .  $]$  represents the ceiling function.

**Theorem 2.2.** [2] For any graph G,  $\gamma_{pr}(\overline{G}) \geq 4$  if and only if diam(G) = 2.

**Theorem 2.3.** [2] If a graph G has no isolated vertices, then  $2 \leq \gamma_{pr}(G) \leq n$ .

**Proposition 2.4.** [2] If a graph G has no isolated vertices, then  $\gamma(G) \leq \gamma_t(G) \leq \gamma_{pr}(G)$ , and  $\gamma_{pr}(G)$  is even.

#### **3. Results**

This section presents the results of the eight different permutations of path and cycle transformation graphs for the paired domination parameter.

#### **3.1. Paired Domination of Transformation Path Graphs**

In this designated section, we intend to acquire and elucidate the outcomes of the paired domination of transformation graphs across all permutations. We shall commence our exposition by focusing on the specific  $P_n^{+++}$  form of permutations.

**Theorem 3.1.** Let  $n > 6$  and  $P_n$  be a path graph. Then,

$$
\gamma_{pr}(P_n^{+++}) = \begin{cases} \left[\frac{4n-2}{7}\right], & n \equiv 0,3,4,6 \pmod{7} \\ \left[\frac{4n-2}{7}\right] + 1, & \text{otherwise} \end{cases}
$$

PROOF.

The vertices of the graph  $P_n^{+++}$  can be partitioned into two parts, such as  $V = V_1 \cup V_2$ . To alleviate the burden of complex notation, the vertex sets can be labeled as  $V_1 = \{1, 2, \dots, n\}$  and  $V_2 = \{n + 1, n + 2, \dots, 2n\}$ . Let D be a PD-set of  $P_n^{+++}$ . To begin, we establish an upper bound for  $\gamma_{pr}(P_n^{+++})$  as follows: Let  $S = S_1 \cup S_2$  be a set where  $S_1$  includes the vertices of  $V(P_n^{+++})$ , and  $S_2$  includes the vertices of  $V(L(P_n^{+++}))$  as follows:

$$
S_1 = \bigcup_{i=0}^{\left\lceil \frac{n-2}{7} \right\rceil -1} \{ (7i+2,7i+3) \}
$$

and

$$
S_2 = \bigcup_{i=0}^{\left[\frac{n-6}{7}\right]-1} \{ (n+7i+5, n+7i+6) \}
$$

Here, the cardinality of S is  $|S| = 2 \left| \frac{n-2}{7} \right|$  $\left[\frac{n-2}{7}\right]$  + 2  $\left[\frac{n-6}{7}\right]$  $\frac{-6}{7}$ . *i*. If  $n \equiv 0.3.4 \pmod{7}$ , then  $D = S$ . *ii.* If  $n \equiv 1.2 \pmod{7}$ , then  $D = S \cup \{n, n - 1\}$ . *iii.* If  $n \equiv 5.6 \pmod{7}$ , then  $D = S \cup \{2n - 2, 2n - 1\}$ . Therefore,  $|D| = \frac{4n-2}{7}$  $\left| \frac{a-2}{7} \right|$ , for  $n \equiv 0.3.4 \pmod{7}$ , otherwise  $|D| = \left| \frac{4n-2}{7} \right|$  $\left| \frac{i-2}{7} \right|$  + 1. Thus,

$$
\gamma_{pr}(P_n^{+++}) \le \begin{cases} \left\lceil \frac{4n-2}{7} \right\rceil, & n \equiv 0,3,4 \pmod{7} \\ \left\lceil \frac{4n-2}{7} \right\rceil + 1, & \text{otherwise} \end{cases}
$$

In order to prove the necessary condition, assume that the set  $T = \{a_1, a_2, \dots, a_t\}$  is  $\gamma_{pr}$ -set of  $P_n^{+++}$ . Assume that the elements of T are ordered as  $a_1 < a_2 < \cdots < a_i < \cdots < a_m < a_{m+1} < \cdots < a_j < \cdots < a_t$ . Here,  $a_i$ and  $a_i$  are positive integers,  $1 \le a_i \le n$ , for all  $i \in \{1, \dots, m\}$ , and  $n + 1 \le a_i \le 2n - 1$ , for all  $j \in \{m + 1, m\}$ 1, …, t}. Define a function such as  $f_x = a_{x+2} - a_x$  where  $x \in \{1,2,\dots, t-2\}$  such that  $x \neq m$  and  $x \neq m$ 1. It is needed to prove that  $f_x \le 7$ . Assume that a value of  $f_x$  is greater than or equal to 8. Without loss of generality, choose  $f_1 = 8$ . Therefore, for  $n \ge 14$ , we can create a set

$$
S' = \{2,3,n+5,n+6,n+7,n+8\} \cup \left( \bigcup_{i=0}^{\left[\frac{n-10}{7}\right]-1} \{7i+10,7i+11\} \right) \cup \left( \bigcup_{i=0}^{\left[\frac{n-10}{7}\right]-1} \{n+7i+13,n+7i+14\} \right)
$$

For  $n \equiv 3 \pmod{7}$ ,  $T = S' \cup \{n, n - 1\}$ . Thus,

$$
|T| = |S'| + 2 = 8 + 2\left(\frac{n-10}{7}\right) + 2\left(\frac{n-10}{7}\right) = \frac{4n+16}{7}
$$

However, it contradicts the obtained upper bound  $\frac{4n-2}{7}$  $\frac{z-2}{7}$ . Hence,  $f_x$  should be less than or equal to 7. Then, the upper bound for the total value of  $f_x$  can be gotten as follows:

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} \le 7(t-4)
$$

In order to construct a relationship between the values of t and n, the total value of  $f_x$  can be evaluated as

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = a_m + a_{m-1} + a_t + a_{t-1} - (2n + 16)
$$

where  $a_1 = 2$ ,  $a_2 = 3$ ,  $a_{m+1} = n + 5$ , and  $a_{m+2} = n + 6$ . If  $n \equiv 0 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 28 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n}{7}$ where  $a_m = n - 4$ ,  $a_{m-1} = n - 5$ ,  $a_t = 2n - 1$ , and  $a_{t-1} = 2n - 2$ . Therefore,

$$
|T| = t \ge \frac{4n}{7} = \left\lceil \frac{4n - 2}{7} \right\rceil
$$

Thus, it requires that  $\gamma_{pr}(P_n^{+++}) \geq \left[\frac{4n-2}{7}\right]$  $\left| \frac{n-2}{7} \right|$ , for  $n \equiv 0 \pmod{7}$ . If  $n \equiv 1 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 22 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n+6}{7}$  $\frac{a_{t-1}}{7}$  where  $a_m = n$ ,  $a_{m-1} = n - 1$ ,  $a_t = 2n - 2$ , and  $a_{t-1} = 2n - 3$ . Therefore,

$$
|T| = t \ge \frac{4n + 6}{7} = \left\lceil \frac{4n - 2}{7} \right\rceil + 1
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \geq \left[\frac{4n-2}{7}\right]$  $\left| \frac{n-2}{7} \right|$  + 1, for  $n \equiv 1 \pmod{7}$ . If  $n \equiv 2 \pmod{7}$ , then

> $\sum_{i} f_{x_i}$  $m-2$  $x_1=1$  $+$   $\sum_{1}$   $f_{x_2}$  $t-2$  $x_2 = m+1$  $= 4n - 24 \leq 7(t - 4)$

It is obtained that  $t \geq \frac{4n+4}{7}$  $\frac{n}{7}$  where  $a_m = n$ ,  $a_{m-1} = n - 1$ ,  $a_t = 2n - 3$ , and  $a_{t-1} = 2n - 4$ . Therefore,

$$
|T| = t \ge \frac{4n+4}{7} = \left\lceil \frac{4n-2}{7} \right\rceil + 1
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \ge \left(\frac{4n-2}{7}\right)$  $\left| \frac{n-2}{7} \right|$  + 1, for  $n \equiv 2 \pmod{7}$ .

If  $n \equiv 3 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 26 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n+2}{7}$  $\frac{a_{t-2}}{7}$  where  $a_m = n$ ,  $a_{m-1} = n - 1$ ,  $a_t = 2n - 4$ , and  $a_{t-1} = 2n - 5$ . Therefore,

$$
|T| = t \ge \frac{4n + 2}{7} = \left\lceil \frac{4n - 2}{7} \right\rceil
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \ge \left(\frac{4n-2}{7}\right)$  $\left| \frac{n-2}{7} \right|$ , for  $n \equiv 3 \pmod{7}$ . If  $n \equiv 4 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 30 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n-2}{7}$  $\frac{a_{t-2}}{7}$  where  $a_m = n - 1$ ,  $a_{m-1} = n - 2$ ,  $a_t = 2n - 5$ , and  $a_{t-1} = 2n - 6$ . Therefore,

$$
|T| = t \ge \frac{4n - 2}{7} = \left\lceil \frac{4n - 2}{7} \right\rceil
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \ge \left(\frac{4n-2}{7}\right)$  $\left| \frac{n-2}{7} \right|$ , for  $n \equiv 4 \pmod{7}$ . If  $n \equiv 5 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 24 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n+4}{7}$  $\frac{a_{t+4}}{7}$  where  $a_m = n - 2$ ,  $a_{m-1} = n - 3$ ,  $a_t = 2n - 1$ , and  $a_{t-1} = 2n - 2$ . Therefore,

$$
|T| = t \ge \frac{4n+4}{7} = \left\lceil \frac{4n-2}{7} \right\rceil + 1
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \ge \left(\frac{4n-2}{7}\right)$  $\left| \frac{n-2}{7} \right|$  + 1, for  $n \equiv 5 \pmod{7}$ . If  $n \equiv 6 \pmod{7}$ , then

$$
\sum_{x_1=1}^{m-2} f_{x_1} + \sum_{x_2=m+1}^{t-2} f_{x_2} = 4n - 26 \le 7(t-4)
$$

It is obtained that  $t \geq \frac{4n+2}{7}$  $\frac{a_{t-2}}{7}$  where  $a_m = n - 3$ ,  $a_{m-1} = n - 4$ ,  $a_t = 2n - 1$ , and  $a_{t-1} = 2n - 2$ . Therefore,

$$
|T| = t \ge \frac{4n+2}{7} = \left\lceil \frac{4n-2}{7} \right\rceil
$$

Thus,  $\gamma_{pr}(P_n^{+++}) \geq \left[\frac{4n-2}{7}\right]$  $\left| \frac{n-2}{7} \right|$  + 1 for  $n \equiv 6 \pmod{7}$ .

Hence, from the necessary and sufficient conditions, we obtain the result as

$$
\gamma_{pr}(P_n^{+++}) = \begin{cases} \left[\frac{4n-2}{7}\right], & n \equiv 0,3,4,6 \pmod{7} \\ \left[\frac{4n-2}{7}\right] + 1, & \text{otherwise} \end{cases}
$$

**Theorem 3.2.** Let  $n > 5$  and  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{++-}) = 4$ .

PROOF.

Consider the set of vertices  $\{1, 2, \dots, n\}$  as the vertex set of  $P_n$ . Let S be paired dominating set of  $P_n^{++-}$  and  $S = S_1 \cup S_2$  where  $S_1 = \{u \mid u \in V(P_n)\}$  and  $S_2 = \{e \mid e \in V(L(P_n))\}$ . If the vertex 1 is part of  $S_1$  then all the vertices in  $L(P_n)$  except for (12)  $\in V(L(P_n))$  are dominated. To achieve perfect matching and ensure the domination of the vertex (12) and the remaining vertices  $\{3,4,\dots,n\} \in V(P_n)$ , it is needed to add a vertex  $e =$ (*uv*) to  $S_2$ , which is adjacent to 1, and  $v$  and  $v + 1$  to  $S_1$ . Therefore, {1, (*uv*),  $v$ ,  $v + 1$ } is a paired dominating set for  $P_n^{++-}$  and  $|S| \leq 4$ . However, it is impossible to dominate all the vertices with two pairs of vertices. Then,  $|S| > 2$ . Moreover, since the paired domination number must be an even number according to Proposition 2.4. Then,

$$
\gamma_{pr}(P_n^{++-})=4
$$

◻

**Theorem 3.3.** Let  $n \ge 6$  and  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{+(-+)}) = 2 + \gamma_{pr}(P_{n-4})$ .

PROOF.

Consider the set of vertices  $\{1,2,\dots,n\}$  as the vertex set of  $P_n$ . Let S be the minimum paired dominating set of  $P_n^{++1}$ . It is known that  $|S| \ge 2$  from Theorem 2.3. If the vertices (12) and  $((n-1)(n))$  with the maximum degree in  $P_n^{+ - +}$  are added to S, then all the vertices in  $V(L(P_n))$  are dominated. However, there are some nondominated vertices such as  $\{3,4,\dots,(n-2)\}\in V(P_n^{+(-+)})$ . These non-dominated  $(n-4)$  vertices construct a  $P_{n-4}$  structure. As a result, the size of *S* can be expressed as 2 +  $\gamma_{pr}(P_{n-4})$ . □

**Theorem 3.4.** Let  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{+--}) = 2$ .

PROOF.

Let  $\{1,2,\dots,n\}$  be the vertex set of  $P_n$ . Let S be a paired dominating set of  $P_n^{+--}$ . It is clear that all the vertices in  $P_n^{+--}$  dominated pairedly by  $\{(12), ((n-1)n)\}\)$ . This insight leads to the conclusion that  $|S| \le 2$ . By Theorem 2.3, it is established that  $S$  must have a size of at least 2. Consequently, combining these observations, we can affirm that  $|S| = 2$ .  $\Box$ 

**Theorem 3.5.** Let  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{---}) = 2$ .

PROOF.

Let  $\{1,2,\dots,n\}$  be the vertex set of  $P_n$ . Let S be a paired dominating set of  $P_n^{---}$ . As in the proof of  $P_n^{+--}$ , all the vertices in  $P_n^{---}$  dominated pairedly by  $\{(12), ((n-1)n)\}\)$ . Therefore,  $|S| \le 2$ . It is known that  $|S| \ge 2$ from Theorem 2.3. Consequently,  $|S| = 2$ .  $\Box$ 

**Theorem 3.6.** Let  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{--+}) = 4$ .

PROOF.

Let  $\{1,2,\dots,n\}$  be the vertex set of  $P_n$ . Consider that S be paired dominating set of  $P_n^{--+}$  and  $S = S_1 \cup S_2$ where  $S_1 = \{u \mid u \in V(P_n)\}\$  and  $S_2 = \{e \mid e \in V(L(P_n))\}\$ . If  $1 \in S_1$ , then all the vertices in  $P_n^{--+}$  except for 2 ∈  $V(P_n)$  and the vertex (12) ∈  $V(L(P_n))$  are dominated. In order to get perfect matching and dominate 2 ∈  $V(P_n)$ , it is needed to add a vertex (12) to  $S_2$ , which is adjacent to 1. However, the vertex (23) is the only non-dominated vertex in the graph. Therefore, a pair should be added to the set as  $\{1, (12), v, v + 1\}$ , a paired dominating set for  $P_n^{--+}$ . Hence,  $|S| \leq 4$ . Furthermore, it is not feasible to achieve domination over all the vertices using only two pairs of vertices, leading to the conclusion that  $|S| > 2$ . Considering Proposition 2.4 stating the paired domination number must be an even value,  $\gamma_{pr}(P_n^{++-}) = 4$ .  $\Box$ 

**Theorem 3.7.** Let  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{-+-}) = 2$ .

PROOF.

Consider the set of vertices  $\{1,2,\dots,n\}$  as the vertex set of  $P_n$ . Let S represent the paired dominating set of  $P_n^{-+}$ . All the vertices in  $P_n^{-+}$  are dominated pairedly by  $\{1, n\}$ . As a direct consequence, we deduce that  $|S| \le 2$ . It is established in Theorem 2.3 that  $|S|$  must have a minimum size of 2. Hence, combining these results,  $|S| = 2$ .  $\Box$ 

**Theorem 3.8.** Let  $P_n$  be a path graph. Then,  $\gamma_{pr}(P_n^{-++}) = 2 + \gamma_{pr}(P_{n-5})$ .

PROOF.

Let  $\{1,2,\dots,n\}$  be the vertex set of  $P_n$ , and S be the minimum paired dominating set of  $P_n^{-++}$ . It is known that  $|S| \ge 2$  from Theorem 2.3. Upon introducing the vertices (2) and  $((n-1))$ , which possess the highest degree in G, are added to the set S, then all the vertices in  $V(\overline{P_n})$  are dominated. Nevertheless, there are some nondominated vertices, such as  $\{(34), (45), \ldots, ((n-3)(n-2))\}$  in  $V(L(G))$ . These non-dominated  $(n-5)$ vertices construct a  $P_{n-5}$  structure. Therefore, it is obtained that  $S = 2 + \gamma_{pr}(P_{n-5})$ .  $\Box$ 

#### **3.2. Paired Domination of Transformation Cycle Graphs**

Given the particular structure of the cycle graph, it is noteworthy that the line graph of  $C_n$  exhibits isomorphism to itself. This observation leads to deduce the subsequent results:

**Theorem 3.9.** Let  $C_n$  be a cycle graph. Then,  $\gamma_{pr}(C_n^{++-}) = 4$ .

PROOF.

Let  $\{1,2,3,\dots,n\}$  be the vertex set of  $C_n$ , and S be the paired dominating set of  $C_n^{++-}$ . If S contains the vertex 1 and the vertex (12), then it ensures the domination of all the vertices within the graph. Nevertheless, these two vertices are not adjacent according to the form of the graph. In order to get perfect matching, it is needed to add two more vertex. It holds that the size of set  $S$  exceeds 2. Consequently, we can deduce that the parameter  $\gamma_{pr}(C_n^{++-})$  is equal to 4.  $\Box$ 

**Theorem 3.10.** Let  $C_n$  be a cycle graph and  $n \ge 9$ . Then,  $\gamma_{pr}(C_n^{+++}) = \gamma_{pr}(C_n^{+++}) = 4 + 2 \left[ \frac{n-8}{4} \right]$  $\frac{1}{4}$ . PROOF.

Let  $\{1,2,3,\dots,n\}$  be the vertex set of  $C_n$  and S be the paired dominating set of  $C_n^{++}$ . If S includes the vertices (12), (1n), 4, and  $(n-2)$ , then all the vertices of  $\overline{C}_n$  and the vertices labeled by (12), (23), (34), (45),  $((n-3)(n-2)), ((n-2)(n-1)), ((n-1)(n)),$  and  $(n1)$  in  $L(C_n)$  are dominated pairedly. However, the non-dominated  $P_{n-8}$  the structure remains. Consequently, from [2]

$$
\gamma_{pr}(C_n^{-++}) = \gamma_{pr}(C_n^{+-+}) = 4 + \gamma_{pr}(P_{n-8}) = 4 + 2\left\lceil \frac{n-8}{4} \right\rceil
$$

◻

**Theorem 3.11.** Let  $C_n$  be a cycle graph and  $n \ge 6$ . Then,  $\gamma_{pr}(C_n^{-+-}) = \gamma_{pr}(C_n^{+--}) = 2$ .

PROOF.

Consider the vertex set  $\{1,2,3,\dots,n\}$  as the vertex set of  $C_n$ , and let S denote a paired dominating set of  $C_n^{+-}$ . If S includes (12) and (45) from  $V(L(C_n))$ , then all the vertices in  $C_n^{+--}$  are dominated pairedly. Consequently, it follows that  $\gamma_{pr}(C_n^{++-}) = \gamma_{pr}(C_n^{+--}) = 2$ , establishing the paired domination number for  $C_n^{++-}$  and  $C_n^{+--}$ .  $\Box$ 

**Theorem 3.12.** Let  $C_n$  be a cycle graph and  $n \ge 6$ . Then,  $\gamma_{pr}(C_n^{--+}) = 4$ .

PROOF.

Given the graph  $C_n^{-(+)}$ , the complete graph of  $C_n^{(++)}$ , it can be deduced from Theorem 2.2 that its diameter is 2. Consequently, this implies that  $\gamma_{pr}(C_n^{--+}) \geq 4$ . Considering the vertex set {1,2,3, …, n} as the vertex set of  $C_n$ , and let S be paired dominating set of  $C_n^{-+}$ . If S is included (1) and (3) from  $V(\overline{C_n})$  and (12) and (34) from  $V(L(\mathcal{C}_n))$ , then all the vertices in  $\mathcal{C}_n^{-++}$  are dominated. This observation leads to the conclusion that  $|S| = \gamma_{pr}(C_n^{--+}) \leq 4$ . Hence,  $\gamma_{pr}(C_n^{--+}) = 4$ .  $\Box$ 

**Theorem 3.13.** Let  $C_n$  be a cycle graph. Then,  $\gamma_{pr}(C_n^{---}) = 2$ .

PROOF.

Consider the vertex set  $\{1,2,3,\dots,n\}$  of the  $C_n$ , and let S be a paired dominating set of  $C_n^{--}$ . If S contains the vertices 1 and 4, then it ensures that S dominates  $\overline{C_n}$  and  $L(C_n)$ . Additionally,  $\langle S \rangle$  forms a graph with perfect matching.  $\square$ 

**Theorem 3.14.** Let  $C_n$  be a cycle graph. Then,

$$
\gamma_{pr}(C_n^{+++}) = \begin{cases} \gamma_{pr}(P_n^{+++}) + 2, & n \equiv 4 \pmod{7} \\ \gamma_{pr}(P_n^{+++}), & \text{otherwise} \end{cases}
$$

PROOF.

Let S be the  $\gamma_{pr}$ -set of  $P_n^{+++}$ . Considered similar to the proof of the  $P_n^{+++}$ , S is also the  $\gamma_{pr}$ -set of  $C_n^{+++}$  when  $n \neq 4$  (mod 7). If the vertex set of  $L(C_n)$  labeled as in the proof of  $P_n^{+++}$ , the vertex  $(n1)$  corresponds to the vertex 2*n*. However, the vertex  $(2n) \in V(L(\mathcal{C}_n^{++}))$  cannot be dominated by *S* when  $n \equiv 4 \pmod{7}$ . In such cases, the addition of two extra vertices becomes necessary to ensure paired domination. Hence,

$$
\gamma_{pr}(C_n^{+++}) = \gamma_{pr}(P_n^{+++}) + 2
$$

when  $n \equiv 4 \pmod{7}$ . Henceforth, it is deduced that

$$
\gamma_{pr}(C_n^{+++}) = \begin{cases} \gamma_{pr}(P_n^{+++}) + 2, & n \equiv 4 \pmod{7} \\ \gamma_{pr}(P_n^{+++}), & \text{otherwise} \end{cases}
$$

#### **4. Conclusion**

We have implemented the discussed paired domination parameter from this research on path and cycle structures, two significant classes within the realm of graphs. We have derived outcomes for the transformation graphs encompassing all eight permutations in both these structural contexts. The outcome  $\gamma_{pr}(P_n^{+++})$  is then applied to derive the  $\gamma_{pr}$ ( $C_n^{+++}$ ) result for the cycle structure.

In upcoming research endeavors related to the topic, paired domination values of transformation graphs can be investigated for a general  $G$  graph. Furthermore, an attempt can be made to establish associations with other graph parameters to derive lower or upper-bound values. Since the concept of domination holds a significant position in graph theory, researchers interested in the subject might also explore the influence of transformation graphs on other domination parameters.

#### **Author Contributions**

The author read and approved the final version of the paper.

#### **Conflict of Interest**

The author declares no conflict of interest.

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## **A New Form of Smooth Cubic Surfaces with 9 Lines**

#### **Fatma Karao˘glu**<sup>1</sup> **[ID](https://orcid.org/0000-0003-2020-7829)**

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**Abstract** − A smooth cubic surface has at most 27 lines, with equality if and only if the underlying field is algebraically closed. Only a few cases are possible regarding the number of lines over fields that are not algebraically closed. The next two cases of interest are smooth cubic surfaces with 15 or 9 lines. The author has recently settled the case of 15 lines. In this paper, we address the case of smooth cubic surfaces with 9 lines. We describe a way to create some cubic surfaces with 9 or more lines based on a set of six field elements. Conditions on the six parameters are given under which the surface has exactly 9, 15, or 27 lines. However, the problem of generating all cubic surfaces with 9 lines remains open.

**Keywords** *Cubic surface, parametrization, non-algebraically closed fields*

**Mathematics Subject Classification (2020)** 14G27, 68W30

#### **1. Introduction**

It is well known that a smooth cubic surface over an algebraically closed field has exactly 27 lines [1]. However, the number of lines over a non-algebraically closed field varies. Naturally, the following question arises: How many lines can a smooth cubic surface have over a non-algebraically closed field? The problem of determining these numbers over the fields of  $\mathbb{R}$ ,  $\mathbb{Q}$ ,  $\mathbb{F}_q$  where q is odd, and  $\mathbb{F}_2$  has been considered by several authors [2–5]. In [6], the author gives the possible number of lines of smooth cubic surfaces over  $\mathbb{F}_q$  where q is even. The number of lines of a smooth cubic surface is one of 27, 15, 9, 7, 5, 3, 2, 1, or 0 [3]. The results on cubic surfaces with 15 or 27 lines over a given field are found in  $[7-11]$ , as well as using alternative methods in papers  $[12-14]$ . All the classification results agree with an enumerative formula recently found by Das [15].

In this paper, we focus on smooth surfaces with at least 9 lines over various fields, characteristic 0 or *p*. In [2], Schlafli described 27 lines of the cubic surface, explaining the line intersection properties. Each line intersects ten others and skews to 16. He defined the term "double-six", which has 12 lines with some special properties, and another 15 can be produced by these 12. To give the intersection properties of the lines for 9, we use the same idea of the double-six but for double-three. Smooth cubic surfaces may only have less than 27 lines if the field is nonalgebraically closed. However, over the algebraic closure of that field, the surface will have 27 lines. Therefore, we can use Schlafli labeling to notate 9 lines. We will then prove that the line intersection graph of the smooth cubic surfaces with 9 lines is unique.

<sup>1</sup> fkaraoglu@gtu.edu.tr (Corresponding Author)

<sup>&</sup>lt;sup>1</sup>Department of Mathematics, Faculty of Natural Sciences, Gebze Technical University, Kocaeli, Türkiye
We go back to the original proof by Cayley and Salmon that the surface has 27 lines over the algebraically closed fields as we did in [11]. We see that there is a discriminant condition certain polynomial of degree 5, which can have irreducible factors of degree 3 over the field  $\mathbb{F}$ , which is not algebraically closed. When this happens, we end up with cubic surfaces with 9 lines. Considering the rational lines over a given field  $\mathbb{F}$ , we formulate the conditions that the surface has 9 lines over the given field. We describe smooth cubic surfaces with at least 9 lines using six parameters. Our approach is experimental. We study some examples of smooth cubic surfaces with 9 lines that we obtain using the computer algebra system Orbiter [16]. Once we observe the pattern, we make the computer free proof. The proof is based on the symmetry of the projective group. We use the computer algebra system Maple for the symbolic computations. When we extend the field over the algebraic closure  $\mathbb{F}$ , the surface is complete to 27 lines since these surfaces are smooth. Using our model, we will show examples of cubic surfaces over a field of characteristic zero and *p*. These examples would have 9, 15, or 27 lines depending on whether the special polynomial is irreducible or reducible into two irreducible polynomials or splits completely over the base field.

We give the rational parameterization of points of our new form. To do this, we study the birational map between cubic surfaces and a plane, [9, 17]. There is an exceptional locus, and the birational map is defined outside the exceptional locus bijective. The exceptional locus of the map on the plane is two conics and a line. The exceptional locus of the map on the surface is two skew lines and one transversal line. We give them explicitly.

The smooth cubic surface has  $q^2 + tq + 1$  points where t is between -2 and 7, but 6 is never possible [18]. Studying the birational map helps us to prove that the smooth cubic surfaces with 9 lines have  $q^2 + 4q + 1$  points. If the cubic surface has a double-six, then the surface is smooth and has exactly 27 lines. However, if the surface has a double-three, it does not necessarily have exactly 9 lines. It can have more lines, in which case the surface is singular. Hence, a necessary condition for the smoothness of our new form is needed. We give this condition using the rational parameterization of our new form.

In section 2, we will provide some basic theory about the cubic surfaces with 27 lines since the structure of smooth cubic surfaces with 9 lines is the sub-configuration of the structure of cubic surfaces with 27 lines. In section 3, we show the uniqueness of 9 lines and investigate the configuration. In this section, we also give our new form for smooth cubic surfaces with at least 9 lines using six parameters, and we provide the conditions when the surface has exactly 9, 15, or 27 over the given field. In section 4, we provide examples of various fields, including Q and some finite fields. In section 5, we give the rational parameterization of our new form and the condition when our form is smooth. In section 6, we discuss future work.

#### **2. Preliminaries**

In this section, we provide some background material on cubic surfaces and projective geometry over finite fields. For a deeper treatment, we refer to [11, 19, 20].

A finite field is a field with only a finite number of elements.  $\mathbb{F}_q$  is a finite field of order  $q = p^k$  where  $p$  is a prime number. The characteristic of the field is the smallest *n* such that  $1 + 1 + 1 + \cdots + 1$  $= 0$ . The

| {z } *n times* characteristic of  $\mathbb{F}_q$  is *p*.  $\overline{\mathbb{F}}_p = \mathbb{F}_p$  adjoints all the roots of polynomials over  $\mathbb{F}_p$ .  $\overline{\mathbb{F}}_p$  is an algebraically closed field of characteristic p.  $\overline{\mathbb{F}}_p$  contains every  $\mathbb{F}_{p^e}$ , for all  $e \geq 1$ . Each  $\mathbb{F}_p$  has a unique  $\overline{\mathbb{F}}_p$ .

Let  $\mathbb F$  be a field. A projective space  $PG(n, \mathbb F)$  is a partially ordered set of subspaces of a vector space  $\mathbf{v}(n+1,\mathbb{F})$ . It is a lattice with respect to "join" and "meet". Join is the span of two subspaces. Meet is the intersection of two subspaces. In PG(3, F), a point is denoted by  $P = \mathbf{P}(\alpha_0, \alpha_1, \alpha_2, \alpha_3)$ . A line through the points  $P(\beta_0, \beta_1, \beta_2, \beta_3)$  and  $P(\gamma_0, \gamma_1, \gamma_2, \gamma_3)$  is denoted by

$$
\ell = \mathbf{L} \left[ \begin{array}{ccc} \beta_0 & \beta_1 & \beta_2 & \beta_3 \\ \gamma_0 & \gamma_1 & \gamma_2 & \gamma_3 \end{array} \right]
$$

The plane consists of the non-collinear points  $\mathbf{P}(\alpha_0, \alpha_1, \alpha_2, \alpha_3)$ ,  $\mathbf{P}(\beta_0, \beta_1, \beta_2, \beta_3)$ , and  $\mathbf{P}(\gamma_0, \gamma_1, \gamma_2, \gamma_3)$ and is denoted by

$$
\pi = \mathbf{v}(c_0x_0 + c_1x_1 + c_2x_2 + c_3x_3) = \begin{bmatrix} \alpha_0 & \alpha_1 & \alpha_2 & \alpha_3 \\ \beta_0 & \beta_1 & \beta_2 & \beta_3 \\ \gamma_0 & \gamma_1 & \gamma_2 & \gamma_3 \end{bmatrix}
$$

where  $c_0$ ,  $c_1$ ,  $c_2$ , and  $c_3$  are elements of the field  $\mathbb{F}$ .

A conic is a curve of degree 2 in  $PG(2, \mathbb{F})$ . It is either an irreducible conic, two distinct lines, or a double line. The space of quadratic polynomials in three variables has dimension 6. To determine conic in the associated projective space, 5 linearly independent conditions are required. A cubic curve is a curve of degree three in  $PG(2, \mathbb{F})$ . It is one of the following: an irreducible cubic, a conic, and a line, 3 different lines, or 2 different lines such that one of which is a double or a triple line. To determine a conic in the associated projective space, 9 linearly independent conditions are required. A *k*-arc in  $PG(2, \mathbb{F})$  is a set of *k* points where no three are collinear.

Let  $\pi$  be a plane in PG(3, F), and *Q* be a point on  $\pi$ . Let  $\ell_1$  and  $\ell_2$  be two skew lines in PG(3, F)  $\pi$ . Then, there is a unique transversal line of  $\ell_1$  and  $\ell_2$  through  $Q$ .

The automorphism group  $P\Gamma\mathrm{L}(n+1,\mathbb{F})$  of  $\mathrm{PG}(n,\mathbb{F})$  is the group of bijective mappings that preserve collinearity. The collineation group contains  $PGL(n+1, \mathbb{F})$  as subgroup which is the group of projectivities of  $PG(n, \mathbb{F})$ .  $PGL(4, \mathbb{F})$  is transitive on the points, lines, and planes of  $PG(3, \mathbb{F})$ . In  $PG(n, \mathbb{F})$ , any  $(n+2)$ -arc can be mapped to any other  $(n+2)$ -arc. The pointwise stabilizer of a hyperplane  $\pi$ in the PGL(4, F) is transitive on the set of two skew lines of PG(3, F) not in  $\pi$  which meet the fixed plane  $\pi$  in two points.

Let f be a homogeneous cubic equation in 4 variables over the field  $\mathbb{F}$ . A *cubic surface*  $\mathcal{F}$  in PG(3,  $\mathbb{F}$ ) is the zero set of *f*. For instance,

$$
\mathcal{F} = \mathbf{v}(f) = \mathbf{v}(x_0^3 + x_1^3 + x_2^3 + x_3^3 - (x_0 + x_1 + x_2 + x_3)^3)
$$

The cubic surface is smooth if the following system of equations has no solution:

$$
\begin{cases}\nf(x_0, x_1, x_2, x_3) = 0 \\
\frac{\partial f(x_0, x_1, x_2, x_3)}{\partial x_0} = 0 \\
\frac{\partial f(x_0, x_1, x_2, x_3)}{\partial x_1} = 0 \\
\frac{\partial f(x_0, x_1, x_2, x_3)}{\partial x_2} = 0 \\
\frac{\partial f(x_0, x_1, x_2, x_3)}{\partial x_3} = 0\n\end{cases}
$$

To define a cubic surface F in  $PG(3, \mathbb{F})$ , it is sufficient to specify 19 linearly independent points on it. A line in  $PG(3, \mathbb{F})$  either intersects cubic surfaces in three points, or it is the line of  $\mathcal{F}$ . Therefore, if the 4 points of the line are on the cubic surface, then the line lies on it. A cubic surface intersects a plane in a cubic curve. If the surface is smooth, then that cubic curve is one of the following: an irreducible cubic, a line, an irreducible conic, or 3 different lines. If the cubic surface intersects a plane in 3 different lines, then that plane is called a tritangent plane.

A "double-six" in PG(3*,* F) is the set of 12 lines

$$
\begin{array}{ccccccccc} a_1 & a_2 & a_3 & a_4 & a_5 & a_6 \\ b_1 & b_2 & b_3 & b_4 & b_5 & b_6 \end{array}
$$

such that  $a_i$  intersects  $b_j$  if and only if  $i \neq j$ ,  $a_i$  are pairwise skew, and  $b_i$  are pairwise skew.



**Figure 1.** 19 independent points of cubic surfaces with 27 lines

A double-six determines a unique cubic surface with 27 lines. 15 further lines  $c_{ij}$  are given by  $\lt a_i, b_j >$  $0 \leq a_j, b_i > 2$ . The red points in Figure 1 represent the 19 independent points that determine the cubic surface with 27 lines.

When three lines of the cubic surface are concurrent at a point, then this point is called an Eckardt point. From line intersection properties, only two cases are possible: either *a<sup>i</sup>* , *b<sup>j</sup>* , and *cij* are concurrent where  $i \neq j$  or  $c_{ij}$ ,  $c_{kl}$ , and  $c_{mn}$  are concurrent where *i*, *j*, *k*, *l*, *m*, and *n* are all different. In the first case, we notate the Eckardt point as  $E_{ij}$ , and for the second case,  $E_{ij,kl,mn}$ .

There is a map between the cubic surface in  $PG(3, \mathbb{F})$  and a plane. This map is called Clebsch map in [9]. Here, we refer to [9,11,17,20] and repeat the description of the map. Let F be a cubic surface and  $\pi$  be a plane in PG(3, F). Let  $\ell_1$  and  $\ell_2$  be two skew lines of  $\mathcal F$  not lying on the plane  $\pi$ , and  $P = \mathbf{P}(X)$  be a point of F which is neither on  $\ell_1$  nor on  $\ell_2$ . There is a unique line  $\ell$  through P which is the transversal to  $\ell_1$  and  $\ell_2$ . The line  $\ell$  meets  $\pi$  in a unique point  $Q = \mathbf{P}(Y)$ . Let  $Q$  be the image of *P*. Therefore,

$$
\begin{array}{c} \Phi : \mathcal{F} \to \pi \\ P \mapsto Q \end{array}
$$

There is a unique line through *Q* that is transversal to  $\ell_1$  and  $\ell_2$ . This line intersect  $\mathcal F$  in 3 points. Two of them are on  $\ell_1$  and  $\ell_2$ , let the third one to be *P*. The inverse  $\Phi^{-1}$  of this map moves *Q* to *P*.



**Figure 2.** Birational map from cubic surfaces with 27 lines to 6-arc not on a conic

Some properties of this map are as follows: Consider that the surface  $\mathcal F$  has 27 lines. Each line of the half double-six of  $\mathcal F$  maps to a single point in  $\pi$  under  $\Phi$ . These six points form a 6-arc not on a conic in the plane, see Figure 2. Outside these six lines, the map is bijective.

The following lemma is elementary. For a proof of the lemma, see Proposition 7.3 in [21]. For the sake of completeness, we include the reference of the previous paper of the author [11].

**Lemma 2.1.** [11,21] Let  $\mathcal F$  be a smooth cubic surface with at least one line. The number of tritangent planes through a line of  $\mathcal F$  is one of 0, 1, 2, 3, and 5 but never 4.

**Lemma 2.2.** [1] If two lines of a smooth cubic surface intersect, then they span a tritangent plane.

We introduce the following notation. Two tritangent planes are called disjoint if their line of intersection does not belong to F.

**Lemma 2.3.** [22] Any two disjoint tritangent planes of  $\mathcal F$  determine a third.

The three tritangent planes in Lemma 2.3 give rise to 9 lines of  $\mathcal F$ . These 9 lines give rise to three further tritangent planes [11, 20].

A trihedral pair consists of two sets of three tritangent planes, which pairwise intersect in 9 lines of  $\mathcal{F}$  [22].

**Lemma 2.4.** Let  $\mathcal F$  be a smooth cubic surface with at least 9 lines over the field  $\mathbb F$ . There exist at least four tritangent planes of  $\mathcal{F}$ .

Proof*.*

As 9 lines of F cannot be pairwise skew, there exist two lines *ℓ*<sup>1</sup> and *ℓ*<sup>2</sup> which intersect. From Lemma 2.2, there is a third line  $\ell_3$  such that  $\ell_1$ ,  $\ell_2$ , and  $\ell_3$  form a tritangent plane. Any line not contained in a hyperplane intersects the hyperplane at a point. In the case of a tritangent plane and a line of the surface, this point of intersection must be on one of the three lines of the tritangent plane. Hence, each  $m_i$  such that  $i \in \{1, \ldots, 6\}$  intersect one of the  $\ell_j$  such that  $j \in \{1, 2, 3\}$ . This gives rise to 6 pairs  $(\ell_i, m_i)$  of intersecting lines. By Lemma 2.2, these 6 pairs create at least 3 tritangent planes different from the tritangent plane through  $\ell_1, \ell_2$ , and  $\ell_3$ .  $\Box$ 

**Lemma 2.5.** Let  $\mathcal F$  be a smooth cubic surface with at least 9 lines over the field  $\mathbb F$ . Lemma 2.4 guarantees that there are at least 4 tritangent planes. If these 4 tritangent planes of  $\mathcal F$  intersect pairwise in a line of the surface, then they all intersect in the same line of  $\mathcal{F}$ .

#### PROOF.

Let  $\pi_1$ ,  $\pi_2$ ,  $\pi_3$ , and  $\pi_4$  be the tritangent planes arising from Lemma 2.4. Let  $\pi_1$  and  $\pi_2$  intersect in the line  $a_1$  of F. Without loss of generality, we may assume that  $\pi_1$  is spanned by  $a_1$ ,  $b_2$ , and  $c_{12}$ , and  $\pi_2$  is spanned by  $a_1$ ,  $b_3$ , and  $c_{13}$ . The lines  $b_2$  and  $c_{12}$  are skew to the lines  $b_3$  and  $c_{13}$ . Therefore, if the third tritangent plane  $\pi_3$  intersects  $\pi_1$  and  $\pi_2$  in lines of F, this line must be  $a_1$ . The same holds true for  $\pi_4$ .  $\square$ 

**Lemma 2.6.** Let  $\mathcal F$  be a smooth cubic surface with at least 9 lines over the field  $\mathbb F$ . Then, there is at least one pair of disjoint tritangent planes.

Proof*.*

From Lemma 2.4, there exist 4 tritangent planes of  $\mathcal{F}$ . To show that a disjoint pair of tritangent planes exist, we assume the opposite. From Lemma 2.5, these planes are all through the same line. By Lemma 2.1, there exists a fifth tritangent plane in this pencil. This configuration gives rise to 11 lines of the surface. Because of Lemma 2.2, any further line would create another tritangent plane not passing through *a*1, contradicting Lemma 2.5, and it is not possible that a smooth cubic surface to have 11 lines.  $\Box$ 

# **3. Construction of a Smooth Cubic Surface with 9 Lines**

This section proves the uniqueness of the line intersection graph of a smooth cubic surface with 9 lines and defines double-three. Besides, is provides a new form of smooth cubic surfaces with at least 9 lines involving 6 parameters.

The result of Cayley (and Salmon) in [1] is strengthened for smooth cubic surfaces with 15 lines in [11]. In this section, we will strength this result for cubic surfaces with 9 lines.

**Theorem 3.1.** The line intersection graph of a smooth cubic surface with exactly 9 lines over  $\mathbb{F}$  is unique.

#### Proof*.*

Let F be a smooth cubic surface with 9 lines over  $\mathbb F$ . By Lemma 2.6, there is a pair of disjoint tritangent planes. Because of Lemma 2.3, there is a unique third tritangent plane which is also disjoint to others. The 9 lines obtained in this way give rise to 3 more tritangent planes. Hence, there exists a trihedral pair of F. In addition, there are 2 tritangent planes through each line of F. Moreover, each line in the trihedral pair intersects 4 others and is skew to 4. Therefore, it is unique.  $\Box$ 

Let  $\ell$  be a line of F and  $\pi(\mu)$  be the plane through  $\ell$ .  $\pi(\mu)$  intersects F in a conic  $C(\mu)$  and the line *ℓ*. Let  $\mathcal{Q}(\mu)$  be a quadratic polynomial which represents  $C(\mu)$ . We define  $\Delta(\mu)$  as the discriminant of the quadratic polynomial  $\mathcal{Q}(\mu)$  as in the proof of Lemma 2.1.

**Theorem 3.2.** Let  $\mathbb{F}$  be a non-algebraically closed field, and  $\mathcal{F}$  be a smooth cubic surface with at least one line. The smooth cubic surface F has exactly 9 lines over F if and only if  $\Delta(\mu)$  consists of an irreducible polynomial of degree 3 and 2 linear factors over F.

#### Proof*.*

If  $\Delta(\mu)$  has two linear factors over the field F, then there exists 2 tritangent planes through  $\ell$ . With a similar argument in the proof of Theorem 11 in [11], we start with a fixed tritangent plane. There is one more tritangent plane through each of the three lines of the fixed plane. This gives 3 tritangent planes different from the one we started with. Each tritangent plane gives rise to 2 new lines. Counting all lines gives  $1 + 2 + 3 \cdot 1 \cdot 2 = 9$  lines. If F has exactly 9 lines, there are 2 tritangent planes through each line of F from Theorem 3.1. Therefore, there are exactly 2 distinct solutions for  $\mu \in \mathbb{F}$  in the  $\Delta(\mu)$ . Hence,  $\Delta(\mu)$  has exactly two distinct linear factors.  $\Box$ 

To give the intersection properties of the lines for smooth cubic surfaces with exactly 9 lines, we used the same idea of double-six but for double-three.

**Definition 3.3.** Let  $\mathbb{F}$  be a field. A *double-three* in  $PG(3,\mathbb{F})$  is the set of 6 lines

$$
\begin{array}{ccc}\na_1 & a_2 & a_3 \\
b_1 & b_2 & b_3\n\end{array}
$$

such that each line is skew to the ones in the same column or row and meets others. One row from the array is called half double-three.

Once a smooth cubic surface include a double-three, then it follows with 3 more lines which arise from the intersection of the following planes:

$$
c_{12} = \langle a_1, b_2 \rangle \cap \langle a_2, b_1 \rangle
$$
  

$$
c_{13} = \langle a_1, b_3 \rangle \cap \langle a_3, b_1 \rangle
$$

and

$$
c_{23} = \langle a_2, b_3 \rangle \cap \langle a_3, b_2 \rangle
$$

The intersection table of 9 lines can be seen in Table 1. We insert 1 in the table if the lines intersect; otherwise, 0.

	$\boldsymbol{a}_1$	$a_2$	$a_3$	$b_1$	$b_2$	$b_3$	$c_{12}$	$c_{13}$	$c_{23}$
$\boldsymbol{a}_1$		$\boldsymbol{0}$	$\boldsymbol{0}$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf 1$	$\mathbf 1$	$\mathbf{1}$	0
$a_2$	0		$\boldsymbol{0}$	$\mathbf{1}$	$\boldsymbol{0}$	$\mathbf 1$	$\mathbf{1}$	$\boldsymbol{0}$	$\mathbf{1}$
$a_3$	$\overline{0}$	$\boldsymbol{0}$		$\mathbf{1}$	$\mathbf{1}$	$\boldsymbol{0}$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf{1}$
$\mathfrak{b}_1$	$\theta$	$\mathbf{1}$	$\mathbf{1}$		$\boldsymbol{0}$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf{1}$	$\boldsymbol{0}$
$b_2$	$\mathbf{1}$	$\boldsymbol{0}$	$\mathbf{1}$	$\boldsymbol{0}$		$\boldsymbol{0}$	$\mathbf 1$	$\boldsymbol{0}$	$\mathbf{1}$
$b_3$	$\mathbf{1}$	$\mathbf{1}$	$\boldsymbol{0}$	$\boldsymbol{0}$	$\boldsymbol{0}$		0	$\mathbf{1}$	$\mathbf{1}$
$c_{12}$	$\mathbf{1}$	$\mathbf{1}$	$\boldsymbol{0}$	$\mathbf{1}$	$\,1$	$\boldsymbol{0}$		$\boldsymbol{0}$	$\boldsymbol{0}$
$c_{13}$	$\mathbf{1}$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf{1}$	$\boldsymbol{0}$	$1\,$	$\boldsymbol{0}$		$\theta$
$c_{23}$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf 1$	$\boldsymbol{0}$	$\mathbf{1}$	$\mathbf 1$	$\boldsymbol{0}$	$\boldsymbol{0}$	

**Table 1.** Pairwise intersection table of the 9 lines

The following theorem presents a new form of a smooth cubic surfaces with 9 lines.

**Theorem 3.4.** Let  $\mathbb{F}$  be a field with at least 4 elements. Let  $a, c \in \mathbb{F} \setminus \{0, 1\}, b, d \in \mathbb{F} \setminus \{0, -1\},$  and  $f, g \in \mathbb{F} \setminus \{0\}$  such that  $b \neq d$  and  $f \neq g$ . Let  $\mathcal{F}_{a,b,c,d,f,g} = \mathbf{v}(\mathfrak{f}_{a,b,c,d,f,g})$  be the variety over  $\mathbb F$  given by the equation f*a,b,c,d,f,g*

$$
d_{002}x_0^2x_2 + d_{012}x_0x_1x_2 + d_{013}x_0x_1x_3 + d_{022}x_0x_2^2 + d_{023}x_0x_2x_3
$$
  
+
$$
d_{112}x_1^2x_2 + d_{113}x_1^2x_3 + d_{122}x_1x_2^2 + d_{123}x_1x_2x_3 + d_{133}x_1x_3^2 = 0
$$
\n(1)

where

$$
d_{002} = bg\kappa_3
$$
  
\n
$$
d_{012} = g(a + b)\kappa_2 - f(bf + a)\kappa_1
$$
  
\n
$$
d_{023} = -abg\kappa_3
$$
  
\n
$$
d_{122} = -cf\kappa_4
$$
  
\n
$$
d_{023} = -abg\kappa_3
$$
  
\n
$$
d_{124} = -cf\kappa_4
$$
  
\n
$$
d_{125} = df\kappa_4
$$
  
\n
$$
d_{126} = adf\kappa_4
$$
  
\n
$$
d_{127} = -cf\kappa_4
$$
  
\n
$$
d_{128} = cdf\kappa_4
$$
  
\n
$$
d_{129} = (d - 1)cf\kappa_4
$$
  
\n
$$
\kappa_5 = abd + acd + bcd + ab
$$
  
\n
$$
\kappa_6 = abc + abd + bcd + ab
$$

 $\kappa_3 \neq 0, \ \kappa_4 \neq 0, \ a\kappa_1 + bc \neq 0, \text{ and } a\kappa_2 + bcf \neq 0$ 

Let  $\mathbf{g}_{a,b,c,d,f,g}$  be the polynomial

$$
\mathbf{g}_{a,b,c,d,f,g} = A_3 \mu^3 + A_2 \mu^2 + A_1 \mu + A_0 \tag{2}
$$

in  $\mu$  such that

$$
A_0 = -bcdf(ac + ad + cd + a)(ag - bf + bg - a)
$$
  
\n
$$
A_1 = g^2 A_{12} + g(c^2 A_{112} + cA_{111} + A_{110}) + A_{10}
$$
  
\n
$$
A_2 = a^2(fA_{221} + A_{220}) - abc(f^2 A_{212} + fA_{211} + A_{210}) + A_{20}
$$
  
\n
$$
A_3 = (f - g)(ac + ad + bc + a)(adg + bcf + ac + ag)
$$

and

$$
A_{12} = -a(d+1)b(abd + acd + bcd + ab)
$$
  
\n
$$
A_{220} = g(dg + c + g)(bc + 2bd + cd + 2b)
$$
  
\n
$$
A_{112} = -(a+b)(bd^2f + abd - abf - adf - 2bdf)
$$
  
\n
$$
A_{212} = (c+d+1)(b+d-1)
$$
  
\n
$$
A_{111} = a(d+1)(ab^2f - ab^2 + abf + adf + b^2f + bdf)
$$
  
\n
$$
A_{210} = -2cdg + 2cd + 2cg + dg - c + g
$$
  
\n
$$
A_{210} = -g(dg + c + g)(b + d)
$$
  
\n
$$
A_{221} = -(c + d + 1)(2bdg + bc + 2bg + cd + cg - c)
$$
  
\n
$$
A_{20} = -b^2c^2f(f - g)(2d - 1)
$$
  
\n
$$
A_{10} = cf(bf + a)(abcd + abd^2 + bcd^2 - abc - acd - ad^2 - 2bcd - ab - ad)
$$

Assume that the surface  $\mathcal{F}_{a,b,c,d,f,g}$  is smooth over F. The surface  $\mathcal{F}_{a,b,c,d,f,g}$  has at least 9 lines, six of which form a double-three. The conditions on the exact number of lines of  $\mathcal{F}_{a,b,c,d,f,g}$  depends on the polynomial **g***a,b,c,d,f,g* of degree three.

*i.* If the polynomial  $\mathbf{g}_{a,b,c,d,f,g}$  is irreducible over the field  $\mathbb{F}$ , then the surface  $\mathcal{F}_{a,b,c,d,f,g}$  has exactly 9 lines.

*ii.* If  $\mathbf{g}_{a,b,c,d,f,g}$  is reducible into one irreducible quadratic polynomial and one linear over the field  $\mathbb{F}$ , then the surface  $\mathcal{F}_{a,b,c,d,f,g}$  has exactly 15 lines.

*iii.* If  $\mathbf{g}_{a,b,c,d,f,g}$  splits completely to 3 linear factors over the field  $\mathbb{F}$ , then the surface  $\mathcal{F}_{a,b,c,d,f,g}$  has exactly 27 lines.

In Table 2, the parametrization of the 9 lines of  $\mathcal{F}_{a,b,c,d,f,q}$  can be observed.



Proof*.*

The group  $PGL(4, \mathbb{F})$  is transitive on the planes. Therefore, we can start from any hyperplane. Hence, we may pick  $\pi = \mathbf{v}(x_3)$ . Consider that we want to construct a cubic surface F with 9 lines, including

the double-three. Considering the sub-configuration of Schlafli configuration and the labeling for 9 lines, we may choose the labels that  $a_1$ ,  $a_2$ ,  $a_3$ ,  $b_1$ ,  $b_2$ ,  $b_3$ ,  $c_{12}$ ,  $c_{13}$ , and  $c_{23}$ . To determine this surface F, we have to specify 19 independent conditions. Nine of them should be on the plane  $\pi$  since the cubic surface intersects  $\pi$  in a cubic curve. We can assume that the coordinates of 4 of those 9 points are  $P_1 = \mathbf{P}(1,0,0,0), P_2 = \mathbf{P}(1,0,0,0), P_3 = \mathbf{P}(1,0,0,0),$  and  $P_4 = \mathbf{P}(1,0,0,0)$  since the projective group of the plane is transitive on the quadrangles. Since  $PGL(4, F)$  is transitive on the lines, we can pick the first line of the surface  $\mathcal F$  without any constraints. Let assume that the first line of  $\mathcal F$  is

$$
c_{12} = \mathbf{L} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}
$$

Since  $c_{12}$  lies on the surface F, there are 2 more independent points on the intersection  $c_{12} \cap \mathcal{F}$ . Assume that  $P_7 = \mathbf{P}(1,1,0,0)$  and  $P_8 = \mathbf{P}(-1,1,0,0)$  are the points of F as well as  $c_{12}$  is transitive. It is known that the pointwise stabilizer of the hyperplane  $\pi$  in the PGL(4, F) is transitive on the set of two skew lines of  $PG(3, \mathbb{F})$  not in  $\pi$  which meet the fixed plane  $\pi$  in two points  $P_1$  and  $P_2$ . Therefore, we are free to choose two skew lines not on the plane  $\pi$  through  $P_1$  and  $P_2$ . Let

$$
b_1 = \mathbf{L} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}
$$
 and  $b_2 = \mathbf{L} \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 1 \end{bmatrix}$ 

are the two skew lines.  $b_1$  are through  $P_1$  and  $P_5 = P(0,0,0,1)$  and  $b_2$  are through  $P_2$  and  $P_6 =$ **P**(0,0, -1, 1). It is known that there is a unique line through  $P_2$  and transversal to  $b_1$  and  $b_2$ . We call this line is a line of  $\mathcal F$ . Because of the configuration of 9 lines, we can label that line as  $a_3$  and it is the line through  $P_5$  and  $P_6$ ,

$$
a_3 = \mathbf{L} \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 1 \end{bmatrix}
$$

From line intersection properties, we know that there exists a line  $c_{23}$  which meets  $b_2$  and  $a_3$  and skew to  $c_{12}$ , i.e., not in  $\pi$ . When  $a_3$ ,  $b_2$ , and  $c_{23}$  are not concurrent which means the Eckardt point  $E_{32}$ does not exist, we may set the line  $c_{23}$  through  $P_9 = P(0, c, -1, 1)$  and  $P_{10} = P(0, 0, d, 1)$ . Since  $P_2$ ,  $P_6$ , and  $P_9$  are three distinct points from our assumption, we have  $0 \neq c \neq 1$ .

From line intersection properties, we know that there exists a line  $c_{13}$  which meets  $b_1$  and  $a_3$  and skew to  $c_{12}$ , i.e., not in  $\pi$ . When  $a_3$ ,  $b_1$ , and  $c_{13}$  are not concurrent which means the Eckardt point  $E_{31}$  does not exist, we may set the line  $c_{13}$  through  $P_{11} = P(a, 0, 0, 1)$  and  $p' = P(0, 0, b, 1)$ . Since  $P_1, P_5$ , and  $P_{11}$  are three distinct points from our assumption, we have  $0 \neq a \neq 1$ . Therefore, two lines of  $\mathcal F$  are

$$
c_{13} = \mathbf{L} \begin{bmatrix} a & 0 & 0 & 1 \\ 0 & 0 & b & 1 \end{bmatrix} \quad \text{and} \quad c_{23} = \mathbf{L} \begin{bmatrix} 0 & c & -1 & 1 \\ 0 & 0 & d & 1 \end{bmatrix}
$$

Since  $P_3$ ,  $P_5$ ,  $P_6$ ,  $P_{10}$ , and  $p'$  are five distinct points on the line  $a_3$ , we have  $b, d \notin \{0, -1\}$  and  $b \neq d$ . Since there are already 4 independent points on  $a_3$ , we do not count the point *p*'. The plane  $\pi$  intersects the line  $c_{13}$  at the point  $P_{13} = \mathbf{P}(-a, 0, b, 0)$  and intersects the line  $c_{23}$  at the point  $P_{12} = P(0, -c, d+1, 0)$ . From line intersection properties, we know that there exists a line *b*<sub>3</sub> which meets  $c_{13}$  and  $c_{23}$  and skew to  $a_3$ ,  $b_1$ ,  $b_2$ , and  $c_{12}$ . The line  $b_3$  cannot intersect the line through  $P_1$  and  $P_3$  and the line through  $P_2$  and  $P_3$  since those lines already have 3 points of  $\mathcal F$  each. The new line  $b_3$  is either  $P_4$  or  $P_{19} = \mathbf{P}(f, g, 1, 0)$ . For this new form, we consider the  $b_3$  is through  $P_4$ . We know that the line  $c_{13}$  and  $c_{23}$  are not in  $\pi$ , and they are skew. Hence, the transversal line  $b_3$  to  $c_{13}$  and  $c_{23}$  through  $P_4$ is uniquely determined. It intersects  $c_{13}$  at the point  $P_{14} = \mathbf{P}(ac(b-d), 0, \kappa_6, a\kappa_1 + bc)$  and intersects  $c_{23}$  at the point  $P_{15} = P(0, -ac(b-d), \kappa_5, a\kappa_1 + bc)$  where  $\kappa_1 = c + d + 1$ ,  $\kappa_5 = abd + acd + bcd + ab$ ,

and  $\kappa_6 = abc + abd + bcd + ab$ . Thus,

$$
b_3 = \mathbf{L} \begin{bmatrix} 1 & 1 & 1 & 0 \\ c(b-d)a & 0 & \kappa_6 & a\kappa_1 + bc \end{bmatrix}
$$

The line  $c_{13}$  already has 4 independent points of  $\mathcal F$  which are  $p'$ ,  $P_{11}$ ,  $P_{13}$ , and  $P_{14}$ . Hence, it is a line of F. The line  $c_{23}$  already has 4 independent points of F which are  $P_{12}$ ,  $P_{10}$ ,  $P_9$ , and  $P_{15}$ . Thus, it is a line of  $\mathcal{F}$ . The lines  $b_1$ ,  $b_2$ , and  $b_3$  have 3 points of  $\mathcal{F}$  each. Therefore, we force the points  $P_{16} = \mathbf{P}(1,0,0,1), P_{17} = \mathbf{P}(0,1,-1,1), \text{ and } P_{18} = a \cdot P_{15} + P_{14}$  to be on the surface *F*. The following 18 points are chosen to force the lines

$$
b_1, b_2, b_3, a_3, c_{12}, c_{13}, \text{ and } c_{23}
$$

to be on the cubic surface. More conditions arise from the fact that the cubic surface intersects the plane  $\mathbf{v}(x_3)$  in a cubic curve which consists of the line  $c_{12}$  and an irreducible conic. Besides, we have considered the 8 points on *π*

$$
P_1
$$
,  $P_2$ ,  $P_3$ ,  $P_4$ ,  $P_7$ ,  $P_8$ ,  $P_{12}$ , and  $P_{13}$ 

We need to force one further point of the surface to lie on this plane. We may pick the point  $P_{19} = \mathbf{P}(f, g, 1, 0)$ . The conic through  $P_3$ ,  $P_4$ ,  $P_{12}$ ,  $P_{13}$ , and  $P_{19}$  is irreducible. This gives the restriction on the parameters *a*, *b*, *c*, *d*, *f*, and *g* such that  $a, c, f, g \neq 0$ ,  $f \neq g$ ,  $\kappa_3 \neq 0$ ,  $\kappa_4 \neq 0$ ,  $a\kappa_1 + bc \neq 0$ , and  $a\kappa_2 + bc \neq 0$ .

The points  $P_1, \ldots, P_{19}$  define the cubic surface uniquely since they are linearly independent. Using Maple, we compute the equation  $f_{a,b,c,d,f,g}$  as in Equation 1 which is the unique equation of a cubic surface  $\mathcal{F}_{a,b,c,d,f,q}$  defined by these 19 points. This surface also involves a double-three. By Theorem 3.1, we know that there is only one way to complete the configuration of the lines. By Definition 3.3, further lines will be defined as  $a_1$  and  $a_2$ . These further lines  $a_1$  and  $a_2$  can be computed using Maple. All these 9 lines can be seen in Table 2.



**Figure 3.** Configuration of 19 points

We check whether the cubic surface  $\mathcal{F}_{a,b,c,d,f,g}$  has any further lines. The only possible way would be the surface having 15 or 27 lines. This depends on the discriminant condition. By Theorem 3.2 and Theorem 11 in [11], we know that it depends on the factors of the discirimant. Therefore, we need to compute it. As in the proof of Theorem 3.2, we start with a fixed tritangent plane  $\pi = \mathbf{v}(x_3)$ . Consider the line

$$
c_{12} = \mathbf{L} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix}
$$

and the planes

 $$ 

through this line  $c_{12}$ . Each of these planes intersects  $\mathcal{F}_{a,b,c,d,f,q}$  in  $c_{12}$  and in a conic  $C(\mu)$ . Substituting  $x_2 = \mu x_3$  into the equation of Equation 1, we find the conic equation  $C(\mu)$  in  $x_0, x_1$ , and  $x_3$ . Hence, we calculate the discriminant  $\Delta(\mu)$  of  $C(\mu)$ . It has a factor of the polynomial  $\mathbf{g}_{a,b,c,d,f,g}$  of degree 3 as in the Equation 2. Similarly, considering the planes through  $a_1$  and  $b_2$ , two further polynomials  $f_{a,b,c,d,f,g}$ and **h***a,b,c,d,f,g* of degree three arise. These polynomials are not stated due to space restrictions. For a given field  $\mathbb{F}$ , if the polynomial  $\mathbf{g}_{a,b,c,d,f,g}$  is irreducible, then  $\mathbf{f}_{a,b,c,d,f,g}$  and  $\mathbf{h}_{a,b,c,d,f,g}$  are irreducible, and vice versa.

There are 2 tritangent planes of  $\mathcal{F}_{a,b,c,d,f,g}$  through each created 9 lines. Hence, there are either one or 3 more tritangent planes through each line. Then, these further lines would arise from the new tritangent planes.

If  $\mathbf{g}_{a,b,c,d,f,g}$  is reducible into one linear factor and one irreducible quadratic, then the surface  $\mathcal{F}_{a,b,c,d,f,g}$ has 15 lines over  $\mathbb F$  from Part 2 of Theorem 11 in [11]. The further 6 lines arise as following: Each of  $c_{12}$ ,  $a_1$ , and  $b_2$  lies on one further tritangent planes, and each of which gives two lines to the surface. This gives  $3 \cdot 1 \cdot 2 = 6$  further lines of  $\mathcal{F}$ .

If  $\mathbf{g}_{a,b,c,d,f,g}$  is reducible into three linear factors, then the surface  $\mathcal{F}_{a,b,c,d,f,g}$  has 27 lines over F from Part 1 of Theorem 11 in [11]. The further 18 lines arise as following: Each of  $c_{12}$ ,  $a_1$ , and  $b_2$  lies on three further tritangent planes, each of which gives two lines to the surface. This gives  $3 \cdot 3 \cdot 2 = 18$ further lines of  $\mathcal{F}$ .

If the polynomial  $\mathbf{g}_{a,b,c,d,f,g}$  is irreducible over  $\mathbb{F}$ , then the surface  $\mathcal{F}_{a,b,c,d,f,g}$  has exactly 9 lines over this field by Theorem 3.2.  $\Box$ 

**Remark 3.5.** Figure 3 summarises the content of Theorem 3.4 by giving the configuration of 19 points on the cubic surface and forcing the surface including the double-three.

**Remark 3.6.** The cubic surface  $\mathcal{F}_{a,b,c,d,f,q}$  cannot have 2 specific Eckardt points. Therefore, we cannot create all possible smooth cubic surfaces with 9 lines using this form.

**Remark 3.7.** The line  $b_3$  of  $\mathcal{F}_{a,b,c,d,f,g}$  could be through  $P_{19}$  instead  $P_4$ . Considering this way would give another form of such surfaces.

#### **4. Illustrative Examples**

This section exemplifies the aforesaid form over finite fields and fields of characteristic zero.

**Example 4.1.** The surface  $\mathcal{F}_{-1,1,-1,2,-1,1}$  given by the equation

 $-2x_0^2x_2 - 2x_0x_1x_2 - 2x_0x_1x_3 + 2x_0x_2^2 - 2x_0x_2x_3 + 3x_1^2x_2 + 3x_1^2x_3 - x_1x_2^2 + x_1x_2x_3 + 2x_1x_3^2 = 0$ 

is smooth over C. The polynomial

$$
\mathbf{g}_{-1,1,-1,2,-1,1} = -6\mu^3 + 10\mu^2 - 16\mu + 16
$$

is irreducible over  $\mathbb Q$ . Hence, the surface  $\mathcal F_{-1,1,-1,2,-1,1}$  has exactly 9 lines over  $\mathbb Q$ . The polynomial **g**−1*,*1*,*−1*,*2*,*−1*,*<sup>1</sup> has one real and two complex roots. Therefore, by Theorem 3.4, it has exactly 15 lines over  $\mathbb R$  and 27 over  $\mathbb C$ .

**Example 4.2.** Over  $\mathbb{F}_5$ , the cubic surface  $\mathcal{F}_{2,1,4,2,4,1}$  given by the equation

$$
3x_0^2x_2 + 3x_1^2x_2 + 3x_1^2x_3 + x_0x_2^2 + 4x_1x_2^2 + 2x_1x_3^2 + 4x_0x_1x_2 + 3x_0x_1x_3 + 4x_0x_2x_3 + x_1x_2x_3 = 0
$$

is smooth. Since the polynomial

$$
\mathbf{g}_{2,1,4,2,4,1} = 2\mu^3 + 3\mu^2 + 4\mu + 3
$$

is reducible into  $(\mu + 1)$  and  $(2\mu^2 + \mu + 3)$  over  $\mathbb{F}_5$ , the surface  $\mathcal{F}_{2,1,4,2,4,1}$  has exactly 15 lines over this field by Theorem 3.4.

Let  $\tau \in \mathbb{F}_{25}$  satisfy the equation  $\tau^2 + \tau + 2 = 0$ . The polynomial

$$
\mathbf{g}_{2,1,4,2,4,1} = 2\mu^3 + 3\mu^2 + 4\mu + 3
$$
  
=  $(\mu + 1)(2\mu^2 + \mu + 3)$   
=  $(\mu + 1)(\mu + 4\tau + 1)(\mu + \tau + 2)$ 

is reducible into three linear factors over  $\mathbb{F}_{25}$ . Therefore, by Theorem 3.4, the surface  $\mathcal{F}_{2,1,4,2,4,1}$  has exactly 27 lines over this field.

**Example 4.3.** Over  $\mathbb{F}_5$ , the cubic surface  $\mathcal{F}_{4,1,4,2,4,1}$  given by the equation

$$
3x_0^2x_2 + 3x_1^2x_2 + 3x_1^2x_3 + 2x_0x_2^2 + 4x_1x_2^2 + 2x_1x_3^2 + 3x_0x_1x_2 + 3x_0x_1x_3 + 3x_0x_2x_3 + x_1x_2x_3 = 0
$$

is smooth. Since the polynomial

$$
\mathbf{g}_{4,1,4,2,4,1} = 3\mu^3 + 3\mu + 2
$$

is irreducible over  $\mathbb{F}_5$ , the surface  $\mathcal{F}_{4,1,4,2,4,1}$  has exactly 9 lines over this field by Theorem 3.4. Over  $\mathbb{F}_{25}$ , this polynomial is also irreducible. Therefore, the cubic surface  $\mathcal{F}_{4,1,4,2,4,1}$  has exactly 9 lines over this field.

Let  $\psi \in \mathbb{F}_{125}$  satisfy the equation  $\psi^3 + \psi^2 + 2 = 0$ . The polynomial

$$
\mathbf{g}_{4,1,4,2,4,1} = 3\mu^3 + 3\mu + 2
$$
  
=  $(\mu + 4\psi^2 + 3\psi + 3)(\mu + 4\psi^2 + 2\psi + 1)(\mu + 2\psi^2 + 1)$ 

is reducible into three linear factors over  $\mathbb{F}_{125}$ . Therefore, the surface  $\mathcal{F}_{4,1,4,2,4,1}$  has exactly 27 lines over this field by Theorem 3.4.

#### **5. Rational Parametrization of the New Form**

This section provides a parametrization of the rational points at lines of the cubic surface  $\mathcal{F}_{a,b,c,d,f,g}$ given in Section 3. The proof of the following theorem is based on the birational map between the cubic surface and a plane. In Subsection 4.2 of [11], rational parametrization of a smooth cubic surface with 15 lines is given explicitly. The proof here and the remarks are similar to the proof of Theorem 8 and its following remarks in [11] but for the cubic surface with at least 9 lines  $\mathcal{F}_{a,b,c,d,f,q}$ .

**Theorem 5.1.** Let  $\mathbb{F}$  be a field with at least 4 elements. Let  $\mathcal{F}_{a,b,c,d,f,g}$  be the variety given in Theorem 3.4. Assume that  $\mathcal{F}_{a,b,c,d,f,g}$  is smooth. Let  $P = \mathbf{P}(x_0, x_1, x_2, x_3)$  be a point on  $\mathcal{F}_{a,b,c,d,f,g}$ and  $Q = P(y_0, y_1, y_2)$  be a point in a plane so that P and Q are the images of each other under a birational map

$$
\Phi_{a,b,c,d,f,g} : \mathcal{F}_{a,b,c,d,f,g} \longrightarrow \mathrm{PG}(2,\mathbb{F})
$$

$$
\mathbf{P}(x_0, x_1, x_2, x_3) \mapsto \mathbf{P}(y_0, y_1, y_2)
$$

Then,  $(x_0, x_1, x_2, x_3)$  can be expressed as

 $\overline{ }$ 

$$
\begin{cases}\nx_0 = -g(a+b)\kappa_3y_0^2y_1 + ga(b+1)\kappa_3y_0^2y_2 - (1+d)f\kappa_4y_0y_1^2 - fc(1+d)\kappa_4y_0y_1y_2 \\
x_1 = abg\kappa_3y_0y_1y_2 - cdf\kappa_4y_1^2y_2 - bg\kappa_3y_0^2y_1 - f\kappa_1\kappa_4y_0y_1^2 \\
x_2 = abg\kappa_3y_0y_2^2 - cdf\kappa_4y_1y_2^2 - bg\kappa_3y_0^2y_2 - f\kappa_1\kappa_4y_0y_1y_2 \\
x_3 = gb\kappa_3y_0^2y_2 + ag\kappa_3y_2^2y_0 - cf\kappa_4y_2^2y_1 + d_{012}y_0y_1y_2 - (1+d)f\kappa_4y_1^2y_2\n\end{cases}
$$
\n(3)

and  $(y_0, y_1, y_2)$  can be expressed as

$$
\begin{cases}\ny_0 = x_0 x_2 \\
y_1 = x_1 x_2 + x_1 x_3 \\
y_2 = x_2^2 + x_2 x_3\n\end{cases}
$$
\n(4)

up to a nonzero scalar multiple.

#### Proof*.*

No generality is lost by picking the plane as  $\mathbf{v}(x_3)$ , and the two skew lines  $b_1$  and  $b_2$  of  $\mathcal{F}_{a,b,c,d,f,g}$ 

$$
b_1 = \mathbf{L} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}
$$
 and  $b_2 = \mathbf{L} \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 1 \end{bmatrix}$ 

Let  $P = \mathbf{P}(x_0, x_1, x_2, x_3)$  be a point on  $\mathcal{F}_{a,b,c,d,f,g}$  but neither on  $b_1$  nor on  $b_2$ . To find the image *Q* of *P* under  $\Phi_{a,b,c,d,f,g}$ , first we compute the transversal line  $\ell$  of  $b_1$  and  $b_2$  through *P*. Then, we find the unique point *Q* where  $\ell$  meets  $\mathbf{v}(x_3)$  as follows:

$$
\ell = \mathbf{L} \begin{bmatrix} x_0 & 0 & 0 & x_2 + x_3 \\ x_0x_2 & x_1x_2 + x_1x_3 & x_2^2 + x_2x_3 & 0 \end{bmatrix}
$$

and

$$
\Phi_{a,b,c,d,f,g}(\mathbf{P}(x_0,x_1,x_2,x_3)) = \mathbf{P}(y_0,y_1,y_2) = \mathbf{P}(x_0x_2,x_1x_2+x_1x_3,x_2^2+x_2x_3)
$$

whose coordinates are given in System 4.

F.

Conversely, let  $Q = P(y_0, y_1, y_2, 0)$  be a point in the plane  $\mathbf{v}(x_3)$ . To find the image *Q* of *P* under  $\Phi_{a,b,c,d,f,g}^{-1}$ , first we compute the transversal line  $\ell'$  of  $b_1$  and  $b_2$  through *Q* as follows:

$$
\ell' = \mathbf{L} \begin{bmatrix} y_0 & 0 & 0 & y_2 \\ y_0 & y_1 & y_2 & 0 \end{bmatrix}
$$

Then, we find the three intersection points where  $\ell'$  meets  $\mathcal{F}_{a,b,c,d,f,g}$  by substituting the point

$$
\mathbf{P}(y_0, 0, 0, y_2) + t \cdot \mathbf{P}(y_0, y_1, y_2, 0)
$$

of  $\ell'$  into the equation of  $\mathcal{F}_{a,b,c,d,f,g}$ . This gives a polynomial of degree three in *t*, for some  $t \in \mathbb{F}$ . Solving this polynomial in *t* gives 3 solutions as follows:

$$
t_1 = 0
$$

$$
t_2 = -1
$$

and

$$
t_3 = \frac{abgs_3y_0y_2 - cdf\kappa_4y_1y_2 - bg\kappa_3y_0^2 - f\kappa_1\kappa_4y_0y_1}{gb\kappa_3y_0^2 + ag\kappa_3y_2y_0 - cf\kappa_4y_2y_1 + d_{012}y_0y_1 - (1+d)f\kappa_4y_1^2}
$$

The point  $\mathbf{P}(y_0, 0, 0, y_2)$  arising from  $t_1$  is the point where  $\ell'$  meets the line  $b_1$ . The point  $\mathbf{P}(0, -y_1, -y_2, y_2)$ arising from  $t_2$  is the point where  $\ell'$  meets the line  $b_2$ . Let  $P = \mathbf{P}(x_0, x_1, x_2, x_3)$  be the third point of intersection of  $\ell'$  with  $\mathcal{F}_{a,b,c,d,f,g}$ . Hence,

```
x_0 = y_0 + t_3 \cdot y_0x_1 = t_3 \cdot y_1x_2 = t_3 \cdot y_2
```
and

$$
x_3=y_2
$$

as in System 3. Then,

$$
\Phi_{a,b,c,d,f,g}^{-1}(\mathbf{P}(y_0,y_1,y_2)) = \mathbf{P}(x_0,x_1x_2,x_3)
$$

 $\Box$ 

**Remark 5.2.** Let  $\mathcal{F}_{a,b,c,d,f,g}$  be the smooth cubic surface over the field  $\mathbb{F}$  as described in Theorem 3.4, and  $\Phi_{a,b,c,d,f,g}$  be the birational map from  $\mathcal{F}_{a,b,c,d,f,g}$  to the plane embedded in PG(3*,* F) as the hyperplane  $\mathbf{v}(x_3)$ , described in Theorem 5.1. Consider the lines  $a_1$ ,  $a_2$ , and  $a_3$  of  $\mathcal{F}_{a,b,c,d,f,g}$ . The map  $\Phi_{a,b,c,d,f,g}$  sends the line  $a_1$  to the point  $\mathbf{P}(1,0,0)$ , the line  $a_2$  to the point  $\mathbf{P}(0,1,0)$ , and the line  $a_3$ to the point  $\mathbf{P}(0,0,1)$ .

**Remark 5.3.** The exceptional locus of  $\Phi_{a,b,c,d,f,g}$  on the surface consists of 3 lines of  $\mathcal{F}_{a,b,c,d,f,g}$  which are  $b_1, b_2$ , and  $c_{12}$ . The exceptional locus on the plane can be found explicitly by applying  $\Phi_{a,b,c,d,f,g}^{-1}$ to  $P(x_0, x_1, x_2, x_3)$  and then by applying  $\Phi_{a,b,c,d,f,g}$  to  $P(y_0, y_1, y_2)$ . It consists of one line and two conics in the plane, which are the line L' through  $P(1, 0, 0)$  and  $P(0, 1, 0)$  and two conics

 $D_1 = abg\kappa_3y_0y_2 - cdf\kappa_4y_1y_2 - bg\kappa_3y_0^2 - f\kappa_1\kappa_4y_0y_1$ 

and

$$
D_2 = -f(1+d)\kappa_4y_1^2 + ag(b+1)\kappa_3y_0y_2 - cf(1+d)\kappa_4y_1y_2 + (-f\kappa_1\kappa_4 + d_{012})y_0y_1
$$

The points of the lines  $b_1$  and  $b_2$  are mapped to the points of the conics  $D_1$  and  $D_2$  and the points of  $c_{12}$  are mapped to the points of *L*' under  $\Phi_{a,b,c,d,f,g}$ .

**Theorem 5.4.** The cubic surface  $\mathcal{F}_{a,b,c,d,f,g} = \mathbf{v}(f_{a,b,c,d,f,g})$  as described in Theorem 3.4 is non-singular over F if and only if the certain four sextic curves in  $(2, \mathbb{F})$  never intersect at a point of  $\mathcal{F}_{a,b,c,d,f,q}$ .

Proof*.*

The idea of the proof is similar to the proof of Theorem 9 in [11] but for  $\mathcal{F}_{a,b,c,d,f,g}$ . The partial derivatives of f*a,b,c,d,f,g* as follows:

$$
\frac{\partial f_{a,b,c,d,f,g}}{\partial x_0} = 2d_{002}x_0x_2 + d_{012}x_1x_2 + d_{013}x_1x_3 + d_{022}x_2^2 + d_{023}x_2x_3
$$

$$
\frac{\partial f_{a,b,c,d,f,g}}{\partial x_1} = d_{012}x_0x_2 + d_{013}x_0x_3 + 2d_{112}x_1x_2 + 2d_{113}x_1x_3 + d_{122}x_2^2 + d_{123}x_2x_3 + d_{133}x_3^2
$$
  

$$
\frac{\partial f_{a,b,c,d,f,g}}{\partial x_2} = d_{002}x_0^2 + d_{012}x_0x_1 + 2d_{022}x_0x_2 + d_{023}x_0x_3 + d_{112}x_1^2 + 2d_{122}x_1x_2 + d_{123}x_1x_3
$$

and

$$
\frac{\partial f_{a,b,c,d,f,g}}{\partial x_3} = d_{013}x_0x_1 + d_{023}x_0x_2 + d_{113}x_1^2 + d_{123}x_1x_2 + 2d_{133}x_1x_3
$$

Let  $P = \mathbf{P}(x_0, x_1, x_2, x_3)$  be a point of  $\mathcal{F}_{a,b,c,d,f,g}$ . By substituting the coordinates of the point  $P(y_0, y_1, y_2)$  as given in Theorem 5.1 into the four partial derivatives of  $f_{a,b,c,d,f,g}$ , we get four polynomials  $S_1, S_2, S_3$ , and  $S_4$  of degree 6 in three variables  $y_0, y_1$ , and  $y_2$ . We did not present these polynomials explicitly here so that they do not take up much space. The zeros of these polynomials form curves  $S_1$ ,  $S_2$ ,  $S_3$ , and  $S_4$  of degree six in PG(2, F). The point  $P = P(x_0, x_1, x_2, x_3)$  is a singular point if and only if it appears on all the partial derivatives of f*a,b,c,d,f,g* if and only if the curves  $S_1, S_2, S_3$ , and  $S_4$  intersect at *P*. Therefore, the surface is non singular if and only if the curves  $\mathcal{S}_1, \mathcal{S}_2, \mathcal{S}_3$ , and  $\mathcal{S}_4$  never intersect in a point of  $\mathcal{F}_{a,b,c,d,f,g}$ .  $\Box$ 

#### **6. Conclusion**

In this paper, the theorems are proved with a computer-free proof. However, at the beginning of the work, computers are used to produce data. Then, the data is analyzed to make a conjecture. To follow this experimental approach, one needs to have good data. To have good data, one must have the right computer software for computations. We use Orbiter [16], which is an open source. The idea behind that paper has been found this way. This paper shows the uniqueness of the line intersection graph of smooth cubic surfaces with 9 lines. The properties of line intersections are defined. A new form of smooth cubic surfaces  $\mathcal{F}_{a,b,c,d,f,g}$  with at least 9 lines is created. The conditions specify when the surface has 9, when 15, and when 27 lines. This form is exemplified over several fields. Besides, the rational parametrization of points and lines of  $\mathcal{F}_{a,b,c,d,f,q}$ .

The Remarks 3.6 and 3.7 will be considered as a continuation of this work. All possible smooth cubic surfaces with 9 lines can be created by considering Remark 3.6 and 3.7. Studying the sub-configuration of Eckardt points of smooth cubic surfaces with 9 lines and investigating which cases are not covered by the form in this paper are worth reaching that aim. Moreover, the solution can be found for the other cases to generalize. Regarding Remark 3.7, the points  $P_{14}, P_{15}$ , and  $P_{18}$  need to be reset. Maple can be used to construct another form, covering all possible cases. Once these two remarks are considered, the classification problem of smooth cubic surfaces with 9 lines over the small finite fields can be considered. Therefore, the result herein can be verified by the enumerative formula of Das [15] using the Orbit Stabilizer Theorem. Besides, one may generalize certain cases and find a family over Q.

Even though it was a popular topic in the 19th century, it still maintains its current and interesting nature. For recent related studies, see [23, 24].

#### **Author Contributions**

The author read and approved the final version of the paper.

#### **Conflicts of Interest**

The author declares no conflict of interest.

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# New Theory

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# **A Novel Sub-Type Mean Estimator for Ranked Set Sampling with Dual Auxiliary Variables**

**Eda Gizem Koçyiğit<sup>1</sup>**

**Article Info** *Received*: 18 Aug 2023 *Accepted*: 22 Sep 2023 *Published*: 30 Sep 2023 doi[:10.53570/jnt.1346020](https://doi.org/10.53570/jnt.1346020) proposed estimator. This empirical validation is based on numerical computations using Monte Carlo Research Article **Abstract** − This research introduces a novel sub-estimator designed to estimate the population mean under ranked set sampling, motivated by the new concept of a recently introduced sub-ratio estimator. The mathematical formulas of the proposed estimator's mean square error and bias are presented and theoretically contrasted with an analogous estimator found in the existing best sub-estimator literature. In addition to the theoretical analysis, empirical evidence is provided to validate the superiority of the simulations, encompassing synthetic and real data applications. The results underscore the effectiveness of the proposed estimator. Finally, this study discusses the need for further research.

**Keywords** *Efficiency, Monte Carlo simulation, ranked set sampling, sub-estimator, survey sampling*

**Mathematics Subject Classification (2020)** 62D05, 94A17

# **1. Introduction**

Ranked set sampling (RSS) represents a recommended option to simple random sampling (SRS), recognized for its capacity to yield more cost-efficient, time-saving, and efficient outcomes in comparison to SRS [1-3]. The RSS technique hinges on the availability of an auxiliary variable (z), ideally easily accessible and correlated with the study variable  $(q)$ . This auxiliary variable is pivotal in the ranked process and sample selection procedure. However, once the sample selection is finalized, the estimation process exclusively pertains to the study variable. While it is feasible to incorporate the auxiliary variable into the estimation phase through various estimator types (e.g., ratio, product, or regression type estimators using auxiliary variables) [4-13], such approaches often necessitate knowledge of population parameters. The undeniable impact of utilizing auxiliary variable information in enhancing efficiency is well established. It is widely recognized that employing multiple auxiliary variables has the potential to yield even greater efficiency gains than using a single auxiliary variable [14-17]. However, obtaining these parameters for all variables is frequently impractical, limiting the applicability of such estimators on applications. **1De Waster Character Controller Controller Controller Controller Statistics**, Faculty, Theory Experiment of Statistics, Faculty, Theory and Statistics, Faculty of Science, Action (Science, Action (Science), Andre Control

A novel method has been proposed, displaying that the auxiliary variable, a crucial component of the RSS method, can be effectively employed in the estimation phase without requiring population parameters. This breakthrough not only broadens the scope of utility for such estimators but also elevates the overall efficiency of the estimation process. These estimators developed for the RSS method are called sub-estimators [18]. The primary objective of this study is to introduce an estimator that surpasses the existing population mean estimators found in the literature in terms of efficiency. For this purpose, a novel sub-estimator using two

<sup>1</sup> eda.kocyigit@deu.edu.tr (Corresponding Author)

auxiliary variables without their population parameters is proposed to estimate the population mean. Once the theoretical foundations of the proposed estimator have been established, the subsequent aim is to validate its efficiency through numerical investigations.

#### **2. Estimators in Literature**

In the RSS method, a simple random sample of  $m<sup>2</sup>$  of the observations z is initially drawn and subsequently divided into  $m$  sets. Within each set, a ranking is established based on  $z$ , leading to the selection of measurements for the  $q$  corresponding to the units situated along the diagonal. If the required sample size cannot be achieved with m, the process is repeated c times, and the desired  $n = mc$  is obtained. The fundamental mean estimator of the RSS method and its mean square error (MSE) obtained after the cycle c:

$$
\hat{g}_{RSS} = \frac{1}{mc} \sum_{j=1}^{c} \sum_{i=1}^{m} g_{[i,i];j}
$$

and

$$
MSE(\hat{g}_{RSS}) = \overline{G}^2 (\gamma C_g^2 - \omega_g^2)
$$

where  $g_{[i];j}$  is the observation value g of  $i^{th}$  ranked in  $i^{th}$  set and  $j^{th}$  cycle,  $\gamma = 1/mc$ ,  $C_g^2$  is the coefficient of variation of  $g$ ,  $\omega_{g(i)}^2 = \frac{1}{m^2}$  $\frac{1}{m^2 c \overline{G}^2} \sum_{i=1}^m \tau_{g(i)}^2$ ,  $\tau_{g(i)}^2 = \mu_{g(i)} - \overline{G}$ , and  $\mu_{g(i)}$  represents the mean of the  $i^{th}$  order statistics of  $g$ . The RSS sub-ratio estimator and its MSE are described as the following equations:

$$
\hat{g}_{KK1} = \frac{\overline{g}_{RSS}}{\overline{z}_{RSS}} \overline{Z}_{SUB}
$$

and

$$
MSE(\hat{g}_{KK1}) = \overline{G}^{2} [\gamma C_{Z_{SUB}}^{2} - \omega_{Z_{SUB}(i)}^{2} + \gamma C_{g}^{2} - \omega_{g(i)}^{2} - 2(\gamma C_{g_{Z_{SUB}}} - \omega_{g_{Z_{SUB}}(i)})]
$$
(1)

where  $\overline{z}_{RSS} = \frac{1}{m}$  $\frac{1}{mc}\sum_{j=1}^{c} \sum_{i=1}^{m} z_{[i,i];j}$  is the ranked set sample mean of z,

$$
\overline{Z}_{SUB} = \frac{1}{m^2 c} \sum_{j=1}^{c} \sum_{k=1}^{m} \sum_{i=1}^{m} Z_{[i,k];j}
$$

 $C_{Z_{SUB}}^2 = \frac{S_{ZSUB}^2}{Z_{ZUS}}$  $\frac{\sigma_{SUB}}{\overline{Z}_{SUB}^2}$  is the coefficient of variation of the  $Z_{SUB}$ ,

$$
C_{gz_{SUB}} = \rho C_g C_{z_{SUB}}
$$

$$
\omega_{z_{SUB}(i)}^2 = \frac{1}{m^2 c \overline{Z}_{SUB}^2} \sum_{i=1}^m \tau_{z_{SUB}(i)}^2
$$

$$
\omega_{gz_{SUB}(i)} = \frac{1}{m^2 c \overline{GZ}_{SUB}} \sum_{i=1} \tau_{gz_{SUB}(i)}^2
$$

and  $\tau_{gz_{SUB}}^2(i)$  represents the cross-product of the deviation [18]. The RSS sub-exponential ratio type estimator and its MSE are formulated as [18]:

$$
\hat{g}_{KK2} = \overline{g}_{RSS} \exp\left(\frac{\overline{Z}_{SUB} - \overline{z}_{RSS}}{\overline{Z}_{SUB} + \overline{z}_{RSS}}\right)
$$

and

$$
MSE(\hat{g}_{KK2}) = \overline{G}^2 \left[ \gamma C_{Z_{SUB}}^2 - \omega_{Z_{SUB}(i)}^2 + \frac{1}{4} (\gamma C_g^2 - \omega_{g(i)}^2) - (\gamma C_{g_{Z_{SUB}}} - \omega_{g_{Z_{SUB}(i)}}) \right]
$$
(2)

The sub-regression type estimator and its MSE are given as [19]:

$$
\hat{g}_{KR} = \overline{g}_{RSS} + \hat{b}(\overline{Z}_{SUB} - \overline{z}_{RSS})
$$

and

$$
MSE(\hat{g}_{KR})_{SUB}^2 = \left[ \gamma C_g^2 - \omega_g^2 - \frac{(\gamma C_{gz_{SUB}} - \omega_{gz_{SUB}})^2}{\gamma C_{z_{SUB}}^2 - \omega_{z_{SUB}}^2} \right]_{\text{min}}
$$
(3)

where  $\hat{b} = R_{SUB} \frac{\gamma c_{gzSUB} - \omega_{gzSUB}}{\gamma c^2 - \omega^2}$  $\frac{C_{gz_{SUB}} - \omega_{gz_{SUB}}}{\gamma c_{z_{SUB}}^2 - \omega_{z_{SUB}}^2}$  and  $R_{SUB} = \frac{G}{Z_{SUB}}$  $rac{u}{\overline{z}_{SUB}}$ .

#### **3. Proposed Estimator**

Building upon the foundation laid by the sub-ratio, sub-exponential, and sub-regression type estimators [18,19], as well as the multiple auxiliary variable estimators [14-17], we introduce a novel sub-estimator for the population mean under ranked set sampling with two auxiliary variables with the following formulation:

$$
\hat{g}_{PRO} = \left[ \overline{g}_{RSS} \exp\left(\frac{\overline{Z}_{SUB} - \overline{z}_{RSS}}{\overline{Z}_{SUB} + \overline{z}_{RSS}}\right) \exp\left(\frac{\overline{X}_{SUB} - \overline{x}_{RSS}}{\overline{X}_{SUB} + \overline{x}_{RSS}}\right) \right]
$$

where  $x$  is the second auxiliary variable,

$$
\overline{x}_{RSS} = \frac{1}{mc} \sum_{j=1}^{c} \sum_{i=1}^{m} x_{[i,i];j}
$$

and

$$
\overline{X}_{SUB} = \frac{1}{m^2 c} \sum_{j=1}^{c} \sum_{k=1}^{m} \sum_{i=1}^{m} X_{[i,k];j}
$$

To derive the MSE of the proposed estimator, we introduce the following definitions:

$$
\overline{g}_{RSS} = \overline{G}(e_g + 1), \quad \overline{z}_{RSS} = \overline{Z}_{SUB}(e_g + 1), \quad \text{and} \quad \overline{x}_{RSS} = \overline{X}_{SUB}(e_x + 1) \tag{4}
$$

where

$$
E(e_g) = E(e_z) = E(e_x) = 0
$$
  
\n
$$
E(e_g^2) = \gamma C_g^2 - \omega_g^2 = V_{200}
$$
  
\n
$$
E(e_z^2) = \gamma C_{z_{SUB}}^2 - \omega_{z_{SUB}}^2 = V_{020}
$$
  
\n
$$
E(e_x^2) = \gamma C_{x_{SUB}}^2 - \omega_{x_{SUB}}^2 = V_{002}
$$
  
\n
$$
E(e_g e_z) = \gamma C_{g_{Z_{SUB}}} - \omega_{g_{Z_{SUB}}} = V_{110}
$$
  
\n
$$
E(e_g e_x) = \gamma C_{g_{X_{SUB}}} - \omega_{g_{X_{SUB}}} = V_{101}
$$
  
\n
$$
E(e_z e_x) = \gamma C_{z_{SUB}x_{SUB}} - \omega_{z_{SUB}x_{SUB}} = V_{011}
$$
  
\n
$$
E(e_g e_z e_x) = \gamma C_{g_{ZSUB}x_{SUB}} - \omega_{g_{ZSUB}x_{SUB}} = V_{111}
$$

By utilizing the provided definitions in Equation 4, we can express the estimator given in Equation 5 in linear form using Taylor series and second-degree approximation:

$$
\hat{g}_{PRO} = \left[ \overline{G}(e_g + 1) \exp\left(\frac{\overline{Z}_{SUB} - \overline{Z}_{SUB}(e_z + 1)}{\overline{Z}_{SUB} + \overline{Z}_{SUB}(e_z + 1)}\right) \exp\left(\frac{\overline{X}_{SUB} - \overline{X}_{SUB}(e_x + 1)}{\overline{X}_{SUB} + \overline{X}_{SUB}(e_x + 1)}\right) \right]
$$
\n
$$
= \overline{G}(e_g + 1) \exp\left(\frac{-e_g}{e_g + 2}\right) \exp\left(\frac{-e_x}{e_x + 2}\right)
$$
\n
$$
= \overline{G}(e_g + 1) \exp\left[\frac{-e_g}{2}\left(\frac{e_g}{2} + 1\right)^{-1}\right] \exp\left[\frac{-e_x}{2}\left(\frac{e_x}{2} + 1\right)^{-1}\right]
$$
\n
$$
= \overline{G}(e_g + 1) \left(1 - \frac{e_g}{2} + \frac{3e_g^2}{8}\right) \left(1 - \frac{e_x}{2} + \frac{3e_x^2}{8}\right)
$$
\n
$$
= \overline{G} + \overline{G}e_g \left(1 - \frac{e_g}{2} + \frac{3e_g^2}{8} - \frac{e_x}{2} + \frac{3e_x^2}{8} + \frac{e_2e_x}{4}\right)
$$

Hence,

$$
\hat{g}_{PRO} = \overline{G} + \overline{G} \left( e_g - \frac{e_z}{2} - \frac{e_x}{2} + \frac{3e_z^2}{8} + \frac{3e_x^2}{8} + \frac{e_z e_x}{4} - \frac{e_g e_z}{2} - \frac{e_g e_z}{2} + \frac{e_g e_z e_x}{4} \right)
$$
(6)

When subtracting from Equation 5 and then taking the expected value, the bias of the estimator is determined as follows:

$$
B(\hat{g}_{PRO}) = \overline{G}^{2} \left( \frac{3V_{020}}{8} + \frac{3V_{002}}{8} + \frac{V_{011}}{4} - \frac{V_{101}}{2} - \frac{V_{110}}{2} + \frac{V_{111}}{4} \right)
$$

After subtracting  $\overline{G}$  from Equation 6 and subsequently squaring and taking the expected value, upon substituting the equations provided in Equation 5 into their corresponding positions, we arrive at the MSE equation of the estimator as follows:

$$
MSE(\hat{g}_{PRO}) = \overline{G}^2 \left( V_{200} + \frac{V_{020}}{4} + \frac{V_{002}}{4} - V_{110} - V_{101} + \frac{V_{011}}{2} \right)
$$
(7)

1.

For theoretical comparisons, we can express Equations 1–3 in the following forms:

$$
MSE(\hat{g}_{KK1}) = \overline{G}^{2}(V_{200} + V_{020} - 2V_{110})
$$

$$
MSE(\hat{g}_{KK2}) = \overline{G}^{2}(V_{200} + \frac{1}{4}V_{020} - V_{110})
$$

and

$$
MSE(\hat{g}_{KR})\overline{G}^{2}\left(V_{200} + \frac{V_{110}^{2}}{V_{020}}\right)_{min}
$$

Subsequently, by employing Equation 7, we can perform concluding the disparities between the estimators:

*i.* If 
$$
MSE(\hat{g}_{PRO}) < MSE(\hat{g}_{KK1})
$$
 then  $\frac{V_{002} - 4V_{101} + 2V_{011}}{3V_{020} - 4V_{110}} < 1$ .  
\n*ii.* If  $MSE(\hat{g}_{PRO}) < MSE(\hat{g}_{KK2})$  then  $\frac{V_{002} + 2V_{011}}{4V_{110}} < 1$ .  
\n*iii.* If  $MSE(\hat{g}_{PRO}) < MSE(\hat{g}_{KR})$  then  $-\frac{V_{020}(V_{002} - 4V_{110} - 4V_{101} + 2V_{011})}{V_{020}^2 + 4V_{110}^2} >$ 

## **4. Numerical Study**

This section encompasses the numerical computations performed on synthetic and real data sets using the existing and proposed estimator. All calculations are carried out utilizing the R programming language in numerical studies.

# **4.1. Monte Carlo Simulation**

Trivariate random observations  $(G, Z, X)$  are generated from a trivariate normal distribution characterized by parameters:  $\mu_g = \mu_z = \mu_x \approx 5$ , and  $\sigma_g = \sigma_z = \sigma_x \approx 1$ , correlation coefficients  $\rho_{gz} = \rho_{gx} = \rho_{zx} \approx 1$ 0.7, 0.8, and 0.9, along with population sizes  $N = 100$  and 1000. In the existing literature, it is recommended that the set size " $m$ " for the RSS method does not exceed 5, and it is often chosen as  $3, 4$ , or 5. Therefore, 100000 RSS samples were drawn from these populations with a set size of  $m \in \{3,4,5\}$  and cycle  $c \in \{1,2,3\}$ {1,2,3,4}. The estimators' values were computed based on these selected RSS samples.

MSE values for the estimators are computed using Equation 8, while the relative efficiency (RE) values are obtained via Equation 9. In the context of the RE comparison, the reference variable is designated as the RSS basic mean estimator ( $\hat{g}_{RSS}$ ). The results of these comparisons are tabulated in Tables 1-3, each corresponding to distinct parameter combinations. These tables encapsulate valuable insights into the performance of the estimators under varying conditions.

$$
MSE(\hat{g}_h) = \sum_{j=1}^{100000} \frac{(\hat{g}_{hj} - \bar{G})^2}{100000}, \quad h \in \{RSS, KK1, KK2, KR, PRO\}
$$
 (8)

$$
RE(\hat{g}_l) = \frac{MSE(\hat{g}_{RSS})}{MSE(\hat{g}_l)}, \quad l \in \{KK1, KK2, KR, PRO\}
$$
\n<sup>(9)</sup>

#### **Table 1.** RE results for  $\rho_{gz} = \rho_{gx} = \rho_{zx} \approx 0.7$



Boldfaced values indicate the "best" performances.



Boldfaced values indicate the "best" performances.

#### **Table 3.** RE results for  $\rho_{gz} = \rho_{gx} = \rho_{zx} \approx 0.9$



Boldfaced values indicate the "best" performances.

# **4.2. Real Data Application**

This section extends the simulation study employed in the preceding segment to utilize real rather than synthetic data. The dataset originates from the compilation of data across 81 provinces in Türkiye [20]. The variables used within this context are as follows:  $\bm{G}$ : The population of Türkiye in the year 2021;  $\bm{Z}$ : The number of registered vehicles in the year 2017, and  $X$ : The tally of traffic accidents involving fatalities or injuries in 2017. The underlying population parameters are briefly summarized as follows:  $\overline{G} = 1045435, \overline{Z} = 15400.6,$  $\overline{X}$  = 2255.173,  $\sigma_q$  = 1914343,  $\sigma_z$  = 50392.4,  $\sigma_x$  = 2662.278,  $\rho_{gz}$  = 0.97,  $\rho_{gx}$  = 0.87, and  $\rho_{zx}$  = 0.77. 100000 RSS samples were drawn from this population ( $N = 81$ ), considering various set sizes  $m \in \{3,4,5\}$ , and cycle  $c \in \{1,2,3,4\}$ . After calculating the estimator values from these samples, RE values were computed



utilizing Equations 8 and 9. The outcomes of these computations are presented in Table 4, offering insights into the performance of the estimators within the context of real data.

Boldfaced values indicate the "best" performances.

## **5. Conclusion**

This study aims to introduce a new sub-estimator that surpasses the existing alternatives available in the literature regarding efficiency. A novel sub-estimator is proposed to achieve this objective, which utilizes two auxiliary variables without relying on population parameters. Once the theoretical framework of this novel estimator is established, the subsequent step involves substantiating its efficiency through numerical investigations.

The outcomes of the numerical studies indicate that, across all scenarios, the proposed estimator consistently outperforms other estimators in terms of effectiveness, as evidenced by simulation results derived from trivariate normal distributions. Considering the simulation study conducted on real data, the proposed estimator is the optimal choice in most cases, except for instances where  $m = 3$ .

In the simulation study and the real data application, a discernible pattern emerged, underscoring that the effectiveness of the estimator exhibited a positive correlation with the parameter " $m$ ", denoting the number of ranked sets employed in the sampling process. Specifically, as the value of "m" increased, there was a notable enhancement in the estimator's performance.

For prospective research, exploring the proposed estimator's behavior in the context of skewed distributions and under varied sampling methodologies is recommended. Furthermore, an intriguing avenue for exploration involves devising sub-estimator adaptations under the ranked set sampling (RSS) of the proposed estimators for other sampling methods, subsequently assessing their efficiencies.

# **Author Contributions**

The author read and approved the final version of the paper.

# **Conflict of Interest**

The author declares no conflict of interest.

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# **An Investigation into LRS Bianchi I Universe in Brans-Dicke Theory**

# **Halife Cağlar**  $\bullet$

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Research Article

 $\text{doi:10.53570}/\text{jnt.1352470}$  density of the matter. When BD scalar field  $\Phi$  is assumed constant, the attained solutions are **Abstract** − In this study, the homogenous and anisotropic locally rotationally symmetric (LRS) Bianchi type-I universe filled with the bulk viscous and the string cloud matter has been investigated in the Brans-Dicke (BD) scalar-tensor theory. The modified Einstein field equations of the constructed model have been solved by using the relation of the scale factors  $A = B<sup>m</sup>$  and considering the deceleration parameter to be  $q = m - 1$ . It is found that the string does not survive for the model due to the obtained zero string energy density ( $\rho_s = 0$ ) and agrees with some studies in recent years. Moreover, it is possible to say that string matter may be turned into phantom energy over time, depending on the obtained negative rest energy reduced to General Relativity (GR) solutions for the static vacuum. Thus, the constructed model has not allowed the reduction of the BD solutions to GR solutions, including the matter distributions. In addition, some expansion models for choosing a value of constant *m* have been obtained and provided. All the results have been analyzed in detail.

**Keywords** *LRS Bianchi I universe, Brans-Dicke theory, bulk viscous, string cloud, deceleration parameter*

**Mathematics Subject Classification (2020)** 83C05, 83C15

# **1. Introduction**

Gravitation theories are based on Newton's theory, successful in determining the orbits of planets and other celestial bodies. However, this theory was insufficient to explain some cosmic issues. On the other hand, Einstein's General Relativity (GR) theory [1] was the first theory that succeeded in explaining gravity geometrically. But Einstein's theory has been suggested for a static universe. After studies of Friedman [2], Lemaitre [3], and Hubble [4], when considering the expanding universe model, a function determining the expansion of the universe has been needed to add to the field equations in Einstein's GR theory. Einstein has suggested the cosmological constant  $\Lambda$  as the added function to GR theory. Hence, the GR theory has become the most valid gravitational theory, which provides tests, such as the perihelion progression of planets, the gravitational deflection, and the redshift of light. Firstly, scientists have discussed decelerating expanding universe until some observation, such as Supernova IA observation [5], Cosmic Microwave Background radiation (CMB) [6], Wilkinson Microwave Anisotropy Probe (WMAP) [7], etc. Such studies have proven the expanding universe model by accelerating, and then researchers have been focused on studies on this expansion model. Especially alternative gravitation theories, such as Lyra theory [8], Barber's Theory [9], Creation Field

<sup>1</sup>halife@comu.edu.tr (Corresponding Author)

<sup>&</sup>lt;sup>1</sup>Department of Electricity and Energy, Biga Vocational School, Çanakkale Onsekiz Mart University, Çanakkale, Türkiye

Theory [10], Yılmaz Theory [11], and Brans-Dicke theory [12] have been investigated with some cosmic matter distribution to explain accelerating expansion of the universe.

In addition, cosmic strings formed during the phase transition have been suggested to cause variations in the density of the particles, which cause the formation of the galaxies [13, 14]. Because of this assumption, it is believed that cosmic strings are important matter structures to investigate and determine the universe's evolution from early to late. Scientists have recently studied the Brans-Dicke (BD) theory with cosmic string for various space-time models. Delice [15] has investigated BD theory for the generalized cylindrical symmetric universe with cosmic string. Reddy and Rao [16] have researched the axially symmetric metric in the presence of a string dust cloud for BD theory. Besides, in BD theory, Rao and Neelima [17] have analyzed the axially symmetric universe model with bulk string matter distribution. Vidya et al. [18] have solved equations of BD theory for the Bianchi type-III universe with bulk string matters. Sharma and Singh [19] have gotten the solutions of Bianchi II universe for the massive string in BD alternative theory. Trivedi and Bhabor [20] have investigated five-dimensional Bianchi III space-time with dark energy fluid and cosmic string in BD theory.

Moreover, Mahanta et al. [21] have calculated plane-symmetric metric with strange quark matter by attaching to string cloud and bulk viscous matter in BD theory. Chakraborty et al. [22] have analyzed BD theory for the Bianchi type-III universe in the presence of perfect fluid matter. Adhav et al. [23] have gotten exact solutions of the Bianchi III cosmological model for vacuum in BD gravity. Çağlar et al. [24] have investigated the locally rotationally symmetric (LRS) Bianchi type I universe in BD theory by assuming strange quark matter (SQM) coupled with string cloud as matter distributions. Cağlar and Aygün [25] have gotten the exact solutions of higher dimensional Friedman–Walker Friedmann-Robertson-Walker (FRW) universe in the presence of SQM with bulk viscous and cosmic string matter in Lyra theory. Furthermore, higher dimensional FRW universe filled with SQM attached to cosmic string matter for Self Creation Cosmology and BD theory, respectively, have been examined by Cağlar and Aygün  $[26, 27]$ .

Motivated by the studies mentioned above, the LRS Bianchi type-I universe has been studied in BD theory by assuming cosmic string with bulk viscous as matter distribution. This paper is organized as follows: In the second section, the universe model has been constructed, and the modified field equations of the model have been solved. In the last section, all the obtained solutions have been concluded in detail.

#### **2. MEFE's of LRS Bianchi-I Universe for Brans-Dicke Theory**

The modified Einstein field equations (MEFE's) in BD theory [12] have been suggested as:

$$
R_{ij} - \frac{1}{2} R g_{ij} = -\frac{8\pi}{\Phi} T_{ij} - \frac{\omega}{\Phi^2} (\Phi_{,i} \Phi_{,j} - \frac{1}{2} g_{ij} \Phi_{,k} \Phi^{,k}) - \frac{1}{\Phi} (\Phi_{i;j} - g_{ij} \Phi^{,k}_{;k}))
$$
(1)

Here,  $R, R_{ij}$ , and  $\omega$  are Ricci scalar, Ricci tensor, and the coupling constant of BD theory, respectively. It is possible to reduce BD theory to GR theory when  $\omega \to \infty$  and  $\Phi = constant$  [28]. In addition,  $\Phi$ symbolizes BD scalar field and it is given as follows [29]:

$$
\Phi = \Phi_{;k}^{\ k} = \frac{8\pi}{3 + 2\omega}T\tag{2}
$$

Here and after, comma (,) and semicolon (;) represent partial derivative and covariant derivative, respectively. Furthermore,  $T_{ij}$  is a tensor characterized stress energy of the matter in Equation 1. In this study, it has been assumed that the string cloud matter with bulk viscous filled the universe and this matter distribution is written as

$$
T_{ij} = \rho u_i u_j - \rho_s x_i x_j - \eta u_{;k}^k (u_i u_j - g_{ij})
$$
\n
$$
\tag{3}
$$

Here,  $\rho$  and  $\rho_s$  symbolizes the rest energy density of matter and the string energy density, respectively. Moreover, the bulk viscous coefficient has been represented as *η*. Besides, *u<sup>i</sup>* defines 4-velocity of the particle and  $x_i$  is named unit space-like vector determined the string direction [21, 30]. Further, the LRS Bianchi type-I space-time line element can be written as

$$
ds^{2} = dt^{2} - A(t)^{2}dx^{2} - B(t)^{2}(dy^{2} + dz^{2})
$$
\n(4)

Here,  $A(t)$  and  $B(t)$  are time-depending scale factors of the universe [31,32]. Some kinematic quantities, such as spatial volume  $(V)$ , Hubble parameter  $(H)$ , expansion scalar  $(\theta)$ , shear scalar  $(\sigma^2)$ , and deceleration parameter (*q*), for LRS Bianchi type-I universe are provided as follows, respectively:

$$
V = a^3 = AB^2 \tag{5}
$$

$$
H = \frac{a_t}{a} = \frac{A_t}{3A} + \frac{2B_t}{3B} \tag{6}
$$

$$
\theta = \frac{A_t}{3A} + \frac{2B_t}{B}
$$

$$
\sigma^2 = \frac{1}{3} \left[ \frac{A_t}{A} - \frac{B_t}{B} \right]^2
$$

and

$$
q = \frac{d}{dt} \left( \frac{1}{H} \right) - 1 = \frac{2(AB_t - A_t B)^2 - 3AB(A_{tt}B + 2AB_{tt})}{(2AB_t + A_t B)^2}
$$
(7)

Here and after, the lower index *t* represents the derivative with respect to cosmic time *t*. Besides, *a* provided in Equations 5 and 6 symbolizes the average scale factor [33]. By using Equations 1-4, the modified Einstein field equations of LRS Bianchi-I space-time for string cloud coupled with bulk viscous matter have been obtained in BD theory as follows:

$$
\frac{B_t^2}{B^2} + \frac{2B_{tt}}{B} = \frac{1}{\Phi} \left[ 8\pi \eta \left( \frac{A_t}{A} + \frac{2B_t}{B} \right) - \frac{\omega \Phi_t^2}{2\Phi} - \frac{2B_t \Phi_t}{B} - \Phi_{tt} \right]
$$
(8)

$$
\frac{B_t^2}{B^2} + \frac{2A_t B_t}{AB} = \frac{1}{\Phi} \left[ 8\pi \rho + \frac{\omega \Phi_t^2}{2\Phi} - \frac{A_t \Phi_t}{A} - \frac{2B_t \Phi_t}{B} \right]
$$
(9)

$$
\frac{A_{tt}}{A} + \frac{B_{tt}}{B} + \frac{A_t B_t}{AB} = \frac{1}{\Phi} \left[ 8\pi \eta \left( \frac{A_t}{A} + \frac{2B_t}{B} \right) - \Phi_t \left( \frac{A_t}{A} - \frac{B_t}{B} \right) - \frac{\omega \Phi_t^2}{2\Phi} - \Phi_{tt} \right]
$$
(10)

$$
\frac{A_{tt}}{A} + \frac{B_{tt}}{B} + \frac{A_t B_t}{AB} = \frac{1}{\Phi} \left[ 8\pi \eta \left( \frac{A_t}{A} + \frac{2B_t}{B} \right) - \Phi_t \left( \frac{A_t}{A} - \frac{B_t}{B} \right) + 8\pi \rho_s - \frac{\omega \Phi_t^2}{2\Phi} - \Phi_{tt} \right]
$$
(11)

Field equations obtained as in Equations 8-11 have six unknown quantities such as  $A, B, \Phi, \rho, \rho_s$ , and *η*. Therefore, two approximations have been used to solve field equations. The first one is the relationship of the scale factors came from a proportion of the shear scalar and the expansion scalar given as follows:

$$
A = B^k \tag{12}
$$

Here, *k* is an arbitrary constant [24,34,35]. The second assumption is a constant form of deceleration parameter by using Equation 7 given as follows:

$$
q = m - 1 \tag{13}
$$

where the arbitrary constant *m* determines the expansion model of the universe. Deceleration parameter  $q$  specifies the fate of the universe model as follows  $[15, 24, 36]$ :

*i*. When *m* < 0, one gets *q* < −1 named super-exponential expansion

- *ii.* When  $m = 0$ , one gets  $q = -1$  named exponential expansion
- *iii.* When  $0 < m < 1$ , one gets  $-1 < q < 0$  named accelerating power law expansion
- *iv.* When  $m = 1$ , one gets  $q = 0$  named constant expansion
- *v.* When  $m > 1$ , one gets  $q > 0$  named decelerating expansion

By using Equation 13 with Equation 12, the scale factor *B* has been calculated as

$$
B = (s_1 t + s_2)^{\frac{3}{m(k+2)}}\tag{14}
$$

Then, by using Equation 14 in Equation 12, the scale factor *A* has been attained as follows:

$$
A = (s_1 t + s_2)^{\frac{3k}{m(k+2)}} \tag{15}
$$

Considering the scale factors' values in Equations 14 and 15 in field Equations 8–11, the BD scalar field  $\Phi$ , the rest energy density  $\rho$ , the string energy density  $\rho_s$ , and the bulk viscous coefficient  $\eta$  have been obtained as follows, respectively:

$$
\Phi = s_3(s_1t + s_2)^{\frac{m-3}{m}}
$$

$$
\rho = \frac{\left[ (k+2)^2 \left( 6\omega + 6m - \omega m^2 - 9\omega \right) - 18k(k+2) - 54 \right] s_3 s_1^2}{16\pi m^2 (k+2)^2 (s_1t + s_2)^{\frac{m+3}{m}}}
$$

$$
\rho_s = 0
$$

and

$$
\eta = \frac{\left[ (k+2)^2 \left( \omega m^2 - 6\omega - 6m + 9\omega \right) + 18k(k+2) + 54 \right] s_1 s_3}{48\pi m^2 (k+2)^2 (s_1 t + s_2)^{\frac{3}{m}}}
$$

Furthermore, kinematic quantities, such as spatial volume  $(V)$ , Hubble parameter  $(H)$ , expansion scalar  $(\theta)$ , and shear scalar  $(\sigma^2)$  for the constructed model are calculated as follows, respectively:

$$
V = (s_1t + s_2)^{\frac{1}{m}}
$$

$$
H = \frac{s_1}{m(s_1t + s_2)}
$$

$$
\theta = \frac{3s_1}{m(s_1t + s_2)}
$$

and

$$
\sigma^{2} = \frac{3(k-1)^{2}s_{1}^{2}}{m^{2}(k+2)^{2}(s_{1}t+s_{2})^{2}}
$$

All the obtained solutions of constructed model have been investigated in the next section.

#### **3. Discussions**

In this paper, the bulk viscous coupled with the string cloud matter distribution has been investigated in BD theory for LRS Bianchi type-I space-time. By considering the scale factor values given in Equations 14 and 15 within Equation 4, the new line element of the LRS Bianchi-I universe can be rewritten as

$$
ds^{2} = dt^{2} - (s_{1}t + s_{2})^{\frac{6k}{m(k+2)}}dx^{2} - (s_{1}t + s_{2})^{\frac{6}{m(k+2)}}(dy^{2} + dz^{2})
$$
\n(16)

It is obviously seen that *s*<sup>1</sup> and *k* in Equation 16 are important constants to determine meaningful the universe model and must be  $s_1 \neq 0$ ,  $k \neq -2$ , and  $k \neq 0$ . Besides, it is well know that the universe is expanding with acceleration. When  $0 < m < 1$ , the deceleration parameter value has given accelerating power law expansion universe. At this condition, the evolution of scale factor  $A(t)$  has been shown in Figure 1. As seen in the figure, the expansion of the constructed universe model is as

expected in given situations.



**Figure 1.** Evolutions of the scale factor  $A(t)$  for the accelerating expanding universe  $(s_1 = s_2 = k = 1$ and  $m = 0.5$ )

Moreover, the string energy density  $\rho_s$  vanishes for constructed model. Thus, it can be said that the string matter does not survive in the model LRS Bianchi type-I universe for BD theory with the bulk viscous and the string cloud. This finding agrees with the studies of Zel'dovich [13], Çağlar and Aygün [25–27], Sahoo and Mishra [37], Kiran and Reddy [38], etc. In addition, the obtained all the solutions give the negative energy density of particle. Hence, it is possible that the cosmic string may have disappeared by time and turned into a phantom [39]. Rest energy density and bulk viscous coefficient have been shown in Figure 2.



Figure 2. Evolutions of the rest energy density and the bulk viscous coefficient for the accelerating expanding universe  $(s_1 = s_2 = s_3 = k = 1, \omega = 2, \text{ and } m = 0.5)$ 

Furthermore, the BD scalar field Φ has been obtained time dependent and shown in Figure 3 for the accelerating expanding universe. It is observed that the scalar field of the theory is inversely proportional to time, intense at the beginning, and loses its effect over time.



**Figure 3.** Evolutions of the BD scalar field for the accelerating expanding universe  $(s_1 = s_2 = s_3 = 1$ and  $m = 0.5$ )

Considering the kinematic quantities, it is obtained that the universe has expansion with accelerating. This expansion is observed in Figures 4 and 5. Further, this result agrees with some studies [4–7].



**Figure 4.** Evolutions of spatial volume for accelerating expanding universe  $(s_1 = s_2 = s_3 = 1$  and  $m = 0.5$ 



**Figure 5.** Evolutions of the Hubble parameter, the expansion scalar, and the shear scalar for the accelerating expanding universe  $(s_1 = s_2 = s_3 = 1, k = 2, \text{ and } m = 0.5)$ 

The obtained all the solutions have been calculated depending on some integer values of arbitrary constant *m* to investigate expansion type of the universe. By assuming  $m = 2$ ,  $m = 1$ , and  $m = 0$ , one obtains the results of the constructed model for decelerating expansion, the constant expansion, and the exponential expansion universe, respectively. Additionally, these results have been provided in Table 1.

Quantities	Decelerating Expansion $q = 1$	Constant Expansion $q = 0$	Exponential Expansion $q = -1$
Scale Factor A	$(s_1t + s_2)^{\frac{3k}{2k+4}}$	$(s_1t + s_2)^{\frac{3k}{k+2}}$	$s_A^k e^{s_5kt}$
Scale Factor $B$	$(s_1t + s_2)^{\frac{3}{2k+4}}$	$(s_1t+s_2)^{\frac{3}{k+2}}$	$s_4e^{s_5t}$
BD Scalar Field $\Phi$	$\frac{s_3}{\sqrt{s_1t+s_2}}$	$\frac{s_3}{(s_1t+s_2)^2}$	$\frac{s_6}{e^{(k+2)s_5t}}$
Rest Energy Density $\rho$	$- \frac{\left[(\omega \!+\! 6)k^2 \!+\! (4\omega \!-\! 12)k \!+\! 4\omega \!+\! 6\right]s_3 s_1^2}{96\pi (k\!+\!2)^2\left(s_1 t\!+\!s_2\right)^{\frac{3}{2}}}$	$-\frac{\left[ (\omega+3)2k^2 + (4\omega+3)2k + 8\omega + 15 \right]s_3s_1^2}{8\pi (k+2)^2(s_1t+s_2)^4}$	$-\frac{\left[ (\omega +2) k^2 + 4 (\omega +1) k + 4 \omega +6) \right] s_6 s_5^2}{16 \pi e^{(k+2)s_5 t}}$
Bulk Viscous Coefficient $\eta$	$\frac{\left[ (\omega+6)k^2+(4\omega-12)k+4\omega+6\right] s_1 s_3}{96\pi (k+2)^2 (s_1 t+s_2)^{\frac{3}{2}}}$	$\frac{[(\omega+3)2k^2+(4\omega+3)2k+8\omega+15]s_1s_3}{24\pi(k+2)^2(s_1t+s_2)^3}$	$\frac{\left[(\omega +2) \, k^2+4 (\omega +1) \, k+4 \omega +6)\right]s_5s_6}{16 \pi \, e^{(k+2)s_5t}}$
Spatial Volume V	$\sqrt{s_1t+s_2}$	$s_1t+s_2$	$s_4e^{\frac{(k+2)s_5t}{3}}$
Hubble Parameter $H$	$\frac{s_1}{2(s_1t+s_2)}$	$\frac{s_1}{(s_1t+s_2)}$	$\frac{(k+2) s_5}{3}$
Expansion Scalar $\theta$	$\frac{3s_1}{2(s_1t+s_2)}$	$\frac{3s_1}{(s_1t+s_2)}$	$(k+2)s_5$
Shear Scalar $\sigma^2$	$\frac{3(k-1)^2s_1^2}{4(k+2)^2(s_1t+s_2)^2}$	$\frac{3(k-1)^2s_1^2}{(k+2)^2(s_1t+s_2)^2}$	$\frac{(k-1)^2s_5^2}{3}$

**Table 1.** Results of Constructed Model for the Deceleration Parameters

It is possible to get GR solutions of the LRS Bianchi type-I space-time with the bulk string cloud matter by assuming  $\Phi = constant$  in BD theory. To get the constant BD scalar field,  $s_3$  must be zero in the model results. In this situation, the energy density and the bulk viscous coefficient became zero, and it gives the static vacuum solutions. Thus, the obtained results do not allow GR solutions of the constructed model by reducing BD solutions with  $\Phi = constant$ .

# **4. Conclusion**

This study investigated BD theory in the LRS Bianchi type-I space-time for bulk viscous and string cloud matter distribution. In the future, it is worth investigating LRS Bianchi I universe filled with bulk string matter or various cosmic matter in alternative gravitation theories such as  $f(G)$ ,  $f(Q)$ . and  $f(Q, T)$ .

# **Author Contributions**

The author read and approved the final version of the paper.

#### **Conflicts of Interest**

The author declares no conflict of interest.

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# **An Alternative Method for Determination of the Position Vector of a Slant Helix**

 $\mathbf{G}$ izem Güzelkardeşler $^{1}$   $\mathbf{D}$  , Burak Şahiner $^{2}$   $\mathbf{D}$ 

**Article Info** *Received*: 7 Sep 2023 *Accepted*: 22 Sep 2023 *Published*: 30 Sep 2023 doi[:10.53570/jnt.1356697](https://doi.org/10.53570/jnt.1356697) helices for chosen curvature and torsion functions as an application of the proposed method. Research Article **Abstract** − In this paper, we provide an alternative method to determine the position vector of a slant helix with the help of an alternative moving frame. We then construct a vector differential equation in terms of the principal normal vector of a slant helix using an alternative moving frame. By solving this vector differential equation, we determine the position vector of the slant helix. Afterward, we obtain parametric representations of some examples of slant Finally, we discuss the method and whether further research should be conducted or not.

**Keywords** *Alternative moving frame, intrinsic equations, position vector, slant helix*

**Mathematics Subject Classification (2020)** 53A04, 34A05

# **1. Introduction**

In differential geometry, the theory of curves is one of the main study areas. The theory of curves is generally studied with the well-known Frenet frame. Many geometric properties of differentiable curves can be defined with the help of this frame. In addition, the determination of the characterization of some special curves can be achieved by the curvatures of the Frenet frame. Among these special curves, various types of helices, including general helices, circular helices, and slant helices, are the curves that attract the most attention from researchers. A general helix (formerly called cylindrical helix) is defined by the property that all tangent vectors along the curve make a constant angle with a fixed direction. A necessary and sufficient condition for a curve to be a general helix is that the ratio of torsion to curvature is a constant [1]. There are numerous uses for general helices in different branches of science, such as biology, fractal geometry, computer-aided geometric design, engineering, and architecture [2–5]. After a long time since the concept of general helix hs been introduced, a new curve called slant helix has been defined with a similar idea. A slant helix is a curve whose principal normal vectors make a constant angle with a fixed straight line, which is the axis of the slant helix [6]. A necessary and sufficient condition for a curve to be a slant helix is that the ratio  $\frac{\kappa^2}{\kappa^2+\kappa^2}$  $\frac{\kappa^2}{(\kappa^2+\tau^2)^{3/2}}\left(\frac{\tau}{\kappa}\right)$  $(\frac{\tau}{\kappa})'$  is constant, where  $\kappa$  and  $\tau$  are curvature and torsion functions of the curve, respectively [6].

According to the fundamental theorem for curves, given two continuous functions of one parameter, a space curve can be determined uniquely up to rigid motion for which the two functions are its curvature and torsion [7]. The problem of determining the position vector of this curve is known as

<sup>&</sup>lt;sup>1</sup> gizemguzelkardesler@ktu.edu.tr; <sup>2</sup>burak.sahiner@cbu.edu.tr (Corresponding Author)

 $1$ Department of Mathematics, Faculty of Science, Karadeniz Technical University, Trabzon, Türkiye

<sup>&</sup>lt;sup>2</sup>Department of Mathematics, Faculty of Science and Letters, Manisa Celal Bayar University, Manisa, Türkiye

solving natural or intrinsic equations [1]. Although the problem has a solution for an arbitrary curve in Galilean space, it remains an open problem for Euclidean and Minkowski spaces [8, 9]. For the problem of determining the position vector of a general helix, a method depending on solving a vector differential equations constructed with the help of the Frenet frame has been proposed in [10]. Using this method, the parametric representation of the position vector of a general helix with the known curvature and torsion functions has been found. Then, a similar method has been used to determine a slant helix's position vector in [11]. After that, similar techniques have been applied for determining position vectors of some special curves with the help of various moving frames, such as the Frenet frame, the type-2 Bishop frame, the Darboux frame, and the alternative moving frame, in various spaces, Euclidean space, Minkowski space, and Galilean space [9, 12–19].

In this paper, we propose an alternative method to the existing methods in the literature to determine the position vector of a slant helix. We first rewrite the derivative formulae of the alternative moving frame to obtain a simpler differential equation. Then, we construct a vector differential equation in terms of the principal normal vector with the help of the new derivative formulae of the alternative moving frame. By solving this vector differential equation, we find the principal normal vector of the slant helix and thus determine the position vector of the curve. Finally, applying the proposed method, we obtain parametric representations of some examples of slant helices for chosen curvature and torsion functions.

#### **2. Preliminaries**

Let  $\alpha = \alpha(s)$  be a unit speed Frenet curve in  $E^3$ , that is,  $\langle \alpha'(s), \alpha'(s) \rangle = 1$  and  $\alpha''(s) \neq 0$ . The Frenet frame along the curve  $\alpha$  consists of three mutually orthonormal vectors defined by

$$
\boldsymbol{T}(s) = \alpha'(s), \quad \boldsymbol{N}(s) = \frac{1}{\|\boldsymbol{T}'(s)\|} \boldsymbol{T}'(s), \quad \text{and} \quad \boldsymbol{B}(s) = \boldsymbol{T}(s) \times \mathbf{N}(s) \tag{1}
$$

where  $T(s)$ ,  $N(s)$ , and  $B(s)$  are called tangent vector, principal normal vector, and binormal vector, respectively. The derivative formulae of the Frenet frame also known as Frenet formulae can be provided as follows:

$$
\begin{bmatrix}\n\boldsymbol{T}'(s) \\
\boldsymbol{N}'(s) \\
\boldsymbol{B}'(s)\n\end{bmatrix} = \begin{bmatrix}\n0 & \kappa(s) & 0 \\
-\kappa(s) & 0 & \tau(s) \\
0 & -\tau(s) & 0\n\end{bmatrix} \begin{bmatrix}\n\boldsymbol{T}(s) \\
\boldsymbol{N}(s) \\
\boldsymbol{B}(s)\n\end{bmatrix}
$$
\n(2)

where  $\kappa$  and  $\tau$  are called curvature and torsion functions of the curve  $\alpha$ , respectively. These functions, also called Frenet curvatures, are defined by

$$
\kappa(s) = ||\mathbf{T}'(s)||
$$
 and  $\tau(s) = -\langle \mathbf{B}'(s), \mathbf{N}(s) \rangle$ 

In Euclidean 3-space, apart from the Frenet frame, many moving frames has been defined to study the differential geometric properties of curves. One of them is the alternative moving frame. The alternative moving frame consists of three mutually orthonormal vectors. These vectors are the unit principal normal vector  $N(s)$ , also included in the Frenet frame, the unit vector  $C(s)$  defined by  $C(s) = \frac{N'(s)}{\|N'(s)\|}$ , and the unit vector  $W(s)$  defined by  $W(s) = N(s) \times C(s)$ , also normalized instantaneous rotation vector of the Frenet frame [20, 21]. The derivative formulae of the alternative moving frame are as follows [20]:

$$
\begin{bmatrix}\n\mathbf{N}'(s) \\
\mathbf{C}'(s) \\
\mathbf{W}'(s)\n\end{bmatrix} = \begin{bmatrix}\n0 & f(s) & 0 \\
-f(s) & 0 & g(s) \\
0 & -g(s) & 0\n\end{bmatrix} \begin{bmatrix}\n\mathbf{N}(s) \\
\mathbf{C}(s) \\
\mathbf{W}(s)\n\end{bmatrix}
$$
\n(3)
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where the functions

$$
f = \sqrt{\kappa^2 + \tau^2} \tag{4}
$$

and

$$
x^2 = (\tau)^{\prime}
$$

$$
g = \frac{\kappa^2}{\kappa^2 + \tau^2} \left(\frac{\tau}{\kappa}\right)'
$$
\n(5)

are called the first and the second alternative curvatures of the curve  $\alpha$ , respectively. From Equalities 4 and 5, we have the followings [22]:

$$
\kappa(s) = f(s)\cos\left(\int g(s)ds\right) \tag{6}
$$

and

$$
\tau(s) = f(s) \sin\left(\int g(s)ds\right) \tag{7}
$$

The alternative curvatures play a major role in the characterizations of curves. The following theorem supports this idea.

**Theorem 2.1.** [20] Let  $\alpha$  be a curve in  $E^3$  with alternative curvatures  $f$  and  $g$ . The curve  $\alpha$  is a slant helix if and only if the function

$$
\sigma = \frac{g}{f} \tag{8}
$$

is a constant.

# **3. Determination of the Position Vector of a Slant Helix**

In this section, the problem of determining a slant helix's position vector is solved using a method based on the alternative moving frame. To achieve this, we first rewrite the derivative formulae of the alternative moving frame with a new parameter. Then, we construct a vector differential equation in terms of the principal normal vector by using these new formulae. By solving this vector differential equation, we obtain the position vector of the slant helix in parametric form. Before constructing the vector differential equation, it will be more useful to use new derivative formulae obtained by the parameter transformation  $\theta = \int f(s)ds$  instead of the derivative formulae given in Equality 3 where *f*(*s*) is the first alternative curvature. According to this new parameter *θ*, the derivative formulae of the alternative moving frame become the following form [19]:

$$
\begin{bmatrix}\n\mathbf{N}'(\theta) \\
\mathbf{C}'(\theta) \\
\mathbf{W}'(\theta)\n\end{bmatrix} = \begin{bmatrix}\n0 & 1 & 0 \\
-1 & 0 & \sigma(\theta) \\
0 & -\sigma(\theta) & 0\n\end{bmatrix} \begin{bmatrix}\n\mathbf{N}(\theta) \\
\mathbf{C}(\theta) \\
\mathbf{W}(\theta)\n\end{bmatrix}
$$
\n(9)

such that  $\sigma(\theta) = \frac{g(\theta)}{f(\theta)}$ .

**Theorem 3.1.** Let  $\alpha = \alpha(s)$  be a unit speed slant helix and  $\alpha = \alpha(\theta)$  be another parametric representation of this curve where  $\theta = \int f(s)ds$ . The principal normal vector  $\mathbf{N}(\theta)$  satisfies the following vector differential equation

$$
\mathbf{N}'''(\theta) + (1 + \sigma^2)\mathbf{N}'(\theta) = 0\tag{10}
$$

where  $\sigma(\theta) = \frac{g(\theta)}{f(\theta)}$ .

Proof*.*

From the second equation of Equality 9, the vector  $W(\theta)$  can be written as

$$
\mathbf{W}(\theta) = \frac{1}{\sigma(\theta)} \left( \mathbf{C}'(\theta) + \mathbf{N}(\theta) \right) \tag{11}
$$

Differentiating the first equation of Equality 9,

$$
\mathbf{N}''(\theta) = \mathbf{C}'(\theta) \tag{12}
$$

Substituting Equality 12 into Equality 11,

$$
\mathbf{W}(\theta) = \frac{1}{\sigma(\theta)} \left( \mathbf{N}''(\theta) + \mathbf{N}(\theta) \right) \tag{13}
$$

By differentiating Equality 13 and by using the first and third equations of Equality 9,

$$
\left(\frac{1}{\sigma(\theta)}\left(\mathbf{N}''(\theta) + \mathbf{N}(\theta)\right)\right)' + \sigma(\theta)\mathbf{N}'(\theta) = 0\tag{14}
$$

According to Theorem 2.1,  $\sigma(\theta)$  is a constant. Thus, Equality 14 becomes Equality 10 which completes the proof.  $\square$ 

Equality 10 is a third-order vector differential equation with constant coefficients. The principal normal vector  $N$  can be found by solving this equation. The following theorem provides a solution for the problem of determining the position vector of a slant helix given the curvature and torsion functions.

**Theorem 3.2.** Let  $\kappa(s) > 0$  and  $\tau(s)$  be two differentiable functions on an interval *I*. On this interval, the position vector  $\alpha(s) = (\alpha_1(s), \alpha_2(s), \alpha_3(s))$  of a slant helix for which *s*,  $\kappa(s)$  and  $\tau(s)$ become the arc-length, the curvature and the torsion, respectively, can be determined up to rigid motions in Euclidean 3-space as follows:

$$
\begin{cases}\n\alpha_1(s) = \frac{1}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) \cos \left( \sqrt{1+\sigma^2} \int f(s) ds \right) ds \right) ds \\
\alpha_2(s) = \frac{1}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) \sin \left( \sqrt{1+\sigma^2} \int f(s) ds \right) ds \right) ds \\
\alpha_3(s) = \frac{\sigma}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) ds \right) ds\n\end{cases}
$$
\n(15)

where  $\sigma = \frac{g}{f}$  $\frac{g}{f}$  and *f* and *g* are the alternative curvatures defined in Equalities 4 and 5, respectively. Proof*.*

Let  $\alpha = \alpha(s)$  be a unit speed slant helix in  $E^3$  and  $\alpha = \alpha(\theta)$  be another parametrization of the same where  $\theta = \int f(s)ds$ . We can write the principal normal vector  $\mathbf{N}(\theta)$  of the curve  $\alpha(\theta)$  with the standard basis of  $E^3$  as  $\mathbf{N}(\theta) = N_1(\theta)\mathbf{e_1} + N_2(\theta)\mathbf{e_2} + N_3(\theta)\mathbf{e_3}$ . We can choose the axis of the slant helix to be parallel to  $e_3$  without losing generality. Note that selections of different axes will produce the same slant helix up to rigid motion in Euclidean 3-space. Since the vector  $\mathbf{N}(\theta)$  makes a constant angle with the axis, then

$$
\langle \mathbf{N}(\theta), \mathbf{e}_3 \rangle = N_3(\theta) = n \tag{16}
$$

where *n* is a constant real number. It is clear that  $N_3(\theta) = n$  satisfies Equality 10. Additionally, since  $N(\theta)$  is a unit vector,

$$
N_1^2(\theta) + N_2^2(\theta) = 1 - n^2 \tag{17}
$$

From the general solution of Equality 17, the components  $N_1$  and  $N_2$  can be written as

$$
N_1(\theta) = \sqrt{1 - n^2} \cos(t(\theta))
$$
 and  $N_2(\theta) = \sqrt{1 - n^2} \sin(t(\theta))$ 

where *t* is a function of  $\theta$ . The components  $N_1$  and  $N_2$  should satisfy Equality 10. Substituting  $N_1(\theta)$ and  $N_2(\theta)$  into Equality 10, we obtain the following equalities:

$$
(3t't'')\cos t + \left(-(t')^3 + t''' + (1+(\sigma)^2)t'\right)\sin t = 0\tag{18}
$$

$$
(3t't'')\sin t - \left(-(t')^3 + t''' + (1+(\sigma)^2)t'\right)\cos t = 0\tag{19}
$$

From Equalities 18 and 19, we have the following equalities:

$$
3t''t' = 0\tag{20}
$$

$$
t'' - (t')^{3} + (1 + \sigma^{2})t' = 0
$$
\n(21)

From Equality 20,

$$
t(\theta) = c_1 \theta + c_2 \tag{22}
$$

where  $c_1$  and  $c_2$  are constants. If we change the parameter  $t \to t + c_2$ , then

$$
t(\theta) = c_1 \theta \tag{23}
$$

Substituting Equality 23 into Equality 21,

$$
c_1(-c_1^2 + 1 + \sigma^2) = 0 \tag{24}
$$

Solving Equality 24,

$$
c_1 = \pm \sqrt{1 + \sigma^2} \tag{25}
$$

Substituting Equality 25 into Equality 23, the function  $t(\theta)$  can be obtained as

$$
t = \pm \sqrt{1 + \sigma^2} \theta
$$

Therefore, the vector  $\mathbf{N}(\theta)$  can be found as follows:

$$
\mathbf{N}(\theta) = \left(\sqrt{1 - n^2} \cos\left(\sqrt{1 + \sigma^2} \theta\right), \pm \sqrt{1 - n^2} \sin\left(\sqrt{1 + \sigma^2} \theta\right), n\right) \tag{26}
$$

Moreover, the constant *n* in Equality 26 can be written in terms of  $\sigma$ . Differentiating Equality 16 and using the first equation of Equality 9,

$$
\langle \mathbf{C}(\theta), \mathbf{e}_3 \rangle = 0 \tag{27}
$$

Differentiating Equality 27 and using the second equation of Equality 9,

$$
\langle \, \textbf{W}(\theta), \textbf{e}_3 \rangle = \frac{n}{\sigma}
$$

Since the axis of the slant helix is a unit vector,

$$
n = \pm \frac{\sigma}{\sqrt{1 + \sigma^2}}\tag{28}
$$

Substituting Equality 28 into Equality 26, the principal normal vector  $\mathbf{N}(\theta)$  of the slant helix can be written in terms of the parameter  $\theta$  as

$$
\mathbf{N}(\theta) = \frac{1}{\sqrt{1+\sigma^2}} \left( \cos\left(\sqrt{1+\sigma^2}\theta\right), \pm \sin\left(\sqrt{1+\sigma^2}\theta\right), \pm \sigma \right)
$$

or in terms of the arc-length parameter *s* as

$$
\mathbf{N}(s) = \frac{1}{\sqrt{1+\sigma^2}} \left( \cos\left(\sqrt{1+\sigma^2} \int f(s)ds\right), \pm \sin\left(\sqrt{1+\sigma^2} \int f(s)ds\right), \pm \sigma \right) \tag{29}
$$

Moreover, from Equalities 1 and 2, the curve  $\alpha$  can be written as

$$
\alpha(s) = \int \left( \int \kappa(s) \mathbf{N}(s) ds \right) ds
$$

By using Equality 4 and the parameter transformation  $\theta = \int f(s)ds$ ,

$$
\alpha(\theta) = \int \frac{1}{f(\theta)} \left( \int \cos(\sigma\theta) \mathbf{N}(\theta) d\theta \right) d\theta
$$

If the curve  $\alpha$  is written again in terms of the arc length parameter *s*,

$$
\alpha(s) = \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) \mathbf{N}(s) ds \right) ds \tag{30}
$$

By choosing the positive sign in Equality 29 and substituting it into Equality 30, the position vector  $\alpha(s) = (\alpha_1(s), \alpha_2(s), \alpha_3(s))$  of the slant helix can be found as

$$
\begin{cases}\n\alpha_1(s) = \frac{1}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) \cos \left( \sqrt{1+\sigma^2} \int f(s) ds \right) ds \right) ds \\
\alpha_2(s) = \frac{1}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) \sin \left( \sqrt{1+\sigma^2} \int f(s) ds \right) ds \right) ds \\
\alpha_3(s) = \frac{\sigma}{\sqrt{1+\sigma^2}} \int \left( \int f(s) \cos \left( \sigma \int f(s) ds \right) ds \right) ds\n\end{cases}
$$
\n(31)

Given the curvature and torsion functions, the functions  $f, g$ , and  $\sigma$  can be determined using Equalities 4 and 5 and Theorem 2.1. Substituting f and  $\sigma$  into Equality 31, the position vector of the slant helix can be obtained.  $\quad \Box$ 

### **4. Illustrative Examples**

In this section, we obtain parametric representations of some examples of slant helices for chosen some special curvature and torsion functions.

**Example 4.1.** If  $\kappa = \cos(s)$  and  $\tau = \sin(s)$ , then  $f = 1$ ,  $g = 1$ , and  $\sigma = 1$  from Equalities 4, 5, and 8. From Equality 15, the position vector  $\alpha(s)$  of the slant helix can be found as

$$
\alpha(s) = \left(-\frac{3\sqrt{2}}{2}\cos(\sqrt{2}s)\cos(s) - 2\sin(\sqrt{2}s)\sin(s), -\frac{3\sqrt{2}}{2}\sin(\sqrt{2}s)\cos(s) + 2\cos(\sqrt{2}s)\sin(s), -\frac{\sqrt{2}}{2}\cos(s)\right)
$$



**Figure 1.** Slant helix with  $\kappa = \cos(s)$  and  $\tau = \sin(s)$ 

**Example 4.2.** Let curvature and torsion functions be given as  $\kappa = \frac{1}{\kappa}$  $\frac{1}{(1+s^2)^{3/2}}$  and  $\tau = \frac{s}{(1+s^2)}$  $\frac{s}{(1+s^2)^{3/2}}$ respectively. From Equalities 4, 5, and 8,  $f = \frac{1}{1+r}$  $\frac{1}{1+s^2}$ ,  $g = \frac{1}{1+s^2}$  $\frac{1}{1+s^2}$ , and  $\sigma = 1$ . Using Equality 15, the position vector of the slant helix can be obtained in the parametric representation as

$$
\alpha(s) = \left(\frac{\sqrt{2}}{2}\sqrt{1+s^2}\cos\left(\sqrt{2}\arctan(s)\right), -\frac{\sqrt{2}}{2}\sqrt{1+s^2}\sin\left(\sqrt{2}\arctan(s)\right), \frac{\sqrt{2}}{2}\sqrt{1+s^2}\right)
$$
\n  
\n1.00  
\n0.80  
\n0.81  
\n0.82  
\n0.83  
\n0.84  
\n0.85  
\n0.87  
\n0.89  
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\n0.82  
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\n0.80  
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**Figure 2.** Slant helix with  $\kappa = \frac{1}{(1+\epsilon)}$  $\frac{1}{(1+s^2)^{3/2}}$  and  $\tau = \frac{s}{(1+s^2)^{3/2}}$  $(1+s^2)^{3/2}$ 

**Example 4.3.** Let curvature and torsion functions be given as  $\kappa = \frac{\cos(\sqrt{s})}{2\sqrt{s}}$  $\frac{\sin(\sqrt{s})}{2\sqrt{s}}$  and  $\tau = \frac{\sin(\sqrt{s})}{2\sqrt{s}}$  $rac{\mu(\nabla^S)}{2\sqrt{s}}$ , respectively. The functions *f*, *g*, and  $\sigma$  can be found as  $f = \frac{1}{2}$  $\frac{1}{2\sqrt{s}}, g = \frac{1}{2\sqrt{s}}$  $\frac{1}{2\sqrt{s}}$ , and  $\sigma = 1$ , respectively. Using Equality 15, the position vector  $\alpha(s)$  of the slant helix can be expressed in the parametric representation as

√

$$
\alpha(s) = \left(-\frac{\sqrt{2}}{2}\cos\left(\sqrt{2s}\right)\sin\left(\sqrt{s}\right) + \sin\left(\sqrt{2s}\right)\cos\left(\sqrt{s}\right), -\frac{\sqrt{2}}{2}\sin\left(\sqrt{2s}\right)\sin\left(\sqrt{s}\right) - \cos\left(\sqrt{2s}\right)\cos\left(\sqrt{s}\right), \frac{\sqrt{2}}{2}\sin\left(\sqrt{s}\right)
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# **5. Conclusion**

In this paper, we give a method for the determination of the position vector of a slant helix. This method basically includes two important steps. The first step is constructing a vector differential equation in terms of the principal normal vector of the curve with the help of the alternative moving frame. The second step is determining the position vector of the slant helix by solving the equation in the first step. Thanks to this method, the problem of determining the position vector of a curve given the curvature and torsion functions is solved for a slant helix case. As an application of the method, some examples of slant helices in parametric form are obtained for given special curvature and torsion functions. It is thought that the obtained examples of slant helices will contribute to the variety of examples of slant helices in the literature.

The problem of determining the position vector of a slant helix, solved in this paper, is also discussed in [11] and [19]. In [11], a method based on solving a vector differential equation in terms of the principal normal vector constructed with the help of Frenet formulae is presented. Since the characterization of a slant helix cannot be used directly in this vector differential equation, this method for determining the position vector of a slant helix involves much more complicated mathematical operations than the method in the present paper. In [19], the authors give a method based on the alternative moving frame to solve the same problem. Unlike the method in the present paper, the vector differential equation is constructed in terms of the vector  $C$  of the alternative moving frame in the method given in [19]. Moreover, this method has been developed to determine a slant helix's position vector in Minkowski 3-space, not Euclidean 3-space. The necessity of finding all the vectors of the alternative moving frame to determine the position vector of the slant helix is an important disadvantage of this method. In light of all these comparisons, it is hoped that the method in the present paper will be an important alternative to the methods in the literature for the problem of determining the position vector of a slant helix.

## **Author Contributions**

All the authors equally contributed to this work. This paper is derived from the first author's master's thesis supervised by the second author. They all read and approved the final version of the paper.

# **Conflicts of Interest**

All the authors declare no conflict of interest.

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