

A Study On The Hyper Darboux Lines In A Generalised Weyl Space

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In the previous works, some properties of the Hyper Darboux lines which are the generalized form of Darboux lines defined in 3- dimensional Euclid space were defined in Riemannian and Kaehlerian space [6], [7]. In this work, some properties of a Hyper Darboux line of various orders in GW_n hyperspace of GW_{n+1} space are obtained. In addition to this, a Hyper Darboux line of order zero in GW_n hyperspace of GW_{n+1} space is considered and the equations involving the second fundamental tensor of the subspace are deduced.

1. INTRODUCTION

An n- dimensional manifold GW_n is said to be a generalized Weyl space if it has an asymmetric conformal metric tensor g_{ij} and asymmetric connection ∇_k satisfying the compatibility condition given by the equation

$$\nabla_k g_{ij} = 2T_k g_{ij} \quad (1.1)$$

where T_k denotes a covariant vector field and ∇_k denotes the usual covariant derivative.

Under a renormalization of the fundamental tensor of the form $\tilde{g}_{ij} = \lambda^2 g_{ij}$ the covariant vector T_k is transformed by the law $\tilde{T}_k = T_k + \partial_k \ln \lambda$, where λ is a scalar function defined on GW_n .

Let L^i_{jk} denote the coefficients of the asymmetric connection ∇_k . So, a generalized Weyl space is shortly written as $GW_n(L^i_{jk}, g_{ij}, T_k)$.

The main properties of $GW_n(L^i_{jk}, g_{ij}, T_k)$ can be expressed as follows

$$g_{ij} = g_{(ij)} + g_{[ij]} \quad (1.2)$$

$$\nabla_k g_{(ij)} = 2g_{(ij)} T_k \quad (1.3)$$

$$\nabla_k g_{[ij]} = 2g_{[ij]} T_k \quad (1.4)$$

$$g_{(ik)} g^{(kl)} = \delta_i^l \quad (1.5)$$

$$\nabla_k g^{(ij)} = -2T_k g^{(ij)} \quad (1.6)$$

where $g_{(ij)}$ and $g_{[ij]}$ denote symmetric and antisymmetric part of g_{ij} , respectively.

The symmetric part of connection coefficients L_{jk}^i are given as ([1], [2], [3])

$$L_{jk}^i = W_{jk}^i = \begin{bmatrix} i \\ jk \end{bmatrix} - (\delta_j^i T_k + \delta_k^i T_j - g_{jk} g^{mi} T_m) \tag{1.7}$$

where $\begin{bmatrix} i \\ jk \end{bmatrix}$ are second kind Christoffel symbols defined by

$$\begin{bmatrix} i \\ jk \end{bmatrix} = \frac{1}{2} g^{(ir)} \left[\frac{\partial g_{(jr)}}{\partial x^k} + \frac{\partial g_{(kr)}}{\partial x^j} - \frac{\partial g_{(jk)}}{\partial x^r} \right]. \tag{1.8}$$

A quantity A is called a satellite of weight $\{p\}$ of the tensor g_{ij} , if it admits a transformation of the form

$$\tilde{A} = \lambda^p A. \tag{1.9}$$

The prolonged covariant derivative of a satellite A of the tensor g_{ij} of weight $\{p\}$ is defined by

$$\dot{\nabla}_k A = \nabla_k A - p T_k A. \tag{1.10}$$

Let $C : x^i = x^i(s)$ be curve in GW_n . The generalized covariant derivative along the curve C of the tensor field T is defined by

$$\frac{\delta T}{\delta s} = \xi_{(t)}^k \nabla_k T \tag{1.11}$$

where $\xi_{(t)}^k$ are the components of the tangent vector of the curve C .

The Frenet equations of C may be written as [4].

$$\frac{\delta \xi_{(\alpha)}^i}{\delta s} = \kappa_{(\alpha)} \xi_{(\alpha+1)}^i - \kappa_{(\alpha-1)} \xi_{(\alpha-1)}^i \quad (\alpha = 1, 2, \dots, n; \kappa_{(0)} = \kappa_{(n)} = 0) \tag{1.12}$$

In the above equation $\xi_{(\alpha)}^i (\alpha = 2, \dots, n)$ denote the α -th curvature of weight $\{-1\}$ normalized by the condition

$$g_{ij} \xi_{(\alpha)}^i \xi_{(\alpha)}^j = 1 \tag{1.13}$$

of the curve C and $\kappa_{(\alpha)} (\alpha = 1, \dots, n-1)$ denote the α -th curvature of weight $\{-1\}$ of the curve C .

Let an n-dimensional hypersurface GW_n given by the equations $y^\alpha = y^\alpha(x^i) (\alpha = 1, \dots, n+1; i = 1, 2, \dots, n)$ be immersed in a generalized Weyl space GW_{n+1} .

The prolonged covariant derivative of the satellite A , relative to GW_{n+1} and GW_n are related by

$$\dot{\nabla}_k A = x_k^\gamma \dot{\nabla}_\gamma A \tag{1.14}$$

where $x_k^\gamma = \frac{\partial y^\gamma}{\partial x^k}$.

The components of any vector U relative to GW_{n+1} and GW_n are related by

$$U^\alpha = x_i^\alpha U^i. \tag{1.15}$$

The prolonged covariant derivative x_i^α is given by

$$\dot{\nabla}_j x_i^\alpha = w_{ij} N^\alpha + A_{ij}^h x_h^\alpha \tag{1.16}$$

where w_{ij} are the components of second fundamental form of GW_n defined by

$$w_{ij} = g_{(\alpha\beta)} N^\beta \dot{\nabla}_j x_i^\alpha \tag{1.17}$$

and A_{ij}^h are defined by

$$A_{ij}^h = g_{(\alpha\beta)} \dot{\nabla}_j x_i^\alpha x_j^\beta g^{(h\gamma)}. \tag{1.18}$$

The components q^α and p^i of the first curvature vectors of the curve $C : x^i = x^i(s)$ with respect to GW_{n+1} and GW_n are given by [5]

$$q^\alpha = \frac{\dot{\delta} \xi_{(1)}^\alpha}{\delta s} = \kappa_{(n)} N^\alpha + p^i x_i^\alpha + I^h x_h^\alpha \tag{1.19}$$

where

$$\kappa_{(n)} = w_{ij} \frac{dx^i}{ds} \frac{dx^j}{ds} \tag{1.20}$$

and I^h are the components of intrinsic curvature vector of the curve C in the hypersurface defined by [5]

$$I^h = A_{ij}^h \frac{dx^i}{ds} \frac{dx^j}{ds}. \tag{1.21}$$

The prolonged covariant derivative of the unit normal is given by

$$\dot{\nabla}_i N^\alpha = -g^{(jk)} w_{ij} B_k^\alpha \tag{1.22}$$

2.HYPER DARBOUX LINES OF ORDER H

Let us take an n -dimensional GW_n defined by the equations $y^\alpha = y^\alpha(x^i)$ ($\alpha = 1, \dots, n+1; i = 1, \dots, n$) which is immersed in a generalized Weyl Space GW_{n+1} .

Let us take $C : x^i = x^i(s)$ which is a curve (not a geodesic of the enveloping space) of the hypersurface .

The components $\xi_{(0)}^\alpha, \xi_{(1)}^\alpha$ and $\xi_{(r)}^\alpha$ are of the unit tangent vectors, of the principal normal vector and of the $(r-1)$ -th binormal vectors at every point of the curve, respectively. These vectors define an orthogonal system unit vectors at every point of the curve. We consider a congruence of curves given by unit vector field λ in GW_{n+1} as

$$\lambda^\alpha = q^i x_i^\alpha + r N^\alpha \tag{2.1}$$

where N^a are the components of the unit normal vector of the hyper space.

Definition 2.1: If the surface spanned by the vectors $\xi_{(0)}^\alpha$ and $R_{(h+1)}\xi_{(h+1)}^\alpha + R_{(h+2)}\frac{\dot{\delta} R_{(h+1)}}{\delta S}\xi_{(h+2)}^\alpha$ ($R_{(\alpha)} \equiv \frac{1}{\kappa_{(\alpha)}}$ and $\kappa_{(\alpha)}$ is the curvature of the α -th order) contains the vector λ^a . The curve C is said to be a hyper D-line of order h ($0 \leq h \leq (n+1) - 3$).

From this definition, for a hyper darbox line of order h we can write

$$\lambda^a = p_{(h)} \left[R_{(h+1)}\xi_{(h+1)}^a + R_{(h+2)}\frac{\dot{\delta} R_{(h+1)}}{\delta S}\xi_{(h+2)}^a \right] + q_{(h)}\xi_{(0)}^a. \tag{2.2}$$

From (1.12) we have

$$\frac{\dot{\delta}^2 \xi_{(r)}^a}{\delta S^2} = -\frac{\dot{\delta} \kappa_{(r)}}{\delta S}\xi_{(r-1)}^a - \kappa_{(r)}\frac{\dot{\delta} \xi_{(r-1)}^a}{\delta S} + \frac{\dot{\delta} \kappa_{(r+1)}}{\delta S}\xi_{(r+1)}^a + \kappa_{(r+1)}\frac{\dot{\delta} \xi_{(r+1)}^a}{\delta S} \tag{2.3}$$

Using the (1.12) in this equation:

$$\begin{aligned} \frac{\dot{\delta}^2 \xi_{(r)}^a}{\delta S^2} &= -\frac{\dot{\delta} \kappa_{(r)}}{\delta S}\xi_{(r-1)}^a - \kappa_{(r)}(-\kappa_{(r-1)}\xi_{(r-2)}^a + \kappa_{(r)}\xi_{(r)}^a) + \\ &\quad + \frac{\dot{\delta} \kappa_{(r+1)}}{\delta S}\xi_{(r+1)}^a + \kappa_{(r+1)}(-\kappa_{(r+1)}\xi_{(r)}^a + \kappa_{(r+2)}\xi_{(r+2)}^a) \\ &= \kappa_{(r)}\kappa_{(r-1)}\xi_{(r-2)}^a - \frac{\dot{\delta} \kappa_{(r)}}{\delta S}\xi_{(r-1)}^a - (\kappa_{(r)}^2 + \kappa_{(r+1)}^2)\xi_{(r)}^a + \\ &\quad + \frac{\dot{\delta} \kappa_{(r+1)}}{\delta S}\xi_{(r+1)}^a + \kappa_{(r+1)}\kappa_{(r+2)}\xi_{(r+2)}^a. \end{aligned} \tag{2.4}$$

This equation gives

$$g_{ab} \left(\frac{\dot{\delta}^2 \xi_{(r)}^a}{\delta S^2} \right) \left[R_{(h+1)}\xi_{(h+1)}^b + R_{(h+2)}\frac{\dot{\delta} R_{(h+1)}}{\delta S}\xi_{(h+2)}^b \right] = 0. \tag{2.5}$$

Multiplying λ^a by $g_{ab}\xi_{(0)}^b$ and $g_{ab}\frac{\delta^2 \xi_{(h)}^b}{\delta S^2}$ and using the last equation, we get

$$g_{ab}\lambda_{(h)}^a \xi_{(0)}^b = q, \tag{2.6}$$

$$g_{ab}\frac{\delta^2 \xi_{(h)}^b}{\delta S^2} \lambda^a = q g_{(h)ab} \frac{\delta^2 \xi_{(h)}^b}{\delta S^2} \xi_{(0)}^a \tag{2.7}$$

and using the definition $\lambda_{(h)} = g_{ab}\lambda^a \xi_{(h)}^b$ we obtain

$$q_{(h)} = \lambda_{(h)} \tag{2.8}$$

and

$$g_{ab}\frac{\delta^2 \xi_{(h)}^b}{\delta S^2} \lambda^a = \lambda_{(h-2)}\kappa_{(h)}\kappa_{(h-1)} - \lambda_{(h-1)}\frac{\delta \kappa_{(h)}}{\delta S} + \tag{2.9}$$

$$-\lambda_{(h)}\left(\kappa_{(h)}^2 + \kappa_{(h+1)}^2\right) + \lambda_{(h+1)}\frac{\delta \kappa_{(h+1)}}{\delta S} + \lambda_{(h+2)}\kappa_{(h+1)}\kappa_{(h+2)} = 0.$$

The above equation is valid for $h = 3, \dots, (n + 1) - 3$.

In the case of $h = 2$ it is written as

$$g_{ab}\frac{\delta^2 \xi_{(2)}^b}{\delta S^2} \lambda^a = \lambda_{(0)}g_{ab}\xi_{(0)}^a \frac{\delta^2 \xi_{(2)}^b}{\delta S^2} = \kappa_{(2)}\kappa_{(1)}\lambda_{(0)} \tag{2.10}$$

and so

$$-\frac{\delta \kappa_{(2)}}{\delta S} \lambda_{(1)} - \left(\kappa_{(2)}^2 + \kappa_{(3)}^2\right) \lambda_{(2)} + \frac{\delta \kappa_{(3)}}{\delta S} \lambda_{(3)} + \kappa_{(3)}\kappa_{(4)}\lambda_{(4)} = 0. \tag{2.11}$$

In the case of $h = 1$ it can be written that

$$g_{ab}\frac{\delta^2 \xi_{(1)}^b}{\delta S^2} \lambda^a = q g_{(0)ab} \frac{\delta^2 \xi_{(1)}^b}{\delta S^2} \xi_{(0)}^a = \lambda_{(0)}g_{ab}\xi_{(0)}^a \frac{\delta^2 \xi_{(1)}^b}{\delta S^2} \tag{2.12}$$

and so

$$-\left(\kappa_{(1)}^2 + \kappa_{(2)}^2\right) \lambda_{(1)} + \frac{\delta \kappa_{(2)}}{\delta S} \lambda_{(2)} + \kappa_{(2)}\kappa_{(3)}\lambda_{(3)} = 0. \tag{2.13}$$

In the case of $h = 0$ it can be written that

$$g_{ab}\frac{\delta^2 \xi_{(0)}^b}{\delta S^2} \lambda^a = q g_{(0)ab} \frac{\delta^2 \xi_{(0)}^b}{\delta S^2} \xi_{(0)}^a = \lambda_{(0)}g_{ab}\xi_{(0)}^a \frac{\delta^2 \xi_{(0)}^b}{\delta S^2} \tag{2.14}$$

and so

$$\frac{\dot{\delta} \kappa_{(1)}}{\delta S} \lambda_{(1)} + \kappa_{(1)} \kappa_{(2)} \lambda_{(2)} = 0. \tag{2.15}$$

These equations show the hyper D-lines of various orders.

Theorem 2.1: If the congruence λ^a is along the h -th binormal $\xi_{(h+1)}^a$ of a curve then the necessary and sufficient to be a hyper D-line of order h is its being a curve of constant $(h+1)$ -th curvature.

Proof: Let the congruence λ^a is along the h -th binormal $\xi_{(h+1)}^a$ of a curve. That is $\lambda^a = c \cdot \xi_{(h+1)}^a$, where c is a constant. From the equations (2.6) and (2.8), we get

$$\lambda_{(h)} = g_{ab} \lambda^a \xi_{(h)}^b = c g_{ab} \xi_{(h+1)}^a \xi_{(h)}^b = q_{(h)} = 0 \tag{2.16}$$

and so, $\lambda_{(i)} = c \delta_i^{h+1}$.

If this curve is a hyper D-line of order h , from (2.9) it is

$$\begin{aligned} \lambda_{(h-2)} \kappa_{(h)} \kappa_{(h-1)} - \lambda_{(h+1)} \frac{\dot{\delta} \kappa_{(h)}}{\delta S} - \lambda_{(h)} (\kappa_{(h)}^2 + \kappa_{(h+1)}^2) + \\ + \lambda_{(h+1)} \frac{\dot{\delta} \kappa_{(h+1)}}{\delta S} + \lambda_{(h+2)} \kappa_{(h+1)} \kappa_{(h+2)} = 0 \end{aligned} \tag{2.17}$$

From the last equation, it is found that

$$\kappa_{(h+1)} = \text{const.} \tag{2.18}$$

Conversely, if the $\kappa_{(h+1)}$ is a constant then $\lambda_{(h+1)} \frac{\dot{\delta} \kappa_{(h+1)}}{\delta S} = 0$ may be written and the equation (2.9) is satisfied and so, the curve is a hyper D-line.

Theorem 2.2: If the congruence λ^a is along the $(h-2)$ th binormal $\xi_{(h-1)}^a$ of a curve then the necessary and sufficient condition to be a hyper D-line of order h ($h \geq 2$) is its being a curve of constant h -th curvature.

Proof: Let the congruence λ^a be along the $(h-2)$ th binormal $\xi_{(h-1)}^a$ of a curve. That is $\lambda^a = c \cdot \xi_{(h-1)}^a$, where c is a constant. From the equation (2.6) and (2.8), we get

$$\lambda_{(h)} = g_{ab} \lambda^a \xi_{(h)}^b = c g_{ab} \xi_{(h-1)}^a \xi_{(h)}^b \tag{2.19}$$

and so, $\lambda_{(i)} = c \delta_i^{h-1}$

If this curve is a hyper D-line of order h ($h \geq 2$), it is

$$\begin{aligned} \lambda_{(h-2)}\kappa_{(h)}\kappa_{(h-1)} - \lambda_{(h-1)} \frac{\dot{\delta} \kappa_{(h)}}{\delta S} - \lambda_{(h)} \left(\kappa_{(h)}^2 + \kappa_{(h+1)}^2 \right) + \\ + \lambda_{(h+1)} \frac{\dot{\delta} \kappa_{(h+1)}}{\delta S} + \lambda_{(h+2)}\kappa_{(h+1)}\kappa_{(h+2)} = 0 \end{aligned} \tag{2.20}$$

From (2.19) and the last equation, it is found that

$$\kappa_{(h)} = \text{const.} \tag{2.21}$$

Conversely, if the $\kappa_{(h)}$ is a constant then $\lambda_{(h-1)} \frac{\dot{\delta} \kappa_{(h)}}{\delta S} = 0$ may be written and the equation is satisfied, and so, the curve is a hyper D-line.

HYPER D-LINES

A hyper D-line of order zero is called the hyper D-line of the hyperspace. This curve is represented by

$$\frac{\dot{\delta} \kappa_{(1)}}{\delta S} \lambda_{(1)} + \kappa_{(1)}\kappa_{(2)}\lambda_{(2)} = 0 \tag{3.1}$$

Theorem 3.1: If the congruence lies along the first binormal $\xi_{(2)}^a$ then the necessary and the sufficient condition that the curve be a hyper D- line is that it be the curve of zero torsion ($\kappa_{(2)} = 0$).

Proof: Let the congruence λ^a be along the first binormal $\xi_{(2)}^a$ of a curve. So, we have

$\lambda^a = c \cdot \xi_{(2)}^a$. From the equations (2.6) and (2.8), we get

$$\lambda_{(h)} = g_{ab} \lambda^a \xi_{(2)}^b = c g_{ab} \xi_{(2)}^a \xi_{(2)}^b \tag{3.2}$$

And so, $\lambda_{(i)} = c \delta_i^2$.

If this curve is a hyper D-line, it is

$$\frac{\dot{\delta} \kappa_{(1)}}{\delta S} \lambda_{(1)} + \kappa_{(1)}\kappa_{(2)}\lambda_{(2)} = 0 \tag{3.3}$$

From (3.2) and the last equation, it is found that $\kappa_{(1)}\kappa_{(2)} = 0$. Since $\kappa_{(1)} \neq 0$, $\kappa_{(2)}$ is equal to zero.

Conversely, if the $\kappa_{(2)}$ is zero then $\frac{\dot{\delta} \kappa_{(1)}}{\delta S} \lambda_{(1)} = 0$ may be written and the equation (3.1) is satisfied. And so, the curve is a hyper D-line.

We shall deduce the equation involving the second fundamental tensors of the hypersurface.

From (1.11) and (1.16) we get

$$\frac{\delta \xi^a}{\delta s} = (p^k + I^k)x_k^a + (w_{ij}\xi^i\xi^j)N^a \tag{3.4}$$

where p^k and I^k are given by

$$p^k = (\dot{\nabla}_j \xi^k)\xi^j, \quad I^k = A_{ij}^k \xi^i \xi^j, \tag{3.5}$$

respectively, and the equation (3.6):

$$\begin{aligned} \frac{\dot{\delta}^2 \xi^a}{\delta s^2} = & \left\{ \frac{\dot{\delta}(p^m + I^m)}{\delta s} + (p^k + I^k)A_{kj}^m \xi^j - w_{ij}w_{kl}\xi^i \xi^j \xi^k g^{lm} \right\} x_m^a + \\ & + \left\{ p^k(2w_{ki} + w_{ik})\xi^i + I^k \xi^j w_{kj} + (\dot{\nabla}_k w_{ij})\xi^i \xi^j \xi^k \right\} N^a. \end{aligned}$$

By putting the last equation into the equation $g_{ab} \frac{\dot{\delta}^2 \xi^a}{\delta s^2} \lambda^b = \lambda_{(0)} g_{ab} \xi^b \frac{\dot{\delta}^2 \xi^a}{\delta s^2}$ and using $g_{ab} N^a N^b = 1$ and $g_{ab} N^a x_i^b = 0$ it is found that

$$\begin{aligned} \frac{\dot{\delta}(p^m + I^m)}{\delta s} q_m + (p^k + I^k)A_{kj}^l \xi^j q_l - q^p w_{ij} w_{kp} \xi^i \xi^j \xi^k + \\ + r[w_{kj} I^k \xi^j + (\dot{\nabla}_k w_{ij})\xi^i \xi^j \xi^k + (2w_{ij} + w_{ji})p^i \xi^j] + \\ - \lambda_{(0)} g_{sm} \xi^s \frac{\dot{\delta}(p^m + I^m)}{\delta s} - \lambda_{(0)} (p^k + I^k) g_{sm} \xi^s \xi^j A_{kj}^m + (w_{ij} \xi^i \xi^j)^2 \lambda_{(0)} = 0 \end{aligned}$$

We have the first two Frenet's formulae for the subspace

$$p^m = \frac{\dot{\delta} \xi^m}{\delta s} = \kappa_{(1)} \xi_1^m, \quad \frac{\dot{\delta} \xi_1^m}{\delta s} = -\kappa_{(1)} \xi_1^m + \kappa_{(2)} \xi_2^m \tag{3.8}$$

where ξ_1^m and ξ_2^m are the principal and binormal vectors and $\kappa_{(1)}$, $\kappa_{(2)}$ are the first and second curvatures with respect to the subspace. From the last equations, by using the fact $\lambda_{(0)} = g_{ab} \lambda^a \xi^b = g_{ab} (q^i x_i^a + r N^a) \xi^b = g_{ab} q^i \xi^j = q_{(0)}$ we get

$$\frac{\dot{\delta} p^m}{\delta s} q_m = -\kappa_{(1)}^2 q_{(0)} + \frac{\dot{\delta} \kappa_{(1)}}{\delta s} q_{(1)} + \kappa_{(1)} \kappa_{(2)} q_{(2)} \tag{3.9}$$

and

$$\begin{aligned} -\kappa_{(1)}^2 q_{(0)} + \frac{\dot{\delta} \kappa_{(1)}}{\delta s} q_{(1)} + \kappa_{(1)} \kappa_{(2)} q_{(2)} + \frac{\dot{\delta} I^m}{\delta s} q_m + p^m A_{mj}^l \xi^j q_l + I^m A_{mj}^l \xi^j q_l + \\ - q^p w_{ij} w_{kp} \xi^i \xi^j \xi^k + r[w_{kj} I^k \xi^j + (\dot{\nabla}_k w_{ij})\xi^i \xi^j \xi^k + w_{ij} p^k \xi^j + w_{ij} p^i \xi^j + w_{ij} p^j \xi^i] + \\ + \lambda_{(0)} \kappa_{(1)} \xi_1^k g_{sm} \xi^s \xi^j A_{kj}^m - \lambda_{(0)} I^k g_{sm} \xi^s \xi^j A_{kj}^m + \lambda_{(0)} (w_{ij} \xi^i \xi^j)^2 = 0 \end{aligned}$$

where $q_{(1)}$ and $q_{(2)}$ are the projections of q^i in the directions ξ_1^m and ξ_2^m , respectively. This represents the hyper D-line of the subspace. The equations have been expressed in terms of the second fundamental tensors and the curvatures (of the curve) with respect to the subspace.

REFERENCES

1. Norden, A., "*Affinely Connected Spaces*", GRFML, Moscow,(1976) (In Russian)
2. Murgescu, V., "Sur les Espaces a Connection Affine", A Tenseur Recurrent, Bul. Inst. Pol de Jassy, VII(XII) 1-2, 65, (1962).
3. Murgescu, V., "Espaces de Weyl generalises", Bui. Inst. Pol de Jassy, (1970).
4. Zeren, Leyla., "*Frenet Formulas for Curves in a Generalized Weyl Space*" Ganita vol.51, No. 2, 149-164, (2000).
5. Demirbükler, H., Zeren Akgun, L., "Hypernormal Curves on the Generalized Weyl Space", Journal of Engineering and Natural Science No.1/2005,68-72, (2005).
6. Sing, U., P., "On Hyper Darboux Lines in Riemannian Subspaces" Tensor, N. S., vol.20, 1-4, (1969).
7. Sing, A., K., "On Hyper Darboux Lines in Kaehlerian Hyperspaces" Acta Ciencia Indica, Vol. XII, No:3, 183-190 , 4445-4452, (1986).