

On the decomposition of curvature tensor in recurrent conformal Finsler space

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M.S. Knebelman [1] has developed conformal geometry of generalised metric spaces. The projective tensor and curvature tensors in conformal Finsler spaces were discussed by R.B.Misra [3,4]. M.Gama [6] has decomposed recurrent curvature tensor in an areal space of submetric class. The decomposition of recurrent curvature tensor in Finsler space was studied by B.B.Sinha and S.P.Singh [5]. The purpose of the present paper is to decompose the recurrent conformal curvature tensor and study the properties of conformal decomposition tensors.

1. Introduction

Let us consider an n -dimensional Finsler space F_n in which two distinct metric functions are represented by $F(x, \dot{x})$ and $\bar{F}(x, \dot{x})$. Then the corresponding metric tensors $g_{ij}(x, \dot{x})$ and $\bar{g}_{ij}(x, \dot{x})$ are called conformal if they are proportional to each other. These quantities are related as under:

$$(1.1) \quad (a) \quad \bar{F}(x, \dot{x}) = e^\sigma F(x, \dot{x}) \quad (b) \quad \bar{g}_{ij}(x, \dot{x}) = e^{2\sigma} g_{ij}(x, \dot{x}) \\ (c) \quad \bar{g}^{ij}(x, \dot{x}) = e^{-2\sigma} g^{ij}(x, \dot{x}),$$

where $\sigma = \sigma(x)$ is atmost a point function as shown by Knebelman [1]. The space equipped with quantities $\bar{F}(x, \dot{x}), \bar{g}(x, \dot{x})$ etc is called a conformal Finsler space usually denoted by \bar{F}_n

1) The numbers in brackets refer to the references given at the end of the paper

2) $\partial_j = \frac{\partial}{\partial x^j}$ and $\dot{\partial}_j = \frac{\partial}{\partial \dot{x}^j}$

In a Finsler space F_n with entities $F(x, \dot{x}), g_{ij}(x, \dot{x})$ etc, the covariant derivatives of a vector $T^i(x, \dot{x})$ with respect to x^j in the sense of Cartan and Berwald are given by [2],

$$(1.2) \quad T^i_{|j} \stackrel{def}{=} \partial_j T^i - (\dot{\partial}_m T^i) G^m_j + T^m \Gamma^*_{mj}{}^i$$

and

$$(1.3) \quad T^i_{(j)} \stackrel{def}{=} \partial_j T^i - (\dot{\partial}_m T^i) G^m_{(j)} + T^m G^i_{m(j)},$$

where

$$(1.4) \quad G^i_j(x, \dot{x}) = G^i_{mj}(x, \dot{x}) \dot{x}^m = \Gamma^*_{mj}{}^i(x, \dot{x}) \dot{x}^m$$

The connection coefficients are homogeneous functions of degree zero in \dot{x}^i . Considering quantities

$$(1.5) \quad \sigma_m(x) \stackrel{def}{=} \partial_m \sigma, \quad B^{ij}(x, \dot{x}) \stackrel{def}{=} \frac{1}{2} F^2 g^{ij} - \dot{x}^i \dot{x}^j,$$

$$U^i_{jk}(x, \dot{x}) \stackrel{def}{=} 2\sigma_{(j} \delta^i_{k)} - \sigma_m \left\{ g^{im} g_{jk} - 2C^i_{r(j} \dot{\partial}_k) B^{rm} + g^{ir} C_{jks} \dot{\partial}_r B^{sm} \right\},$$

where $C^i_{mj}(x, \dot{x}) = g^{ih} C_{mhj}(x, \dot{x})$, one can have conformal connection coefficients :

$$(1.6) \quad \bar{\Gamma}^*_{jk}{}^i(x, \dot{x}) = \Gamma^*_{jk}{}^i(x, \dot{x}) + U^i_{jk}(x, \dot{x}),$$

$$(1.7) \quad \bar{G}^i_{jk}(x, \dot{x}) = G^i_{jk}(x, \dot{x}) - \sigma_m \dot{\partial}_j \dot{\partial}_k \dot{\partial}_h B^{im},$$

where the functions $B^{im}(x, \dot{x})$ are homogeneous of degree two in \dot{x}^i .

The variation in $G^i_{jkh}(x, \dot{x})$ under the conformal change (1.1) is given by

$$(1.8) \quad \bar{G}^i_{j'ch}(x, \dot{x}) = G^i_{jkh}(x, \dot{x}) - \sigma_m \dot{\partial}_j \dot{\partial}_k \dot{\partial}_h B^{im},$$

where

$$\bar{G}^i_{jkh}(x, \dot{x}) \stackrel{def}{=} \dot{\partial}_h G^i_{jk}.$$

The curvature tensors $K^i_{j'jk}, H^i_{j'jk}$ and H^i_{kh} transform under the conformal change (1.1) as

$$(1.9) \quad \bar{K}^i_{j'jk} = K^i_{j'jk} + 2U^i_{r[j|k]} + 2\left[\sigma_m \left\{ \dot{\partial}_s (\Gamma^*_{r|j}{}^i + U^i_{r|j}) \right\} \dot{\partial}_k \right] B^{sm} + U^i_{s[k} U^s_{j]r}$$

$$(1.10) \quad \begin{aligned} \bar{H}^i_{\ rjk} = & H^i_{\ rjk} + 2\sigma_{m[(j)}\dot{\partial}_k\dot{\partial}_r B^{im} - 2\sigma_m\dot{\partial}_r(\dot{\partial}_{[j}B^{im})_{(k)}] - (\dot{\partial}_{[j}B^{is})G^m_{k]rs} \} \\ & + 2\sigma_m\sigma_s\dot{\partial}_r\{(\dot{\partial}_{[j}B^{pm})\dot{\partial}_k\dot{\partial}_p B^{is} \} \end{aligned}$$

and

$$(1.11) \quad \begin{aligned} \bar{H}^i_{\ kh} = & H^i_{\ kh} - 2\sigma_m\dot{\partial}_{[k}B^{im})_{(h)}] + 2\sigma_{m[(k)}\dot{\partial}_h]B^{im} \\ & + 2\sigma_m\sigma_r(\dot{\partial}_{[k}B^{sm})\dot{\partial}_h\dot{\partial}_s B^{ir} \end{aligned}$$

respectively.

The Bianchi identities satisfied by $\bar{H}^i_{\ jk}$, $\bar{H}^i_{\ rjk}$ and $\bar{K}^i_{\ rjk}$ are [4]

$$(1.12) \quad \bar{H}^i_{\ [jk(\bar{h})]} = 0$$

$$(1.13) \quad \begin{aligned} H^i_{\ r[jk(\bar{h})]} = & \{H^p_{[kj} - \sigma_{q[(j)}\dot{\partial}_k B^{pq} + \sigma_{q[(k)}\dot{\partial}_j B^{pq}\}G^i_{h]pr} + \sigma_m\{(\dot{\partial}_{[j}B^{pm})_{(k)}] - (\dot{\partial}_{[k}B^{pm})_{(j)}\}G^i_{h]pr} \\ & + \{H^p_{[jk} + \sigma_{q[(j)}\dot{\partial}_k B^{pq} - \sigma_{q[(k)}\dot{\partial}_j B^{pq}\}\dot{\partial}_h\dot{\partial}_p\dot{\partial}_r B^{im}] - \sigma_m\sigma_l\{G^i_{pr[lj}\{(\dot{\partial}_k B^{st})\dot{\partial}_h]B^{pm} - \\ & - (\dot{\partial}_h B^{st})\dot{\partial}_k\dot{\partial}_s B^{pm} \} + (\dot{\partial}_p\dot{\partial}_r\dot{\partial}_{[j}B^{it})\{(\dot{\partial}_k B^{pm})_{(h)}] - (\dot{\partial}_h B^{pm})_{(k)}\} \} + \\ & + \sigma_m\sigma_q\sigma_l(\dot{\partial}_p\dot{\partial}_r\dot{\partial}_{[j}B^{iq})\{(\dot{\partial}_k B^{st})\dot{\partial}_h\dot{\partial}_s B^{pm} - (\dot{\partial}_h B^{st})\dot{\partial}_k\dot{\partial}_s B^{pm} \} \} \end{aligned}$$

and

$$(1.14) \quad \begin{aligned} \bar{K}^i_{\ r[jk(\bar{h})]} + & (\dot{\partial}_p\Gamma^*_{r[j}H^p_{kh}) + (\dot{\partial}_p U^i_{r[j}H^p_{kh}) + \dot{\partial}_p(\Gamma^*_{r[j} + U^i_{r[j}])\{(\dot{\partial}_h B^{pm})_{(k)}] - (\dot{\partial}_k B^{pm})_{(h)}\}\sigma_m \\ & + \{(\dot{\partial}_{[k}B^{sm})\dot{\partial}_h\dot{\partial}_s B^{pi} - (\dot{\partial}_h B^{sm})\dot{\partial}_k\dot{\partial}_s B^{pi}\}\sigma_m\sigma_l + \sigma_{[m|(k)}\dot{\partial}_h]B^{pm} - \sigma_{|m|(h)}\dot{\partial}_k]B^{pm} (\} = 0 \end{aligned}$$

respectively. The notations (\bar{k}) and $|\bar{h}$ represent covariant differentiation in the sense of Berwald and Cartan respectively in conformal Finsler space \bar{F}_n .

2. Decomposition of conformal curvature tensor $\bar{H}^i_{\ rjk}$

The recurrent conformal curvature tensor $\bar{H}^i_{\ rjk}$ is characterised by the condition

$$(2.1) \quad \bar{H}^i_{\ rjk(\bar{m})} = \bar{V}^i_{\ m}\bar{H}^i_{\ rjk}, \quad \bar{H}^i_{\ rjk} \neq 0,$$

where the barred index enclosed within the round bracket denotes the covariant derivative with respect to x^s for the connection parameters

$\bar{G}_{jk}^i(x, \dot{x})$. The covariant vector $\bar{V}_m(x, \dot{x})$ is called the conformal recurrence vector. The conformal space equipped with such curvature tensor is called recurrent conformal Finsler space and we denote it by \bar{F}_n^* .

We consider the decomposition of the recurrent conformal curvature tensor in the form

$$(2.2) \quad \bar{H}_{rjk}^i = \bar{X}^i \bar{\Phi}_{rjk},$$

where $\bar{\Phi}_{rjk}$ is a homogeneous conformal decomposition tensor and \bar{X}^i is a non-zero conformal vector such that

$$(2.3) \quad \bar{X}^i \bar{V}_i = 1$$

Sinha and Singh [5] have decomposed the recurrent curvature tensor H_{rjk}^i in similar manner

$$(2.4)(a) \quad H_{rjk}^i = X^i \Phi_{rjk},$$

where the decomposition vector X^i also satisfies the condition

$$(2.5) \quad X^i V_i = 1.$$

Transvecting (2.4)(a) by \dot{x}^r and noting $H_{rjk}^i \dot{x}^r = H_{jk}^i$ [2],

We obtain

$$(2.4)(b) \quad H_{jk}^i = X^i \Phi_{jk},$$

where $\Phi_{jk} = \Phi_{rjk} \dot{x}^r$.

The decomposition tensor Φ_{rjk} satisfies the identities

$$(2.6) \quad \Phi_{rjk} = -\Phi_{rkj}$$

and

$$(2.7) \quad \Phi_{rjk} + \Phi_{jkr} + \Phi_{kjr} = 0.$$

We notice that the decomposition vector X^i and the recurrence vector V_i are transformed conformally as under :

$$(2.8) \quad \bar{X}^i = e^{-\sigma} X^i$$

and

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$$(2.9) \quad \bar{V}_i = e^\sigma V_i$$

respectively.

Applying equations (2.2) and (2.4)(a) in equation (1.10), we obtain

$$(2.10) \quad \bar{X}^i \bar{\Phi}_{rjk} = X^i \Phi_{rjk} + 2\sigma_{m[(j)} \dot{\partial}_k \dot{\partial}_r B^{im} - 2\sigma_m \left\{ \dot{\partial}_r \left(\dot{\partial}_{[j} B^{im} \right)_{(k)} \right\} - \left(\dot{\partial}_{[j} B^{is} \right) G_{k]rs}^m \right\} \\ + 2\sigma_m \sigma_s \dot{\partial}_r \left\{ \left(\dot{\partial}_{[j} B^{pm} \right) \dot{\partial}_k \dot{\partial}_p B^{is} \right\}.$$

Transvecting (2.10) by \bar{V}_i and using the relations (2.3), (2.5), and (2.9), it yields

$$(2.11) \quad \bar{\Phi}_{rjk} = e^\sigma \left[\Phi_{rjk} + 2\sigma_{m[(j)} \dot{\partial}_k \dot{\partial}_r B^{im} V_i - 2\sigma_m \left\{ \dot{\partial}_r \left(\dot{\partial}_{[j} B^{im} \right)_{(k)} \right\} - \left(\dot{\partial}_{[j} B^{is} \right) G_{k]rs}^m \right\} V_i \\ + \sigma_m \sigma_s \dot{\partial}_r \left\{ \left(\dot{\partial}_{[j} B^{pm} \right) \dot{\partial}_k \dot{\partial}_p B^{is} \right\} V_i \right],$$

which represents the conformal transformation of the decomposition tensor Φ_{rjk} under the change (1.1)

Thus we state

Theorem 2.1 : Under the decomposition (2.2), the conformal decomposition tensor Φ_{rjk} is expressed in the form (2.11).

Interchanging the indices j and k in the equation (2.11), we get

$$(2.12) \quad \bar{\Phi}_{rjk} = -\bar{\Phi}_{rki}$$

in view of (2.6).

The cyclic permutation of the indices r, j, k in the equation (2.11) yields the identity

$$(2.13) \quad \bar{\Phi}_{[rjk]} = e^\sigma V_i \left[\sigma_m \dot{\partial}_{[r} \left(\dot{\partial}_k B^{im} \right)_{(j)} - \sigma_m \dot{\partial}_{[r} \left(\dot{\partial}_j B^{im} \right)_{(k)} \right] \\ + \sigma_m \sigma_s \dot{\partial}_{[r} \left\{ \left(\dot{\partial}_j B^{pm} \right) \dot{\partial}_k \dot{\partial}_p B^{is} \right\} - \sigma_m \sigma_s \dot{\partial}_{[r} \left\{ \left(\dot{\partial}_k B^{pm} \right) \dot{\partial}_j \dot{\partial}_p B^{is} \right\} \right]^{3)}$$

by applying (2.7) and the symmetry property of G_{rjk}^i .

Hence we state

Theorem 2.2 : Under the decomposition (2.2), the conformal decomposition tensor $\bar{\Phi}_{rjk}$ satisfies the identities (2.12) and (2.13).

Transvecting the equation (2.2) by \dot{x}^r , we have

$$(2.14) \quad \bar{H}^i_{jk} = \bar{X}^i \bar{\Phi}_{jk},$$

where

$$(2.15) \quad \bar{\Phi}_{jk} = \bar{\Phi}_{rjk} \dot{x}^r$$

and \bar{H}^i_{jk} satisfies the relation (1.1).

Applying the equation (2.5), (2.9) and (2.14) in the equation (1.11), it assumes the form

$$(2.16) \quad \bar{\Phi}_{kh} = e^\sigma V_i \left[H^i_{kh} - 2\sigma_m \left\{ \dot{\partial}_{[k} B^{im} \right\}_{(h)} \right] + 2\sigma_{m[(k)} \dot{\partial}_{h]} B^{im} + 2\sigma_m \sigma_r \left(\dot{\partial}_{[k} B^{sm} \right) \dot{\partial}_{h]} \dot{\partial}_s B^{ir} \right]$$

Interchanging the indices k and h in the equation (2.16), we get

$$(2.17) \quad \bar{\Phi}_{kh} = -\bar{\Phi}_{hk}$$

by virtue of the relation $H^i_{kh} = -H^i_{hk}$ [2].

In view of the equations (2.4)(b) and (2.5), the equation (2.16) becomes

$$(2.18) \quad \bar{\Phi}_{kh} = e^\sigma \left[\Phi_{kh} - 2\sigma_m \left\{ \dot{\partial}_{[k} B^{im} \right\}_{(h)} \right] V_i + 2\sigma_{m[(k)} \dot{\partial}_{h]} B^{im} V_i + 2\sigma_m \sigma_r \left(\dot{\partial}_{[k} B^{sm} \right) \dot{\partial}_{h]} \dot{\partial}_s B^{ir} V_i \right],$$

which gives the conformal transformation of the decomposition tensor Φ_{kh} under the change (1.1)

Accordingly, we have

Theorem 2.3 : Under the decomposition (2.2) and (2.14) the conformal decomposition tensor $\bar{\Phi}_{kh}$ is expressed in the form (2.18)

In view of the equations (2.2), (2.3), (2.4), (2.9) and (2.14), the identities (1.12) and (1.13) assume the forms

$$(2.19) \quad \bar{\Phi}_{[jk(h)} = 0$$

and

$$\begin{aligned}
 (2.20) \quad \bar{\Phi}_{r[jk(\bar{h})]} &= e^\sigma V_i \{ X^p \Phi_{[kj} - \sigma_{q[(j)} \dot{\partial}_k B^{pq} + \sigma_{q[(k)} \dot{\partial}_j B^{pq} \} G_{h]pr}^i \\
 &+ e^\sigma V_i \{ \{ (\dot{\partial}_{[j} B^{pm})_{(k)} - (\dot{\partial}_{[k} B^{pm})_{(j)} \} G_{h]pr}^i + \{ X^p \Phi_{[jk} + \sigma_{q[(j)} \dot{\partial}_k B^{pq} \\
 &- \sigma_{q[(k)} \dot{\partial}_j B^{pq} \} \dot{\partial}_h \dot{\partial}_p \dot{\partial}_r B^{im} \} - e^\sigma V_i \sigma_m \sigma_t \{ G_{pr[j}^i \{ (\dot{\partial}_k B^{st}) \dot{\partial}_h \dot{\partial}_s B^{pm} \\
 &- (\dot{\partial}_h B^{st}) \dot{\partial}_k \dot{\partial}_s B^{pm} \} + (\dot{\partial}_p \dot{\partial}_r \dot{\partial}_{[j} B^{it}) \{ (\dot{\partial}_k B^{pm})_{(h)} - (\dot{\partial}_h B^{pm})_{(k)} \} \} \\
 &+ e^\sigma V_i \sigma_m \sigma_q \sigma_t \{ \dot{\partial}_p \dot{\partial}_r \dot{\partial}_{[j} B^{iq} \} \{ (\dot{\partial}_k B^{st}) \dot{\partial}_h \dot{\partial}_s B^{pm} - (\dot{\partial}_h B^{st}) \dot{\partial}_k \dot{\partial}_s B^{pm} \}
 \end{aligned}$$

respectively . Hence we have

Theorem 2.4 : Under the decomposition (2.2) and (2.14), the conformal decomposition tensors $\bar{\Phi}_{jk}$ and $\bar{\Phi}_{rjk}$ satisfy the Bianchi identities (2.19) and (2.20) respectively .

Differentiating (2.2) covariantly with respect to x^h in the sense of Berwald, we get

$$(2.21) \quad \bar{H}_{rjk(\bar{h})}^i = \bar{X}_{(\bar{h})}^i \bar{\Phi}_{rjk} + \bar{X}^i \bar{\Phi}_{rjk(\bar{h})} .$$

Applying (2.1) and (2.2) in the above equation, we find

$$(2.22) \quad \bar{X}^i \bar{\Phi}_{rjk} \bar{V}_h = \bar{X}_{(\bar{h})}^i \bar{\Phi}_{rjk} + \bar{X}^i \bar{\Phi}_{rjk(\bar{h})} .$$

Let us assume that the conformal vector \bar{X}^i is covariant constant, then (2.22) reduces to

$$(2.23) \quad \bar{\Phi}_{rjk(\bar{h})} = \bar{V}_h \bar{\Phi}_{rjk}$$

Conversely, if the above equation is true, the equation (2.22) yields

$$\bar{X}_{(\bar{h})}^i \bar{\Phi}_{rjk} = 0$$

Since $\bar{\Phi}_{rjk}$ is non-zero conformal decomposition tensor ,it implies

$$(2.25) \quad \bar{X}_{(\bar{h})}^i = 0 ,$$

which shows that \bar{X}^i is covariant constant in the space \bar{F}_n^* .

Transvecting (2.23) by \dot{x}^r and using (2.15), we obtain

$$(2.26) \quad \overline{\Phi}_{jk|\bar{h}} = \overline{V}_h \overline{\Phi}_{jk}$$

Thus we state

Theorem 2.5 : In a recurrent conformal Finsler space \overline{F}_n^* , the necessary and sufficient condition for the conformal decomposition tensor field $\overline{\Phi}_{ijk}$ and $\overline{\Phi}_{jk}$ to be recurrent is that the conformal vector field \overline{X}^i is covariant constant in the sense of Berwald .

Decomposition of conformal curvature tensor \overline{K}_{ijk}^i

In this section we consider the decomposition of the conformal curvature tensor \overline{K}_{ijk}^i . In similar manner, the recurrent conformal curvature tensor \overline{K}_{ijk}^i is characterised by

$$(3.1) \quad \overline{K}_{ijk|\bar{h}}^i = \overline{V}_h \overline{K}_{ijk}^i, \quad \overline{K}_{ijk}^i \neq 0,$$

where the symbol $|\bar{h}$ denotes the covariant derivative with respect to x^i for the connection coefficients $\overline{\Gamma}_{jk}^i(x, \dot{x})$. The non-zero covariant vector $\overline{V}_h(x, \dot{x})$ is called the conformal recurrence vector .

We decompose the recurrent conformal curvature tensor \overline{K}_{ijk}^i in the following manner :

$$(3.2) \quad \overline{K}_{ijk}^i = \overline{X}^i \overline{\Psi}_{ijk}$$

where $\overline{\Psi}_{ijk}$ is conformal decomposition tensor and \overline{X}^i is conformal vector which satisfies the relation (2.3).

In a recurrent Finsler space, if the curvature tensor K_{ijk}^i is decomposed as

$$(3.3) \quad K_{ijk}^i = X^i \Psi_{ijk}, \quad K_{ijk}^i \neq 0$$

where the decomposition vector X^i satisfies the relation (2.5), then the decomposition vector X^i and the recurrence vector V_i are transformed conformally in the form (2.8) and (2.9) respectively .

In view of the identities $K^i_{rjk} = -K^i_{rjk}$ and $K^i_{[rjk]} = 0$ [2], the decomposition tensor $\Psi_{r,jk}$ satisfies the identities

$$(3.4) \quad \Psi_{rjk} = -\Psi_{rkj},$$

$$(3.5) \quad \Psi_{[rjk]} = 0.$$

Using the equations (3.2) and (3.3) in the equation (1.9), it assumes the form

$$(3.6) \quad \bar{X}^i \bar{\Psi}_{rjk} = X^i \Psi_{rjk} + 2U^i_{r[j|k]} + 2[\sigma_m \{ \dot{\partial}_s (\Gamma^*_{r[j} + U^i_{r[j}]) \dot{\partial}_k \} B^{sm}] + U^i_{s[k} U^s_{j]r}].$$

Transvecting the above equation by \bar{V}_i and applying the equations (2.3), (2.5) and (2.9), it becomes

$$(3.7) \quad \bar{\Psi}_{rjk} = e^\sigma \Psi_{rjk} + 2e^\sigma U^i_{r[j|k]} + 2e^\sigma [\sigma_m \{ \dot{\partial}_s (\Gamma^*_{r[j} + U^i_{r[j}]) \dot{\partial}_k \} B^{sm} + U^i_{s[k} U^s_{j]r}],$$

which gives the conformal transformation of the decomposition tensor Ψ_{rjk} under the change (1.1).

Accordingly, we have

Theorem 3.1: Under the decomposition (3.2), the decomposition tensor $\bar{\Psi}_{rjk}$ is expressed in the form (3.7).

Interchanging the indices j and k in the equation (3.7) and noting (3.4), we get

$$(3.8) \quad \bar{\Psi}_{rjk} = -\bar{\Psi}_{rkj}.$$

Also the cyclic permutation of the indices r, j, k in the equation (3.7) yields

$$(3.9) \quad \bar{\Psi}_{[rjk]} = 0$$

in view of (3.5).

Thus we have

Theorem 3.2: Under the decomposition (3.2), the conformal decomposition tensor $\bar{\Psi}_{rjk}$ satisfies the identities (3.8) and (3.9).

Applying (2.3), (2.4)(b), (2.5), (2.9) and (3.2) in the Bianchi identity (1.14), it assumes the form

$$\begin{aligned}
 (3.10) \quad & \bar{\Psi}_{r[jk|\bar{h}]} + e^\sigma V_i (\dot{\partial}_p \Gamma_r^i [j]) X^p \Phi_{kh} + e^\sigma V_i (\dot{\partial}_p U_r^i [j]) X^p \Phi_{kh} \\
 & + e^\sigma V_i \{ \dot{\partial}_p (\Gamma_r^i [j] + U_r^i [j]) \} \{ [(\dot{\partial}_h B^{pm})_{(k)}] - (\dot{\partial}_k B^{pm})_{(h)} \} \sigma_m \\
 & + \{ (\dot{\partial}_k B^{sm}) \dot{\partial}_h \dot{\partial}_s B^{pi} - (\dot{\partial}_h B^{sm}) \dot{\partial}_k \dot{\partial}_s B^{pi} \} \sigma_m \sigma_i \\
 & + \sigma_{|m|(k)} \dot{\partial}_h B^{pm} - \sigma_{|m|(h)} \dot{\partial}_k B^{pm}] = 0 .
 \end{aligned}$$

Hence we state

Theorem 3.3 : Under the decomposition (3.2) the conformal decomposition tensor $\bar{\Psi}_{ijk}$ satisfies the Bianchi identity (3.10).

Taking covariant differentiation of (3.2) with respect to x^h in the sense of Cartan, it yields

$$(3.11) \quad \bar{K}_{ijk|\bar{h}}^i = \bar{X}_{|\bar{h}}^i \bar{\Psi}_{ijk} + \bar{X}^i \bar{\Psi}_{ijk|\bar{h}} .$$

In view of the equation (3.1), it becomes

$$(3.12) \quad \bar{V}_h \bar{K}_{ijk}^i = \bar{X}_{|\bar{h}}^i \bar{\Psi}_{ijk} + \bar{X}^i \bar{\Psi}_{ijk|\bar{h}} .$$

Using the decomposition (3.2) in the above equation, we get

$$(3.13) \quad \bar{X}^i \bar{V}_h \bar{\Psi}_{ijk} = \bar{X}_{|\bar{h}}^i \bar{\Psi}_{ijk} + \bar{X}^i \bar{\Psi}_{ijk|\bar{h}} .$$

Let us assume that the conformal vector \bar{X}^i is covariant constant, that is, $\bar{X}_{|\bar{h}}^i = 0$. Then the equation (3.13) reduces to

$$(3.14) \quad \bar{\Psi}_{ijk|\bar{h}} = \bar{V}_h \bar{\Psi}_{ijk} .$$

Conversely, if the above relation is true, then from the equation (3.13), we find

$$(3.15) \quad \bar{X}_{|\bar{h}}^i \bar{\Psi}_{ijk} = 0 .$$

Since $\bar{\Psi}_{ijk}$ is non-zero, it implies

$$(3.16) \quad \bar{X}_{|\bar{h}}^i = 0,$$

which implies that the conformal vector \bar{X}^i is covariant constant.

We have

Theorem 2.4 : Under the decomposition (3.2), the necessary and sufficient condition for the conformal decomposition tensor field $\bar{\Psi}_{ijk}$ to be recurrent is that the conformal vector \bar{X}^i is covariant constant in the sense of Cartan.

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