

ON HYPER DARBOUX LINES IN FINSLER SUBSPACE

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The hyper DARBOUX lines (hyper D-lines) on a surface of a 3-dimensional Euclidean space have been defined and studied by PRVANOVITCH [2]. SINGH [3] investigated some properties of these curves for a Riemannian subspace. The object of this paper is to obtain the differential equation of the hyper D-lines of a FINSLER subspace and deduce some properties of union hyper D-lines.

1. Fundamental formulae. Let a subspace F_m given by the equations $x^i = x^i(u^\alpha)$ ($i = 1, \dots, n; \alpha = 1, \dots, m$) be immersed in an n -dimensional FINSLER space F_n . Consider any arbitrary curve $C: u^\alpha = u^\alpha(s)$ (or $x^i = x^i(s)$) of the subspace. The components x'^i (or $\frac{dx^i}{ds}$) and u'^α (or $\frac{du^\alpha}{ds}$) of the unit tangent vector of C are related by $x'^i = B_\alpha^i u'^\alpha$ where $B_\alpha^i = \frac{\partial x^i}{\partial u^\alpha}$.

The metric tensors $g_{ij}(x, x')$ and $g_{\alpha\beta}(u, u')$ of F_n and F_m are such that

$$(1.1) \quad g_{\alpha\beta}(u, u') = g_{ij}(x, x') B_\alpha^i B_\beta^j$$

and the tensor

$$(1.2) \quad C_{ijk}(x, x') = \frac{1}{2} \frac{\partial g_{ij}(x, x')}{\partial x'^k}$$

satisfies the identities

$$(1.3) \quad C_{ijk}(x, x') x'^i = C_{ijk}(x, x') x'^j = C_{ijk}(x, x') x'^k = 0.$$

There exist $(n-m)$ vectors, called secondary normal vectors $n_{(\mu)}^{*i}$ ($\mu = m+1, \dots, n$), which satisfy the relations

$$(1.4) \quad a) \quad g_{ij}(x, x') n_{(\mu)}^{*j} B_\alpha^i = n_{(\mu)i}^* B_\alpha^i = 0,$$

$$b) \quad g_{ij}(x, n_{(\mu)}^*) n_{(\mu)}^{*i} n_{(\mu)}^{*j} = 1,$$

$$c) \quad g_{ij}(x, x') n_{(\mu)}^{*i} n_{(\nu)}^{*j} = \delta_{\mu\nu} \psi_{(\mu)}$$
 (no summation on μ).

The δ -derivative of a vector field V^i of F_n along the curve C has been defined by RUND [3] as

$$(1.5) \quad \frac{\delta V^i}{\delta s} = \frac{dV^i}{ds} + P_{jk}^{*i} V^j \frac{dx^k}{ds}$$

where P_{jk}^{*i} are the connection coefficients of the embedding space.

The δ -derivative of the metric tensor $g_{ij}(x, x')$ along the curve C is given by RUND [4]

$$(1.6) \quad \frac{\delta}{\delta s} g_{ij}(x, x') = 2C_{ijk}(x, x') \frac{\delta x'^k}{ds} = C_{ijk}^* \frac{dx^k}{ds}$$

where

$$(1.7) \quad C_{ijk}^*(x, x') = g_{ij;k}(x, x').$$

If $q^i \left(= \frac{\delta x'^i}{\delta s} \right)$ and $p^a \left(= \frac{\delta u'^a}{\delta s} \right)$ are the components of the first curvature vectors of the curve with respect to F_n and F_m respectively, we have

$$(1.8) \quad q^i = p^a B_a^i + \sum_{\mu} \Omega_{(\mu)\alpha\beta}^* (u, u') \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \eta_{(\mu)}^{*i}$$

where $\Omega_{(\mu)\alpha\beta}^*$ are the components of the secondary second fundamental tensors of the subspace.

2. **Frenet's formulae.** After putting $n_{(0)}^i = \frac{dx^i}{ds}$ and $n_{(1)}^i$ as a unit vector along q^i , we have the following equations

$$(2.1) \quad q^i = \frac{\delta n_{(0)}^i}{\delta s} = K_{(1)} n_{(1)}^i$$

where $K_{(1)}$ is the first curvature of the curve. The vector $n_{(1)}^i$ is consistent with the conditions

$$(2.2) \quad g_{ij}(x, x') n_{(0)}^i n_{(1)}^j = 0 \quad \text{and} \quad g_{ij}(x, x') n_{(0)}^i n_{(0)}^j = 1.$$

The δ -derivative {refer (1b.5)} of the vector $n_{(1)}^i$ gives

$$(2.3) \quad \frac{\delta}{\delta s} n_{(1)}^i = \frac{d}{ds} n_{(1)}^i + P_{jk}^{*i} n_{(1)}^j \frac{dx^k}{ds}.$$

Decomposing this vector in the surface determined by $n_{(0)}^i$ and $n_{(1)}^i$ and orthogonal to this surface we have

$$(2.4) \quad \frac{\delta}{\delta s} n_{(1)}^i = a n_{(0)}^i + b n_{(1)}^i + c n_{(2)}^i$$

where the vector $n_{(2)}^i$ satisfies

$$(2.5) \quad g_{ij}(x, x') n_{(p)}^i n_{(2)}^j = \delta_2^p \quad (p = 0, 1, 2).$$

A simple calculation based on the equations (1.3), (1.6) (2.1), (2.4) and (2.5) gives

$$a = -K_{(1)}; b = -\frac{1}{2} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j.$$

Substituting the values of a, b in (2.4) and putting $C = K_{(2)}$ (the second curvature of the curve) we get

$$(2.6) \quad \frac{\delta}{\delta s} n_{(1)}^i = -K_{(1)} n_{(0)}^i + K_{(2)} n_{(2)}^i - \frac{1}{2} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j n_{(1)}^i.$$

The equations (2.1) and (2.6) represents the first two FRENET's formulae for a curve in the FINSLER space.

3. Hyper D-line. Consider a congruence of curves given by the vector field λ^i such that through each point of F_m there passes exactly one curve of the congruence. At the point of the subspace we have

$$(3.1) \quad \lambda^i = t^{*\alpha} B_\alpha^i + \sum_{\mu} C_{(\mu)}^* n_{(\mu)}^i.$$

The curve C is said to be a hyper D-line of the subspace if the surface spanned by the vectors $n_{(0)}^i$ and $R_{(1)} n_{(1)}^i + R_{(2)} \frac{dR_{(1)}}{ds} n_{(2)}^i$ contains the vector λ^i . ($R_{(1)} = \frac{1}{K_{(1)}}$ and $R_{(2)} = \frac{1}{K_{(2)}}$ are the radii of curvatures of the first and second order respectively.) From (2.1) and (2.6) we have

$$(3.2) \quad \frac{\delta^2}{\delta s^2} n_{(0)}^i = \frac{dK_{(1)}}{ds} n_{(1)}^i - K_{(1)}^2 n_{(0)}^i + K_{(1)} K_{(2)} n_{(2)}^i - \frac{1}{2} K_{(1)} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j n_{(1)}^i.$$

This equation gives

$$(3.3) \quad g_{ij}(x, x') \left(\frac{\delta^2}{\delta s^2} n_{(0)}^i \right) \left(R_{(1)} n_{(1)}^j + R_{(2)} n_{(2)}^j \frac{dR_{(1)}}{ds} \right) = -\frac{1}{2} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j.$$

For the hyper D-line, we have

$$(3.4) \quad \lambda^i = A \left(R_{(1)} n_{(1)}^i + R_{(2)} \frac{dR_{(1)}}{ds} n_{(2)}^i \right) + B n_{(0)}^i.$$

Multiplying this equation by $g_{ij}(x, x') n_{(0)}^j$, $g_{ij}(x, x') n_{(1)}^j$ and $g_{ij}(x, x') \frac{\delta^2}{\delta s^2} n_{(0)}^j$ respectively and putting $\lambda_{(h)} = g_{ij}(x, x') \lambda^i n_{(h)}^j$ ($h = 0, 1, 2$) we get $B = \lambda_{(0)}$, $A = \lambda_{(1)} K_{(1)}$ and

$$(3.5) \quad g_{ij}(x, x') \lambda^i \frac{\delta^2}{\delta s^2} n_{(0)}^j = -\frac{1}{2} \lambda_{(1)} K_{(1)} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j \\ + \lambda_{(0)} g_{ij}(x, x') n_{(0)}^i \left(\frac{\delta^2}{\delta s^2} n_{(0)}^j \right).$$

which with the help of equation (3.2) gives

$$(3.6) \quad \frac{dK(1)}{ds} \lambda_{(1)} + K(1) K(2) \lambda_{(2)} = 0.$$

This equation represents the equation of hyper D-line of the subspace.

Theorem (3.1). If the congruence λ^i is along the first curvature vector of the curve then the necessary and sufficient condition that it be a hyper D-line is that it be a curve of constant first curvature.

Proof. The proof follows from equation (3.6), the definition of $\lambda_{(1)}$ and $\lambda_{(2)}$ and the fact that the vectors $n_{(1)}^i$ and $n_{(2)}^i$ are orthogonal.

4. Hyper D-line involving secondary second fundamental tensors. Writing the δ -differential of equation (1.8) and using [1]

$$(4.1) \quad n_{(\mu);\gamma}^{*i} = -\psi_{(\mu)} B_{\alpha}^i g^{\alpha\delta} \Omega_{(\mu)\delta\gamma}^* - C_{jhk}^* g^{ji} B_{\gamma}^k n_{(\mu)}^{*h} + \sum_{\lambda} N_{(\lambda)\gamma}^{(\mu)} n_{(\lambda)}^{*i}$$

we obtain

$$(4.2) \quad \begin{aligned} \frac{\delta^2}{\delta s^2} n_{(0)}^i &= \left[\frac{\delta p^{\alpha}}{\delta s} - \sum_{\mu} \psi_{(\mu)} \Omega_{(\mu)\theta\beta}^* \frac{du^{\theta}}{ds} \frac{du^{\beta}}{ds} \Omega_{(\mu)\theta\gamma}^* g^{\alpha\delta} \frac{du^{\gamma}}{ds} \right] B_{\alpha}^i \\ &+ \sum_{\mu} \left[3 \Omega_{(\mu)\alpha\beta}^* p^{\alpha} \frac{du^{\beta}}{ds} + \Omega_{(\mu)\alpha\beta;\gamma}^* \frac{du^{\alpha}}{ds} \frac{du^{\beta}}{ds} \frac{du^{\gamma}}{ds} \right. \\ &+ \left. \sum_{\lambda} \Omega_{(\lambda)\alpha\beta}^* \frac{du^{\alpha}}{ds} \frac{du^{\beta}}{ds} N_{(\lambda)\gamma}^{(\mu)} \frac{du^{\gamma}}{ds} \right] n_{(\mu)}^{*i} \\ &- \sum_{\mu} \Omega_{(\mu)\alpha\beta}^* \frac{du^{\alpha}}{ds} \frac{du^{\beta}}{ds} C_{ijk}^* g^{ji} B_{\gamma}^k n_{(\mu)}^{*h} \frac{du^{\gamma}}{ds}. \end{aligned}$$

Substituting the values of $\frac{\delta^2}{\delta s^2} n_{(0)}^i$ from (4.2) and λ^i from (3.1) in the equation (3.5) we get

$$\begin{aligned}
 (4.3) \quad g_{\alpha\beta}(u, u') \frac{\delta p^\alpha}{\delta s} t^{*\beta} - \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta}^* t^{*\alpha} \frac{du^\beta}{ds} - \sum_{\mu} K_{(\mu)}^* C_{ihk}^* t^{*i} B_{\epsilon}^i n_{(\mu)}^{*h} \frac{dx^k}{ds} \\
 + \sum_{\mu} \psi_{(\mu)} C_{(\mu)}^* \left[3 \Omega_{(\mu)\alpha\beta}^* p^\alpha \frac{du^\beta}{ds} + \Omega_{(\mu)\alpha\beta;\gamma}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \frac{du^\gamma}{ds} + \sum_{\lambda} K_{(\mu)}^* N_{(\mu)\gamma}^{(\lambda)} \frac{du^\gamma}{ds} \right] \\
 - \sum_{\mu} \sum_{\nu} C_{(\nu)}^* K_{(\mu)}^* n_{(\nu)}^{*j} C_{ihk}^* n_{(\mu)}^{*h} \frac{dx^k}{ds} + \frac{1}{2} \lambda_{(1)} K_{(1)} C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j \\
 - \lambda_{(0)} \left[g_{\alpha\beta} \left[\frac{\delta p^\alpha}{\delta s} \frac{du^\beta}{ds} - \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^{*2} \right] \right] = 0.
 \end{aligned}$$

where we have used

$$(4.4) \quad K_{(\mu)}^* = \Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds}$$

and the fact that $n_{(0)}^i$ is orthogonal to $n_{(\mu)}^{*i}$.

The first two FRENET's formulae for the subspace are

$$(4.5) \quad p^\alpha = \frac{\delta}{\delta s} \left(\frac{du^\alpha}{ds} \right) = k_{(1)} \xi_{(1)}^\alpha,$$

$$(4.6) \quad \frac{\delta \xi_{(1)}^\alpha}{\delta s} = -k_{(1)} \xi_{(0)}^\alpha + k_{(2)} \xi_{(2)}^\alpha - \frac{1}{2} C_{\beta\gamma\delta}^* \frac{du^\delta}{ds} \xi_{(1)}^\alpha \xi_{(1)}^\gamma \xi_{(1)}^\beta$$

where $C_{\beta\gamma\delta}^* = g_{\beta\gamma;\delta}$, $\xi_{(0)}^\alpha = \frac{du^\alpha}{ds}$, $\xi_{(1)}^\alpha$ and $\xi_{(2)}^\alpha$ are the first and second curvature vectors and $k_{(1)}$, $k_{(2)}$ are the first and second curvatures with respect to the subspace. These equations give

$$(4.7) \quad \frac{\delta p^\alpha}{\delta s} = -k_{(1)}^2 \xi_{(0)}^\alpha + \xi_{(1)}^\alpha \frac{dk_{(1)}}{ds} + k_{(1)} k_{(2)} \xi_{(2)}^\alpha - \frac{1}{2} k_{(1)} C_{\beta\gamma\delta}^* \frac{du^\delta}{ds} \xi_{(1)}^\alpha \xi_{(1)}^\gamma \xi_{(1)}^\beta.$$

Substituting the value of $\frac{\delta p^\alpha}{\delta s}$ in (4.3) and using the fact

$$\lambda_{(0)} = g_{ij}(x, x') \lambda^i n_{(0)}^j = g_{\alpha\beta} t^{*\alpha} \frac{du^\beta}{ds} = t_{(0)}^*,$$

$$K_{(1)} \lambda_{(1)} = g_{ij}(x, x') \lambda^i q^j = g_{\alpha\beta} t^{*\alpha} p^\beta + \sum_{\mu} \psi_{(\mu)} C_{(\mu)}^* K_{(\mu)}^*$$

we obtain

$$\begin{aligned}
 (4.8) \quad & \frac{dk(1)}{ds} t_{(1)}^* + k(1) k(2) t_{(2)}^* - \frac{1}{2} k(1) C_{\beta\gamma\delta}^* \frac{du^\delta}{ds} t_{(1)}^* \xi_{(1)}^\gamma \xi_{(1)}^\beta \\
 & - \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta}^* t^{*\alpha} \frac{du^\beta}{ds} - \sum_{\mu} K_{(\mu)}^* C_{ihk}^* t^{*i} B_{\epsilon}^i n_{(\mu)}^{*h} \frac{dx^k}{ds} \\
 & + \sum_{\mu} \psi_{(\mu)} C_{(\mu)}^* \left[3 \Omega_{(\mu)\alpha\beta}^* p^\alpha \frac{du^\beta}{ds} + \Omega_{(\mu)\alpha\beta;\gamma}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \frac{du^\gamma}{ds} + \sum_{\lambda} K_{(\lambda)}^* N_{(\mu)\gamma}^{(\lambda)} \frac{du^\gamma}{ds} \right] \\
 & - \sum_{\mu} \sum_{\nu} C_{(\nu)}^* K_{(\mu)}^* n_{(\nu)}^{*i} C_{ihk}^* n_{(\mu)}^{*h} \frac{dx^k}{ds} + \frac{1}{2} (k_{(1)} t_{(1)}^* + \sum_{\mu} \psi_{(\mu)} C_{(\mu)}^* K_{(\mu)}^*) \\
 & C_{ijk}^* \frac{dx^k}{ds} n_{(1)}^i n_{(1)}^j + t_{(0)}^* \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^{*2} = 0
 \end{aligned}$$

where

$$t_{(1)}^* = g_{\alpha\beta} t^{*\alpha} \xi_{(1)}^\beta ; \quad t_{(2)}^* = g_{\alpha\beta} t^{*\alpha} \xi_{(2)}^\beta.$$

The equation (4.8) represents the hyper D-line of the subspace. The equation has been expressed in the secondary second fundamental tensors.

5. Hyper D-lines and union curves. The union curves of the subspace have been studied by SINGH [4]. For these curves we have

$$\lambda^i = A n_{(0)}^i + B \frac{\delta}{\delta s} n_{(0)}^i.$$

Using (3.1), (1.8) and (5.4), we obtain

$$\begin{aligned}
 t^{*\alpha} &= A \frac{du^\alpha}{ds} + B p^\alpha, \quad C_{(\mu)}^* = B \Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds}, \\
 t_{(0)}^* &= A, \quad t_{(1)}^* = B k_{(1)}, \quad t_{(2)}^* = 0.
 \end{aligned}$$

Substituting these values in (4.8) we get, after some simplification,

$$\begin{aligned}
 (5.1) \quad & k(1) \frac{dk(1)}{ds} + 2 \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta}^* p^\alpha \frac{du^\beta}{ds} + \sum_{\mu} \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta;\gamma}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \frac{du^\gamma}{ds} \\
 & + \sum_{\mu} \sum_{\gamma} \psi_{(\mu)} K_{(\mu)}^* K_{(\gamma)}^* N_{(\gamma)\gamma}^{(\lambda)} \frac{du^\gamma}{ds} - \sum_{\mu} K_{(\mu)}^* C_{ihk}^* p^i B_{\epsilon}^i n_{(\mu)}^{*h} \frac{dx^k}{ds} \\
 & - \sum_{\mu} \sum_{\nu} K_{(\mu)}^* K_{(\nu)}^* n_{(\nu)}^{*i} C_{ihk}^* n_{(\mu)}^{*h} \frac{dx^k}{ds} - \frac{1}{2} C_{\beta\gamma\delta}^* \frac{du^\delta}{ds} p^\gamma p^\beta \\
 & + \frac{1}{2} C_{ihk}^* \frac{dx^k}{ds} q^i q^h = 0.
 \end{aligned}$$

Further since [1]

$$(5.2) \quad \psi_{(\gamma)} N_{(\lambda)\gamma}^{(\mu)} + \psi_{(\mu)} N_{(\mu)\gamma}^{(\lambda)} = \frac{\partial \psi_{(\mu)}}{\partial \mu^\gamma} \delta_\mu^\lambda - C_{ihk}^* n_{(\lambda)}^{*h} n_{(\mu)}^{*i} B_\beta^k$$

we have

$$(5.3) \quad \sum_\lambda \sum_\mu K_{(\mu)}^* K_{(\lambda)}^* \psi_{(\mu)} N_{(\mu)\gamma}^{(\lambda)} \frac{du^\gamma}{ds} = \frac{1}{2} \sum_\mu K_{(\mu)}^{*2} \frac{d\psi_{(\mu)}}{ds} - \frac{1}{2} \sum_\mu \sum_\lambda K_{(\mu)}^* K_{(\lambda)}^* C_{ihk}^* n_{(\mu)}^{*i} n_{(\lambda)}^{*h} \frac{dx^k}{ds}$$

A simple calculation based on the equations (1.1), (1.8) and (1.4) yields

$$(5.4) \quad \frac{1}{2} C_{ihk}^* \frac{dx^k}{ds} q^i q^h - \frac{1}{2} C_{\alpha\beta\gamma}^* \frac{du^\gamma}{ds} p^\alpha p^\beta = \frac{1}{2} \sum_\mu \sum_\nu C_{ihk}^* \frac{dx^k}{ds} n_{(\mu)}^{*i} n_{(\nu)}^{*h} K_{(\mu)}^* K_{(\nu)}^* + \sum_\mu K_{(\mu)}^* C_{ihk}^* \frac{dx^k}{ds} B_\beta^h p^\beta n_{(\nu)}^{*i}$$

With the help of equations (5.3) and (5.4), the equation (5.1) reduces to

$$(5.5) \quad k_{(1)} \frac{dk_{(1)}}{ds} + 2 \sum_\mu \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta}^* p^\alpha \frac{du^\beta}{ds} + \sum_\mu \psi_{(\mu)} K_{(\mu)}^* \Omega_{(\mu)\alpha\beta\gamma}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \frac{du^\gamma}{ds} + \frac{1}{2} \sum_\mu K_{(\mu)}^{*2} \frac{d\psi_{(\mu)}}{ds} - \sum_\mu \sum_\nu K_{(\mu)}^* K_{(\nu)}^* C_{ihk}^* \frac{dx^k}{ds} n_{(\mu)}^{*i} n_{(\nu)}^{*h} = 0$$

or

$$(5.6) \quad k_{(1)} \frac{dk_{(1)}}{ds} + \frac{1}{2} \frac{\delta}{\delta s} \left[\sum_\mu \psi_{(1)} \left(\Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \right)^2 \right] - 2 \sum_\mu \sum_\nu K_{(\mu)}^* K_{(\nu)}^* M_{(\mu\nu)k} q^k = 0$$

where

$$(5.7) \quad M_{(\mu\nu)k} = C_{ihk}^* n_{(\mu)}^{*i} n_{(\nu)}^{*h}$$

Assuming

$$(5.8) \quad \sum_\mu \sum_\nu K_{(\mu)}^* K_{(\nu)}^* M_{(\mu\nu)k} q^k = 0$$

and using the fact

$$\frac{\delta}{\delta s} \left[\sum_\mu \psi_{(\mu)} \left(\Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \right)^2 \right] = \frac{d}{ds} \left[\sum_\mu \psi_{(\mu)} \left(\Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \right)^2 \right]$$

we get after integrating (5.6)

$$(5.9) \quad \frac{1}{r_1^2} + \frac{1}{\varrho^2} = \frac{1}{a^2}$$

where $r(1) = \frac{1}{k_{(1)}}$ and $\varrho = \left[\sum_{\mu} \psi_{(\mu)} \left(\Omega_{(\mu)\alpha\beta}^* \frac{du^\alpha}{ds} \frac{du^\beta}{ds} \right)^2 \right]^{-\frac{1}{2}}$ are the radii of geodesic

(with respect to the subspace) and the secondary normal curvatures and $\frac{1}{a^2}$ is the constant of integration.

We have thereby established the following theorems:

Theorem (5.1). The sum of the squares of the geodesic and the secondary normal curvatures in the direction of a union hyper D-line is the same at all those points where the vector $M_{(\mu\nu)k}$ is orthogonal to the first curvature vector q^k :

Theorem (5.2). At all those points where the vector $M_{(\mu\nu)k}$ is orthogonal to q^k , the sum of squares of the geodesic and secondary normal curvatures in the direction of the union hyper D-line is the same relative to every congruence λ^i .

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ÖZET

Üç boyutlu bir Euklid uzayında bulunan bir yüzeye ait hiper DARBOUX çizgileri (hiper D-çizgileri) PRYANOVITCH tarafından tanımlanmış ve incelenmiştir [2]. SINGH, bir RIEMANN uzayının alt uzaylarındaki benzer çizgilerin bazı özelliklerini incelemiştir [3]. Bu yazının gayesi bir FINSLER uzayının bir alt uzayına ait hiper D-çizgilerinin diferansiyel denklemini elde etmek ve bundan, birleşim eğrileri olan hiper D-çizgilerinin bazı özelliklerini bulmaktır.