Investigation of Some Parameters Which Affects Into The Efficiency of Quantum Dot Intermediate Band Solar Cell

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Abstract- In this paper, Quantum dots intermediate band solar cell (QDIBSC) is used to the enhancement of the power conversion efficiency of solar cell. The main advantage of this type is that it preserves large open-circuit voltage with increasing the produced photocurrent in the solar cell. For a best efficiency of one-intermediate band solar cell (one-IBSC), the induced detailed balance efficiency is determined as a function of changing locations of intermediate-band (IB) by using the blackbody. The QDs have the ability to confirm the chosen higher efficiency by assigning the appropriate values of quantum dot width size (QDW) and barrier thickness (BT). It means that the best location of IB in solar cell is realized. The results show that to obtain the maximum power conversion efficiency of QDIBSC, the QDWs and BTs for nanostructured model (Al0.4Ga0.6As/In0.42Ga0.58As) are limited by a surface contour. The highest power efficiencies in this located contour are 45.32% and 62.81% for QDWs = (1.60 nm, 1.64 nm) and BTs = (1.98 nm, 1.94 nm) for 1 sun and maximum light concentrations; respectively.

Keywords-Quantum Dots Width, Barrier Thickness, Intermediate Band, Efficiency, Nanostructured Solar Cell.

1. Introduction

Throughout the past decades, the whole world started to feel the seriousness of the deficiency in the fossil fuel. Therefore, a large number of researchers around the world began to research and study the alternative energy to compensate the loss in traditional energy. Solar energy is one of the important options for alternative energy and improving this technology can clearly reduce the problem of energy in the world [1-5]. Consequently, a number of researchers at Nano-power Research Laboratories and others focus on the improvement of the power conversion efficiency of solar cells by using nanostructures with suitably adjusted properties [6-8]. The new proposed concepts are called thirdgeneration solar cell, such as hot carrier cells, tandem cells, and intermediate-band cells [4, 7, 9-12]. The intermediate energy bands are established within a forbidden gap of a bulk semiconductor material when a superlattice low-dimensional nanostructure (quantum well, wire, and dot (QW, QR, and QD respectively)) of other semiconductors material injected within a bulk semiconductor material [13-15]. This structure is called heterostructure and it allows for the proper operation of IBSC. For QW and QR heterostructures, the photogenerated charge carriers will lose energy through thermalization due to the high density of states available in the in-plane and longitudinal direction; respectively [16]. Also, the Fermi level will not split into the number of bands, but it will split into two levels similar to the traditional solar cell. Moreover, the lowest intermediate band will act as the conduction band and the heterostructure will behave like the traditional solar cell [17-18]. For QDs heterostructure, it will provide zero density of states between the excited bands due to the discrete energy spectrum; thus the thermalization is reduced. Each band is thermally isolated; therefore the Fermi energy level will split into the number of bands and charge concentration in each band is described by its own chemical potential as required for improving IBSC operation [16, 19]. Fig. 1(a) shows QDIBSC with the barrier material surrounding the QDs. QDs are made of the materials such as Ge, InAs, InAsN,..etc., they have smaller bandgap energy than the bulk materials such as Si, GaAs, GaAsN,..etc [4, 6, 7, 10, 11]. The superlattice structure is an array of closely

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spaced QDs, therefore a sufficient amount of wavevectors overlapped to form the intermediate energy bands as shown in Fig. 1(b) [8, 12]. These bands are critical in the nanostructured solar cell because they allow the absorption of low energy photons in addition to the normal absorbed photons processes. Therefore, the efficiency is improved, as a result of increasing the generated photocurrent [20]. The organization of this paper includes the following parts. Part II includes the related work for QDIBSC. The analyses of the mathematical model are described in Part III. The one-intermediate band solar cell is discussed in Part IV. The quantum dot intermediate band solar cell at different QDs width sizes and barrier thickness is described in Part V. Finally, Part VI is the conclusions of this article.



Fig. 1. (a) Schematic diagram of the QDs inside a bulk semiconductor of larger bandgap (b) Intermediate energy bands of an array of QDs

2. Related Work

Since 1997, IBSC have received the attention of a large number of research groups due to the potential of the IB to increase the power conversion efficiency of solar cell devices from 41% to 63% under maximum light concentration [21]. The key working principle of the IBSC have been demonstrated in a InAs/GaAs quantum dot solar cell [22], though so far to the knowledge of the author no net increase of the efficiency has been shown by an IB device respect to a device with the same characteristics but lacking the IB. Reviews of the topic can be found in references [23, 24]. Many authors report strong enhancement of the photocurrent [25] due to subbandgap absorption, and even large increase in efficiency compared to reference devices [2]. Sablon et al. [2] fabricated InAs/GaAs quantum dot solar cells and reported a large increase in current without significant voltage degradation. This field has expanded a lot, as much to propose materials that show an IB without the insertion of nanostructures [26, 27]. Other models proposed also include space-charge and other recombination effects to best fit the experimental data [28, 29]. Recent progress has also been made with the growth of InAs/GaAs quantum dot solar cells with MOCVD [4] and even transfer to flexible substrates by epitaxial lift-off [30]. A possible limitation to this technology comes from strain-induced defects that can generate in the OD layers [31], and eventually limit the number of layers that can be grown. Higher densities of QDs in strain compensated materials have been proposed and realized [32, 33]. An interesting feature of strain-compensated QD multilayer's is the self-ordering effect observed in multilayered samples [34]. This ordering may lead to the formation of a mini-band that would serve as the perfect IB as theorized in [21].

3. Mathematical Model Analyses

The principle of a mathematical model to describe the operation of a solar cell is built on using the sun as a blackbody. The generalization of Planck's radiation law for blackbody and solar radiation is [5, 14, 16]

$$S_R = \frac{2\pi hc^2}{\lambda^5} \frac{1}{e^{((hc/\lambda-\mu)/kT)} - 1}$$
(1)

Where S_R is a solar spectrum, λ is wavelength of light, h is Planck's constant, c is the speed of light, μ is the chemical potential, k is Boltzmann's constant, and T is the temperature of a blackbody. From the Roosbroeck-Shockley equation, the flux of photons (φ) absorbed in or emitted from the semiconductor is given by [19]

$$\varphi(\lambda_a, \lambda_b, T, \mu) = \int_{\lambda_a}^{\lambda_b} S_R \, d\lambda \tag{2}$$

Where λ_a and λ_b are the lower and upper wavelengths of the light at the energy limits of the photo flux. When handling the sun as an ideal blackbody and the solar absorption of a semiconductor is limited by the bandgap energy E_g , the flux of photons absorbed at the sun's surface (φ_s) is

$$\varphi_s(0,\lambda_g,T_s,0) = \int_0^{\lambda_g} \frac{2\pi hc^2}{\lambda^5} \frac{1}{e^{((hc/\lambda-\mu)/kT_s)} - 1} d\lambda$$
(3)

Where λ_g is the wavelength of light at the bandgap energy of the solar semiconductor, T_s is the temperature of the sun, 6000K. The photon flux absorbed by the solar cell, φ_{ssc} , is shown in Fig. 2 and denoted by:

$$\varphi_{ssc} = \varphi_s \times f_s \tag{4}$$

Where $f_s = 2.1646 \times 10^{-5}$ is called geometric parameter [16].

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Similarly, the solar cell may be considered as an ideal blackbody at temperature $T_c = 300$ K. Then, a photon emission from the solar cell with a flux φ_{sc} is emitted when an electron-hole pair is directly recombined. The flux φ_{sc} is the same as in Equation 3, but the surface sun temperature, T_s is replaced by the surface solar cell temperature, T_c . According to the Shockley-Queisser model, the current density (*j*) of the solar cell device can be written as [35] $j = q[c_s f_s \varphi_s(0, \lambda_{a_l} T_{c_l} 0) + (1 - c_s f_s) \varphi_{sc}(0, \lambda_{a_l} T_{c_l} 0) -$

$$\varphi_{sc}(0,\lambda_g, T_c, \mu)] \tag{5}$$

Where c_s is the light concentration on a solar cell and it is called the number of suns, i.e., 1 sun means that $c_s = 1$ at the

surface of the Earth's atmosphere and 46198 suns mean that $c_s = 1/f_s$ or maximum light concentration at the surface of the sun's. The chemical potential of solar photons (μ) is equal to the applied bias (qV) which it is equal to the number of splitting in the quasi-Fermi energy levels, where q is electron charge and V is the applied voltage of the solar cell. This equation represents the detailed balance formula; it directly provides the current-voltage characteristics of a solar cell. Also, the solar power conversion efficiency can be directly calculated.



Fig. 2. Solar spectrum versus wavelength at blackbody

4. One-Intermediate Band Solar Cell

The main challenge in improving the performance of the traditional solar cells is the case where the increasing output current decreases the output voltage and vice versa. A number of articles discussed how to handle this problem by making a tradeoff between the materials used in the manufacturing of solar cells. The traditional solar cell would only be able to absorb photons with energy equal to or greater than the energy between conduction and valence bands (E_{CV}). For the nanostructured solar cell, new bands are located at an intermediate level (E_l) between conduction and valence bands. These bands increase the absorption of photons that lead to the increased performance of this design when compared to the traditional solar cell. To study the mathematical model of this type, the following assumptions are necessary [5, 12, 36, 37]:

- i. There is no overlap of energy transitions for a given photon energy; i.e., if a photon may energetically induce a transition from one band to another, then all photons of the same energy will only cause that particular transition.
- ii. Electrons enter the IB only when transferred from the valence band (VB) and they leave only to the conduction band (CB).

iii. The difference in chemical potentials between any two bands causes only the difference between the quasi-Fermi levels of these bands.

In the following, the study will concern a one-IBSC taking into account the considered assumptions. It has three transitions of electrons between bands. (1) The standard effect of an electronic transition across the traditional bandgap E_{CV} generates the current density (j_{CV}) . (2) The effect of the intermediate to conduction band transition (E_{CI}) produces the current density (j_{CI}) . (3) The transition between valence and intermediate bands (E_{IV}) will produce a current density (j_{IV}) . One can notice that j_{CI} must be equal to the current density j_{IV} . The two energy intermediate band transitions, E_{CI} and E_{IV} are independent of each other; however, the traditional bandgap E_{CV} is a function of these bands. Thus, the three generated current densities are

$$j_{CV} = q[c_s f_s \varphi_s(0, \lambda_{CV}, T_s, 0) + (1 - c_s f_s) \varphi_{sc}(0, \lambda_{CV}, T_c, 0) - \varphi_{sc}(0, \lambda_{CV}, T_c, \mu)]$$

$$j_{CI} = q[c_s f_s \varphi_s(\lambda_{CI}, \lambda_{IV}, T_s, 0) + (1 - c_s f_s) \varphi_{sc}(\lambda_{CI}, \lambda_{IV}, T_c, 0) - \varphi_{sc}(\lambda_{CI}, \lambda_{IV}, T_c, \mu_{CI})]$$

$$(6) \qquad j_{IV} = q[c_s f_s \varphi_s(\lambda_{IV}, \lambda_{CV}, T_s, 0) + (1 - c_s f_s) \varphi_{sc}(\lambda_{IV}, \lambda_{CV}, T_s, 0) + (1 - c_s f_s) \varphi_{sc}(\lambda_{IV}, \lambda_{CV}, T_c, \mu_{VI})]$$

Since the current densities j_{CI} and j_{IV} must be equal, therefore the total current density $j_T = j_{CV} + j_{CI}$. As in previous analyses [13, 16, 38], the chemical potential of photons is

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equal to the applied bias ($\mu = qV = \mu_{CI} + \mu_{IV}$), which is equal to the number of splitting in the quasi-Fermi energy levels. Fig.3 analyzes the detailed balance efficiency under blackbody solar spectrum at 1 sun and maximum light concentrations.



Fig. 3. Detailed balance efficiency of one-IBSC varies with the place of IB between the CB and VB at blackbody spectrum under 1 sun, and maximum light concentrations

The place of one-IB that achieves the maximum power efficiency can be recognized by changing the limited values of bandgap energies E_{IV} and E_{CI} . The maximum efficiency is reached to 46.74% which corresponds to the optimum intermediate energy band position at $E_{IV} = 0.93$ eV and $E_{CI} = 1.5$ eV; the total bandgap energy $E_{CV} = 2.43$ eV. When the

light concentration is increased to a maximum, the maximum efficiency is increased to 63.13% and the optimum intermediate energy band position in this case will be $E_{IV} = 0.72\text{eV}$ and $E_{CI} = 1.25\text{eV}$; $E_{CV} = 1.97\text{eV}$, these determined values also shown in Table 1.

Table 1.	Summarizes	the maximum	n efficiency,	optimum	bandgap,	and IB	for any	one-IBS	C at 1	sun and	l maximun	ı light
			concentrat	ions for th	e blackbo	dy spec	trum					

	1 sun lig	ght conce	entratior	ı	Max. light concentration					
Solar	Max.	Opt. Bandgaps (eV)			Max.	Opt. Bandgaps (eV)				
Spectrum	Efficiency (%)	Eci	EIV	Ecv	Efficiency (%)	Есі	EIV	Ecv		
Blackbody	46.74	1.50	0.93	2.43	63.13	1.25	0.72	1.97		

One can notice that, in the case of maximum light concentration, the IB energies, E_{CI} and E_{IV} , and then corresponding bandgap energy, E_{CV} , are decreased. Practically, there is a critical note should be taken into consideration that is the order of the energy intermediate bands E_{CI} and E_{IV} does not affect the efficiency of the one-IBSC. Because of the efficiency depends only on the transition energies themselves [16]. For example, the theoretical efficiency is equal to 63.13% under maximum

light concentration when $E_{CI} = 1.25$ eV and $E_{IV} = 0.72$ eV. However, the same result is also achieved when $E_{CI} = 0.72$ eV and $E_{IV} = 1.25$ eV, therefore the order of the transition energies is unimportant. From Fig. 3, it is clear that in the case of the maximum light concentration, not all the values of E_{IV} and E_{CI} are included in the 3D curve, which indicates the difficulty to realize the transition energies condition for the same values of induced currents; $j_{CI} = j_{IV}$. This is the reason not all efficiency values are dominated. For further declaration, the light concentration range is scanned as shown in Fig. 4.



Fig. 4. Maximum efficiency and optimum bandgap energy,

 E_{cv} , versus with entire light concentration for the one-IBSC It confirms the previous results and provides the all possibilities of maximum power efficiency at specified optimum bandgap energy. Of course, the change of incident solar light concentration will need different optimum bandgap energy. Also by investigating Fig. 4, one can recognize that when the solar light concentration increases, the maximum balanced power efficiency for the IBSC increases while the optimum bandgap energy, E_{cv} , decreases. From this behavior, it is clear that the bandgap energy of semiconductor material is inversely proportional to temperature under blackbody spectrum [38].

5. Quantum Dot Intermediate Band Solar Cell

For this type of solar cell, the size, shape, and spacing of the quantum dots in the barrier material are critical to improve the efficiency of the device. Whenever the QDs are arranged uniformly in the barrier material, the IBs are wellplaced and the recombination is reduced [15]. Various techniques such as Metal Organic Chemical Vapor Deposition (MOCVP) and Molecular Beam Epitaxy (MBE) are available to produce highly uniform QD arrays in barrier material [38]. In this paper, the geometry of the QD is assumed cubic and thus it is characterized by one dimensional of cubic. Also, to simplify the analysis no valence band offset is assumed and the only confining potential occurs at the conduction band offset. Under these assumptions, the time-independent Schrödinger equation and the Krong-Penny model are used to calculate the energy bands in the QDs using the effective mass of charge carriers [10, 38, 39- 41]. Al_{1-x}Ga_xAs/In_{1-x}Ga_xAs QDIBSC is

analysed theoretically to show the effect of QDs width size and the spacing between them in the barrier material on the efficiency. Here $Al_{1-x}Ga_xAs$ is the barrier ternary alloy and $In_{1-x}Ga_xAs$ is the QD ternary alloy. The energy bandgaps and effective masses of these alloys are calculated by the following chemical formulas [7, 42, 43]:

$$\begin{split} Eg_{(A11-xGaxAs)} &= 3.099(1-x) + 1.519x - 0.2x(1-x), \\ x_{GaxAs)} &= (1-x)m^*_{(A1As)} + xm^*_{(GaAs)} \end{split}$$

$$\begin{split} Eg_{(InxGa1-xAs)} &= 1.519(1\text{-}x) + 0.417x - 0.588x(1\text{-}x), m^*_{(InxGa1-xAs)} \\ &= (1\text{-}x)m^*_{(GaAs)} + xm^*_{(InAs)} \end{split}$$

Here x (limited by $0 \le x \le 1$) is called the molar concentration for each GaAs and InAs in barrier and QD ternary alloys materials, respectively. There are two important steps in the study of this type of technology. (1) How to compound the semiconductor materials to obtain the required energy bandgaps. This step is studied in details in several researches [43-46]. (2) Studying the effect of QDW and BT on the efficiency of solar cell and also knowing the permissible limits of them into the selected materials to obtain one-IB in the conduction band offset. More attention in this work is given to the second point. The semiconductors compound model used in this study of QDIBSC is Al_{0.4}Ga_{0.6}As/In_{0.42}Ga_{0.58}As [44-46]. The bandgap of ternary alloy In_{0.42}Ga_{0.58}As (0.913eV) dot array material is smaller to that of the bandgap of ternary alloy Al_{0.4}Ga_{0.6}As (2.103eV) barrier or bulk material to induce the one-IB by coupling of confined electronic states in the In_{0.42}Ga_{0.58}As conduction band. The results of bandgap energies for the structure Al_{0.4}Ga_{0.6}As/In_{0.42}Ga_{0.58}As QDIBSC are approximately close to the ideal values determined theoretically from Table 1 as in the previous section. Fig. 5 shows the changes of the one-IB energy when each of QDW and BT changes.



Fig. 5. The relation between width of one-IB energy and each of QDW and BT

It is clear that the larger values of QDW and BT tend to decrease the one-IB bandwidth and the later moves toward the valence band. Therefore, E_{ci} increases and E_{iv} decreases as shown in Figs. 6 and 7.



This work concerns with the one-IB; therefore the QDW and BT in this type of ternary materials are limited with the values in Fig. 5, i.e., $1 \text{ nm} \le \text{QDW} \le 3 \text{ nm}$ and $0.5 \text{ nm} \le \text{BT}$ \leq 2 nm. The detailed balance efficiency for this structure under the above assumptions with changing QDW and BT, and thus E_{iv} and E_{ci} , at 1 sun and maximum light concentrations are plotted in Figs. 8 and 9. At first glance, when comparing the efficiency in Figs. 8, 9 and Table 2, with the results in Fig. 3 and Table 1, it is clear that the efficiency is approximately the same. This is due to the benefit from the ternary compound semiconductor materials. From the results of efficiency in Figs 8 and 9, the maximum power conversion efficiency of Al_{0.4}Ga_{0.6}As/In_{0.42}Ga_{0.58}As QDIBSC as a function of QDW and BT or E_{vi} and E_{ci} is changed from 45.32% to 62.81% at 1 sun and maximum light concentrations, respectively.

When one compares these values with the previous results in the one-IBSC, it is clear that the values are smaller by a little bit. This is natural because we are limited by the selected QDIBSC materials. Also, one can notice that the coverage or matching areas are probably small. So the movement in QDW range that is covered by the defined BT range has become narrower. When QDW decreases towards a lower limit, BT increases to an upper limit. Another thing to notice from these curves is that the maximum balanced efficiencies QDIBSC behave the same trend as in an ideal case, one-IBSC. That means that it realizes current-matching easily for 1 sun compared to the maximum light concentration, which is very difficult. Table 2 summarizes the maximum efficiency and optimum QDW and BT of QDIBSC at 1 sun and maximum light concentrations.



Fig. 7. Changing energy E_{iv} versus QDW and BT







Fig. 9. Detailed balance efficiency of QD solar cell one- IB varies with E_{iv} and E_{ci} at blackbody spectrum under 1 sun and maximum light concentrations

Table 2. Summarizes the maximum efficiency and optimum BT and QDW of QDIBSC at 1 sun and maximum light concentrations for the blackbody spectrum

		1 Sun Con	centration		Max. Concentration					
Solar Spectrum	Max. Efficienc	Opt. BT & QDW (nm)		IB Width	Max. Efficiency	Opt. BT (ni	IB Width			
	y (%)	BT	QDW	(eV)	(%)	BT	QDW	(ev)		
Blackbody	45.32	1.98	1.60	0.0696	62.81	1.94	1.64	0.0684		

6. Conclusion

In this article, the detailed balance efficiency of nanostructured solar cell dependent on Shockley and Queisser design model has been studied under the standard blackbody spectrum. For the nanostructured solar cell, the analyses of detailed balance efficiency showed that it varies with the place of the one-IB within the bandgap energy of the bulk semiconductor. Therefore, the optimum location of one-IB that provides the maximum power conversion efficiency can be determined. Accordingly, the QDIBSC is investigated to achieve the results of one-IBSC. By utilizing the Schrödinger equation and the Krong-penny model, the relation between QDIBSC parameters, (QDW and BT), and detailed balance efficiency for achieving maximum efficiencies are defined. The results have shown that both the location of the one-IB and the maximum power conversion efficiency strongly depend on the considered parameters. The maximum efficiencies are changed from 46.74% and 63.13% for one-IBSC to 45.32% and 62.81% for Al_{0.4}Ga_{0.6}As/In_{0.42}Ga_{0.58}As QDIBSC under the 1 sun and maximum light concentrations; respectively. In the QDIBSC, it is very difficult to obtain the current-matching. In future work, increasing the energy bands to two-IB in QDIBSC to improve the efficiency will be considered.

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