

## LEAST-SQUARES FITS OF FUNDAMENTAL PARAMETERS FOR X-RAY ANALYSES AS A FUNCTION OF Z ( $10 \leq Z \leq 93$ )

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### ABSTRACT

In order to make elemental analyses with X-ray fluorescence technique without standard elements, all of the fundamental parameters (Edge energy, Characteristic energy, Fluorescence yield, Jump factor, Transition probability and so on) corresponding to each element have to be known. In this study, the fundamental parameters of elements for which  $10 \leq Z \leq 93$  have been fitted to polynomials with the best determination coefficients ( $r^2$ ). The results are in good agreement with a previous work for which  $11 \leq Z \leq 83$ , and we extend the upper range of Z from  $Z=83$  to  $Z=93$  with better determination coefficients in the fitting process.

### INTRODUCTION

Numerical values of fundamental parameters have central importance in XRF analyses of elements without standards. The fundamental parameters of an element are the edge energy ( $E_{edge}$ ), the characteristic energy ( $E_{char}$ ), the fluorescence yield ( $w_i$ ), the jump factor ( $1 - 1/r_i$ ), the transition probability ( $p_i$ ) and so on.

In this study the fundamental parameters of elements for which  $10 < Z < 93$ , are fitted to polynomials which are functions of the atomic number Z. The fitted polynomials are so determined that to give the best values for the determination coefficient ( $r^2$ ). Our results are compared with the results of previous works (Poehn et al., 1985, Fink et al., 1966) and the range of Z is extended 11 elements beyond their range.

### THE SHIRAIWA AND FUJINO EQUATION

The elemental weight fraction of samples are determined by making use of the Shiraiwa and Fujino (1966) equation in computerized calculation of concentrations. The equation reads;

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$$I_{i,s} = q Q_i C_i \frac{\mu_i(E_0)}{\mu_s(E_0) + A \mu_s(E_i)} \left[ 1 + \frac{1}{2} Q_j C_j \frac{\mu_j(E_0)}{\mu_i(E_0)} L_0 \frac{\mu_i(E_j)}{\mu_s(E_j)} \right] \quad (1)$$

where

$$L_0 = \frac{\ln [1 + \mu_s(E_0) / \sin \psi_1 \mu_s(E_j)]}{[\mu_s(E_0) / \sin \psi_1 \mu_s(E_1)]} + \frac{\ln [1 + \mu_s(E_i) / \sin \psi_2 \mu_s(E_j)]}{[\mu_s(E_i) / \sin \psi_2 \mu_s(E_j)]}$$

and

$$Q_j = (1 - 1/r_j) w_j f_j$$

In order to save the computer memory location, the fundamental parameters appearing in the eq. (1), are expressed as function of the atomic number Z. In the previous works done by Poehn et al. and Fink et al., the fundamental parameters are fitted to different order of polynomials in Z. In those studies values of the determination coefficient ( $r^2$ ) for their polynomials are subject for discussions as it is commented in the conclusion and the range of the atomic number Z is not as broad as it is in this work. We enlarged the range of Z as  $10 \leq Z \leq 93$  and selected the polynomials for which  $r^2$  has the best values.

### LEAST SQUARES FITS OF FUNDAMENTAL PARAMETERS OF ELEMENTS

Any fundamental parameter F(Z) of elements is expressed as a function of its atomic number Z as the following

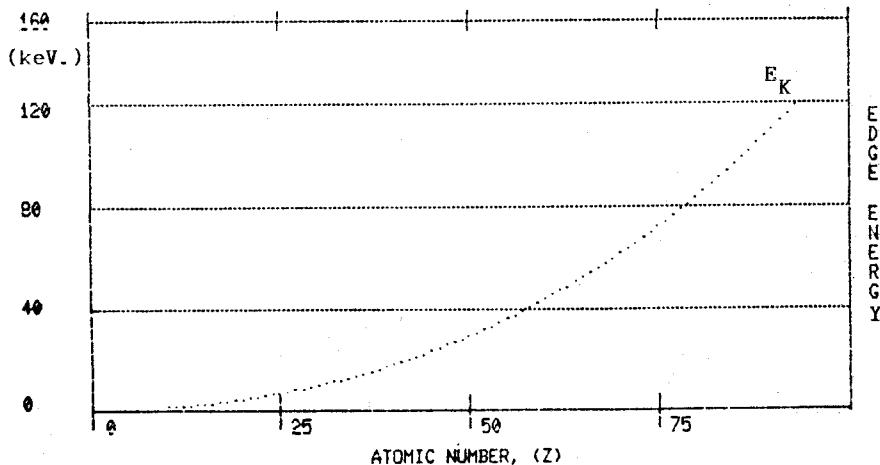
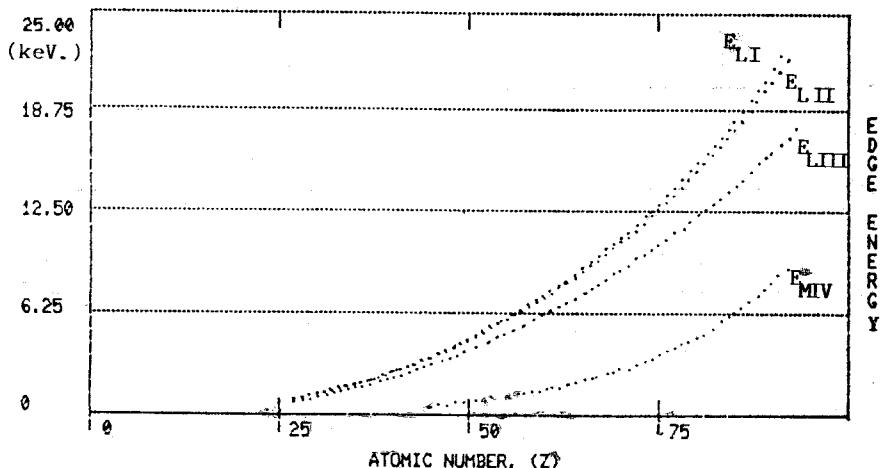
$$F(Z) = \sum_{n=0}^N A_n Z^n \quad (2)$$

where  $A_n$  is the regression coefficient. The regression coefficients in the eq. (2) can be determined by the method of least squares. A computer program, (BASICA-REG) is prepared to determine the coefficients of the polynomials up to the order of  $n=N=7$ . Consequently  $A_n$  values are determined by the computer programme.

For the case of edge energies;  $E_K$ ,  $E_{LI}$ ,  $E_{LII}$ ,  $E_{LIII}$  and  $E_{MIV}$  values are taken from Bearden (1967), and fitted to the polynomials which give the best determination coefficients. The range of the edge energy is  $0.6 \text{ KeV} \leq E_{\text{crit}} \leq 116 \text{ KeV}$ . and the corresponding range of the atomic number is  $10 \leq Z \leq 93$ . The fitted values of the regression coefficients are listed in the Table 1. And the fitted curves of the critical energy as a function of Z are given the Fig. 1 and 2.

Table 1. Fit Parameters of Edge Energies

$E/A_n$	$A_0$	$A_1$	$A_2$	$A_3$	$A_4$	$A_5$	Z
K	-0.288	3.165E-2	7.041E-3	1.412E-6	7.607E-9	7.607E-9	10-93
L <sub>I</sub>	1.096	-8.920E-2	2.834E-3	2.169E-5	-4.070E-7	2.733E-9	27-97
L <sub>II</sub>	0.978	-6.698E-2	1.500E-3	4.932E-5	-6.872E-7	3.794E-9	27.92
L <sub>III</sub>	.995	-6.531E-2	1.092E-3	5.802E-5	-7.856E-7	3.657E-9	27.92
M <sub>V</sub>	-5.259	0.220	-2.290E-3	-1.321E-5	3.542E-7	-3.35;E-10	45.92

FIG. 1. LSF Edge energy,  $E_K$ FIG. 2. LSFs of Edge energies,  $E_{LI}$ ,  $E_{LII}$ ,  $E_{LIII}$ ,  $E_{LMIV}$

For the case of the characteristic energies, values of  $E_{K\alpha}$  and  $E_{L\alpha}$  are calculated from the following equations:

$$E_{K\alpha} = \frac{2E_{K\alpha 1} + E_{K\alpha 2}}{3} \quad (3a)$$

and

$$E_{L\alpha} = \frac{9E_{L\alpha 1} + E_{L\alpha 2}}{10} \quad (3b)$$

for each element separately. Values of the quantities on the right hand side of the eq. (3) are taken from Bertin (1975). The numerical values found from this calculation is fitted to the best fit polynomial which is a function of the atomic number  $Z$ . The fitted values of the regression coefficients are given in the Table 2. And the fitted curve of the characteristic energy as a function of  $Z$  is given in the Fig. 3 for  $E_{K\alpha}$  and  $E_{L\alpha}$  separately.

Table 2. Fit Parameters of Characteristic  $K\alpha$  and  $L\alpha$  Radiations.

$E / A_n$	$A_0$	$A_1$	$A_2$	$A_3$	$A_4$	$Z$
K	0.220	-4.898E-2	1.604E-2	-3.266E-5	3.985E-7	10-90
L111	-0.577	3.109E-2	2.384E-4	1.872E-5	-7.416E-8	30.92

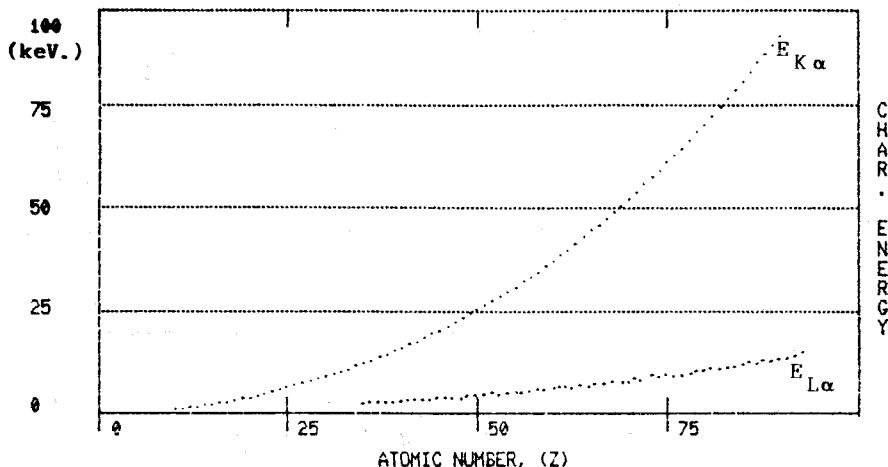


FIG. 3. LSFs of Characteristic  $K\alpha$  and  $L\alpha$  energies

For the case of the jump factor, the fluorescence yield and the transition probability, the experimental values of McMaster et al. (1969), Bambynek et al. (1972), Schreiber and Wims (1982) are used respectively in these calculations. The corresponding fitted values of the regression coefficients are given in the Tables 3-5. The plots of the jump factor and the fluorescence yield are given in the Figs. 4 and 5 respectively.

Table 3. Fit parameters for jump Factor

J/A <sub>n</sub>	A <sub>0</sub>	A <sub>1</sub>	A <sub>2</sub>	A <sub>3</sub>	A <sub>4</sub>	A <sub>5</sub>	A <sub>6</sub>	A <sub>7</sub>	Z
K	1.051	-1.251E-2	1.485E-4	1.411E-5	-6.046E-7	2.858E-9	2.247E-10	-3.065E-12	15-48
L <sub>III</sub>	4.463	-0.349	1.042E-2	-1.174E-4	1.635E-7	-5.929E-9	2.590E-10	-1.839E-12	35.70

Table 4. Fit Parameters of Fluorescence Yields wK and wL

W/A <sub>n</sub>	A <sub>0</sub>	A <sub>1</sub>	A <sub>2</sub>	A <sub>3</sub>	A <sub>4</sub>	A <sub>5</sub>	Z
K	0.293	-6.416E-2	4.296E-3	-7.954E-5	4.770E-7	—	10-55
L <sub>III</sub>	0.148	-8.184E-3	8.361E-5	1.586E-6	-8.156E-9	-4.276E-11	30-90

Table 5. Fit Parameters of Transition Probabilities, PK

P <sub>K</sub> /A <sub>n</sub>	A <sub>0</sub>	A <sub>1</sub>	A <sub>2</sub>	A <sub>3</sub>	A <sub>4</sub>	A <sub>5</sub>	Z
I	-1.224	0.132	-4.372E-3	5.205E-5	—	—	20-32
II	-3.557E-2	0.729E-3	-1.382E-4	1.020E-5	-1.864E-7	7.384E-10	33-44
III	-0.247	1.245E-2	1.190E-4	-6.334E-6	6.432E-8	-1.913E-10	55-85

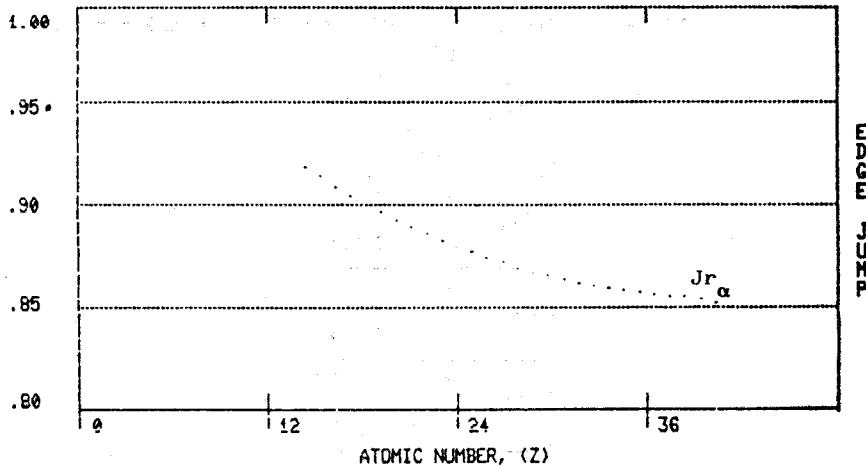


FIG. 4. LSF of J. Jump factor

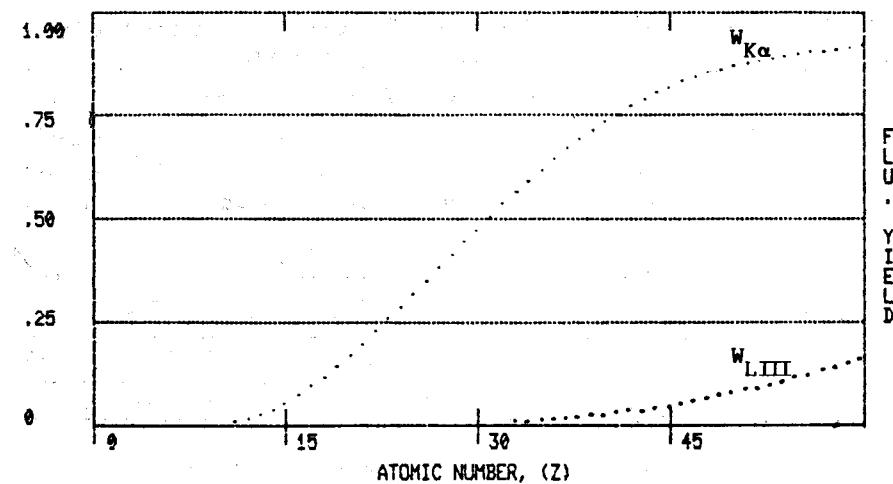


FIG. 5. LSFs of wK and wLIII fluorescence yield

## CONCLUSION

Least-square fits of fundamental parameters are very beneficial in XRF analyses of elements without standards. The determination coefficients of the previous works done by Poehn et al. and Fink et al. range  $0.985 \leq r^2 \leq 0.998$ , while the values of  $r^2$  for this work are almost ideal. Our values for  $r^2$  are tabulated in the Table 6. The range of Z we

Table 6. The Determination Coefficient of Fundamental Parameters

$F_{edge}$	$E_K$	1.000
	$E_{L1}$	1.000
	$E_{LII}$	1.000
	$E_{LIII}$	1.000
	$E_{MIV}$	1.000
$E_{char}$	$E_{K\alpha}$	1.000
	$E_{L\alpha}$	1.000
$(1-1/r)$	K	1.000
	L <sub>111</sub>	0.994
$W$	K	1.000
	L <sub>111</sub>	1.000
$P$	KI	1.000
	KII	0.998
	KIII	0.997

used is  $10 \leq Z \leq 93$ , which covers almost all the elements with which one may perform tubular or radioisotope excitations in XRF analyses.

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