

# COMMUNICATIONS

DE LA FACULTÉ DES SCIENCES  
DE L'UNIVERSITÉ D'ANKARA

Série A<sub>1</sub> : Mathématiques

---

TOME 32

ANNÉE 1983

---

**The Singular Solutions of the Generalized Poly-Axially Symmetric  
Helmholz Equations**

by

**İ. Ethem ANAR**

12

Faculté des Sciences de l'Université d'Ankara  
Ankara, Turquie

# Communications de la Faculté Sciences de l'Université d'Ankara

Comité de Redaction de la Série A,

Cengiz Uluçay – Cevat Kart – H. M. Hacısalihoğlu

Secrétaire de Publication

Ö. Çakar

---

La Revue "Communications de la Faculté des Sciences de l'Université d'Ankara" est un organe de publication englobant toutes les disciplines scientifique représentées à la Faculté de Sciences de l'Université d'Ankara.

La Revue, jusqu'à 1975 à l'exception des tomes I, II, III était composé de trois séries

Série A : Mathématiques, Physique et Astronomie,

Série B : Chimie,

Série C : Sciences Naturelles.

A partir de 1975 la Revue comprend sept séries:

Série A<sub>1</sub> : Mathématiques,

Série A<sub>2</sub> : Physique,

Série A<sub>3</sub> : Astronomie,

Série B : Chimie,

Série C<sub>1</sub> : Géologie,

Série C<sub>2</sub> : Botanique,

Série C<sub>3</sub> : Zoologie.

En principe, la Revue est réservée aux mémoires originaux des membres de la Faculté des Sciences de l'Université d'Ankara. Elle accepte cependant, dans la mesure de la place disponible les communications des auteurs étrangers. Les langues Allemande, Anglaise et Française seront acceptées indifféremment. Tout article doit être accompagné d'un résumé.

Les articles soumis pour publications doivent être remis en trois exemplaires dactylographiés et ne pas dépasser 25 pages des Communications, les dessins et figures portés sur les feuilles séparées devant pouvoir être reproduits sans modifications.

Les auteurs reçoivent 25 extraits sans couverture.

l'Adresse : Dergi Yayın Sekreteri,  
Ankara Üniversitesi,  
Fen Fakültesi,  
Beşevler-Ankara

# The Singular Solutions of the Generalized Poly-Axially Symmetric Helmholtz Equations

I. Ethem ANAR

## SUMMARY

In this article the solution of the generalized poly-axially symmetric Helmholtz equation is obtained on the common intersection of the singular hyperplanes. In the case of spherical region the solution on the common intersection of the singularity hyperplanes is given by a Poisson type integral. We also have obtained solutions of the equations  $h_{\Sigma}^p[u]=0$  and  $h_{\Sigma}^p[u]=f$ .

## I. INTRODUCTION

In this paper we shall investigate the solutions of the generalized poly-axially symmetric Helmholtz equation

$$h_{\Sigma}[U] \equiv \sum_{i=1}^n U_{x_i x_i} + \sum_{j=1}^n \frac{2k_j}{x_j} U_{x_j} + k^2 U = 0. \quad (1.1)$$

The equation (1.1) has the singularity hyperplanes  $x_j = 0$  if  $k_j \neq 0$ . Throughout this paper we assume that  $k_j$  and  $k$  are real positive constants. Let us denote by  $D_0$  the common intersection of the singularity hyperplanes, and consider the portion of the  $n$ -dimensional space, defined as follows

$$Q = \{(x_1, \dots, x_{l-1}, x_1, \dots, x_n) : x_i \geq 0, \dots, x_n \geq 0\}. \quad (1.2)$$

Let  $D$  be a bounded region in  $Q$ , and  $\partial D$  denote the boundary of  $D$ . Let us assume that  $\partial D = D'_0 \cup S$  where  $D'_0 \subset D_0$ . Let  $n$  be the inward normal on  $\partial D$ .

A solution of (1.1) is

$$U = r^{-s} N_s(kr), \quad 2s = n - 2 + 2 \sum_{j=1}^n k_j \quad (1.3)$$

where

$$r = [(x_1 - \xi_1)^2 + \dots + (x_{l-1} - \xi_{l-1})^2 + x_l^2 + \dots + x_n^2]^{\frac{1}{2}}$$

In the representation (1.3)

$P = (x_1, \dots, x_{l-1}, x_l, \dots, x_n) \in D$ ,  $P_0 = (\xi_1, \dots, \xi_{l-1}, 0, \dots, 0) \in D'_0$  and  $N_s$  is a Bessel function of the second-kind of order  $s$ . Let us denote the sphere with radius  $\epsilon$ , centred at a point on  $D'_0$  by  $D\epsilon$  and its boundary by  $\partial D\epsilon$ . Let  $u, v$  be two functions having continuous derivatives of second order in  $D$ . Then the Green's formula for the operator  $h_\Sigma$  can be written as follows,

$$D \int_{j=l}^n x_j^{2k_j} (uh_\Sigma[v] - vh_\Sigma[u]) d\tilde{\mathcal{G}} = - \int_{\partial D} \prod_{j=l}^n x_j^{2k_j} \left( u \frac{\partial v}{\partial n} - v \frac{\partial u}{\partial n} \right) d\sigma. \tag{1.4}$$

Here  $d\tilde{\mathcal{G}}$  and  $d\sigma$  are the volume and surface elements respectively.

### 2. THE SOLUTION OF THE EQUATION $h_\Sigma[u] = f$ ON $D'_0$

By a similar reasoning used in [1] we choose the function  $v$  as a singular solution of (1.3). Let us write the Green's formula (1.4) for the region  $D-D\epsilon$  as follows:

$$D-D\epsilon \int_{j=l}^n x_j^{2k_j} [u(P)h_\Sigma(r^{-s}N_s(kr)) - r^{-s}N_s(kr)h_\Sigma(u(P))] d\tilde{\mathcal{G}} \tag{2.1}$$

$$= - \int_{\partial D + \partial D\epsilon} \prod_{j=l}^n x_j^{2k_j} \left\{ u(P) \frac{\partial}{\partial n} [r^{-s}N_s(kr)] - r^{-s}N_s(kr) \frac{\partial u(P)}{\partial n} \right\} d\sigma.$$

We already know that [3] a Bessel function of the second-kind has the property

$$\frac{d}{dx} [x^{-\nu}N_\nu(kx)] = -k x^{-\nu} N_{\nu+1}(kx) \tag{2.2}$$

and for very small  $x$ 's,

$$N_\nu(x) \sim \begin{cases} -\frac{1}{\pi} 2^\nu \Gamma(\nu) x^{-\nu} ; \nu \neq 0 \\ \frac{2}{\pi} \ln x ; \nu = 0 \end{cases} \quad \nu \neq -1, -2, \dots \tag{2.3}$$

Using (2.2) and (2.3) we obtain

$$\begin{aligned}
 & - \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} u(P) \frac{\partial}{\partial n} [r^{-s} N_s(kr)] d\sigma_\varepsilon \\
 & = - \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} u(P) \frac{\partial}{\partial r} [r^{-s} N_s(kr)] d\sigma_\varepsilon \\
 & = k \int u(P) r^{s+1} N_{s+1}(kr) \prod_{j=1}^{n-1} \sin^{2 \sum_{l=1}^j k_s} \theta_j \prod_{j=1}^n \cos^{2k_j} \theta_j dw
 \end{aligned}$$

If  $\varepsilon$  tends to zero we get

$$- \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} u(P) \frac{\partial}{\partial n} [r^{-s} N_s(kr)] d\sigma_\varepsilon = - u(P_0) w^*, \tag{2.4}$$

where

$$w^* = \frac{k}{\pi} \left( \frac{2}{k} \right)^{s+1} \Gamma(s+1) \int \prod_{j=1}^{n-1} \sin^{2 \sum_{l=1}^j k_s} \theta_j \prod_{j=1}^n \cos^{2k_j} \theta_j dw \tag{2.5}$$

$$d\sigma = r^{n-1} dw, \quad dw = \sin^{n-2} \theta_{n-1} \sin^{n-3} \theta_{n-2} \dots \sin \theta_2 d\theta_{n-1} d\theta_{n-2} \dots d\theta_1 \tag{2.6}$$

The second integral over  $\partial D_\varepsilon$  appearing in the equation (2.1) easily gives

$$\int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} r^{-s} N_s(kr) \frac{\partial u(P)}{\partial n} d\sigma_\varepsilon = 0. \tag{2.7}$$

Let  $u$  be a solution of the equation  $h_\Sigma[u] = f(P)$  and  $f$  be a continuous function defined in  $D$ . Hence the solution of  $h_\Sigma[u] = f(P)$  at the point  $P_0$  can be written as follows:

$$\begin{aligned}
 u(P_0) & = \frac{1}{w^*} \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} \left\{ u(P) \frac{\partial}{\partial n} [r^{-s} N_s(kr)] - r^{-s} N_s(kr) \frac{\partial}{\partial n} u(P) \right\} d\sigma \\
 & + \frac{1}{w^*} \int_D \prod_{j=1}^n x_j^{2k_j} f(P) r^{-s} N_s(kr) d\mathcal{G}. \tag{2.8}
 \end{aligned}$$

Obviously the solution of  $h_\Sigma[u] = 0$  in  $D'_0$  is

$$u(P_0) = \frac{1}{w^*} \int_{\partial D} \prod_{j=l}^n x_j^{2k_j} \left\{ u(P) \frac{\partial}{\partial n} [r^{-s} N_s(kr)] - r^{-s} N_s(kr) \frac{\partial}{\partial n} u(P) \right\} d\sigma \quad (2.9)$$

### 3. A FUNDAMENTAL FUNCTION FOR THE SPHERICAL REGIONS

Let  $D$  be the spherical region with the radius  $R$ , centred at the common intersection of the singularity hyperplanes. Let us define

$$R_1 = (\rho^2 + r_1^2 - 2r_1\rho \cos\theta)^{1/2} \quad R_1^* = (r_1^2 + \frac{R^4}{\rho^2} - 2r_1 \frac{R^2}{\rho} \cos\theta)^{1/2}$$

and

$$R_1 \left| \begin{array}{l} = R_1^* \\ r_1 = R \end{array} \right. \quad R_1^* \left| \begin{array}{l} = R_1^* \\ r_1 = R \end{array} \right.$$

where

$$\rho = \overline{OP}_0, \quad \overline{OP}_0 \cdot \overline{OP}'_0 = R^2, \quad r_1 = \overline{OP}$$

In the spherical region  $D$  the fundamental function  $\gamma(P; P_0)$  of the equation  $h_{\Sigma}[u] = 0$  is

$\gamma(P; P_0) = (R_1 R_1^*)^{-s} N_s(kR_1) N_s(kR_1^*) - (R_1^* R_1)^{-s} N_s(kR_1^*) N_s(kR_1)$  where  $\gamma(P; P_0)$  vanishes on  $\partial D - D_0'$ . If we substitute  $\gamma(P; P_0)$  for the solution (1.3) in the formula (2.1) we obtain

$$\begin{aligned} & - \int_{\partial D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} u(P) \frac{\partial \gamma(P; P_0)}{\partial n} d\sigma_\varepsilon \\ &= - \int_{\partial D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} u(P) \frac{\partial}{\partial R_1} \gamma(P; P_0) d\sigma_\varepsilon \\ &= - \int_{\partial D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} u(P) \frac{\partial}{\partial R_1} \{(R_1 R_1^*)^{-s} N_s(kR_1) N_s(kR_1^*)\} d\sigma_\varepsilon \quad (3.1) \\ &= - w^* (R_1^*)^{-s} N_s(kR_1^*) u(P_0) \end{aligned}$$

as in the previous section. On the other hand we can easily show that

$$\int_{\partial D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} \gamma(P; P_0) \frac{\partial u(P)}{\partial n} d\sigma_\varepsilon = 0. \quad (3.2)$$

Making use of (3.1) and (3.2) in (2.1) we get

$$\int_D \prod_{j=1}^n x_j^{2k_j} \{u(P) h_\Sigma[\gamma(P;P_0)] - \gamma(P;P_0) h_\Sigma[u(P)]\} d\bar{\sigma}$$

$$= - w^* (R_1^*)^{-s} N_s(kR_1^*) u(P_0) \tag{3.3}$$

$$- \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} \left\{ u(P) \frac{\partial}{\partial n} [\gamma(P;P_0)] - \gamma(P;P_0) \frac{\partial u(P)}{\partial n} \right\} d\sigma$$

or

$$u(P_0) = \frac{(R_1^*)^s}{w^* N_s(kR_1^*)} \left\{ \int_D \prod_{j=1}^n x_j^{2k_j} \gamma(P;P_0) h_\Sigma[u(P)] d\bar{\sigma} \right.$$

$$\left. - \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} u(P) \frac{\partial \gamma(P;P_0)}{\partial n} d\sigma \right\}. \tag{3.4}$$

In the case of spherical region (3.4) is the solution of the equation  $h_\Sigma[u] = f$  at the singularity point  $P_0 \in D'$ . If  $h_\Sigma[u] = 0$  then (3.4) reduces to the form

$$u(P_0) = \frac{-(R_1^*)^s}{w^* N_s(kR_1^*)} \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} u(P) \left[ \frac{\partial}{\partial r_1} \gamma(P;P_0) \right] \Big|_{r_1 = R} d\sigma. \tag{3.5}$$

The value of the normal derivative of the fundamental solution on the spherical surface is

$$\frac{\partial \gamma}{\partial r_1} \Big|_{r_1 = R} = k(R^*)^{-(s+1)} (R_1^*)^{-s} \{(\rho - R \cos\theta) N_s(kR^*_1) N_{s+1}(kR_1^*)$$

$$- (R - \rho \cos\theta) N_s(kR_1^*) N_{s+1}(kR^*_1)\}$$

Thus the solution (3.5) becomes

$$u(P_0) = \frac{k(R_1^*)^s}{w^* N_s(kR_1^*)} \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} u(P) (R_1^*)^{-(s+1)} (R_1^*)^{-s} \tag{3.6}$$

$$\times \{ (R - \rho \cos\theta) N_s(kR_1^*) N_{s+1}(kR_1^*) - (\rho - R \cos\theta) N_s(kR_1^*) N_{s+1}(kR_1^*) \} d\sigma$$

#### 4. A SOLUTION OF THE EQUATION $h_{\Sigma}^p[u] = 0$

If we want to obtain a solution of (1.1), which depends only on  $r = [(x_1 - \xi_1)^2 + \dots + (x_{l-1} - \xi_{l-1})^2 + x_l^2 + \dots + x_n^2]^{1/2}$  then the equation becomes

$$\frac{d^2u}{dr^2} + \frac{2s+1}{r} \frac{du}{dr} + k^2u = 0. \quad (4.1)$$

We shall use the method of induction to find a solution of the equation  $h_{\Sigma}^p[u] = 0$ , where  $p$  is a natural number. First let us consider the equation  $h_{\Sigma}^2[u] = 0$ : the system of equation

$$\begin{aligned} h_{\Sigma}[\varphi_1] &= 0 \\ h_{\Sigma}[u] &= \varphi_1 \end{aligned} \quad (4.2)$$

corresponds to the equation

$$h_{\Sigma}^2[u] \equiv h_{\Sigma}(h_{\Sigma}[u]) = 0. \quad (4.3)$$

By (1.3) a solution of the equation  $h_{\Sigma}[\varphi_1] = 0$  is,

$$\varphi_1 = r^{-s} J_s(kr) \quad (4.4)$$

Here  $J_s(kr)$  is the Bessel function of the first-kind. Thus the equation (4.2) takes the form

$$r^2 \frac{d^2u}{dr^2} + (2s+1)r \frac{du}{dr} + k^2 r^2 u = r^2 J_s(kr). \quad (4.5)$$

Now let us look for a solution of (4.5) in the form

$$u = A r^{-s+1} J_{s-1}(kr) \quad (4.6)$$

After the necessary calculations we obtain

$$\{Ak[2s - n - 2 \sum_{j=1}^n k_j] - 1\} r^{-s+2} J_s(kr) = 0.$$

Thus in order to get a solution of the form (4.6) we must choose

$$A = - \frac{1}{2k}.$$

Hence a solution of the equation  $h_{\Sigma}^2[u] = 0$  is



$$u_2 = - \frac{1}{2k} r^{-s+1} J_{s-1}(kr). \quad (4.7)$$

Now, let us suppose that a solution of the equation

$$h_\Sigma h_\Sigma \dots h_\Sigma[u] \equiv h_\Sigma^{p-1}[u] = 0 \quad (4.8)$$

is

$$u_{p-1} = \frac{(-1)^{p-2}}{2^{p-2}(p-2)! k^{p-2}} r^{-s+p-2} J_{s-(p-2)}(kr). \quad (4.9)$$

Since

$$h_\Sigma^p[u] \equiv h_\Sigma^{p-1}(h_\Sigma[u]) = 0. \quad (4.10)$$

we get the system

$$\begin{aligned} h_\Sigma^{p-1}[\varphi_{p-1}] &= 0 \\ h_\Sigma[u] &= \varphi_{p-1}. \end{aligned} \quad (4.11)$$

Substituting (4.9) into (4.11), we get

$$r^2 \frac{d^2u}{dr^2} + (2s+1)r \frac{du}{dr} + k^2 r^2 u = \frac{(-1)^{p-2}}{2^{p-2}(p-2)! k^{p-2}} r^{-s+p} J_{s-(p-2)}(kr) \quad (4.12)$$

If we try to find a solution of (4.12) in the form

$$u = A r^{-s+p-1} J_{s-(p-1)}(kr) \quad (4.13)$$

we obtain

$$A = \frac{(-1)^{p-1}}{(2k)^{p-1} (p-1)!}$$

Hence a solution of the equation  $h_\Sigma^p[u] = 0$  is

$$u_p = \frac{(-1)^{p-1}}{(2k)^{p-1} (p-1)!} r^{-s+p-1} J_{s-(p-1)}(kr). \quad (4.14)$$

## 5. THE SOLUTION OF THE EQUATION $h_\Sigma^p[u] = f$ IN $D'_0$

Following the method used in obtaining (4.14) we can find the solution

$$u = r^{-q} N_q(\mathbf{kr}), \quad q = \mathbf{s} - p + 1 \quad (5.1)$$

of the equation  $h_{\Sigma^p}[u] = 0$ . On the other hand the Green's formula for the operator  $h_{\Sigma^p}$ , [2] is,

$$\begin{aligned} & \int_D \prod_{j=l}^n x_j^{2k_j} (u h_{\Sigma^p}[v] - v h_{\Sigma^p}[u]) d\mathcal{G} \\ &= - \sum_{i=0}^{p-1} \int_{\partial D} \prod_{j=l}^n x_j^{2k_j} \left( h^i_{\Sigma}[u] \frac{\partial h_{\Sigma^p}[v]}{\partial \mathbf{n}} - h^i_{\Sigma}[v] \frac{\partial h_{\Sigma^p}[u]}{\partial \mathbf{n}} \right) d\sigma \quad (5.2) \end{aligned}$$

here  $u, v \in C^{2p}(D)$  and  $u, v \in C^{2p-2}(\partial D)$ . In (5.2) we define  $h_{\Sigma^0}[u] = u$ . We want to use the singular solution (5.1) instead of  $v$ . Let us subtract a neighborhood  $D_\varepsilon$  of the point from the region  $D$ ,  $v$  is singular. Therefore we can apply the formula (5.2) to the functions  $u$  and  $v$  in the region  $D - D_\varepsilon$ . From (5.2) we can write

$$\begin{aligned} & \int_{D-D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} \{ u h_{\Sigma^p}[r^{-q}N_q(\mathbf{kr})] - r^{-q}N_q(\mathbf{kr}) h_{\Sigma^p}[u] \} d\tau \\ & \partial D + \partial D_\varepsilon \int_{j=l}^n x_j^{2k_j} \left\{ u \frac{\partial}{\partial \mathbf{n}} h_{\Sigma^{p-1}}[r^{-q}N_q(\mathbf{kr})] - r^{-q}N_q(\mathbf{kr}) \frac{\partial}{\partial \mathbf{n}} h_{\Sigma^{p-1}}[u] \right\} d\sigma \\ & - \sum_{i=1}^{p-1} \int_{\partial D + \partial D_\varepsilon} \prod_{j=l}^n x_j^{2k_j} \left\{ h_{\Sigma^i}[u] \frac{\partial}{\partial \mathbf{n}} h_{\Sigma^{p-i-1}}[r^{-q}N_q(\mathbf{kr})] - h_{\Sigma^i}[r^{-q}N_q(\mathbf{kr})] \right. \\ & \left. \times \frac{\partial}{\partial \mathbf{n}} h_{\Sigma^{p-i-1}}[u] \right\} d\sigma \quad (5.3) \end{aligned}$$

By induction

$$h_{\Sigma^{p-1}}[r^{-q}N_q(\mathbf{kr})] = (2k)^{p-1} \prod_{m=1}^{p-1} (m-p) r^{-(q+p-1)} N_{p+q-1}(\mathbf{kr}) \quad (5.4)$$

and

$$\frac{\partial}{\partial \mathbf{r}} h_{\Sigma^{p-1}}[r^{-q}N_q(\mathbf{kr})] = - 2^{p-1} k^p \prod_{m=1}^{p-1} (m-p) r^{-(q+p-1)} N_{p+q}(\mathbf{kr}) \quad (5.5)$$

are easily obtained. Thus

$$\begin{aligned}
 & - \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} u(P) \frac{\partial}{\partial r} h_{\Sigma^{p-1}} [r^{-q} N_q(kr)] d\sigma_\varepsilon \\
 & = 2^{p-1} k^p \prod_{m=1}^{p-1} (m-p) \int_C r^{2 \sum_{l=1}^n k_l} u(P) r^{-(p+q-p)} N_{p+q}(kr) \\
 & \quad \times \prod_{j=1}^n \sin^{2 \sum_{l=1}^j k_s} \theta_j \prod_{j=1}^n \cos^{2k_j} \theta_j r^{n-1} dw \\
 & = - W_1 u(P_0),
 \end{aligned} \tag{5.6}$$

where

$$\begin{aligned}
 W_1 & = \frac{1}{\pi} \left( \frac{2}{k} \right)^{q+p} \Gamma(p+q) 2^{p-1} k^p \prod_{m=1}^{p-1} (m-p) \int_C \prod_{j=1}^n \sin^{2 \sum_{l=1}^j k_s} \theta_j \prod_{j=1}^n \cos^{2k_j} \theta_j r^{n-1} dw.
 \end{aligned} \tag{5.7}$$

On the other hand we can easily verify that

$$\int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} r^{-q} N_q(kr) \frac{\partial}{\partial r} h_\Sigma^{p-1} [u] d\sigma_\varepsilon = 0. \tag{5.8}$$

To complete the solution we are left with the calculation of the following:

$$\begin{aligned}
 & - \sum_{i=1}^{p-1} \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} \left\{ h_\Sigma^i [u] \frac{\partial}{\partial r} h_\Sigma^{p-i-1} [r^{-q} N_q(kr)] - h_\Sigma^i [r^{-q} N_q(kr)] \right. \\
 & \quad \left. \times \frac{\partial}{\partial r} h_\Sigma^{p-i-1} [u] \right\} d\sigma_\varepsilon
 \end{aligned} \tag{5.9}$$

From (5.4), we have

$$h_\Sigma^j [r^{-q} N_q(kr)] = (2k)^j \prod_{m=1}^j (m-j-1) r^{-(q+j)} N_{q+j}(kr), \quad 1 \leq j \leq p-1 \tag{5.10}$$

Let us choose the term

$$I_{\mu-1} = \int_{\partial D_\varepsilon} \prod_{j=1}^n x_j^{2k_j} \left\{ h_\Sigma^{\mu-1}[u] \frac{\partial}{\partial r} h_\Sigma^{p-\mu}[r^{-q}N_q(kr)] - h_\Sigma^{\mu-1}[r^{-q}N_q(kr)] \frac{\partial}{\partial r} \right. \\ \left. \times h_\Sigma^{p-\mu}[u] \right\} d\sigma_\varepsilon \quad (5.11)$$

( $\mu = 2, 3, \dots, p$ )

in the sum (5.9). In order to obtain the value of (5.9) it suffices to calculate the integral  $I_{\mu-1}$ . Putting  $j = p - \mu$  in (5.10) we get

$$\frac{\partial}{\partial r} h_\Sigma^{p-\mu}[r^{-q}N_q(kr)] = -k (2k)^{p-\mu} \prod_{m=1}^{p-\mu} (m-p+\mu-1) r^{-(p+q-\mu)} N(kr)_{p+q-\mu+1} \quad (5.12)$$

and from (5.10) we find

$$h_\Sigma^{\mu-1}[r^{-q}N_q(kr)] = (2k)^{\mu-1} \prod_{m=1}^{\mu-1} (m-\mu) r^{-(q+\mu-1)} N(kr)_{q+\mu-1} \quad (5.13)$$

for  $j = \mu-1$ . So (5.11) takes the

$$I_{\mu-1} = \int_{\partial D} r^{2 \sum_{j=1}^n k_j} \prod_{j=1}^{n-1} \sin^{2k_j} \theta_j \prod_{j=1}^n \cos^{2k_j} \theta_j \left\{ -k(2k)^{p-\mu} \prod_{m=1}^{p-\mu} (m-p+\mu-1) r^{-(p+q-\mu)} \right. \\ \left. \times N(kr)_{q+p-\mu+1} h_\Sigma^{\mu-1}[u] - (2k)^{\mu-1} \prod_{m=1}^{\mu-1} (m-\mu) r^{-(q+\mu-1)} N(kr)_{q+\mu-1} \frac{\partial}{\partial r} h_\Sigma^{p-\mu} \right. \\ \left. \times [u] \right\} r^{n-1} dw. \quad (5.14)$$

Taking the limit when  $r \rightarrow 0$  we obtain  $I_{\mu-1} = 0$ . Hence the solution of the equation

$$h_\Sigma^p[u] = f$$

at the singularity point  $P_0 \in D_0'$  is

$$u(P_0) = \frac{1}{W_1} \left\{ \int_D \prod_{j=1}^n x_j^{2k_j} r^{-q} N_q(kr) f(P) d\tilde{\sigma} \right. \\ \left. - \sum_{i=0}^{p-1} \int_{\partial D} \prod_{j=1}^n x_j^{2k_j} (h^i_\Sigma[u(P)]) \frac{\partial}{\partial n} h_\Sigma^{p-i-1}[r^{-q}N_q(kr)] - h^i_\Sigma[r^{-q}N_q(kr)] \frac{\partial}{\partial n} h_\Sigma^{p-i-1}[u(P)] \right\} d\sigma$$

## ÖZET

Bu makalede, genelleştirilmiş çok simetrik Helmholtz denkleminin çözümü, tekil hiperdüzlemlerin arakesiti üzerindeki noktalarda elde edilmiştir. Bölgenin küresel bir bölge olması halinde, bir temel fonksiyon elde ederek tekil hiperdüzlemlerin arakesiti üzerindeki noktalarda çözüm, Poisson integraline benzer biçimde verilmiştir. Ayrıca  $h^p_\Sigma[u]=0$  denkleminin bir çözümü elde edilerek  $h^p_\Sigma[u]=f$  denkleminin tekil çözümü incelenmiştir.

## REFERENCES

1. Çelebi A.O., Some properties for the solutions of  $\sum_{i=1}^n \left( \frac{\partial^2}{\partial x_i^2} + \frac{k_i}{x_i} \frac{\partial}{\partial x_i} \right) u = 0$   
Comm. Fac. Sci. Univ. Ankara Ser. A Tome 19, (1970) pp. 31-39.
2. Çelebi Ü., Boggio's formulas for the operator  $\mathcal{L}$ . (Doctoral thesis 1974, unpublished) Faculty of Science University of Ankara.
3. Watson. G.N.. Theory of Bessel Functions. Cambridge Univ. Press, 1962.