

The Curvature Property of a Linear Dynamical System

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Abstract

In this work a two-dimensional smooth autonomous dynamical system is regarded as a three-dimensional Riemannian manifold and it is shown that the scalar curvature of a linear dynamical system dx/dt = ax + by, dy/dt = cx + dy is non-positive. The manifold is scalar-flat iff b = -c and a = d = 0.

Keywords: Linear dynamical systems, Riemann curvature tensor, scalar curvature

Bir Lineer Dinamik Sistemin Eğrilik Özelliği

Öz

Bu çalışmada iki-boyutlu, düzgün, otonom bir dinamik sistem üç-boyutlu bir Riemann manifoldu olarak değerlendirilmiş ve bir dx/dt = ax + by, dy/dt = cx + dy lineer dinamik sisteminin skaler eğriliğinin pozitif olmadığı gösterilmiştir. Manifold skalerdüzdür ancak ve ancak b = -c ve a = d = 0.

Anahtar Kelimeler: Lineer dinamik sistemler, Riemann eğrilik tensörü, skaler eğrilik.

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1. Introduction

A smooth autonomous dynamical system (SADS) on a twodimensional manifold N = (D; x, y), where D is a connected open set in \mathbb{R}^2 endowed with coordinates (x, y), is given by a system of first order ordinary differential equations

$$\frac{dx}{dt} = f(x, y)$$

$$\frac{dy}{dt} = g(x, y)$$
(1)

such that f and g are smooth functions on N. The system (1) defines a smooth vector field

$$\xi = f(x, y)\partial_x + g(x, y)\partial_y \tag{2}$$

on *N* which is a smooth section of the tangent bundle *TN*, i.e. is a mapping $\xi: N \to TN$ defined to be

$$\xi(x,y) = (x,y,\dot{x} = f(x,y), \dot{y} = g(x,y)).$$

Since the rank of $d\xi$, the differential of ξ , equals 2 on *N*, a SADS may be regarded as a surface in *TN*.

In this work, to capture all the information about the dynamics, we identify a SADS as a three-dimensional manifold $M = I \times N$ with local coordinates (t, x, y). Here t is a coordinate function on an open interval $I \subset \mathbb{R}$ and stands for the time coordinate. This consideration gives rise to define a SADS as a submanifold of the first-order jet bundle $J^1(\mathbb{R}, N)$ of maps $\mathbb{R} \to N$ [1, 2].

This paper is devoted to the investigation of the curvature property of a linear system

$$\frac{dx}{dt} = ax + by$$

$$\frac{dy}{dt} = cx + dy,$$
(3)

where $a, b, c, d \in \mathbb{R}$, within the context of Riemannian geometry. By means of the method of moving frames we evaluate the connection 1-form and the curvature 2-form of the Levi-Civita connection in *TM* compatible with the Riemannian metric which is defined by the sum of squares of the one-forms

$$\begin{aligned} \omega^1 &= \mathrm{d}t \\ \omega^2 &= \mathrm{d}x - (ax+by)\mathrm{d}t \\ \omega^3 &= \mathrm{d}y - (cx+dy)\mathrm{d}t, \end{aligned}$$

and show that the scalar curvature of the connection is nonpositive and the curvature vanishes if and only if b = -c and a = d = 0.

2. Riemannian structure and the curvature

The solutions of (3) are identified with the solutions of the Pfaffian system

$$\omega^2 = 0, \qquad \omega^3 = 0$$

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such that on a solution curve we have $\omega^1 \neq 0$. The column vector $\omega = (\omega^1, \omega^2, \omega^3)^t$, where ^t denotes the transposition, defines a coframe on *M* which is dual to the frame of the vector fields

$$e_1 = \frac{\partial}{\partial t} + f \frac{\partial}{\partial x} + g \frac{\partial}{\partial y}, \qquad e_2 = \frac{\partial}{\partial x}, \qquad e_3 = \frac{\partial}{\partial y}$$

If we introduce the Riemannian metric on M

$$ds^2 = \sum_i \,\omega^i \otimes \omega^i, \tag{4}$$

then *M* becomes a Riemannian manifold and the frame $e = (e_1, e_2, e_3)$ forms an orthonormal frame. Let ∇ be the connection compatible with the Riemannian metric (4). The structure equations for the coframe ω are given by

$$d\omega^{1} = 0$$

$$d\omega^{2} = a\omega^{1} \wedge \omega^{2} + b\omega^{1} \wedge \omega^{3}$$

$$d\omega^{3} = c\omega^{1} \wedge \omega^{2} + d\omega^{1} \wedge \omega^{3}$$

As it is considered in [4], the $\mathfrak{so}(3, \mathbb{R})$ -valued connection form is obtained by solving the system of equations

$$d\omega^i = -\theta^i_j \wedge \omega^j, \qquad \theta^i_j = -\theta^j_i.$$

The unique 1-form

$$\theta = \begin{pmatrix} 0 & \theta_2^1 & \theta_3^1 \\ -\theta_2^1 & 0 & \theta_3^2 \\ -\theta_3^1 & -\theta_3^2 & 0 \end{pmatrix}$$

satisfying the structure equations is obtained by the Cartan's Lemma, where

$$\theta_{2}^{1} = -a\omega^{2} - \frac{1}{2}(b+c)\omega^{3}$$
$$\theta_{3}^{1} = -\frac{1}{2}(b+c)\omega^{2} - d\omega^{3}$$
$$\theta_{3}^{2} = -\frac{1}{2}(b-c)\omega^{1}.$$

The 1-form θ is called $\mathfrak{so}(3, \mathbb{R})$ -valued connection form of the Levi-Civita connection.

The curvature 2-form of the Levi-Civita connection is defined by the anti-symmetric matrix of two-forms:

$$\Omega^i_j = \mathrm{d} heta^i_j + \sum_k \, heta^i_k \wedge heta^k_j, \qquad \Omega^i_j = -\Omega^j_i.$$

In terms of ω^{i} 's, the components of the Riemannian curvature tensor are determined by

$$\Omega_j^i = \sum_{k < l} R_{jkl}^i \omega^k \wedge \omega^l.$$

For the details on a connection in an arbitrary vector bundle we refer to [3]. By a direct calculation we obtain

$$\Omega_{2}^{1} = -\frac{1}{2} \left(2a^{2} + \frac{3}{2}c^{2} + bc - \frac{1}{2}b^{2} \right) \omega^{1} \wedge \omega^{2}$$

-(ab + cd)\omega^{1} \wedge \omega^{3}
$$\Omega_{3}^{1} = -(ab + cd)\omega^{1} \wedge \omega^{2}$$

$$-\frac{1}{2} \left(\frac{3}{2}b^{2} + bc + 2d^{2} - \frac{1}{2}c^{2} \right) \omega^{1} \wedge \omega^{3}$$

$$\Omega_3^2 = \frac{1}{2} \left(\frac{1}{2} (b+c)^2 - 2ad \right) \omega^2 \wedge \omega^3.$$

The independent nonzero components of the Riemann curvature tensor are

$$R_{212}^{1} = -\frac{1}{2} \Big(2a^{2} + c(b+c) - \frac{1}{2}(b^{2} - c^{2}) \Big),$$

$$R_{213}^{1} = -(ab + cd),$$

$$R_{313}^{1} = -\frac{1}{2} \Big(b(b+c) + 2d^{2} + \frac{1}{2}(b^{2} - c^{2}) \Big),$$

$$R_{323}^{2} = \frac{1}{2} \Big(\frac{1}{2}(b+c)^{2} - 2ad \Big).$$

At a point, the sectional curvatures of the two dimensional subspaces spanned by (e_1, e_2) , (e_1, e_3) and (e_2, e_3) are given respectively by R_{212}^1 , R_{313}^1 , and R_{323}^2 . The scalar curvature is defined by the trace $R = \sum_{i,j} R_{jij}^i$ and is found to be $R = 2(R_{212}^1 + R_{313}^1 + R_{323}^2)$. Note that the scalar curvature is an invariant, that is, it does not depend on the choice of an orthonormal frame. It follows that

$$R = -\left(2a^{2} + \frac{3}{2}c^{2} + bc - \frac{1}{2}b^{2} + \frac{3}{2}b^{2} + bc + 2d^{2} - \frac{1}{2}c^{2} - \frac{1}{2}b^{2} - \frac{1}{2}c^{2} - bc + 2ad\right).$$

By arranging the terms we obtain

$$2R = -((b + c)^2 + (2a + d)^2 + 3d^2).$$

As a result we have the following:

Theorem 1 *The Riemannian manifold corresponding to a linear system*

$$\frac{\mathrm{d}x}{\mathrm{d}t} = ax + by$$
$$\frac{\mathrm{d}y}{\mathrm{d}t} = cx + dy,$$

has non-positive scalar curvature. The scalar curvature vanishes if and only if b = -c and a = d = 0.

We say that the Riemannian manifold is flat iff the curvature tensor identically vanishes. Substituting b = -c and a = d = 0 into the components of the Riemannian curvature tensor yields the following:

Corollary 2 The Riemannian manifold corresponding to a linear system

$$\frac{\mathrm{d}x}{\mathrm{d}t} = \lambda y$$
$$\frac{\mathrm{d}y}{\mathrm{d}t} = -\lambda x,$$

where $\lambda \in \mathbb{R}$, is flat.

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References

- Vassiliou, P. J. (2000). Introduction: Geometric approaches to differential equations. In Vassiliou P. J. and Lisle I. G. (Eds.), *Geometric approaches to differential equations* (pp. 1-15). Australian Mathematical Society Lecture Series. 15, Cambridge University Press, Cambridge UK.
- [2]Saunders, D. J. (1989). *The Geometry of Jet Bundles*. Cambridge University Press, Cambridge UK.
- [3] Morita, S. (2001). *Geometry of Differential Forms*. American Mathematical Society, Providence, RI, USA.
- [4]Ok Bayrakdar, Z., Bayrakdar, T. (2019). A geometric description for simple and damped harmonic oscillators, *Turk J Math*, 43: 2540 – 2548.