

RESEARCH ARTICLE

Essential self-adjointness for covariant tri-harmonic operators on manifolds and the separation problem

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Abstract

Consider the tri-harmonic differential expression $L_V^{\nabla} u = (\nabla^+ \nabla)^3 u + V u$, on sections of a hermitian vector bundle over a complete Riemannian manifold (M, g) with metric g, where ∇ is a metric covariant derivative on bundle E over complete Riemannian manifold, ∇^+ is the formal adjoint of ∇ and V is a self adjoint bundle on E. We will give conditions for L_V^{∇} to be essential self-adjoint in $L^2(E)$. Additionally, we provide sufficient conditions for L_V^{∇} to be separated in $L^2(E)$. According to Everitt and Giertz, the differential operator L_V^{∇} is said to be separated in $L^2(E)$ if for all $u \in L^2(E)$ such that $L_V^{\nabla} u \in L^2(E)$, we have $Vu \in L^2(E)$.

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1. Introduction

Assume that (M, g) be a smooth Riemannian manifold without boundary and $\dim M = n$, let M is connected, with metric g, and with Riemannian volume element d_{μ} . Assume that E be a vector bundle over M. The study of essential self-adjointness for differential operators on \mathbb{R}^n has been the central theme of numerous studies, such as [11] and [30]. Gaffney studied essential self-adjointness for differential operators on Riemannian manifolds in [15]. This problem has lead to many works, such as [4, 5, 10, 17, 18, 21, 25]. The study of the separation property for Schrodinger operators on \mathbb{R}^n was studied through Everitt and Giertz, see [15]. The operator $-\Delta + V$ in $L^p(\mathbb{R}^n)$ is separated if the following condition is satisfied for all $u \in L^p(\mathbb{R}^n)$ such that $(\Delta + V) u \in L^p(\mathbb{R}^n)$, we have that $\Delta u \in L^p(\mathbb{R}^n)$ and $Vu \in L^p(\mathbb{R}^n)$. For the separation problem of second and higher order differential operators, see [1,2,6,7,9,13,27,28]. The separation problem of the differential operator $\Delta_M + V$ on $L^2(M)$ where M is a non-compact Riemannian manifold, Δ_M is the scalar laplacian on M and $V \in C^1(M)$, was studied in [23]. Milatovic was studied the separation property for $\Delta_M + V$ in $L^p(M)$ in [24]. The separation problem for

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Schrodinger operators on \mathbb{R}^n goes back to the work of Everitt and Giertz in [14]. Some authors have studied the separation problem for differential operators on Riemannian manifolds, see [2, 24, 25]. The separation property is linked to the self adjointness in $L^2(M)$, see [26]. Separation problem of differential operators has strong links with the essential self-adjointness of the underlying operator. In this article, we will give the conditions for essentially self-adjointness of $\Delta_A^3 + V$ on $C_c^{\infty}(M)$, where Δ_A^3 be the (non-negative) magnetic tri-Laplacian (with a smooth magnetic field A) on a geodescially complete Riemannian manifold M and $0 \leq V \in C^2(M)$. Additionally, we provide sufficient conditions for L_V^{∇} to be separated in $L^2(E)$.

2. General notations

In this article we consider the differential operator $(\nabla^+ \nabla)^3 u + V u$, where ∇ is a metric covariant derivative on a hermitian bundle E over a Riemannian manifold M, ∇^+ its formal adjoint and V is a linear self-adjoint bundle map over E. In general, the symbols $C^{\infty}(E)$ and $C_{c}^{\infty}(E)$ and $\Omega^{1}(M)$ denote sections of E and compactly supported sections of E, and complex-valued smooth 1-forms on M respectively. We call $L^{p}(E), 1 \leq p < \infty$, indicates the space of *p*-integrable sections of *E* with the norm $||u||_p := (\int_M |u(x)|^p d\mu)^{1/p}$. In the special case p = 2, we have the Hilbert space $L^{2}(E)$ and we use (.,.) to denote the corresponding inner product in $L^{2}(M)$ and the pairing (linear in the first and anti-linear in the second slot) between $L^p(M)$ and $L^q(M)$ with 1/p+1/q = 1. For local Sobolev spaces of sections we use the notation $W_{loc}^{k,p}(E)$, with p and k indicating the corresponding L^p spaces and the highest order of derivatives, respectively. For k = 0 we use $L_{loc}^p(E)$. In the case $E = M \times C$, we denote the corresponding function spaces by $C^{\infty}(M)$, $C_c^{\infty}(M)$, $L^p(M)$, $L^{p}_{loc}(M)$ and $W^{k,p}_{loc}(M)$. In this paper, $\nabla : C^{\infty}(E) \to C^{\infty}(T^{*}M \otimes E)$ stands for a smooth metric covariant derivative on E, and $\nabla^{+} : C^{\infty}(T^{*}M \otimes E) \to C^{\infty}(E)$ indicates the formal adjoint of ∇ with respect to (.,.). The covariant derivative ∇ on E induces the covariant derivative ∇^{End} on the bundle of endomorphisms End E, making $\nabla^{End}V$ a section of the bundle $T^*M \otimes (End \ E)$. Working in the space $L^2(E)$ only, we find it convenient to indicate by (.,.) and $\|.\|$ the inner product and the norm in the spaces $L^{2}(E)$ and $L^{2}(T^{*}M \otimes E)$. We study the separation property on $L^{2}(E)$ that can be seen as an extension of the work mentioned in [23]. We define a set $D_{2}^{\nabla} := \{u \in L^{2}(E) : L_{V}^{\nabla}u \in L^{2}(E)\}$, it is not true that for all $u \in D_2^{\nabla}$, we have $(\nabla^+ \nabla)^3 u \in L^2(E)$ and $Vu \in L^2(E)$ simultaneously, using the terminology of Everitt and Giertz [15] we will say that $L_V^{\nabla} = (\nabla^+ \nabla)^3 + V$ is separated in $L^{2}(E)$ when the following statement holds true for all $u \in D_{2}^{\nabla}$, we have $Vu \in L^{2}(E)$, to make the notations less cumbersome, the symbols (.,.) and $\| \cdot \|_p$ will also be used when referring to $L^p(\Lambda^1T^*M)$, the space of p-integrable 1-forms on M, we only consider the space $L^{2}(M)$, we will use $\| \cdot \|$ instead of $\| \cdot \|_{2}$ to indicate the norm. In a magnetic field $A \in \Omega^1(M)$, where A be real-valued form, the operator $d_A : C^{\infty}(M) \to \Omega^1(M)$ stands for the magnetic differential where $d_A u = du + iu A$ where $d : C^{\infty}(M) \to \Omega^1(M)$ be the standard differential and $i = \sqrt{-1}$. We denote the formal adjoint of d_A with respect to (\cdot, \cdot) by d_A^+ , the (non-negative) magnetic Laplacian on M by $\Delta_A := d_A^+ d_A$, and the magnetic tri-Laplacian by $\Delta_A^3 := \left(d_A^+ d_A\right)^3$. We start recalling some abstract terminology concerning m-accretive operators on Banach spaces. A linear operator S on a Banach space \varkappa is called accretive if $\|(\xi + S)u\|_{\varkappa} \ge \xi \|u\|_{\varkappa}$, for all $\xi > 0$ and all $u \in Dom(S)$. In [12], a densely defined accretive operator S is close and its closure S^{\sim} is also accretive. An operator S on \varkappa is called m-accretive if it is accretive and $\xi + S$ is surjective for all $\xi > 0$. An operator S on \varkappa is named essentially m-accretive if it is accretive and S^{\sim} is m-accretive. We use the relation between m-accretivity and self-adjointness of operators on Hilbert spaces which stated in the paper [22] that the operator S is a self-adjoint and non-negative operator if and only if S is symmetric, closed and m-accretive. We mentioned

some results on the essential m-accretivity of operators in $L^2(E)$ used in this paper with L_V^{∇} with $0 \leq V \in L_{loc}^{\infty}(End \ E)$ and $0 \leq v \in L_{loc}^{\infty}(M)$, we define an operator $H_{2,V}^{\nabla}$ as $H_{2,V}^{\nabla}u := L_V^{\nabla}u$ with the domain $D_2^{\nabla} := \{u \in L^2(E) : L_V^{\nabla}u \in L^2(E)\}$ and an operator $H_{2,v}^d$ as $H_{2,v}^du := S_v^du$ for all $u \in D_2^d$, where $D_2^d := \{u \in L^2(M) : S_v^du \in L^2(M)\}$. Atia was studied the separation problem of bi-harmonic differential operators on Riemannian manifolds in [3,4].

3. Essential self-adjointness result for a perturbation of Δ_A^3

Let p(t, x, y) be the heat kernel of M. For a Borel function $f: M \to R$, define

$$F(t) := \sup_{x \in M} \int_0^t \int_M p(s, x, y) |f(y)| d\mu(y) ds,$$

where $f \in D(M)$ if there exists t > 0 such that F(t) < 1. We say that f belongs to the Kato class K(M) of M if $F(t) \to 0$ as $t \to 0+$. So $K(M) \subset D(M)$, see Theorem 7.13 in [16].

Now, we remind our main result, in this section.

Theorem 3.1. Let M is a geodesically complete connected Riemannian manifold. Let $U = U_1 + U_2$ with $0 \le U_1 \in L^2_{loc}(M)$ and $0 \le U_2 \in L^2_{loc}(M) \cap D(M)$. Furthermore, let $W \in W^{3,2}_{loc}(M)$ and $W \ge 0$. Also there exist constants $c_1 \ge 0$ and $c_2 \ge 0$ such that

$$dW(x)|^{2} \le c_{1} + c_{2}W(x), \qquad (3.1)$$

for all almost $x \in M$. Then $\Delta_A^3 + U - W$ is essentially self-adjoint on $C_c^{\infty}(M)$.

Now, we explain some notations used in subsequent results. We define operators $T_{\Delta_A}^{(p)}$ and $T_{\Delta_A^3}^{(p)}$ in $L^p(M)$ where 1 by the formulas

$$T_{\Delta_A}^{(p)}u := \Delta_A u, \ u \in Dom\left(T_{\Delta_A}^{(p)}\right) := \left\{u \in L^p(M) : \Delta_A u \in L^p(M)\right\},\tag{3.2}$$

and

$$T_{\Delta_{A}^{3}}^{(p)}u := \Delta_{A}^{3}u, \ u \in Dom\left(T_{\Delta_{A}^{3}}^{(p)}\right) := \left\{u \in L^{p}(M) : \Delta_{A}^{3}u \in L^{p}(M)\right\}.$$
(3.3)

We now state the result of Okazawa, see [29].

Lemma 3.2. Let S and G be nonnegative symmetric operators in a Hilbert space H with inner product $(.,.)_H$ and norm $\| \cdot \|_H$. We assume D be a linear subspace of H on which S+G is essentially self-adjoint. Assume that the following inequalities hold for all $u \in D$:

$$||Su||_{H} + ||Gu||_{H} \le a_{1} ||u||_{H} + a_{2} ||(S+G)u||_{H}$$
(3.4)

and

 $|\mathrm{Im}\,(Gu, Su)_H| \le \widetilde{a_1} \, \|u\|_H^2 + \widetilde{a_2} \, \|(S+G)\,u\|_H \, \|u\|_H \,, \tag{3.5}$

where $a_1, a_2, \widetilde{a_1}$ and $\widetilde{a_2} \ge 0$ are constants. Then S - G is essentially self-adjoint on D.

The following lemma is a direct consequence of Lemma VI.7 in [20] and Corollary 2.5 in [19]:

Lemma 3.3. Let M be a geodesically complete connected Riemannian manifold with Riemannian volume element $d\mu$. Let $0 \ge U_2 \in L^2_{loc}(M) \cap D(M)$. Then, there exist constants $0 \le \delta < 1$ and $\xi \ge 0$ such that

$$\int_{M} |U_2| |u|^2 d\mu \le \delta \int_{M} |d_A u|^2 d\mu + \xi ||u||^2, \qquad (3.6)$$

for all $u \in W^{3,2}_A(M)$.

We will introduce the following lemma which will be used in the proof of the main theorem in this section,

Lemma 3.4. Let M be a geodesically complete connected Riemannian manifold with Riemannian volume element $d\mu$. Let $U = U_1 + U_2$ with $0 \le U_1 \in L^2_{loc}(M)$ and $0 \ge U_2 \in L^2_{loc}(M) \cap D(M)$. Additionally, let $0 \le W \in W^{3,2}_{loc}(M)$ is a function satisfying (3.1). Then we have

$$2\operatorname{Re}\left((T_{W}+1)u, \left(T_{\Delta_{A}^{3}}+T_{U}+\xi\right)u\right) \geq -c_{3}\left\|d_{A}^{+}d_{A}u\right\|^{2},$$
(3.7)

for all $u \in C_c^{\infty}(M)$, where ξ is as in (3.6) and $c_3 := (c_1 + c_2)/2$.

Proof. We will use the integration by parts and the product rule

$$d_A(fv) = f d_A v + (df) v,$$

where f and v are functions on M, also we define

$$\begin{split} W_A^{1,3}(M) &:= \left\{ u \in L^2(M) : d_A u \in L^2\left(\Lambda^1 T^* M\right) \right\} \\ \mathrm{Re}(T_{\Delta_A^3} u, (T_W + 1)u) &= \mathrm{Re}\left(\Delta_A^3 u, (W + 1)u\right) \\ &= \mathrm{Re}\left(\left(d_A^+ d_A\right)^2 u, (W + 1)u\right) \\ &= \mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}\sqrt{W + 1}u\right)\right) \\ &= \mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, \sqrt{W + 1}d_A\left(\sqrt{W + 1}u\right)\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}d_A\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}d_A\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}d_A\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}d_A\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \frac{1}{2}\,\mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, d_W u\right) \\ &= \mathrm{Re}\left(\frac{d_A\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u\right)}{-d\left(\sqrt{W + 1}\right)\left(d_A^+ d_A\right)^2 u\right)}, d_A\left(\sqrt{W + 1}u\right) \right) \\ &+ \frac{1}{2}\,\mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{Re}\left(d_A\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u\right), d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \frac{1}{2}\,\mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \frac{1}{2}\,\mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \frac{1}{2}\,\mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{Re}\left(\sqrt{W + 1}\left(d_A^+ d_A\right)^2 u, d_A\left(\sqrt{W + 1}u\right)\right) \\ &+ \mathrm{Re}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\ &= \mathrm{RE}\left(d_A\left(d_A^+ d_A\right)^2 u, dW u\right) \\$$

$$\begin{split} &= \sqrt{W+1} \operatorname{Re} \left(u, \left(d_{A}^{+} d_{A} \right)^{3} \left(\sqrt{W+1} u \right) \right) \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, d_{A} \left(\sqrt{W+1} u \right) \right) \\ &+ \frac{1}{2} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \operatorname{Re} \left(\left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, \sqrt{W+1} d_{A} u + d\sqrt{W+1} u \right) \right) \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, \sqrt{W+1} d_{A} u + d\sqrt{W+1} u \right) \\ &+ \frac{1}{2} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, \sqrt{W+1} d_{A} u + d\sqrt{W+1} u \right) \\ &+ \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, \left(W+1 \right)^{1/2} d_{A} u \right) \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, d\sqrt{W+1} u \right) \\ &+ \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u \left(W+1 \right)^{-1/2}, \frac{dW}{2\sqrt{W+1}} u \right) \\ &+ \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{2} \operatorname{Re} \left(dW \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &+ \frac{1}{2} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{4} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{4} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{4} \int_{M} \frac{\left| d_{A}^{+} d_{A} u \right|^{2}}{W+1} | dW|^{2} d\mu \\ &- \frac{1}{2} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{4} \int_{M} \frac{\left| d_{A}^{+} d_{A} u \right|^{2}}{W+1} | dW|^{2} d\mu \\ &- \frac{1}{2} \operatorname{Re} \left(d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u, dW u \right) \\ &= \left\| \left(d_{A}^{+} d_{A} \right)^{3/2} \left(\sqrt{W+1} u \right) \right\|^{2} \\ &- \frac{1}{4} \int$$

We use our assumptions on W, we get $\left(\sqrt{W+1}u\right) \in W_A^{3,2}(M)$, we combine the last equality, (3.1) and (3.6) we get

$$\begin{aligned} &\operatorname{Re}\left((T_{W}+1)u, \left(T_{\Delta_{A}^{3}}+T_{U}\right)u\right) \\ &= \operatorname{Re}\left((T_{W}+1)u, T_{\Delta_{A}^{3}}u\right) + \operatorname{Re}\left((T_{W}+1)u, T_{U}u\right) \\ &= \operatorname{Re}\left(T_{\Delta_{A}^{3}}u, (T_{W}+1)u\right) + \operatorname{Re}\int_{M}U_{1}\left|u\sqrt{W+1}\right|^{2}d\mu \\ &+ \operatorname{Re}\int_{M}U_{2}\left|u\sqrt{W+1}\right|^{2}d\mu \\ &\geq \left\|\left(d_{A}^{+}d_{A}\right)^{3/2}\left(\sqrt{W+1}u\right)\right\|^{2} - \frac{1}{4}\int_{M}\frac{\left|d_{A}^{+}d_{A}u\right|^{2}}{W+1}\left|dW\right|^{2}d\mu \\ &- \delta\left\|\left(d_{A}^{+}d_{A}\right)^{3/2}\left(\sqrt{W+1}u\right)\right\|^{2} - \xi\left\|\sqrt{W+1}u\right\|^{2} \\ &= (1-\delta)\left\|\left(d_{A}^{+}d_{A}\right)^{3/2}\left(\sqrt{W+1}u\right)\right\|^{2} - \frac{1}{4}\int_{M}\frac{\left|d_{A}^{+}d_{A}u\right|^{2}}{W+1}\left|dW\right|^{2}d\mu \\ &- \xi\left\|\sqrt{W+1}u\right\|^{2} \\ &\geq -\xi\left\|\sqrt{W+1}u\right\|^{2} - \frac{c_{1}+c_{2}}{4}\left\|d_{A}^{+}d_{A}u\right\|^{2}. \end{aligned}$$

So we proved the lemma.

3.1. Proof of theorem 3.1

We assume that the hypotheses of Theorem 3.1 are satisfied, we assume that M is a geodesically complete connected Riemannian manifold with Riemannian volume element d_{μ} . Let $U = U_1 + U_2$ with $0 \leq U_1 \in L^2_{loc}(M)$ and $0 \geq U_2 \in L^2_{loc}(M) \cap D(M)$, let $0 \leq W \in W^{3,2}_{loc}(M)$ is a function satisfying (3.1). We get $(T_{\Delta^3_A} + T_U + \xi)_{C_c^{\infty}(M)}$ is a non-negative symmetric operator. By the assumptions on W, it follows that $(T_W + 1)_{C_c^{\infty}(M)}$ is a non-negative symmetric operator. Since $0 \leq (U_1 + W) \in L^2_{loc}(M)$ also $U_2 \in L^2_{loc}(M) \cap D(M)$ by using Theorem X.1 in [21] to conclude that the operator $(T_{\Delta^3_A} + T_U + T_W + 1 + \xi)$ is essentially self-adjoint on $C_c^{\infty}(M)$. By using (3.7) we get

$$\begin{split} & \left\| (T_{\Delta_A^3} + T_U + T_W + 1 + \xi) u \right\|^2 + c_3 \|u\|^2 \\ &= \left((T_{\Delta_A^3} + T_U + T_W + 1 + \xi) u, (T_{\Delta_A^3} + T_U + T_W + 1 + \xi) u \right) \\ &+ c_3 \|u\|^2 \\ &= \left\| \left(T_{\Delta_A^3} + T_U + \xi \right) u \right\|^2 + 2 \operatorname{Re} \left((T_W + 1) u, \left(T_{\Delta_A^3} + T_U + \xi \right) u \right) \\ &+ \| (T_W + 1) u \|^2 + c_3 \|u\|^2 \\ &\geq \left\| \left(T_{\Delta_A^3} + T_U + \xi \right) u \right\|^2 + \| (T_W + 1) u \|^2 \,, \end{split}$$

for all $u \in C_c^{\infty}(M)$. Applying hypothesis (3.4) of Lemma 3.2 is satisfied

with $S = T_{\Delta_A^3} + T_U + \xi$, $G = T_W + 1$ and $D = C_c^{\infty}(M)$, we will use the integration by parts and the product rule, for all $u \in C_c^{\infty}(M)$ we have

$$\operatorname{Im}\left(T_{W}u, T_{\Delta_{A}^{3}}u\right) = \operatorname{Im}\left(Wu, \left(d_{A}^{+}d_{A}\right)^{3}u\right)$$
$$= \operatorname{Im}\left(d_{A}(Wu), d_{A}\left(d_{A}^{+}d_{A}\right)^{2}u\right)$$
$$= \operatorname{Im}\left(Wd_{A}u + dWu, d_{A}\left(d_{A}^{+}d_{A}\right)^{2}u\right)$$
$$= \operatorname{Im}\left(Wd_{A}u, d_{A}\left(d_{A}^{+}d_{A}\right)^{2}u\right) + \operatorname{Im}\left(dWu, d_{A}\left(d_{A}^{+}d_{A}\right)^{2}u\right)$$
$$= \operatorname{Im}\left(dWu, d_{A}\left(d_{A}^{+}d_{A}\right)^{2}u\right).$$
(3.8)

From (3.1) and (3.8) we obtain

$$\begin{aligned} \left| \operatorname{Im} \left(T_{W} u, T_{\Delta_{A}^{3}} u \right) \right| &= \left| \operatorname{Im} \left(dW u, d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u \right) \right| \\ &\leq \frac{1}{2} \int_{M} |dW|^{2} |u|^{2} d\mu + \frac{1}{2} \int_{M} \left| d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u \right|^{2} d\mu \\ &\leq \frac{c_{1}}{2} ||u||^{2} + \frac{c_{2}}{2} \int_{M} W |u|^{2} d\mu + \frac{1}{2} \left\| d_{A} \left(d_{A}^{+} d_{A} \right)^{2} u \right\|^{2} \\ &\leq \frac{c_{1}}{2} ||u||^{2} + \frac{c_{2}}{2} (u, T_{W} u) + \frac{1}{2} \left(u, T_{\Delta_{A}^{5}} u \right) \\ &\leq \frac{c_{1}}{2} ||u||^{2} + \frac{c_{2+1}}{2} \left(u, \left(T_{\Delta_{A}^{5}} + T_{W} \right) u \right) \end{aligned}$$

For all $u \in C_{c}^{\infty}(M)$ we get

$$\left| \operatorname{Im} \left(T_{W} u, \left(T_{\Delta_{A}^{3}} + T_{U} + \xi \right) u \right) \right| = \left| \operatorname{Im} \left(T_{W} u, T_{\Delta_{A}^{3}} u \right) \right| \\ \leq \frac{c_{1}}{2} \| u \|^{2} + \frac{c_{2+1}}{2} \left(u, \left(T_{\Delta_{A}^{5}} + T_{W} \right) u \right).$$
(3.9)

From (3.6) we obtain

$$(u, T_U u) = (u, (U_1 + U_2) u) \ge (u, U_2 u)$$

 $\ge -\delta (u, \Delta_A u) - \xi ||u||^2,$

for all $u \in C_c^{\infty}(M)$ we get

$$\begin{pmatrix} u, (T_{\Delta_{A}^{3}} + T_{U}) u \end{pmatrix} + \xi ||u||^{2} = (u, T_{\Delta_{A}^{3}} u) + (u, T_{U}u) + \xi ||u||^{2} = (u, T_{\Delta_{A}^{3}} u) + \xi ||u||^{2} + (u, Uu) \geq (u, T_{\Delta_{A}^{3}} u) + \xi ||u||^{2} - \delta \int |d_{A}u|^{2} d\mu - \xi ||u||^{2} = (u, T_{\Delta_{A}^{3}} u) - \delta (u, (d_{A}^{+}d_{A})^{3} u) = (1 - \delta) (u, T_{\Delta_{A}^{3}} u),$$
(3.10)

for all $u \in C_{c}^{\infty}(M)$ from (3.9), (3.10) and since

$$\operatorname{Im}\left(u,\left(T_{\Delta_{A}^{3}}+T_{U}+\xi\right)u\right)=0.$$

We get

$$\begin{split} \left| \operatorname{Im} \left((T_W + 1) \, u, \left(T_{\Delta_A^3} + T_U + \xi \right) u \right) \right| &= \left| \operatorname{Im} \left(T_W u, \left(T_{\Delta_A^3} + T_U + \xi \right) u \right) \right| \\ &\leq \frac{c_1}{2} \, \| u \|^2 + \frac{c_{2+1}}{2} \left(u, \left(T_{\Delta_A^5} + T_W \right) u \right) \\ &= \frac{c_1}{2} \, \| u \|^2 + \frac{c_{2+1}}{2} \left(u, T_W u \right) + \frac{c_{2+1}}{2} \left(u, T_{\Delta_A^5} u \right) \\ &\leq \frac{c_1}{2} \, \| u \|^2 + \frac{c_{2+1}}{2} \left(u, T_W u \right) \\ &+ \frac{c_{2+1}}{2 \left(1 - \delta \right)} \left(u, \left(T_{\Delta_A^5} + T_U \right) u \right) \\ &+ \left(\frac{c_{2+1}}{2} \right) \left(\frac{\xi}{1 - \delta} \right) \| u \|^2 \\ &\leq \left(\frac{c_1}{2} + \left(\frac{c_{2+1}}{2} \right) \left(\frac{\xi}{1 - \delta} \right) \right) \| u \|^2 \\ &+ \frac{c_{2+1}}{2 \left(1 - \delta \right)} \left(u, \left(T_{\Delta_A^5} + T_U + T_W + 1 + \xi \right) u \right). \end{split}$$

Hence the assumptions of (3.5) of Lemma 3.2 is satisfied with $S = T_{\Delta_A^3} + T_U + \xi$ and $G = T_W + 1$ and $D = C_c^{\infty}(M)$, Thus by Lemma 3.2 it follows that $S - G = T_{\Delta_A^3} + T_U + \xi - T_W - 1$ is essentially self-adjoint on $C_c^{\infty}(M)$, since $\xi - 1$ is a constant, then $\Delta_A^3 + U - W$ is essentially self-adjoint on $C_c^{\infty}(M)$.

4. The separation problem result

We will introduce our main theorem in this section.

Theorem 4.1. Let (M, g) be a complete connected Riemannian manifold without boundary, let E be a vector bundle over M with a metric covariant derivative ∇ . We assume $V \in C^1$ (End E), $V(x) \ge 0$, for all $x \in M$, where the inequality is understood in the sense of linear operators $E_x \to E_x$ and

$$\left| \left(\nabla^{End} V(x) \right) \right| \le \sigma \left(V(x) \right)^{3/2}, \ 0 \le \sigma < 1.$$

$$(4.1)$$

Then

$$\left\| \left(\nabla^+ \nabla \right)^3 u \right\| + \| V u \| \le C \left[\left\| L_V^{\nabla} u \right\| + \| u \| \right], \tag{4.2}$$

for all $u \in D_2^{\nabla}$, where $C \ge 0$ be a constant, that is L_V^{∇} is separated in $L^2(E)$.

Lemma 4.2. Under the hypothesis of the Theorem 3.1, then the following inequalities are valid for all $u \in C_c^{\infty}(E)$,

$$\left\| \left(\nabla^+ \nabla \right)^3 u \right\| + \| V u \| \le C_1 \left\| L_V^{\nabla} u \right\|, \tag{4.3}$$

and

$$\left\| V^{1/2} \left(\nabla^+ \nabla \right)^{3/2} u \right\| \le C_1 \left\| L_V^{\nabla} u \right\|, \tag{4.4}$$

where $V^{1/2}$ is the square root of the operator $V(x) : E_x \to E_x$ and C_1 is a constant depending on $n = \dim M$, $m = \dim E_x$ and σ .

Proof. By the definition of L_{V}^{∇} , for all $\delta > 0$ and all $u \in C_{c}^{\infty}(E)$ we obtain

$$\begin{split} \left\| L_{V}^{\nabla} u \right\|^{2} &= \left\| \left(\nabla^{+} \nabla \right)^{3} u + V u \right\|^{2} \\ &= \left(\left(\nabla^{+} \nabla \right)^{3} u + V u, \left(\nabla^{+} \nabla \right)^{3} u + V u \right) \\ &= \left\| V u \right\|^{2} + \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} + 2 \operatorname{Re} \left(\left(\nabla^{+} \nabla \right)^{3} u, V u \right) \\ &= \left\| V u \right\|^{2} + \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} + 2 \operatorname{Re} \left(\left(\nabla^{+} \nabla \right)^{3} u, V u \right) + \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} \\ &- \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} \\ &= \left\| V u \right\|^{2} + \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} + (1 - \delta) \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} \\ &+ 2 \operatorname{Re} \left(\left(\nabla^{+} \nabla \right)^{3} u, V u \right) \\ &= \left\| V u \right\|^{2} + \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} + (1 - \delta) \operatorname{Re} \left(\nabla^{+} \nabla \right)^{3} u, L_{V}^{\nabla} u - V u \right) \\ &+ 2 \operatorname{Re} \left(\nabla^{+} \nabla \right)^{3} u, V u \right) \\ &= \left\| V u \right\|^{2} + \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} + (1 - \delta) \operatorname{Re} \left(\nabla^{+} \nabla \right)^{3} u, L_{V}^{\nabla} u \right) \\ &+ (1 + \delta) \operatorname{Re} \left(\left(\nabla^{+} \nabla \right)^{3} u, V u \right). \end{split}$$

By the product rule $\nabla (Vu) = (\nabla^{End}V) u + V\nabla u$, so for all $u \in C_c^{\infty}(E)$, we have

$$\operatorname{Re}(\left(\nabla^{+}\nabla\right)^{3}u, Vu) = \operatorname{Re}(\left(\nabla^{+}\nabla\right)^{2}u, \left(\nabla^{+}\nabla\right)(Vu))$$

$$= \operatorname{Re}(\left(\nabla^{+}\nabla\right)^{2}u, \nabla^{+}(\nabla^{End}V)u + \nabla^{+}V\nabla u)$$

$$= \operatorname{Re}(\left(\nabla^{+}\nabla\right)^{2}u, \nabla^{+}(\nabla^{End}V)u) + \operatorname{Re}(\left(\nabla^{+}\nabla\right)^{2}u, \nabla^{+}V\nabla u)$$

$$= \operatorname{Re}(Z) + W, \qquad (4.5)$$

where $Z := ((\nabla^+ \nabla)^2 u, \nabla^+ (\nabla^{End} V)u)$ and $W := (V^{1/2} (\nabla^+ \nabla)^{3/2} u, V^{1/2} (\nabla^+ \nabla)^{3/2} u)$, then, we obtain

$$(1+\delta)\operatorname{Re}((\nabla^{+}\nabla)^{3}u, Vu) = (1+\delta)\operatorname{Re}Z + (1+\delta)W \ge -(1+\delta)|Z| + (1+\delta)W.$$
(4.6)

By Cauchy-Schwartz $2ab \leq ka^2 + k^{-1}b^2$, where k, a and b are positive real numbers and the condition (4.1) we obtain

$$|Z| \leq (\sigma+1) \int_{M} \left| \left(\nabla^{+} \nabla \right)^{3/2} V^{1/2} u \right|_{(T^{*}M \otimes E)_{x}} |Vu|_{E_{x}} d\mu,$$

$$|Z| \leq \frac{\delta \alpha}{2} \left\| V^{1/2} \left(\nabla^{+} \nabla \right)^{3/2} u \right\|^{2} + \frac{(\sigma+1)^{2}}{2\delta \alpha} \|Vu\|^{2}, \qquad (4.7)$$

for all $\alpha > 0$, we use Cauchy-Schwartz again, we obtain

$$\left|\operatorname{Re}\left(\left(\nabla^{+}\nabla\right)^{3}u, L_{V}^{\nabla}u\right)\right| \leq \left|\left(\left(\nabla^{+}\nabla\right)^{3}u, L_{V}^{\nabla}u\right)\right| \leq \frac{\gamma}{2} \left\|\left(\nabla^{+}\nabla\right)^{3}u\right\|^{2} + \frac{1}{2\gamma} \left\|L_{V}^{\nabla}u\right\|^{2}, \quad (4.8)$$

for all $\gamma > 0$, we obtain

$$\begin{split} \left\| L_{V}^{\nabla} u \right\|^{2} &\geq \| V u \|^{2} + \delta \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} \\ &- \frac{\left(1 + \delta \right) \delta \alpha}{2} \left\| V^{1/2} \left(\nabla^{+} \nabla \right)^{3/2} u \right\|^{2} - \frac{\left(1 + \delta \right) \left(\sigma + 1 \right)^{2}}{2 \delta \alpha} \| V u \|^{2} \\ &+ \left(1 + \delta \right) \left\| V^{1/2} \left(\nabla^{+} \nabla \right)^{3/2} u \right\|^{2} - \frac{\left| 1 - \delta \right| \gamma}{2} \left\| \left(\nabla^{+} \nabla \right)^{3} u \right\|^{2} \\ &- \frac{\left| 1 - \delta \right|}{2 \gamma} \left\| L_{v}^{\nabla} u \right\|^{2}, \end{split}$$

from this, we obtain

$$\left(1 + \frac{|1-\delta|}{2\gamma}\right) \left\|L_V^{\nabla} u\right\|^2 \ge \left(1 - \frac{(1+\delta)(\sigma+1)^2}{2\delta\alpha}\right) \|V u\|^2 + \left(\delta - \frac{|1-\delta|\gamma}{2}\right) \left\|\left(\nabla^+ \nabla\right)^3 u\right\|^2 + \left((1+\delta) - \frac{(1+\delta)\delta\alpha}{2}\right) \left\|V^{1/2} \left(\nabla^+ \nabla\right)^{3/2} u\right\|^2.$$

Now the inequalities (4.3) and (4.4) holds if

$$|1-\delta| < \frac{2\delta}{\gamma}, \delta\alpha < 2 \text{ and } (1+\delta) (\sigma+1)^2 < 4.$$
(4.9)

Since, from $0 \le \sigma < 1$, there exist $\delta > 0$, $\gamma > 0$ and $\alpha > 0$ such that the inequalities (4.10) hold.

4.1. Proof of theorem 4.1

As M be a geodesically complete manifold it is known that $\left(L_V^d|_{C_c^{\infty}(M)}\right)^{\sim}$ in $L^2(M)$, be m-accretive and it coincides with $H_{2.V}^d$. Also, from the assumption on M, the operator $\left(L_V^{\nabla}|_{C_c^{\infty}(E)}\right)^{\sim}$ in $L^2(E)$, is m-accretive and it coincides with $H_{2.V}^{\nabla}$. Both of these statements are proven in [31]. From the strategy of Milatovic employs in [25], then the operator $L_V^{\nabla}|_{C_c^{\infty}(E)}$ is essentially self-adjoint and $\left(L_V^{\nabla}|_{C_c^{\infty}(E)}\right)^{\sim} = H_{2.V}^{\nabla}$. We prove (4.3) and (4.4) for all $u \in D_2^{\nabla} = Dom\left(H_{2.V}^{\nabla}\right)$, from which (4.2) follows directly. Since $H_{2.V}^{\nabla}$ is a closed operator, there exists a sequence $\{u_k\}$ in $C_c^{\infty}(E)$ such that $u_k \to u$ and $L_V^{\nabla}u_k \to$ $H_{2.V}^{\nabla}u$ in $L^2(E)$, by the previous lemma the sequence $\{u_k\}$ satisfies (4.3) and (4.4), hence $\{(\nabla^+\nabla)^3 u_k\}, \{Vu_k\}$ and $\{V^{1/2}(\nabla^+\nabla)^{3/2} u_k\}$ are Cauchy sequences in the space $L^2(E)$. Furthermore, $\{\nabla u_k\}$ is a Cauchy sequence in $L^2(T^*M \otimes E)$ as

$$\|\nabla u_k\|^2 = (\nabla u_k, \nabla u_k) = \left(\nabla^+ \nabla u_k, u_k\right) \le \left\|\nabla^+ \nabla u_k\right\| \|u_k\|,$$

and

$$\left\| \left(\nabla^+ \nabla \right)^{3/2} u_k \right\|^2 = \left(\left(\nabla^+ \nabla \right)^{3/2} u_k, \left(\nabla^+ \nabla \right)^{3/2} u_k \right) \\ = \left(\left(\nabla^+ \nabla \right)^3 u_k, u_k \right) \le \left\| \left(\nabla^+ \nabla \right)^3 u_k \right\| \|u_k\|.$$

We will prove that $(\nabla^+\nabla)^3 u_k \to (\nabla^+\nabla)^3 u, \quad Vu_k \to Vu, \quad V^{1/2} (\nabla^+\nabla)^{3/2} u_k \to V^{1/2} (\nabla^+\nabla)^{3/2} u, \quad \nabla u_k \to \nabla u \text{ and } (\nabla^+\nabla)^{3/2} u_k \to (\nabla^+\nabla)^{3/2} u.$ As the essential self-adjointness of $\nabla^+\nabla|_{C_c^{\infty}(E)}$ and $(\nabla^+\nabla)^3|_{C_c^{\infty}(E)}$ we obtain $(\nabla^+\nabla)^{3/2} u_k \to (\nabla^+\nabla)^{3/2} u$ and $(\nabla^+\nabla)^3 u_k \to (\nabla^+\nabla)^3 u$ in $L^2(E)$. As $\{\nabla u_k\}$ is a Cauchy sequences in $L^2(T^*M \otimes E)$, it follows that ∇u_k convergent

to some elements $z \in L^2(T^*M \otimes E)$, respectively then for all $\Psi \in C_c^{\infty}(T^*M \otimes E)$ we have $0 = (\nabla u_k, \Psi) - (u_k, \nabla^+ \Psi) \rightarrow (z, \Psi) - (u, \nabla^+ \Psi) = (z, \Psi) - (\nabla u, \Psi)$, as $Dom\left(H_{2.V}^{\nabla}\right) \subset W_{loc}^{2,2}(E)$ (see, Lemma 8.8 in [8]). With the convergence relations, we take the limit as $k \rightarrow \infty$ in all terms in (4.3) and (4.4) with u replaced by u_k then (4.3) and (4.4) hold for all $u \in D_2^{\nabla} = Dom\left(H_{2.V}^{\nabla}\right)$.

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