# One Parameter Commutative Octonions 

Göksal Bilgici ${ }^{1}$<br>${ }^{1}$ Department of Elementary Mathematics Education, Education Faculty, Kastamonu University 37200, Kastamonu, tÜRKIYE


#### Abstract

Hyperbolic numbers had been developed in the 19th century. Octonions forms a noncommutative and nonassociative normed division algebra over reals. Octonions have many applications in fields of physics such as quantum logic and string theory. Cayley-Dickson process is applied to quaternions in order to construct octonions and in a sense, we follow a similar process. The aim of this study is to introduce the concept of commutative octonions. We construct this algebra by using some matrix methods. After construction, we give a number of properties of commutative octonions such as fundamental matrices and principal conjugates. We also obtain representation of a commutative octonion as decomposed form, holomorphic and analytic functions of commutative octonions.


Keywords: Commutative octonions; Segre's quaternions; holomorphic functions
2010 Mathematics Subject Classification: Primary 13A99, 11R52; Secondary 32A30

## 1. Introduction

Four-dimensional hyper complex numbers, real quaternions, were introduced by Hamilton in 1843 to extend complex numbers [7].The set of real quaternions are

$$
\mathbb{H}=\left\{x_{0}+x_{1} i+x_{2} j+x_{3} k: i, j, k \notin \mathbb{R} \text { and } x_{0}, x_{1}, x_{2}, x_{3} \in \mathbb{R}\right\}
$$

where the multiplication rules of the elements of the ordered basis $\{1, i, j, k\}$ is

$$
i^{2}=j^{2}=k^{2}=-1,-j i=i j=k,-k j=j k=i \text { and }-i k=k i=j .
$$

One can see that quaternions are not commutative. The Hamilton quaternions form a division algebra and in this aspect, they can be regarded as an extension of complex numbers. Octonions have many applications in field. For instance, we can give M-theory cosmology [6], quantum theory [5, 12], cough monitoring [8], space time coding [13], electromagnetic and gravitational equations [16] as most striking examples. Segre gave a new type of quaternions whose multiplication rule has commutative property [11]. These numbers are called commutative quaternions or Segre's quaternions. The set of commutative quaternions is [3]

$$
\mathrm{Q}=\left\{x_{0}+x_{1} i+x_{2} j+x_{3} k: i, j, k \notin \mathbb{R} \text { and } x_{0}, x_{1}, x_{2}, x_{3} \in \mathbb{R}\right\}
$$

where the versors satisfy

$$
i^{2}=k^{2}=-1, j^{2}=1, j i=i j=k,-k j=j k=i \text { and } i k=k i=-j .
$$

A more general commutative quaternions were given by Catoni et al. [3]. They represented these numbers by

$$
\mathbb{S}=\left\{x_{0}+x_{1} i+x_{2} j+x_{3} k: i, j, k \notin \mathrm{R} \text { and } x_{0}, x_{1}, x_{2}, x_{3} \in \mathbb{R}\right\}
$$

where the versors satisfy the following multiplication rules where $\delta$ is an arbitrary real number.
Table 1: Multiplication rules of the elements of $\{1, i, j, k\}$.

| 1 | $i$ | $j$ | $k$ |
| :--- | :--- | :--- | :--- |
| $i$ | $\delta$ | $k$ | $\delta j$ |
| $j$ | $k$ | 1 | $i$ |
| $k$ | $\delta j$ | $i$ | $\delta$ |

They also examined algebraic properties of this type of quaternions. Following their work [2,3], we construct the commutative octonions. This study can be regarded as an application of [2]. Kosal et al. [9] studied matrices for commutative quaternions and gave some interesting properties.
After discovery of quaternion algebra, Cayley and Graves gave octonion algebra independently. Octonions algebra is constructed by using the Cayley-Dickson method. An octonion o can be written as

$$
o=p+p^{\prime} e
$$

where $p, p^{\prime} \in \mathbb{Q}$ and $e$ is a new imaginary unit, i.e. it is a square root of -1 . Let $o_{1}=p_{1}+p_{1}^{\prime} e$ and $o_{2}=p_{2}+p_{2}^{\prime} e$ be two any octonions. Addition and multiplication of these two octonions are

$$
\begin{aligned}
o_{1}+o_{2} & =p_{1}+p_{1}+\left(p_{1}^{\prime}+p_{2}^{\prime}\right) e \\
o_{1} o_{2} & =\left(p_{1} p_{2}-\overline{p_{2}^{\prime}} p_{1}^{\prime}\right)+\left(p_{2}^{\prime} p_{1}+p_{1}^{\prime} \overline{p_{2}^{\prime}}\right)
\end{aligned}
$$

where $\bar{q}$ is the conjugate of the quaternion $q$. The ordered basis for octonion algebra over $\mathbb{R}$ consists of the elements

$$
e_{0}=1, e_{1}=i, e_{2}=j, e_{3}=k, e_{4}=t, e_{5}=i t, e_{6}=j t, e_{7}=k t
$$

where $t$ is another versor different from $\{1, i, j, k\}$, and any octonion $o$ can be expressed as

$$
o=\sum_{i=0}^{7} a_{i} e_{i}, \quad a_{i} \in \mathbb{R}
$$

Thus, there are eight objects $e_{i} \quad(i=0, \ldots, 7)$ in the ordered basis of octonion algebra. The multiplication rules of the elements of standard basis $\left\{e_{0}, e_{1}, e_{2}, \cdots, e_{7}\right\}$ for octonions algebra can be found in [10]. The octonions division algebra over the real numbers $\mathbb{R}$ is a non-commutative and non-associative algebra.
There are some studies on octonions whose coefficients are well-known integer sequences. We can refer to [1, 14, 15] for this type of studies.

## 2. Commutative Octonions

By following Catoni et al [3], for any real number $\alpha$, we define the commutative octonions with the help of the multiplication rules as follows.

Table 2: Multiplication rules of elements of standard basis $\left\{e_{0}, e_{1}, e_{2}, \cdots, e_{7}\right\}$ for commutative octonions algebra.

| $\cdot$ | 1 | $e_{1}$ | $e_{2}$ | $e_{3}$ | $e_{4}$ | $e_{5}$ | $e_{6}$ | $e_{7}$ |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 1 | 1 | $e_{1}$ | $e_{2}$ | $e_{3}$ | $e_{4}$ | $e_{5}$ | $e_{6}$ | $e_{7}$ |
| $e_{1}$ | $e_{1}$ | $\alpha$ | $e_{3}$ | $\alpha e_{2}$ | $e_{5}$ | $\alpha e_{4}$ | $e_{7}$ | $\alpha e_{6}$ |
| $e_{2}$ | $e_{2}$ | $e_{3}$ | 1 | $e_{1}$ | $e_{6}$ | $e_{7}$ | $e_{4}$ | $e_{5}$ |
| $e_{3}$ | $e_{3}$ | $\alpha e_{2}$ | $e_{1}$ | $\alpha$ | $e_{7}$ | $\alpha e_{6}$ | $e_{5}$ | $\alpha e_{4}$ |
| $e_{4}$ | $e_{4}$ | $e_{5}$ | $e_{6}$ | $e_{7}$ | 1 | $e_{1}$ | $e_{2}$ | $e_{3}$ |
| $e_{5}$ | $e_{5}$ | $\alpha e_{4}$ | $e_{7}$ | $\alpha e_{6}$ | $e_{1}$ | $\alpha$ | $e_{3}$ | $\alpha e_{2}$ |
| $e_{6}$ | $e_{6}$ | $e_{7}$ | $e_{4}$ | $e_{5}$ | $e_{2}$ | $e_{3}$ | 1 | $e_{1}$ |
| $e_{7}$ | $e_{7}$ | $\alpha e_{6}$ | $e_{5}$ | $\alpha e_{4}$ | $e_{3}$ | $\alpha e_{2}$ | $e_{1}$ | $\alpha$ |

These multiplication rules can be obtained by the similar way to the octonion algebra mentioned above. Let $\mathbb{O}$ be the set of commutative octonions, i.e.

$$
\mathbb{O}=\left\{o=\sum_{i=0}^{7} a_{i} e_{i}: a_{0}, a_{1}, \ldots, a_{7} \in \mathbb{R}, e_{0}=1, e_{1}, e_{2}, \ldots, e_{7} \notin \mathbb{R}\right\}
$$

where the versors $e_{0}, e_{1}, \ldots, e_{7}$ satisfy the multiplication rules in Table 2.
Let $o=\sum_{i=0}^{7} c_{i} e_{i} \in \mathbb{O}$, then the characteristic matrix of $o$ is
$\mathrm{N}=\left[\begin{array}{rrrrrrrr}\mathrm{c}_{0} & \alpha \mathrm{c}_{1} & \mathrm{c}_{2} & \alpha \mathrm{c}_{3} & \mathrm{c}_{4} & \alpha \mathrm{c}_{5} & \mathrm{c}_{6} & \alpha \mathrm{c}_{7} \\ \mathrm{c}_{1} & \mathrm{c}_{0} & \mathrm{c}_{3} & \mathrm{c}_{2} & \mathrm{c}_{5} & \mathrm{c}_{4} & \mathrm{c}_{7} & \mathrm{c}_{6} \\ \mathrm{c}_{2} & \alpha \mathrm{c}_{3} & \mathrm{c}_{0} & \alpha \mathrm{c}_{1} & \mathrm{c}_{6} & \alpha \mathrm{c}_{7} & \mathrm{c}_{4} & \alpha \mathrm{c}_{5} \\ \mathrm{c}_{3} & \mathrm{c}_{2} & \mathrm{c}_{1} & \mathrm{c}_{0} & \mathrm{c}_{7} & \mathrm{c}_{6} & \mathrm{c}_{5} & \mathrm{c}_{4} \\ \mathrm{c}_{4} & \alpha \mathrm{c}_{5} & \mathrm{c}_{6} & \alpha \mathrm{c}_{7} & \mathrm{c}_{0} & \alpha \mathrm{c}_{1} & \mathrm{c}_{2} & \alpha \mathrm{c}_{3} \\ \mathrm{c}_{5} & \mathrm{c}_{4} & \mathrm{c}_{7} & \mathrm{c}_{6} & \mathrm{c}_{1} & \mathrm{c}_{0} & \mathrm{c}_{3} & \mathrm{c}_{2} \\ \mathrm{c}_{6} & \alpha \mathrm{c}_{7} & \mathrm{c}_{4} & \alpha \mathrm{c}_{5} & \mathrm{c}_{2} & \alpha \mathrm{c}_{3} & c_{0} & \alpha \mathrm{c}_{1} \\ \mathrm{c}_{7} & \mathrm{c}_{6} & \mathrm{c}_{5} & \mathrm{c}_{4} & \mathrm{c}_{3} & \mathrm{c}_{2} & \mathrm{c}_{1} & \mathrm{c}_{0}\end{array}\right] \equiv\left[\begin{array}{l|l}\Psi & \Omega \\ \hline \Omega & \Psi\end{array}\right]$
where $\Psi$ and $\Omega$ are the following $4 \times 4$ matrices:
$\Psi=\left[\begin{array}{rrrr}c_{0} & \alpha c_{1} & c_{2} & \alpha c_{3} \\ \mathrm{c}_{1} & \mathrm{c}_{0} & \mathrm{c}_{3} & \mathrm{c}_{2} \\ \mathrm{c}_{2} & \alpha \mathrm{c}_{3} & \mathrm{c}_{0} & \alpha \mathrm{c}_{1} \\ \mathrm{c}_{3} & \mathrm{c}_{2} & \mathrm{c}_{1} & \mathrm{c}_{0}\end{array}\right]$ and $\Omega=\left[\begin{array}{rrrr}\mathrm{c}_{4} & \alpha \mathrm{c}_{5} & \mathrm{c}_{6} & \alpha \mathrm{c}_{7} \\ \mathrm{c}_{5} & \mathrm{c}_{4} & \mathrm{c}_{7} & \mathrm{c}_{6} \\ \mathrm{c}_{6} & \alpha \mathrm{c}_{7} & \mathrm{c}_{4} & \alpha \mathrm{c}_{5} \\ \mathrm{c}_{7} & \mathrm{c}_{6} & \mathrm{c}_{5} & \mathrm{c}_{4}\end{array}\right]$.

Similarly, we can express matrices $\Psi$ and $\Omega$ as

$$
\Psi=\left[\begin{array}{cc}
\Psi^{\prime} & \Psi^{\prime \prime} \\
\Psi^{\prime \prime} & \Psi^{\prime}
\end{array}\right] \text { and } \Omega=\left[\begin{array}{cc}
\Omega^{\prime} & \Omega^{\prime \prime} \\
\Omega^{\prime \prime} & \Omega^{\prime}
\end{array}\right]
$$

where

$$
\Psi^{\prime}=\left[\begin{array}{rr}
c_{0} & \alpha c_{1} \\
c_{1} & c_{0}
\end{array}\right], \Psi^{\prime \prime}=\left[\begin{array}{rr}
c_{2} & \alpha c_{3} \\
c_{3} & c_{2}
\end{array}\right]
$$

and

$$
\Omega^{\prime}=\left[\begin{array}{cc}
c_{4} & \alpha c_{5} \\
c_{5} & c_{4}
\end{array}\right], \Omega^{\prime \prime}=\left[\begin{array}{cc}
c_{6} & \alpha c_{7} \\
c_{7} & c_{6}
\end{array}\right]
$$

Determinant of the matrix N is given in the following theorem.
Theorem 2.1. Determinant of the matrix N is

$$
\begin{align*}
\operatorname{det}(\mathrm{N})= & {\left[\varepsilon(0,2,4,6)^{2}-\alpha \varepsilon(1,3,5,7)^{2}\right] \times\left[\varepsilon(0,-2,4,-6)^{2}-\alpha \varepsilon(1,-3,5,-7)^{2}\right] } \\
& \times\left[\varepsilon(0,2,-4,-6)^{2}-\alpha \varepsilon(1,3,-5,-7)^{2}\right] \times\left[\varepsilon(0,-2,-4,6)^{2}-\alpha \varepsilon(1,-3,-5,7)^{2}\right] . \tag{2.2}
\end{align*}
$$

where $\varepsilon(p, q, r, s)=c_{p}+c_{q}+c_{r}+c_{s}$ and $\varepsilon(-p)=-c_{p}$.
Proof. If we evaluate the determinant of the matrix N in Eq. (2.1), we obtain

$$
\begin{aligned}
\operatorname{det}(\mathrm{N}) & =\operatorname{det}\left(\Psi^{2}-\Omega^{2}\right) \\
& =\operatorname{det}\left(\left(\Psi^{\prime}+\Omega^{\prime}\right)^{2}-\left(\Psi^{\prime \prime}+\Omega^{\prime \prime}\right)^{2}\right) \times \operatorname{det}\left(\left(\Psi^{\prime}-\Omega^{\prime}\right)^{2}-\left(\Psi^{\prime \prime}-\Omega^{\prime \prime}\right)^{2}\right) \\
& =\operatorname{det}\left(\Psi^{\prime}+\Omega^{\prime}+\Psi^{\prime \prime}+\Omega^{\prime \prime}\right) \times \operatorname{det}\left(\Psi^{\prime}+\Omega^{\prime}-\Psi^{\prime \prime}-\Omega^{\prime \prime}\right) \times \operatorname{det}\left(\Psi^{\prime}-\Omega^{\prime}+\Psi^{\prime \prime}-\Omega^{\prime \prime}\right) \times \operatorname{det}\left(\Psi^{\prime}-\Omega^{\prime}-\Psi^{\prime \prime}+\Omega^{\prime \prime}\right)
\end{aligned}
$$

Substituting determinants of these matrices into the last equation completes the proof.
Eq. (2.1) also gives the principal conjugations of the commutative octonion $o=\sum_{i=0}^{7} c_{i} e_{i}$ as follows
$o_{1}=\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{3} e_{3}-\left(\mathrm{c}_{4} e_{4}+\mathrm{c}_{5} e_{5}+\mathrm{c}_{6} e_{6}+\mathrm{c}_{7} e_{7}\right)$,
$o_{2}=\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\mathrm{c}_{4} e_{4}+\mathrm{c}_{5} e_{5}-\left(\mathrm{c}_{2} e_{2}+\mathrm{c}_{3} e_{3}+\mathrm{c}_{6} e_{6}+\mathrm{c}_{7} e_{7}\right)$,
$o_{3}=\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\mathrm{c}_{6} e_{6}+\mathrm{c}_{7} e_{7}-\left(\mathrm{c}_{2} e_{2}+\mathrm{c}_{3} e_{3}+\mathrm{c}_{4} e_{4}+\mathrm{c}_{5} e_{5}\right)$,
$o_{4}=c_{0}+c_{2} e_{2}+c_{4} e_{4}+c_{6} e_{6}-\left(c_{1} e_{1}+c_{3} e_{3}+c_{5} e_{5}+c_{7} e_{7}\right)$,
$o_{5}=\mathrm{c}_{0}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{5} e_{5}+\mathrm{c}_{7} e_{7}-\left(\mathrm{c}_{1} e_{1}+\mathrm{c}_{3} e_{3}+\mathrm{c}_{4} e_{4}+\mathrm{c}_{6} e_{6}\right)$,
$o_{6}=\mathrm{c}_{0}+\mathrm{c}_{3} e_{3}+\mathrm{c}_{4} e_{4}+\mathrm{c}_{7} e_{7}-\left(\mathrm{c}_{1} e_{1}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{5} e_{5}+\mathrm{c}_{6} e_{6}\right)$,
and
$o_{7}=\mathrm{c}_{0}+\mathrm{c}_{3} e_{3}+\mathrm{c}_{5} e_{5}+\mathrm{c}_{6} e_{6}-\left(\mathrm{c}_{1} e_{1}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{4} e_{4}+\mathrm{c}_{7} e_{7}\right)$.
From these conjugations, for $\alpha \neq 0$, we can construct the bijective mapping between $\mathrm{c}_{0}, \mathrm{c}_{1}, \ldots, \mathrm{c}_{7} \rightarrow o_{0}=o, o_{1}, \ldots, 0_{7}$ as
$\mathrm{c}_{0}=\frac{\zeta(0,1,2,3,4,5,6,7)}{8}$,
$\mathrm{c}_{1}=e_{1} \frac{\zeta(0,1,2,3,-4,-5,-6,-7)}{8 \alpha}$,
$c_{2}=e_{2} \frac{\zeta(0,1,-2,-3,4,5,-6,-7)}{8}$,
$c_{3}=e_{3} \frac{\zeta(0,1,-2,-3,-4,-5,6,7)}{8 \alpha}$,
$\mathrm{c}_{4}=e_{4} \frac{\zeta(0,-1,2,-3,4,-5,6,-7)}{8}$,
$c_{5}=e_{5} \frac{\zeta(0,-1,2,-3,-4,5,-6,7)}{8 \alpha}$,
$\mathrm{c}_{6}=e_{6} \frac{\zeta(0,-1,-2,3,4,-5,-6,7)}{8}$
$\mathrm{c}_{7}=e_{7} \frac{\zeta(0,-1,-2,3,-4,5,6,-7)}{8 \alpha}$
where $\zeta\left(r_{0}, r_{1}, \cdots, r_{7}\right)=o_{r_{0}}+o_{r_{1}}+\cdots+o_{r_{7}}$ and $\zeta(-r)=-o_{r}$.
From the eigenvalues of the characteristic matrix (2.1), we obtain the norm of a commutative octonion $o=\sum_{i=0}^{7} c_{i} e_{i}$ as
$\|o\|=o \cdot o_{1} \cdot o_{2} \cdots o_{7}=\operatorname{det}(\mathrm{N})$.

By using Eq.(2.1), we can obtain the following matrix expressions of all the versors:
$\mathrm{e}_{0}=\left[\begin{array}{cccccccc}1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1\end{array}\right], \mathrm{e}_{1}=\left[\begin{array}{cccccccc}0 & \alpha & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \alpha & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & \alpha & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \alpha \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0\end{array}\right]$,
$\mathrm{e}_{2}=\left[\begin{array}{llllllll}0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0\end{array}\right], \mathrm{e}_{3}=\left[\begin{array}{llllllll}0 & 0 & 0 & \alpha & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \alpha & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \alpha \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & \alpha & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0\end{array}\right]$,
$\mathrm{e}_{4}=\left[\begin{array}{llllllll}0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0\end{array}\right], \mathrm{e}_{5}=\left[\begin{array}{llllllll}0 & 0 & 0 & 0 & 0 & \alpha & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \alpha \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & \alpha & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \alpha & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0\end{array}\right]$,
$\mathrm{e}_{6}=\left[\begin{array}{llllllll}0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0\end{array}\right], \mathrm{e}_{7}=\left[\begin{array}{llllllll}0 & 0 & 0 & 0 & 0 & 0 & 0 & \alpha \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & \alpha & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & \alpha & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \alpha & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0\end{array}\right]$.

## 3. Some Properties of Commutative Octonions

Except two systems, any system which satisfies associativity and distributivity with respect to the sum, it can satisfy either commutativity or it does not have divisors of zero [2]. One can see that the commutative octonions satisfy the commutative properties from the multiplication rules easily. For the first two properties (distributive and associative), the characteristic matrix in Eq.(2.1) can be used.
A commutative octonion can be represented by four complex numbers with four linearly independent bases. Let $o=\sum_{i=0}^{7} c_{i} e_{i}$ be a commutative octonion. We can derive

$$
\begin{align*}
o & =\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\cdots+\mathrm{c}_{7} e_{7} \\
& =\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{3} e_{3}+\left(\mathrm{c}_{4}+\mathrm{c}_{5} e_{1}+\mathrm{c}_{6} e_{2}+\mathrm{c}_{7} e_{3}\right) e_{4} \\
& =\left(\mathrm{c}_{0}+\mathrm{c}_{1} e_{1}+\mathrm{c}_{2} e_{2}+\mathrm{c}_{3} e_{3}\right)\left(\frac{1+e_{4}}{2}+\frac{1-e_{4}}{2}\right)+\left(\mathrm{c}_{4}+\mathrm{c}_{5} e_{1}+\mathrm{c}_{6} e_{2}+\mathrm{c}_{7} e_{3}\right)\left(\frac{1+e_{4}}{2}-\frac{1-e_{4}}{2}\right) \\
& =q_{1} e_{1}^{\prime}+q_{2} e_{2}^{\prime} \tag{3.1}
\end{align*}
$$

where
$q_{1}=\left(\mathrm{c}_{0}+\mathrm{c}_{4}\right)+\left(\mathrm{c}_{1}+\mathrm{c}_{5}\right) e_{1}+\left(\mathrm{c}_{2}+\mathrm{c}_{6}\right) e_{2}+\left(\mathrm{c}_{3}+\mathrm{c}_{7}\right) e_{3}$,
$q_{2}=\left(\mathrm{c}_{0}-\mathrm{c}_{4}\right)+\left(\mathrm{c}_{1}-\mathrm{c}_{5}\right) e_{1}+\left(\mathrm{c}_{2}-\mathrm{c}_{6}\right) e_{2}+\left(\mathrm{c}_{3}-\mathrm{c}_{7}\right) e_{3}$,
and

$$
e_{1}^{\prime}=\frac{1+e_{4}}{2}, \quad e_{2}^{\prime}=\frac{1-e_{4}}{2}
$$

Thus we can represent a commutative octonion by two commutative quaternions. Here $e_{1}^{\prime}$ and $e_{2}^{\prime}$ are idempotent basis and satisfy

$$
\left(e_{1}^{\prime}\right)^{2}=e_{1}^{\prime}, \quad\left(e_{2}^{\prime}\right)^{2}=e_{2}^{\prime}, \quad e_{1}^{\prime} e_{2}^{\prime}=0
$$

From [3], we know that the commutative quaternions $q_{1}$ and $q_{2}$ can be represented as follows
$q_{1}=x^{(1)} e_{1}^{\prime \prime}+x^{(2)} e_{2}^{\prime \prime}, \quad q_{2}=x^{(3)} e_{1}^{\prime \prime}+x^{(4)} e_{2}^{\prime \prime}$
where

$$
e_{1}^{\prime \prime}=\frac{1+e_{2}}{2}, e_{2}^{\prime \prime}=\frac{1-e_{2}}{2}
$$

and
$x^{(1)}=\left(x_{0}+x_{2}\right)+\left(x_{1}+x_{3}\right) e_{1}$,
$x^{(2)}=\left(x_{0}-x_{2}\right)+\left(x_{1}-x_{3}\right) e_{1}$,
$x^{(3)}=\left(y_{0}+y_{2}\right)+\left(y_{1}+y_{3}\right) e_{1}$,
$x^{(4)}=\left(y_{0}-y_{2}\right)+\left(y_{1}-y_{3}\right) e_{1}$.
Here we set

$$
x_{i}=c_{i}+c_{i+4} \text { and } y_{i}=c_{i}-c_{i+4}(0 \leq i \leq 3) .
$$

By using Eqs.(3.1) and (3.2), we obtain

$$
\begin{align*}
o & =q_{1} e_{1}^{\prime}+q_{2} e_{2}^{\prime} \\
& =\left[x^{(1)} e_{1}^{\prime \prime}+x^{(2)} e_{2}^{\prime \prime}\right] e_{1}^{\prime}+\left[x^{(3)} e_{1}^{\prime \prime}+x^{(4)} e_{2}^{\prime \prime}\right] e_{2}^{\prime} \\
& =x^{(1)} i_{1}+x^{(2)} i_{2}+x^{(3)} i_{3}+x^{(4)} i_{4} \tag{3.3}
\end{align*}
$$

where
$i_{1}=e_{1}^{\prime \prime} e_{1}^{\prime}=\frac{1+e_{2}+e_{4}+e_{6}}{4}, i_{2}=e_{2}^{\prime \prime} e_{1}^{\prime}=\frac{1-e_{2}+e_{4}-e_{6}}{4}, i_{3}=e_{1}^{\prime \prime} e_{2}^{\prime}=\frac{1+e_{2}-e_{4}-e_{6}}{4}, i_{4}=e_{2}^{\prime \prime} e_{2}^{\prime}=\frac{1-e_{2}-e_{4}+e_{6}}{4}$.
Here $i_{k}(k=1, \ldots, 4)$ satisfy
$i_{k}^{2}=i_{k}$ and $i_{k} i_{l}=0($ for $k \neq l)$.
Finally, to summarize above transformations, we can say that a commutative octonion $o=\sum_{i=0}^{7} c_{i} e_{i}$ can be represented as the form in Eq.(3.3) where
$x^{(1)}=\varepsilon(0,2,4,6)+\varepsilon(1,3,5,7) e_{1}$,
$x^{(2)}=\varepsilon(0,-2,4,-6)+\varepsilon(1,-3,5,-7) e_{1}$
$x^{(3)}=\varepsilon(0,2,-4,-6)+\varepsilon(1,3,-5,-7) e_{1}$
$x^{(4)}=\varepsilon(0,-2,-4,6)+\varepsilon(1,-3,-5,7) e_{1}$.
Thus $i_{1}, i_{2}, i_{3}$ and $i_{4}$ are linearly independent and the set of commutative octonions can be shown as a direct sum of four complex number fields. For any positive integer $n$, Eqs. (3.3) and (3.4) give
$o^{n}=\left(x^{(1)}\right)^{n} i_{1}+\left(x^{(2)}\right)^{n} i_{2}+\left(x^{(3)}\right)^{n} i_{3}+\left(x^{(4)}\right)^{n} i_{4}$.
Let $o^{\prime}=x_{1}^{(1)} i_{1}+x_{1}^{(2)} i_{2}+x_{1}^{(3)} i_{3}+x_{1}^{(4)} i_{4}$ and $o^{\prime \prime}=x_{2}^{(1)} i_{1}+x_{2}^{(2)} i_{2}+x_{2}^{(3)} i_{3}+x_{2}^{(4)} i_{4}$ be two commutative octonions. Then by using Eqs. (3.3) and (3.4) again, we have
$o^{\prime} o^{\prime \prime}=x_{1}^{(1)} x_{2}^{(1)} i_{1}+x_{1}^{(2)} x_{2}^{(2)} i_{2}+x_{1}^{(3)} x_{2}^{(3)} i_{3}+x_{1}^{(4)} x_{2}^{(4)} i_{4}$
and
$\frac{o^{\prime}}{o^{\prime \prime}}=\frac{x_{1}^{(1)}}{x_{2}^{(1)}} i_{1}+\frac{x_{1}^{(2)}}{x_{2}^{(2)}} i_{2}+\frac{x_{1}^{(3)}}{x_{2}^{(3)}} i_{3}+\frac{x_{1}^{(4)}}{x_{2}^{(4)}} i_{4}$.
Seven conjugates of a commutative octonion in Eqs.(2.3) - (2.9) can be represented according to the decomposed form (3.3) as follows
Theorem 3.1. For a commutative octonion $o=\sum_{i=0}^{7} c_{i} e_{i}=x^{(1)} i_{1}+x^{(2)} i_{2}+x^{(3)} i_{3}+x^{(4)} i_{4}$, we have
$o_{1}=x^{(3)} i_{1}+x^{(4)} i_{2}+x^{(1)} i_{3}+x^{(2)} i_{4}$,
$o_{2}=x^{(2)} i_{1}+x^{(1)} i_{2}+x^{(4)} i_{3}+x^{(3)} i_{4}$,
$o_{3}=x^{(4)} i_{1}+x^{(3)} i_{2}+x^{(2)} i_{3}+x^{(1)} i_{4}$,
$o_{4}=\overline{x^{(1)}} i_{1}+\overline{x^{(2)}} i_{2}+\overline{x^{(3)}} i_{3}+\overline{x^{(4)}} i_{4}$,
$o_{5}=\overline{x^{(3)}} i_{1}+\overline{x^{(4)}} i_{2}+\overline{x^{(1)}} i_{3}+\overline{x^{(2)}} i_{4}$,
$o_{6}=\overline{x^{(2)}} i_{1}+\overline{x^{(1)}} i_{2}+\overline{x^{(4)}} i_{3}+\overline{x^{(3)}} i_{4}$,
$o_{7}=\overline{x^{(4)}} i_{1}+\overline{x^{(3)}} i_{2}+\overline{x^{(2)}} i_{3}+\overline{x^{(1)}} i_{4}$
where $\bar{x}$ is the complex conjugate of a complex number $x$.
Proof. The proof can be done easily by using Eqs. (2.3) - (2.9), Eqs.(3.5) - (3.8) and the multiplication rule (3.10).

We can give the following results immediately.
Corollary For a commutative octonion $o=\sum_{i=0}^{7} c_{i} e_{i}=x^{(1)} i_{1}+x^{(2)} i_{2}+x^{(3)} i_{3}+x^{(4)} i_{4}$, we have

$$
\begin{align*}
o_{4} & =\left\|x^{(1)}\right\|^{2} i_{1}+\left\|x^{(2)}\right\|^{2} i_{2}+\left\|x^{(3)}\right\|^{2} i_{3}+\left\|x^{(4)}\right\|^{2} i_{4},  \tag{3.19}\\
o_{1} o_{5} & =\left\|x^{(3)}\right\|^{2} i_{1}+\left\|x^{(4)}\right\|^{2} i_{2}+\left\|x^{(1)}\right\|^{2} i_{3}+\left\|x^{(2)}\right\|^{2} i_{4},  \tag{3.20}\\
o_{2} o_{6} & =\left\|x^{(2)}\right\|^{2} i_{1}+\left\|x^{(1)}\right\|^{2} i_{2}+\left\|x^{(4)}\right\|^{2} i_{3}+\left\|x^{(3)}\right\|^{2} i_{4},  \tag{3.21}\\
o_{3} o_{7} & =\left\|x^{(4)}\right\|^{2} i_{1}+\left\|x^{(3)}\right\|^{2} i_{2}+\left\|x^{(2)}\right\|^{2} i_{3}+\left\|x^{(1)}\right\|^{2} i_{4} \tag{3.22}
\end{align*}
$$

where $\|x\|$ is the absolute square of a complex number $x$, i.e. $\|x\|^{2}=x \bar{x}$.

### 3.1. Holomorphic and Analytic Functions

We show an octonion function by
$G \equiv \sum_{i=0}^{7} G_{i}\left(\mathrm{c}_{0}, \mathrm{c}_{1}, \ldots, \mathrm{c}_{7}\right) e_{i}$
where $G_{i}(i=0,1, \cdots, 7)$ are real functions with partial derivatives for the variables, $\mathrm{c}_{0}, \mathrm{c}_{1}, \cdots, \mathrm{c}_{7}$. Catoni et al. [2] introduced the Generalized Cauchy - Riemann - like (GCR) conditions. While there are some methods for calculation of the GCR conditions, Catoni et al. [4] give the following theorem.
Theorem 3.2. [4, p.91] The Jacobian matrix of a hypercomplex function's components has the same form of the characteristic matrix.
By combining this method and the study of Catoni et al. [3] for Segre's commutative quaternions we have the following theorem for holomorphic functions of octonion.
Theorem 3.3. $G$ is called a holomorphic functions of octonion if

1) $G$ is differentiable with non-zero derivatives and not a zero divisor,
2) The GCR conditions for the partial derivatives of components of $G$ are
$G_{0, \mathrm{c}_{0}}=G_{1, \mathrm{c}_{1}}=G_{2, \mathrm{c}_{2}}=G_{3, \mathrm{c}_{3}}=G_{4, \mathrm{c}_{4}}=G_{5, \mathrm{c}_{5}}=G_{6, \mathrm{c}_{6}}=G_{7, \mathrm{c}_{7}}$
$G_{0, \mathrm{c}_{1}}=\alpha G_{1, \mathrm{c}_{0}}=G_{2, \mathrm{c}_{3}}=\alpha G_{3, \mathrm{c}_{2}}=G_{4, \mathrm{c}_{5}}=\alpha G_{5, \mathrm{c}_{4}}=G_{6, \mathrm{c}_{7}}=\alpha G_{7, \mathrm{c}_{6}}$
$G_{0, \mathrm{c}_{2}}=G_{1, \mathrm{c}_{3}}=G_{2, \mathrm{c}_{0}}=G_{3, \mathrm{c}_{1}}=G_{4, \mathrm{c}_{6}}=G_{5, \mathrm{c}_{7}}=G_{6, \mathrm{c}_{4}}=G_{7, \mathrm{c}_{5}}$
$G_{0, \mathrm{c}_{3}}=\alpha G_{1, \mathrm{c}_{2}}=G_{2, \mathrm{c}_{1}}=\alpha G_{3, \mathrm{c}_{0}}=G_{4, \mathrm{c}_{7}}=\alpha G_{5, \mathrm{c}_{6}}=G_{6, \mathrm{c}_{5}}=\alpha G_{7, \mathrm{c}_{4}}$
$G_{0, \mathrm{c}_{4}}=G_{1, \mathrm{c}_{5}}=G_{2, \mathrm{c}_{6}}=G_{3, \mathrm{c}_{7}}=G_{4, \mathrm{c}_{0}}=G_{5, \mathrm{c}_{1}}=G_{6, \mathrm{c}_{2}}=G_{7, \mathrm{c}_{3}}$
$G_{0, \mathrm{c}_{5}}=\alpha G_{1, \mathrm{c}_{4}}=G_{2, \mathrm{c}_{7}}=\alpha G_{3, \mathrm{c}_{6}}=G_{4, \mathrm{c}_{1}}=\alpha G_{5, \mathrm{c}_{0}}=G_{6, \mathrm{c}_{3}}=\alpha G_{7, \mathrm{c}_{2}}$
$G_{0, \mathrm{c}_{6}}=G_{1, \mathrm{c}_{7}}=G_{2, \mathrm{c}_{4}}=G_{3, \mathrm{c}_{5}}=G_{4, \mathrm{c}_{2}}=G_{5, \mathrm{c}_{3}}=G_{6, \mathrm{c}_{0}}=G_{7, \mathrm{c}_{1}}$
$G_{0, \mathrm{c}_{7}}=\alpha G_{1, \mathrm{c}_{6}}=G_{2, \mathrm{c}_{5}}=\alpha G_{3, \mathrm{c}_{4}}=G_{4, \mathrm{c}_{3}}=\alpha G_{5, \mathrm{c}_{2}}=G_{6, \mathrm{c}_{1}}=\alpha G_{7, \mathrm{c}_{0}}$.
Let $G(o)$ be an octonion holomorphic function and its power series in $q$ about 0 be
$G(o)=\sum_{r=0}^{\infty} t_{r} o^{r}$
where $t_{r} \in \mathbb{O}$. From Eq.(3.1), we write

$$
t_{r}=f_{r} e_{1}^{\prime}+g_{r} e_{2}^{\prime}
$$

where $f_{r}, g_{r} \in \mathbb{H}$. Then we obtain

$$
\begin{aligned}
G(q) & =\sum_{r=0}^{\infty}\left(f_{r} e_{1}^{\prime}+g_{r} e_{2}^{\prime}\right)\left(q_{1} e_{1}^{\prime}+q_{2} e_{2}^{\prime}\right)^{r} \\
& =\sum_{r=0}^{\infty}\left(f_{r} e_{1}^{\prime}+g_{r} e_{2}^{\prime}\right)\left(q_{1}^{r} e_{1}^{\prime}+q_{2}^{r} e_{2}^{\prime}\right) \\
& =\sum_{r=0}^{\infty}\left(f_{r} q_{1}^{r} e_{1}^{\prime}+g_{r} q_{2}^{r} e_{2}^{\prime}\right) \\
& =e_{1}^{\prime} \sum_{r=0}^{\infty} f_{r} q_{1}^{r}+e_{2}^{\prime} \sum_{r=0}^{\infty} g_{r} q_{2}^{r}
\end{aligned}
$$

where $q_{1}$ and $q_{2}$ are commutative quaternions given in Eq.(3.1).

## 4. Conclusion

Octonions have many applications in field. We mentioned some of them such as cosmology, quantum theory, etc. Octonions form a non-commutative and non-associative algebra. These hyper-complex numbers are constructed by Cayley-Dickson process over Hamilton quaternions. The aim of this study is to introduce commutative octonions constructed by Cayley-Dickson process over commutative quaternions. We think that several applications in field will be applied by other researchers. It will be very interesting that investigation of commutative octonions whose coefficients are well-known integer sequences such as Fibonacci, Lucas, Pell, etc.

## Article Information

Acknowledgements: The authors would like to express their sincere thanks to the editor and the anonymous reviewers for their helpful comments and suggestions.

Author's contributions: All authors contributed equally to the writing of this paper. All authors read and approved the final manuscript
Conflict of Interest Disclosure: No potential conflict of interest was declared by the author.
Copyright Statement: Authors own the copyright of their work published in the journal and their work is published under the CC BY-NC 4.0 license.

Supporting/Supporting Organizations: No grants were received from any public, private or non-profit organizations for this research.
Ethical Approval and Participant Consent: It is declared that during the preparation process of this study, scientific and ethical principles were followed and all the studies benefited from are stated in the bibliography

Plagiarism Statement: This article was scanned by the plagiarism program. No plagiarism detected.
Availability of data and materials: Not applicable.

## References

[1] Bilgici, G., Unal, Z., Tokeser, U. and Mert, T., On Fibonacci and Lucas generalized octonions, Ars Combinatoria, 138 (2018), 35-44
[2] Catoni, F., Cannata, R., Catoni, V. and Zampetti, P., N-dimensional geometries generated by hypercomplex numbers, Adv. Appl. Clifford Algebr., 15 (2005), 1-25.
[3] Catoni, F., Cannata, R. and Zampetti, P., An introduction to commutative quaternions, Adv. Appl. Clifford Algebr., 16 (2006), 1-28.
[4] Catoni, F., Boccalett, D., Cannata, R., Catoni, V., Nichelatti, E. and Zampetti, P., The Mathematics of Mikowski Space-Time and Introduction to Commutative Hypercomplex Numbers. Birkhauser-Verlag, Basel, 2008.
[5] Freedman, M., Shokrian-Zini, M. and Wang, Z., Quantum computing with octonions, Peking Math. J., 2(3) (2019), 239-273
[6] Gunaydin, M., Kallosh, R., Linde, A. and Yamada, Y., M-theory cosmology, octonions, error correcting codes, J. High Energ. Phys., 2021(1) (2021) 1-60.
[7] Hamilton, W.R., Lectures on Quaternions, Hodges and Smith, Dublin, 1853.
[8] Klco, P., Kollarik, M. and Tatar, M., Novel computer algorithm for cough monitoring based on octonions, Respiratory Physiology \& Neurobiology, 257 (2018), 36-41.
[9] Kosal, H.H. and Tosun, M., Commutative quaternion matrices, Adv. Appl. Clifford Algebr., 24 (2014), 769-779.
10] Okubo, S., Introduction to Octonion and Other Non-Associative Algebras in Physics, Cambridge University Press, London, 1995.
[11] Segre, C., The real representations of complex elements and extension to bicomplex system, Math. Ann., 40 (1892), 413-467.
[12] Singh, T.P., Octonions, trace dynamics and noncommutative geometry-A case for unification in spontaneous quantum gravity, Zeitschrift für Naturforschung A, 75(12) (2020), 1051-1062.
[13] Srivastava, G., Gupta, R. Kumar, R. and Le, D.N., Space-time code design using quaternions, octonions and other non-associative structures, International Journal of Electrical and Computer Engineering Systems 10(2) (2019), 91-95
[14] Tokeser, U., Mert, T., Unal, Z. and Bilgici, G., On Pell and Pell-Lucas generalized octonions, Turkish Journal of Mathematics and Computer Sciences 13(2) (2021), 226-233.
[15] Tokeser, U., Mert, T. and Dundar, Y., Some properties and Vajda theorems of split dual Fibonacci and split dual Lucas octonions, AIMS Math., 7(5) (2022), 8645-8653
[16] Weng, Z.H., Frequencies of astrophysical jets and gravitational strengths in the octonion spaces, International Journal of Modern Physics D., 31 (4) (2022), 2250024-1-2250024-16.

