



## SOME BOUNDS FOR THE $k$ -GENERALIZED DIGAMMA FUNCTION

Hesham MOUSTAFA<sup>1</sup>, Mansour MAHMOUD<sup>2</sup> and Ahmed TALAT<sup>3</sup>

<sup>1</sup>Mathematics Department, Mansoura University, Mansoura 35516, EGYPT

<sup>2</sup>Mathematics Department, King Abdulaziz University, P.O. Box 80203, Jeddah 21589, SAUDI ARABIA

<sup>3</sup>Mathematics and Computer Sciences Department, Port Said University, Port Said, EGYPT

ABSTRACT. We presented some monotonicity properties for the  $k$ -generalized digamma function  $\psi_k(h)$  and we established some new bounds for  $\psi_k^{(s)}(h), s \in \mathbb{N} \cup \{0\}$ , which refine recent results.

### 1. INTRODUCTION

The ordinary Gamma function is given by [1]:

$$\Gamma(h) = \lim_{s \rightarrow \infty} \frac{s! s^{h-1}}{h(h+1)(h+2) \cdots (h+(s-1))}, \quad h > 0$$

was discovered by Euler when he generalized the factorial function to non integer values. The digamma function is the logarithmic derivative of the ordinary gamma function and is given by [1]:

$$\psi(1+h) = -\gamma + \sum_{s=1}^{\infty} \frac{h}{s(h+s)}, \quad h > -1$$

where  $\gamma = \lim_{m \rightarrow \infty} \left( \sum_{s=1}^m \frac{1}{s} - \log m \right) \simeq 0.577$  is the Euler-Mascheroni constant. In 2006, Kirchoff applied the polygamma functions in the field of physics [3] and many series involving polygamma functions appeared in Feynman calculations [8]. In 2021, Wilkins and Hromadka [16] use the digamma function, as well as new

2020 *Mathematics Subject Classification.* 33B15, 26A48, 26D07.

*Keywords.* Gamma function, digamma function, polygamma function, completely monotonic function, asymptotic expansion, inequalities.

<sup>1</sup> heshammoustafa14@gmail.com-Corresponding author; 0000-0002-2792-6239

<sup>2</sup> mansour@mans.edu.eg; 0000-0002-5918-1913

<sup>3</sup> a\_t\_amer@yahoo.com; 0000-0001-7702-8093.

variants of the digamma function, as a new family of basis functions in mesh-free numerical methods for solving partial differential equations. polygamma functions are used to approximate the values of many special functions and have many applications in physics, statistics and applied mathematics [14].

Many mathematicians studied the completely monotonic (CM) of some functions including the digamma function to deduce some of its bounds. An infinitely differentiable function  $L(h)$  on  $\mathbb{R}^+$  is CM if  $(-1)^s L^{(s)}(h) \geq 0$  for  $s \in \mathbb{N} \cup \{0\}$ . A theorem [15, Theorem 12b] stated the sufficient condition for  $L(h)$  being CM on  $\mathbb{R}^+$  as:

$$L(h) = \int_0^\infty e^{-hy} dv(y),$$

where  $v(y)$  is non-decreasing and the integral converges for  $h \in \mathbb{R}^+$ .

In 2006, Muqattash and Yahdi [10] presented the following inequality:

$$\ln h < \psi(1+h) < \ln(1+h), \quad h \in \mathbb{R}^+. \quad (1)$$

In 2011, Batir [2] presented the following inequalities:

$$\ln \left( h^2 + h + e^{-2\gamma} \right) \leq 2\psi(h+1) < \ln \left( h^2 + h + \frac{1}{3} \right), \quad h \in [0, \infty) \quad (2)$$

$$\ln \left( \frac{2h+2}{e^{\frac{2}{h+1}} - 1} \right) < 2\psi(h+1) \leq \ln \left( \frac{2h + (e^2 - 1)e^{-2\gamma}}{e^{\frac{2}{1+h}} - 1} \right), \quad h \in [0, \infty) \quad (3)$$

and

$$\left( \frac{1+2h}{2} \right) e^{-2\psi(1+h)} < \psi'(1+h) < \left( \frac{\pi^2 e^{-2\gamma} + 6h}{6} \right) e^{-2\psi(h+1)}, \quad h \in (0, \infty). \quad (4)$$

In 2014, Guo and Qi [5] refined the inequality (1) by

$$\ln(h+1/2) < \psi(1+h) < \ln(e^{-\gamma} + h), \quad h \in \mathbb{R}^+. \quad (5)$$

Diaz and Pariguan [4] presented the  $k$ -generalized gamma function as:

$$\Gamma_k(h) = \lim_{s \rightarrow \infty} \frac{s! k^s (sk)^{\frac{h}{k}-1}}{h(k+h)(2k+h) \cdots ((s-1)k+h)}, \quad k, h \in \mathbb{R}^+.$$

Mansour [7] determined the  $\Gamma_k$  by a combination of some functional equations. The  $k$ -analogue of the digamma function is introduced by [11]

$$\psi_k(h) = \frac{-1}{k} (\gamma - \ln k) - 1/h - \sum_{s=1}^{\infty} \left( \frac{1}{sk+h} - \frac{1}{sk} \right), \quad k, h \in \mathbb{R}^+$$

and it has the following relations for  $h, k \in \mathbb{R}^+$  and  $s \in \mathbb{N} \cup \{0\}$

$$k\psi_k(h) - \psi\left(\frac{h}{k}\right) = \ln k, \quad \psi_k^{(s)}(k+h) = \frac{(-1)^s s!}{h^{s+1}} + \psi_k^{(s)}(h) \text{ and } \psi_k'(k) = \frac{\pi^2}{6k^2}. \quad (6)$$

In 2018, Nantomah, Nisar and Gehlot [12] introduced the following integral formulas:

$$\psi_k(h) = \int_0^\infty \left( \frac{2e^{-y} - e^{-ky}}{ky} - \frac{e^{-hy}}{1 - e^{-ky}} \right) dy, \quad h, k > 0 \quad (7)$$

and

$$\psi_k^{(s)}(h) = (-1)^{s+1} \int_0^\infty y^s \left( \frac{e^{-hy}}{1 - e^{-ky}} \right) dy, \quad h, k > 0; s \in \mathbb{N}. \quad (8)$$

Yin, Huag, Song and Dou [19] deduced the following inequality:

$$0 \leq \psi'_k(h) - \frac{1}{kh} \leq \frac{1}{h^2}, \quad h, k \in \mathbb{R}^+. \quad (9)$$

In 2020, Yildirim [17] deduced the following inequality:

$$-\frac{k}{12h^2} < \psi_k(h+k) - \frac{1}{k} \ln h - \frac{1}{2h} < 0, \quad h, k \in \mathbb{R}^+. \quad (10)$$

In 2021, Moustafa, Almuashi and Mahmoud [9] presented the following asymptotic formulas for  $k > 0$ :

$$\psi_k(h) \sim \frac{1}{k} \ln h - \frac{1}{2h} - \sum_{m=1}^{\infty} \frac{k^{2m-1} B_{2m}}{(2m) h^{2m}}, \quad h \rightarrow \infty \quad (11)$$

and for  $s \in \mathbb{N}$ ,

$$\psi_k^{(s)}(h) \sim \frac{(-1)^{s-1} (s-1)!}{kh^s} - \frac{(-1)^s s!}{2h^{s+1}} + (-1)^{s+1} \sum_{m=1}^{\infty} \frac{(s+2m-1)! k^{2m-1} B_{2m}}{(2m)! h^{2m+s}}, \quad h \rightarrow \infty \quad (12)$$

and they also deduced the inequalities:

$$\frac{1}{k} \ln h + \frac{1}{h} - \frac{k}{2} \psi'_k(h) < \psi_k(k+h) < \frac{\ln h}{k} + 1/h - \frac{k}{2} \psi'_k\left(\frac{k+3h}{3}\right), \quad h, k > 0 \quad (13)$$

where the upper bound of (13) refines upper bound of (10) for all  $h > \frac{k}{3}$ , and for  $s \in \mathbb{N}$ ,  $h, k \in \mathbb{R}^+$

$$\frac{(s-1)!}{kh^s} + \frac{(-1)^s k}{2} \psi_k^{(s+1)}\left(h + \frac{k}{3}\right) < (-1)^{s+1} \psi_k^{(s)}(h) < \frac{(s-1)!}{kh^s} + \frac{(-1)^s k}{2} \psi_k^{(s+1)}(h). \quad (14)$$

Notes: All of the  $k$ -digamma function results allow us to make new conclusions about the classical digamma function or new proofs for some of its established conclusions when  $k$  tends to one, and likewise for the  $k$ -gamma function [6, 18, 20]. For extra information about  $\Gamma_k$  and  $\psi_k$  functions, see [4, 7, 9, 19] and the related references therein.

We will introduce two CM functions involving  $\psi_k(h)$  and  $\psi'_k(h)$  functions. Some new bounds for  $\psi_k^{(s)}(h)$  functions ( $s \in \mathbb{N} \cup \{0\}$ ) will be deduced, which generalize and refine some recent results. Also, we will study the monotonicity of two functions

containing the  $k$ -generalized digamma function and consequently, we will deduce some new best bounds for  $\psi_k^{(s)}(h)$  functions ( $s \in \mathbb{N} \cup \{0\}$ ).

2. AUXILIARY RESULTS

In [13], the following corollary was introduced:

**Corollary 1.** *Assume that  $S$  is a function defined on  $h > h_0$ ,  $h_0 \in \mathbb{R}$  with  $\lim_{h \rightarrow \infty} S(h) = 0$ . Then for  $\omega \in \mathbb{R}^+$ ,  $S(h) > 0$ , if  $S(h + \omega) < S(h)$  for  $h > h_0$  and  $S(h) < 0$ , if  $S(h + \omega) > S(h)$  for  $h > h_0$ .*

Using the monotonicity properties, we can conclude the following results:

**Lemma 1.**

$$\ln\left(\frac{h^2 + 3h + 3}{3}\right) < h, \quad \forall h \in \mathbb{R}^+, \tag{15}$$

$$\ln\left(\frac{1}{e^{-\gamma} + h} + 1\right) < \frac{1}{1+h}, \quad \forall h > \frac{e^{2\gamma} - e^\gamma - 1}{e^\gamma(2 - e^\gamma)} \simeq 1.00313 \tag{16}$$

and

$$\frac{2}{1+h} < \ln\left(\frac{2(1+h)}{h^2 + h + e^{-2\gamma}} + 1\right), \quad \forall h > \frac{1}{\sqrt{e^{2\gamma}(-3 + e^{2\gamma})}} - 1 \simeq 0.352938. \tag{17}$$

*Proof.* Let the function  $L(h) = \ln\left(\frac{h^2+3h+3}{3}\right) - h$  and then  $L'(h) = \frac{-h(1+h)}{3+3h+h^2} < 0$  for all  $h > 0$  and then  $L(h)$  is decreasing on  $(0, \infty)$  with  $\lim_{h \rightarrow 0^+} L(h) = 0$  and then  $L(h) < 0$  for all  $h > 0$  which proves (15). Secondly, we let the function  $C(h) = \ln\left(\frac{1}{e^{-\gamma}+h} + 1\right) - \frac{1}{1+h}$ . Then

$$C'(h) = \frac{1 + e^\gamma - e^{2\gamma} + e^\gamma(2 - e^\gamma)h}{(1+h)^2(1 + e^\gamma h)(1 + e^\gamma(1+h))} > 0, \quad h > h_1 = \frac{-1 - e^\gamma + e^{2\gamma}}{e^\gamma(2 - e^\gamma)} \simeq 1.00313.$$

Then  $C(h)$  is increasing on  $(h_1, \infty)$  with  $\lim_{h \rightarrow \infty} C(h) = 0$  and this proves (16). By the same way, we obtain (17). □

**Lemma 2.** *For  $k \in \mathbb{R}^+$ , we have*

$$k\psi_k(k+h) < \gamma + \ln\left(\frac{3}{\pi^2}\right) + \ln(2h+k), \quad \forall h > \frac{k}{2}. \tag{18}$$

*Proof.* Let the function  $N_k(h) = \ln\left(\frac{3}{\pi^2}\right) + \gamma + \ln(2h + k) - k\psi_k(k + h)$ . Then  $N'_k(h) = \frac{2}{k+2h} - k\psi'_k(k + h)$  and by using (6), we obtain

$$N'_k(h) - N'_k(k + h) = \frac{k^3}{(3k + 2h)(k + 2h)(k + h)^2} > 0, \quad h, k > 0.$$

Using the asymptotic formula (12), we have  $\lim_{h \rightarrow \infty} N'_k(h) = 0$  and then Corollary 1 gives us that  $N'_k(h) > 0$  for  $h > 0$  and  $k > 0$ . Then, we have  $N_k(h)$  is increasing on  $\mathbb{R}^+$  and by using (6) again, we get  $N_k(h) > N_k\left(\frac{k}{2}\right) \simeq 0.0430254 > 0$  for all  $h > \frac{k}{2}$  and  $k > 0$ .  $\square$

**Lemma 3.** For  $k > 0$ , we have

$$e^{k\psi_k(h+2k)} < e^{k\psi_k(h+k)} + k, \quad \forall h > 0. \tag{19}$$

*Proof.* Set the function  $B_k(h) = e^{k\psi_k(h+2k)} - e^{k\psi_k(h+k)} - k$ . Then by using (6), we get

$$\frac{B_k(h+k) - B_k(h)}{e^{k\psi_k(h+k)}} = e^{\frac{k}{h+k} + \frac{k}{h+2k}} - 2e^{\frac{k}{h+k}} + 1 \doteq D_k(h).$$

Then

$$\frac{(h+k)^2}{ke^{\frac{k}{h+k}}} D'_k(h) = 2 - \frac{(5k^2 + 6kh + 2h^2)e^{\frac{k}{h+2k}}}{(2k+h)^2} \doteq f_k(h).$$

Then  $f'_k(h) = \frac{k^3 e^{\frac{k}{h+2k}}}{(2k+h)^4} > 0$  for all  $h, k > 0$  and hence  $f_k(h)$  is increasing on  $\mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} f_k(h) = 0$ . Then  $f_k(h) < 0$  for  $h, k > 0$  and then  $D_k(h)$  is decreasing on  $\mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} D_k(h) = 0$ . Then  $B_k(h+k) - B_k(h) > 0$  for  $h, k > 0$ . Using the asymptotic formula (11), we have  $\lim_{h \rightarrow \infty} B_k(h) = 0$  and then Corollary 1 gives us that  $B_k(h) < 0$  for all  $h, k > 0$ .  $\square$

**Lemma 4.** For  $k > 0$ , we have

$$e^{2k\psi_k(h+2k)} > e^{2k\psi_k(h+k)} + 2k(h+k), \quad \forall h > 0. \tag{20}$$

*Proof.* Set the function  $m_k(h) = e^{2k\psi_k(h+2k)} - e^{2k\psi_k(h+k)} - 2k(h+k)$ . Then

$$m_k(h+k) - m_k(h) = e^{2k\psi_k(h+3k)} - 2e^{2k\psi_k(h+2k)} + e^{2k\psi_k(h+k)} - 2k^2 \doteq t_k(h).$$

Then by using (6), we get

$$\frac{t_k(h+k) - t_k(h)}{e^{2k\psi_k(h+k)}} = e^{\frac{2k}{h+k} + \frac{2k}{h+2k} + \frac{2k}{h+3k}} - 3e^{\frac{2k}{h+k} + \frac{2k}{h+2k}} + 3e^{\frac{2k}{h+k}} - 1 \doteq s_k(h).$$

Then

$$\begin{aligned} \frac{(h+k)^2 s'_k(h)}{2ke^{\frac{2k}{h+k}}} &= -\frac{(49k^4 + 96k^3h + 72k^2h^2 + 24kh^3 + 3h^4)e^{\frac{2k}{h+2k} + \frac{2k}{h+3k}}}{(2k+h)^2(3k+h)^2} \\ &+ \frac{3(5k^2 + 6kh + 2h^2)}{(2k+h)^2 e^{\frac{-2k}{h+2k}}} - 3 \doteq u_k(h) \end{aligned}$$

Then

$$\frac{(2k + h)^4 u'_k(h)}{2k(3k^2 + 3kh + h^2)e^{\frac{2k}{h+2k}}} = \frac{379k^6 + 959k^5h + 1049k^4h^2 + 626k^3h^3 + 213k^2h^4}{(3k^2 + 3kh + h^2)(3k + h)^4 e^{\frac{-2k}{h+3k}}} + \frac{39kh^5 + 3h^6}{(3k^2 + 3kh + h^2)(3k + h)^4 e^{\frac{-2k}{h+3k}}} - 3 \doteq w_k(h).$$

Then

$$w'_k(h) = \frac{-2k^5 \left( 129k^4 + 240k^3h + 172k^2h^2 + 56kh^3 + 7h^4 \right) e^{\frac{2k}{h+3k}}}{(3k + h)^6 (3k^2 + 3kh + h^2)^2} < 0, \quad h, k > 0$$

and hence  $w_k(h)$  is decreasing on  $(0, \infty)$  with  $\lim_{h \rightarrow \infty} w_k(h) = 0$ . Then  $w_k(h) > 0$  for  $h, k \in \mathbb{R}^+$  and then  $u_k(h)$  is increasing on  $\mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} u_k(h) = 0$ . Then  $s_k(h)$  is decreasing on  $\mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} s_k(h) = 0$ . Then  $t_k(h + k) - t_k(h) > 0$  for  $h, k \in \mathbb{R}^+$ . Using the asymptotic formula (11), we have  $\lim_{h \rightarrow \infty} t_k(h) = 0$  and then Corollary 1 gives us that  $t_k(h) < 0$  for all  $h, k \in \mathbb{R}^+$ . Then  $m_k(h + k) - m_k(h) < 0$  for  $h, k \in \mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} m_k(h) = 0$  and then  $m_k(h) > 0$  for all  $h, k \in \mathbb{R}^+$ .  $\square$

### 3. SOME CM MONOTONIC FUNCTIONS

**Theorem 1.** Assume that  $h, k > 0$ . Then the function

$$U_{\beta,k}(h) = \psi'_k(h) - \frac{2}{kh} + \frac{2}{k^2} \ln \left( 1 + \frac{\beta k}{h} \right)$$

is CM on  $\mathbb{R}^+$  if and only if  $\beta \geq \frac{1}{2}$ .

*Proof.*

$$U'_{\beta,k}(h) = \psi''_k(h) + \frac{2}{kh^2} + \frac{2}{k^2} \left( \frac{1}{h + \beta k} - \frac{1}{h} \right)$$

and by using (8) and the identity  $\frac{1}{h^l} = \frac{1}{(l-1)!} \int_0^\infty y^{l-1} e^{-hy} dy$  for  $h > 0$ , (see [1]), we have

$$U'_{\beta,k}(h) = \int_0^\infty \frac{2e^{-hy}}{k^2(e^{ky} - 1)} \phi_k(y) dy,$$

where

$$\phi_k(y) = (e^{ky} - 1)(e^{-\beta ky} - 1) + ky(e^{ky} - 1) - \frac{(ky)^2}{2} e^{ky}.$$

Let  $\beta \geq \frac{1}{2}$ . Then

$$\begin{aligned} e^{\frac{ky}{2}} \phi_k(y) &\leq e^{ky} - e^{\frac{3ky}{2}} + e^{\frac{ky}{2}} - 1 + ky \left( e^{\frac{3ky}{2}} - e^{\frac{ky}{2}} \right) - 1/2(ky)^2 e^{\frac{3ky}{2}} \\ &= \sum_{m=1}^\infty \frac{n(m)}{2^{m+2} (m+2)!} (ky)^{m+2} \end{aligned}$$

where

$$\begin{aligned} n(m) &= 2^{m+2} - 3^{m+2} + 1 + 2(m+2)(3^{m+1} - 1) - 2(m+2)(m+1)3^m \\ &= -2(m+2)\left((m-1)2^m + 1\right) - \sum_{s=1}^m \frac{(m+2)(m+1)(m-s)}{(m+2-s)} \binom{m}{s} 2^{m+1-s} \\ &< 0 \end{aligned}$$

and then  $-U'_{\beta,k}(h)$  is CM on  $\mathbb{R}^+$  and hence  $U_{\beta,k}(h)$  is decreasing on  $\mathbb{R}^+$ . Using the asymptotic formula (12), we have  $\lim_{h \rightarrow \infty} U_{\beta,k}(h) = 0$  and then  $U_{\beta,k}(h) > 0$ . Then  $U_{\beta,k}(h)$  is CM on  $\mathbb{R}^+$  for  $\beta \geq \frac{1}{2}$ . On the other side, if  $U_{\beta,k}(h)$  is CM, then by using again the asymptotic formula (12), we get  $\lim_{h \rightarrow \infty} h U_{\beta,k}(h) = \frac{2\beta-1}{k} \geq 0$  and hence  $\beta \geq \frac{1}{2}$ . □

**Theorem 2.** Assume that  $h, k > 0$  and  $\lambda \in \mathbb{R}$ . Then the function

$$F_{\lambda,k}(h) = \psi_k(h+k) - \frac{1}{k} \ln(h+\lambda k)$$

is CM on  $\mathbb{R}^+$  if and only if  $\lambda \leq \frac{1}{2}$ . Also, the function  $-F_{\lambda,k}(h)$  is CM on  $\mathbb{R}^+$  if  $\lambda \geq 1$ .

*Proof.*

$$F'_{\lambda,k}(h) = -\frac{1}{k(h+\lambda k)} + \psi'_k(h+k) = \int_0^\infty \frac{e^{-hy}}{k(e^{ky}-1)} \varphi_k(y) dy,$$

where

$$\varphi_k(y) = ky - e^{-\lambda ky} (e^{ky} - 1).$$

Let  $\lambda \leq \frac{1}{2}$ , then we obtain

$$\begin{aligned} e^{\frac{ky}{2}} \varphi_k(y) &\leq 1 + ky e^{\frac{ky}{2}} - e^{ky} \\ &= -\sum_{l=2}^\infty \frac{(2^l - l - 1)(ky)^{1+l}}{2^l (1+l)!} \\ &= -\sum_{l=2}^\infty \frac{\left(\sum_{s=2}^l \binom{l}{s}\right) (ky)^{1+l}}{2^l (1+l)!} < 0 \end{aligned}$$

and consequently,  $-F'_{\lambda,k}(h)$  is CM on  $\mathbb{R}^+$  for  $\lambda \leq \frac{1}{2}$  and hence  $F_{\lambda,k}(h)$  is decreasing on  $\mathbb{R}^+$ . Using the asymptotic formula (11), we obtain  $\lim_{h \rightarrow \infty} F_{\lambda,k}(h) = 0$  and then  $F_{\lambda,k}(h) > 0$ . Hence  $F_{\lambda,k}(h)$  is CM on  $\mathbb{R}^+$  for  $\lambda \leq \frac{1}{2}$ . On the other hand, if  $F_{\lambda,k}(h)$  is CM, then by using again the asymptotic formula (11), we obtain  $\lim_{h \rightarrow \infty} h F_{\lambda,k}(h) =$

$\frac{1}{2} - \lambda \geq 0$  and then  $\lambda \leq \frac{1}{2}$ . Now for  $\lambda \geq 1$ , we have  $e^{ky} \varphi_k(y) \geq \sum_{l=1}^\infty \frac{l(ky)^{l+1}}{(l+1)!} > 0$

and consequently,  $F'_{\lambda,k}(h)$  is CM on  $\mathbb{R}^+$  for  $\lambda \geq 1$  and hence  $F_{\lambda,k}(h)$  is increasing on  $\mathbb{R}^+$  with  $\lim_{h \rightarrow \infty} F_{\lambda,k}(h) = 0$  and then  $F_{\lambda,k}(h) < 0$ . Then  $-F_{\lambda,k}(h)$  is CM on  $\mathbb{R}^+$  for  $\lambda \geq 1$ . □

4. SOME INEQUALITIES FOR THE  $\psi_k$  AND  $\psi_k^{(s)}$  FUNCTIONS

Let us mention some important consequences of Theorems 1 and 2.

**Corollary 2.** *Let  $a \in (0, \infty)$ . Then we have*

$$\frac{1}{kh} - \frac{1}{k^2} \ln\left(\frac{ak+h}{h}\right) < \frac{\psi'_k(h)}{2}, \quad k, h \in \mathbb{R}^+ \tag{21}$$

with the best possible constant  $a = \frac{1}{2}$ .

*Proof.* The inequality (21) at  $a = \frac{1}{2}$  follows from  $U_{\frac{1}{2},k}(h) > 0$  in Theorem 1 and the inequality (21) is equivalent that  $h U_{a,k}(h) > 0$  which yields  $a \geq \frac{1}{2}$  as stated when we proved Theorem 1. Then  $a = \frac{1}{2}$  is the best in (21), since the logarithmic function is strictly increasing on  $\mathbb{R}^+$ . □

**Remark 1.** *Using the identity  $\ln(1+h) < h$  for all  $h > -1$ , (see [1]) yields the lower bound of (21) refines the lower bound of (9) for all  $h, k > 0$ .*

**Corollary 3.** *Let  $a \in (0, \infty)$  and  $s = 1, 2, 3, \dots$ . Then we have*

$$\frac{2s!}{kh^{s+1}} + \frac{2(s-1)!}{k^2} \left( \frac{1}{(h+ak)^s} - \frac{1}{h^s} \right) < (-1)^s \psi_k^{(s+1)}(h), \quad h, k \in (0, \infty) \tag{22}$$

with the best possible constant  $a = \frac{1}{2}$ .

*Proof.* The inequality (22) at  $a = \frac{1}{2}$  follows from  $(-1)^s U_{\frac{1}{2},k}^{(s)}(h) > 0$  in Theorem 1 and the inequality (22) is equivalent that  $h^{s+1} (-1)^s U_{a,k}^{(s)}(h) > 0$ . Using the asymptotic expansion (12), we have  $\lim_{h \rightarrow \infty} h^{s+1} (-1)^s U_{a,k}^{(s)}(h) = \frac{s!}{k} (2a-1) \geq 0$  and hence  $a \geq \frac{1}{2}$ . Using the decreasing property of the function  $\frac{1}{h^s}$  on  $(0, \infty)$  for  $s = 1, 2, 3, \dots$ , we deduce that  $a = \frac{1}{2}$  is the best possible constant in (22). □

**Corollary 4.** *Let  $a \in [0, \infty)$ . Then we have*

$$\ln(h+ak) < k\psi_k(k+h) < \ln(k+h), \quad k, h \in \mathbb{R}^+ \tag{23}$$

with the best possible constant  $a = \frac{1}{2}$ .

*Proof.* The inequality (23) at  $a = \frac{1}{2}$  is deduced from  $F_{\frac{1}{2},k}(h) > 0$  and  $F_{1,k}(h) < 0$  in Theorem 2. The left-hand side of (23) is equivalent that  $h F_{a,k}(h) > 0$  and this gives  $a \leq \frac{1}{2}$  as stated when we proved Theorem 2. Then  $a = \frac{1}{2}$  is the best in (23). □

**Remark 2.** • *Letting  $k = 1$  and  $a = 0$  in (23), we obtain (1).*



- Using (21), we deduce that the lower bound of (23) refines the lower bound of (13) for every  $k, h \in \mathbb{R}^+$ .

**Corollary 5.** Let  $a \in [0, \infty)$  and  $s = 1, 2, 3, \dots$ . Then we have

$$\frac{s!}{h^{s+1}} + \frac{(s-1)!}{k(h+k)^s} < (-1)^{s+1} \psi_k^{(s)}(h) < \frac{s!}{h^{s+1}} + \frac{(s-1)!}{k(h+ak)^s}, \quad h, k \in (0, \infty) \quad (24)$$

with the best possible constant  $a = \frac{1}{2}$ .

*Proof.* The inequality (24) at  $a = \frac{1}{2}$  is deduced from  $(-1)^s F_{\frac{1}{2}, k}^{(s)}(h) > 0$  and  $(-1)^s F_{1, k}^{(s)}(h) < 0$  in Theorem 2. The right-hand side of (24) is equivalent that  $h^{1+s} (-1)^s F_{a, k}^{(s)}(h) > 0$ . Using the asymptotic expansion (12), we have

$$\lim_{h \rightarrow \infty} h^{s+1} (-1)^s F_{a, k}^{(s)}(h) = s! \left( \frac{1}{2} - a \right) \geq 0$$

and hence  $a \leq \frac{1}{2}$ . Then  $a = \frac{1}{2}$  is the best possible constant in (24). □

**Remark 3.** Using (22), we deduce that the upper bound of (24) refines the upper bound of (14) for every  $s \in \mathbb{N}$  and  $h, k > 0$ .

**Lemma 5.** For  $k > 0$ , the function

$$T_k(h) = e^{k\psi_k(h+k)} - h \quad (25)$$

is strictly decreasing convex on  $(-k, \infty)$  with  $\lim_{h \rightarrow \infty} T_k(h) = \frac{k}{2}$  and  $\lim_{h \rightarrow 0} T_k(h) = ke^{-\gamma}$ .

*Proof.* Using (6), we have  $\lim_{h \rightarrow 0} T_k(h) = ke^{-\gamma}$ . Differentiating (25) yields

$$T'_k(h) = -1 + k\psi'_k(h+k)e^{k\psi_k(h+k)}$$

and

$$\frac{T''_k(h)}{k e^{k\psi_k(h+k)}} = k \left[ \psi'_k(h+k) \right]^2 + \psi''_k(h+k) \doteq S_k(h).$$

Applying (6), we get

$$\frac{(k+h)^2}{2k} \left[ S_k(k+h) - S_k(h) \right] = -\psi'_k(k+h) - \frac{2h^2 + 4kh + 2k^2 - 1}{2(h+k)^2} \doteq Q_k(h).$$

Applying (6) again, we get

$$Q_k(k+h) = Q_k(h) + \frac{A_k(k+h)}{2(k+h)^2(2k+h)^2},$$

where

$$A_k(h) = k^2 + 2kh + 2h^2 > 0, \quad h, k \in \mathbb{R}^+$$

and then  $Q_k(k+h) > Q_k(h)$  for all  $h > -k$  and by using the asymptotic formula (12), we have  $\lim_{h \rightarrow \infty} Q_k(h) = 0$  and then Corollary 1 gives us  $Q_k(h) < 0$  for every  $h > -k$ . Consequently, we have  $S_k(k+h) < S_k(h)$  for all  $h > -k$  and by using the

asymptotic expansion (12), we have  $\lim_{h \rightarrow \infty} S_k(h) = 0$  and then  $T_k''(h) > 0$  for every  $h > -k$ . Then  $T_k'(h)$  is strictly increasing on  $(-k, \infty)$ . By using the asymptotic formulas (11) and (12), we have

$$\lim_{h \rightarrow \infty} T_k'(h) = 0 \text{ and } \lim_{h \rightarrow \infty} T_k(h) = k/2.$$

Then  $T_k'(h) < \lim_{h \rightarrow \infty} T_k'(h) = 0$  and this finishes the proof. □

And consequently, we have the following Corollary:

**Corollary 6.** *Set  $a$  and  $b$  be positive real numbers. Then we have*

$$\ln(h + ak) < k\psi_k(h + k) < \ln(h + bk), \quad h, k \in \mathbb{R}^+ \tag{26}$$

where  $a = \frac{1}{2}$  and  $b = \frac{1}{e^\gamma} \simeq 0.56$  being the best.

**Remark 4.** • Letting  $k = 1$  in (26), we obtain (5).

• The upper bound of (26) refines the upper bound of (23) for all  $h, k > 0$ .

**Lemma 6.** *For  $h \geq 0$  and  $k \in \mathbb{R}^+$ ,*

$$\ln\left(\frac{c k}{e^{\frac{k}{k+h}} - 1}\right) \leq k\psi_k(k + h) < \ln\left(\frac{d k}{e^{\frac{k}{k+h}} - 1}\right), \tag{27}$$

where the constants  $c = e^{-\gamma}(e - 1) \simeq 0.965$  and  $d = 1$  are the best possible.

*Proof.* Set

$$f_k(h) = T_k(k + h) - T_k(h) = -k + e^{k\psi_k(k+h)}\left(e^{\frac{k}{k+h}} - 1\right), \quad h \geq 0 \text{ and } k > 0.$$

Since  $T_k'(h)$  is strictly increasing on  $(-k, \infty)$ , then  $f_k(h)$  is strictly increasing on  $[0, \infty)$  and by using (6) and the asymptotic expansion (11), we get

$$f_k(0) = -k + ke^{-\gamma}(e - 1) \leq -k + e^{k\psi_k(k+h)}\left(e^{\frac{k}{k+h}} - 1\right) < \lim_{h \rightarrow \infty} f_k(h) = 0$$

and this gives (27). □

**Remark 5.** *Using (16), we deduce that the upper bound of (27) refines the upper bound of (26) for all  $h > 1.00313k$  and  $k > 0$ .*

**Lemma 7.** *For  $h > 0$  and  $k > 0$ , we have*

$$g e^{-k\psi_k(h+k)} < k\psi'_k(k + h) < r e^{-k\psi_k(k+h)}, \tag{28}$$

where the constants  $g = \frac{\pi^2 e^{-\gamma}}{6} \simeq 0.924$  and  $r = 1$  are the best possible.

*Proof.* By using the increasing property of  $T_k'(h)$  on  $(-k, \infty)$ , we have

$$T_k'(0) = -1 + k\psi'_k(k)e^{k\psi_k(k)} < -1 + k\psi'_k(k + h)e^{k\psi_k(k+h)} < \lim_{h \rightarrow \infty} T_k'(h) = 0.$$

Using (6) yields

$$\frac{\pi^2 e^{-\gamma}}{6} < k\psi'_k(k + h)e^{k\psi_k(k+h)} < 1,$$

which finishes the proof. □

**Remark 6.** Using (23) yields  $e^{-k\psi_k(k+h)} < \frac{2}{2h+k}$  for every  $h, k > 0$  and then the upper bound of (28) refines the upper bound of (24) at  $s = 1$  for all  $h, k > 0$ .

**Lemma 8.** For  $h > 0$  and  $k > 0$ ,

$$1 - e^{-\frac{k}{h+k}} + \frac{k^2}{(h+k)^2} < k^2\psi'_k(h+k) < e^{\frac{k}{h+k}} - 1. \tag{29}$$

*Proof.* By applying the mean value theorem to  $T_k$  on the interval  $[h, h+k]$ , we obtain

$$\frac{-T_k(h) + T_k(k+h)}{k} = T'_k(h + \alpha_h), \quad 0 < \alpha_h < k.$$

By using the increasing property of  $T'_k(h)$  on  $(-k, \infty)$ , we obtain

$$T'_k(h) < T'_k(h + \alpha_h) < T'_k(k+h), \quad 0 < \alpha_h < k.$$

Combining the last two relations yields

$$T'_k(h) < \frac{-T_k(h) + T_k(k+h)}{k} < T'_k(k+h)$$

and this gives us (29). □

**Remark 7.** Using (19), we deduce that the upper bound of (29) refines the upper bound of (28) for every  $h, k \in \mathbb{R}^+$ .

**Lemma 9.** For  $k > 0$ , the function

$$W_k(h) = e^{2k\psi_k(h+k)} - h^2 - hk \tag{30}$$

is strictly increasing concave in  $(-k, \infty)$  with  $\lim_{h \rightarrow \infty} W_k(h) = \frac{k^2}{3}$  and  $\lim_{h \rightarrow 0} W_k(h) = k^2e^{-2\gamma}$ .

*Proof.* Using (6), we have  $\lim_{h \rightarrow 0} W_k(h) = k^2e^{-2\gamma}$ . Differentiating (30) yields

$$W'_k(h) = -2h - k + 2k\psi'_k(h+k)e^{2k\psi_k(k+h)},$$

$$\frac{1}{2}W''_k(h) = -1 + ke^{2k\psi_k(k+h)} \left[ \psi''_k(k+h) + 2k(\psi'_k(k+h))^2 \right]$$

and

$$\frac{1}{2ke^{2k\psi_k(k+h)}}W'''_k(h) = \psi'''_k(k+h) + 6k\psi'_k(k+h)\psi''_k(k+h) + 4k^2(\psi'_k(k+h))^3 \doteq V_k(h).$$

Applying (6), we get

$$\begin{aligned} \frac{(h+k)^2}{2k} [V_k(k+h) - V_k(h)] &= -3\psi''_k(k+h) + \frac{6(h+2k)}{(h+k)^2}\psi'_k(h+k) \\ &\quad - 6k(\psi'_k(h+k))^2 - \frac{11k^2 + 12kh + 3h^2}{k(h+k)^4} \doteq U_k(h). \end{aligned}$$

Similarly, we get

$$\begin{aligned} \frac{(h+k)^2(h+2k)^2}{6k(3k^2+3kh+h^2)} [U_k(h+k) - U_k(h)] &= \psi'_k(h+k) \\ &- \frac{114k^5 + 298k^4h + 321k^3h^2}{6k(k+h)^2(2k+h)^2(3k^2+3kh+h^2)} \\ &- \frac{178k^2h^3 + 51kh^4 + 6h^5}{6k(k+h)^2(2k+h)^2(3k^2+3kh+h^2)} \\ &\doteq H_k(h). \end{aligned}$$

And finally, we get

$$H_k(k+h) - H_k(h) = - \frac{k^4 P_k(k+h)}{3(k+h)^2(2k+h)^2(3k+h)^2(3k^2+3kh+h^2)(7k^2+5kh+h^2)}$$

where

$$P_k(h) = 12k^4 + 36k^3h + 46k^2h^2 + 28kh^3 + 7h^4 > 0, \quad h, k \in \mathbb{R}^+$$

and then  $H_k(h+k) < H_k(h)$  for all  $h > -k$  and by using the asymptotic formula (12), we have  $\lim_{h \rightarrow \infty} H_k(h) = 0$  and then Corollary 1 gives us  $H_k(h) > 0$  for every  $h > -k$ . Consequently, we obtain  $U_k(h+k) > U_k(h)$  for all  $h > -k$  with  $\lim_{h \rightarrow \infty} U_k(h) = 0$  and then  $U_k(h) < 0$  for every  $h > -k$  and similarly, we get  $V_k(h) > 0$  for all  $h > -k$ . Then  $W'_k(h)$  is strictly increasing on  $(-k, \infty)$ . By using the asymptotic formulas (11) and (12), we have

$$\lim_{h \rightarrow \infty} W''_k(h) = \lim_{h \rightarrow \infty} W'_k(h) = 0 \text{ and } \lim_{h \rightarrow \infty} W_k(h) = \frac{k^2}{3}.$$

Then  $W''_k(h) < 0$  for all  $h > -k$  and then  $W'_k(h)$  is strictly decreasing on  $(-k, \infty)$ . Hence  $W'_k(h) > \lim_{h \rightarrow \infty} W'_k(h) = 0$  and this completes the proof.  $\square$

And consequently, we have the following Corollary:

**Corollary 7.** *Set  $a, b \in \mathbb{R}^+$  and  $k > 0$ . Then we have*

$$\frac{1}{2k} \ln(h^2 + hk + ak^2) \leq \psi_k(h+k) < \frac{1}{2k} \ln(h^2 + hk + bk^2), \quad h \in [0, \infty) \quad (31)$$

where the constants  $a = e^{-2\gamma} \simeq 0.315$  and  $b = \frac{1}{3}$  are the best possible.

**Remark 8.**  $\bullet$  Putting  $k = 1$  in (31) yields (2).

- $\bullet$  Using (15), we deduce that the upper bound of (31) refines the upper bound of (10) for  $h, k > 0$ .
- $\bullet$  For  $k > 0$ , the upper and lower bounds of (31) refine the upper and lower bounds of (26) for  $h > \left(\frac{\frac{1}{3} - e^{-2\gamma}}{2e^{-\gamma} - 1}\right)k \simeq 0.147224 k$  and  $h > 0$  respectively.

**Lemma 10.** For  $h \geq 0$  and  $k > 0$ , we have

$$\frac{1}{2k} \ln \left( \frac{2hk + ck^2}{e^{\frac{2k}{h+k}} - 1} \right) < \psi_k(h+k) \leq \frac{1}{2k} \ln \left( \frac{2hk + dk^2}{e^{\frac{2k}{h+k}} - 1} \right), \tag{32}$$

where  $c = 2$  and  $d = e^{-2\gamma}(e^2 - 1) \simeq 2.014$  are the best possible.

*Proof.* Set

$$M_k(h) = W_k(k+h) - W_k(h) = e^{2k\psi_k(k+h)} \left( e^{\frac{2k}{h+k}} - 1 \right) - 2hk - 2k^2, \quad h \geq 0, \quad k > 0.$$

Since  $W'_k(h)$  is strictly decreasing on  $(-k, \infty)$ , then  $M_k(h)$  is strictly decreasing on  $[0, \infty)$  and by using (6) and the asymptotic expansion (11), we get

$$M_k(0) = k^2 e^{-2\gamma}(e^2 - 1) - 2k^2 \geq e^{2k\psi_k(h+k)} \left( e^{\frac{2k}{h+k}} - 1 \right) - 2hk - 2k^2 > \lim_{h \rightarrow \infty} M_k(h) = 0$$

and this gives (32). □

**Remark 9.** • Letting  $k = 1$  in (32), we obtain (3).

- Using (17), we deduce that the lower bound of (32) refines the lower bound of (31) for  $h > 0.352938k$ .

**Lemma 11.** For  $h > 0$  and  $k > 0$ , we have

$$\left( \frac{h}{k} + a \right) e^{-2k\psi_k(h+k)} < \psi'_k(h+k) < \left( \frac{h}{k} + b \right) e^{-2k\psi_k(h+k)}, \tag{33}$$

where the constants  $a = \frac{1}{2}$  and  $b = \frac{\pi^2 e^{-2\gamma}}{6} \simeq 0.519$  are the best possible.

*Proof.* Using the decreasing property of  $W'_k(h)$  on  $(-k, \infty)$  yields

$$W'_k(0) > 2k\psi'_k(h+k)e^{2k\psi_k(h+k)} - 2h - k > \lim_{h \rightarrow \infty} W'_k(h) = 0.$$

Using (6), we have

$$\frac{\pi^2 e^{-2\gamma} k}{3} > 2k\psi'_k(h+k)e^{2k\psi_k(h+k)} - 2h > k,$$

which finishes the proof. □

**Remark 10.** • Putting  $k = 1$  in (33) gives (4).

- Using (18), we deduce that the lower bound of (33) refine the lower bound of (28) for  $h > \frac{k}{2}$  and  $k > 0$ .

**Lemma 12.** For  $h > 0$  and  $k > 0$ ,

$$\begin{aligned} \frac{1}{2k^2} \left( e^{\frac{2k}{h+k}} - 1 - k^2 e^{-2k\psi_k(h+k)} \right) &< \psi'_k(h+k) \\ &< \frac{1}{2k^2} \left( \frac{2k^2}{(h+k)^2} + 1 - e^{-\frac{2k}{h+k}} + k^2 e^{-2k\psi_k(h+2k)} \right). \end{aligned} \tag{34}$$

*Proof.* By applying the mean value theorem to  $W_k$  on the interval  $[h, h+k]$ , we get

$$\frac{W_k(h+k) - W_k(h)}{k} = W'_k(h + \beta_h), \quad 0 < \beta_h < k.$$

Using the decreasing property of  $W'_k(h)$  on  $(-k, \infty)$  yields

$$W'_k(h+k) < \frac{W_k(h+k) - W_k(h)}{k} < W'_k(h)$$

and this gives us (34).  $\square$

**Remark 11.** Using (20), we deduce that the lower bound of (34) refine the lower bound of (33) for  $h, k \in \mathbb{R}^+$ .

## 5. CONCLUSION

The main conclusions of this paper are stated in Theorems 1 and 2 and Lemmas 5 and 9. The authors proved the CM and the monotonicity properties of four functions containing the  $k$ -generalized digamma and polygamma functions, derived some new bounds for  $\psi_k^{(s)}(h)$  functions ( $s \in \mathbb{N} \cup \{0\}$ ). These bounds refine some recent results.

**Author Contribution Statements** The authors contributed equally to this article.

**Declaration of Competing Interests** The authors declare that they have no competing interest.

## REFERENCES

- [1] Abramowitz, M., Stegun, I. A., Handbook of Mathematical Functions with Formulas, Graphs and Mathematical Tables, Dover, New York, 1965.
- [2] Batir, N., Sharp bounds for the psi function and harmonic numbers, *Math. Inequal. Appl.*, 14(4) (2011), 917-925. <http://files.ele-math.com/abstracts/mia-14-77-abs.pdf>
- [3] Coffey, M. W., One integral in three ways: moments of a quantum distribution, *J. Phys. A: Math. Gen.*, 39 (2006), 1425-1431. <https://doi.org/10.1088/0305-4470/39/6/015>
- [4] Díaz, R., Pariguan, E., On hypergeometric functions and  $k$ -Pochhammer symbol, *Divulg. Mat.*, 15(2) (2007), 179-192. <https://doi.org/10.48550/arXiv.math/0405596>
- [5] Guo, B.-N., Qi, F., Sharp inequalities for the psi function and harmonic numbers, *Analysis*, 34(2) (2014), 201-208. DOI 10.1515/anly-2014-0001.
- [6] Kokologiannaki, C. G., Krasniqi, V., Some properties of the  $k$ -gamma function, *Le Matematiche*, 68(1) (2013), 13-22. DOI 10.4418/2013.68.1.2
- [7] Mansour, M., Determining the  $k$ -generalized gamma function  $\Gamma_k(x)$  by functional equations, *Int. J. Contemp. Math. Sciences*, 4(21) (2009), 653-660. <http://www.m-hikari.com/ijcms-password2009/21-24-2009/mansourIJCMS21-24-2009.pdf>
- [8] Miller, A. R., Summations for certain series containing the digamma function, *J. Phys. A: Math. Gen.*, 39 (2006), 3011-3020. DOI 10.1088/0305-4470/39/12/010
- [9] Moustafa, H., Almuashi, H., Mahmoud, M., On some complete monotonicity of functions related to generalized  $k$ -gamma function, *J. Math.*, 2021 (2021), 1-9. <https://doi.org/10.1155/2021/9941377>

- [10] Muqattash, I., Yahdi, M., Infinite family of approximations of the digamma function, *Math. Comput. Modelling*, 43(11-12) (2006), 1329-1336. <https://doi.org/10.1016/j.mcm.2005.02.010>
- [11] Nantomah, K., Iddrisu, M. M., The  $k$ -analogue of some inequalities for the gamma function, *Electron. J. Math. Anal. Appl.*, 2(2) (2014), 172-177.
- [12] Nantomah, K., Nisar, K. S., Gehlot, K. S., On a  $k$ -extension of the Nielsen's (beta)-function, *Int. J. Nonlinear Anal. Appl.*, 9(2) (2018), 191-201. <http://dx.doi.org/10.22075/ijnaa.2018.12972.1668>
- [13] Qi, F., Guo, S.-L., Guo, B.-N., Complete monotonicity of some functions involving polygamma functions, *J. Comput. Appl. Math.*, 233 (2010), 2149-2160. <https://doi.org/10.1016/j.cam.2009.09.044>
- [14] Qiu, S.-L., Vuorinen, M., Some properties of the gamma and psi functions with applications, *Math. Comp.*, 74(250) (2005), 723-742. DOI 10.1090/S0025-5718-04-01675-8
- [15] Widder, D. V., The Laplace Transform, Princeton University Press, Princeton, 1946.
- [16] Wilkins, B. D., Hromadka, T. V., Using the digamma function for basis functions in mesh-free computational methods, *Engineering Analysis with Boundary Elements*, 131 (2021), 218-227. <https://doi.org/10.1016/j.enganabound.2021.06.004>
- [17] Yildirim, E., Monotonicity properties on  $k$ -digamma function and its related inequalities, *J. Math. Inequal.*, 14(1) (2020), 161-173. <https://doi.org/10.7153/jmi-2020-14-12>
- [18] Yildirim, E., Ege, I., On  $k$ -analogues of digamma and polygamma functions, *J. Class. Anal.*, 13(2) (2018), 123-131. <https://doi.org/10.7153/jca-2018-13-08>
- [19] Yin, L., Huag, L. G., Song, Z. M., Dou, X. K., Some monotonicity properties and inequalities for the generalized digamma and polygamma functions, *J. Inequal. Appl.*, 1 (2018), 249. <https://doi.org/10.1186/s13660-018-1844-2>
- [20] Yin, L., Zhang, J., Lin, X., Complete monotonicity related to the  $k$ -polygamma functions with applications, *Ad. Diff. Eq.*, (2019), 1-10. <https://doi.org/10.1186/s13662-019-2299-6>