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# The comparison of the methods used for determining of Schottky diode parameters in a wide temperature range 

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#### Abstract

The current-voltage ( $I-V$ ) data of $\mathrm{Ni} / n$-GaAs Schottky diodes with 50 nm Schottky metal thickness has been measured in the temperature range of 60 K to 320 K . The important contact parameters of $\mathrm{Ni} / n-\mathrm{GaAs}$ Schottky diodes have been obtained by using conventional $I-V$ method, Norde method, generalized Norde method, and Cheung functions for each temperature. Then, the results have been compared each other.


Keywords: Schottky diode, conventional $I-V$ method, Cheung method, Norde method, generalized Norde method, temperature dependence

## Schottky diyot parametrelerini belirlemede kullanılan metotların geniş bir sıcaklık aralığı için kıyaslanması

## ÖZ

50 nm Schottky kontak kalınlığına sahip Ni/ $n$-GaAs Schottky diyotlarının akım-gerilim (I-V) verileri 60 K'den 320 K'e kadar olan geniş bir sıcaklık aralığında ölçüldü. Ni/ $n$-GaAs Schottky diyotlarının önemli kontak parametreleri geleneksel I-V metodu, Norde metodu, genelleştirilmiş Norde metodu ve Cheung fonksiyonları kullanılarak her bir sıcaklık değeri için ayrı ayrı elde edildi. Daha sonra sonuçlar birbirleriyle kıyasland.

Anahtar Kelimeler: Schottky diyot, geleneksel I-V metot, Cheung metot, Norde metot, genelleştirilmiş Norde metot, sıcaklık bağlılığ 1

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## 1. INTRODUCTION

Several methods based on thermionic emission theory have been derived for determination of Schottky diode parameters by using forward bias current-voltage ( $I-V$ ) data [1-8]. The conventional I-V method is the most popular method to obtain the ideality factor and Schottky barrier height (BH) [1,2]. This method loses its reliability when lnI-V plot has a narrow linear region because of the high series resistance [3]. Norde has proposed a method which overcomes this problem. The reliable values of Schottky BH and series resistance can be determined by means of this method. However, this method is not suitable to determine the contact parameters of non-ideal Schottky diode since it assumes that the contact is abrupt (i.e. ideal) between metal and semiconductor [3,9]. For non-ideal Schottky diodes, Norde's method has been generalized by Bohlin. The method enables to calculate Schottky BH , ideality factor, and series resistance from one I-V measurement [4]. Furthermore, Cheung functions have been presented as another way to find series resistance, Schottky BH and $n$ from the downward curvature region of the forward bias $\ln I-V$ plot [8].
In this study, the forward bias $I-V$ measurements of $\mathrm{Ni} / n$-GaAs Schottky diodes (SDs) which were fabricated by magnetron DC sputtering system were taken in a wide temperature range of 60-320 K. The common analytical tools have been used for data analysis and the contact parameters obtained by using different methods have been compared to each other. In literature there is no comparison of well-known methods used to find the contact parameters in a wide temperature range. This study presents information about the temperature dependence of the contact parameters obtained by different methods and try to explain the inconsistency between the methods.

## 2. EXPERIMENTAL

N-type GaAs (100) wafer with $1.46 \times 10^{16} \mathrm{~cm}^{-3}$ donor concentration atoms has been used to fabricate the Schottky diodes. Before the Schottky and ohmic metallization, the wafer was exposed to wet chemical cleaning which has the following steps [10,11]:

1) D ipping in trichloroethylene, methanol, and acetone for 3.0 minutes (to remove the organic contamination)
2) Etching in a $\mathrm{H}_{2} \mathrm{SO}_{4}: \mathrm{H}_{2} \mathrm{O}_{2}: \mathrm{H}_{2} \mathrm{O}$ (5:1:1) solution for 1.0 minute (to remove the surface damage layer and undesirable impurities),
3) Etching in a $\mathrm{H}_{2} \mathrm{O}: \mathrm{HCl}$ (1:1) solution (to remove the metallic contamination and thin oxide layer).

Generally, removing of organic contamination realizes before cleaning of the inorganic contamination because the presence of organic layer can prevent the acidic solutions to reach to the wafer surface [12]. After each process, the wafer was rinsed in deionized water with $18 \mathrm{M} \Omega$ resistivity. The drying process of the wafer was realized with the high purity nitrogen gas and then the wafer was taken to the deposition chamber, instantly. Ohmic contact was realized by evaporation of indium (In) at $10^{-5} \mathrm{Torr}$ and then the wafer was annealed at $380{ }^{\circ} \mathrm{C}$ for 3 minutes in flowing $\mathrm{N}_{2}$ for low resistance ohmic contact. The Schottky metallization of nickel ( Ni ) was realized by magnetron DC sputtering. Ni contacts with 50 nm thicknesses have a circular shape with a diameter of 1.5 mm . A Keithley 2400 SourceMeter and a Leybold Heraeus closed-cycle helium cryostat were employed to take the temperature dependent $I-V$ measurements of the diodes. A Windaus MD850 electronic thermometer and a copper constantan thermocouple have been used to control sample temperature.

## 3. RESULTS AND DISCUSSION

### 3.1. Conventional $I-V$ method

The relationship between current and voltage for Schottky diodes if the dominant mechanism is thermionic emission can be given as [1]:
$I=A A^{*} T^{2} \exp \left(-\frac{q \Phi_{b}^{\text {eff }}}{k T}\right)\left[\exp \left(\frac{q\left(V-I R_{s}\right.}{k T}\right)-1\right]$
Here, $A$ is the contact area, $A^{*}$ is the effective Richardson constant ( $8.16 \mathrm{~A} / \mathrm{cm}^{2} \mathrm{~K}^{2}$ for $n$-GaAs [1]), $k$ is the Boltzmann's constant, $T$ is the absolute temperature in $\mathrm{K}, q$ is the elementary charge, $\Phi_{b}{ }^{\text {eff }}$ is the effective BH , and $R_{s}$ is the series resistance. A bias dependent Schottky BH and other effects which cause the deviation from thermionic emission are included by addition of an ideality factor ( $n$ ) to Eq. (1). The relationship
between $\Phi_{b}^{\text {eff }}$ and $n$ is given by the following expression $[1,10]$,
$\Phi_{b}^{e f f}=\Phi_{b o}+\left(1-\frac{1}{n}\right) V$
Here, $\Phi_{b 0}$ is BH under zero bias. For $V-I R_{s}>3 k T / q$, Eq. (1) can be written as the following equation by taking into Eq. 2.
$I=I_{0} \exp \left(\frac{q V}{n k T}\right)$
$I_{0}=A A^{*} T^{2} \exp \left(-\frac{q \Phi_{b 0}}{k T}\right)$
At low and intermediate voltage region, the plot of $\ln I-V$ is linear and the extrapolating of this plot to $\ln I$ axis gives the reverse saturation current $\left(I_{0}\right)$. The ideality factor can be determined from the slope of $\ln I-V$ plot in according to the following equation.
$n=\frac{q}{k T} \frac{\mathrm{~d} V}{\mathrm{~d} \ln I}$
The $\ln I-V$ characteristics of $\mathrm{Ni} / n$-GaAs Schottky diodes as a function of temperatures are given in Figure 1. It can be said that the dominant mechanism for current transport across $\mathrm{Ni} / n-\mathrm{GaAs}$ contacts is thermionic emission because the $\ln I-V$ characteristics are linear over a wide range of current values [11]. For each temperature, the existence of the series resistance which limits the current flow in the diodes causes a downward curvature region in the forward $\ln I-V$ curves [12]. This effect can be shown in Figure1. The ideality factor and the zero-bias BH values determined from $\ln I-V$ curves are listed in Table 1.

Table 1. The ideality and the zero-bias barrier height values obtained by using conventional $I-V$ method

| Temperature (K) | $\boldsymbol{n}$ | $\boldsymbol{\Phi}_{\text {bo }}(\mathbf{I}-\boldsymbol{V})(\boldsymbol{e V})$ |
| :---: | :---: | :---: |
| 320 | 1.015 | 0.720 |
| 300 | 1.026 | 0.723 |
| 280 | 1.033 | 0.727 |
| 260 | 1.056 | 0.735 |
| 240 | 1.048 | 0.747 |
| 220 | 1.101 | 0.739 |
| 200 | 1.121 | 0.739 |
| 180 | 1.113 | 0.748 |
| 160 | 1.147 | 0.735 |
| 140 | 1.192 | 0.709 |
| 120 | 1.321 | 0.656 |
| 100 | 1.415 | 0.615 |
| 80 | 1.461 | 0.588 |
| 60 | 1.934 | 0.457 |

In our previous work [13], the temperature dependence of $\ln I-V$ curves and the parameters obtained from $I-V$ data was studied in detail. As can be seen from the table, the ideality factor and zero-bias BH values change between 60 K and 180 K and the values show very little fluctuations after 180 K. Real metal-semiconductor interfaces where include low and high barrier height regions have an inhomogeneties barrier height distribution [16]. At low temperatures, since the electrons can only across the regions where include low barriers, the zero-bias BH value decreases and the ideality factor value increases as the temperature decreases. As the temperature increases, the electrons are able to surmount regions including high barrier heights. Therefore, effect of the barrier height inhomogeneties on $I-V$ characteristic loses its importance and ideality factor and zerobias height values show very little fluctuations at high temperatures $[15,16]$.


Figure 1. Forward bias $I-V$ characteristics of $\mathrm{Ni} / n-\mathrm{GaAs} / \mathrm{In}$ Schottky diode for different temperatures

### 3.2. Cheung functions

The series resistance of a diode is the main reason of downward curvature region in an $\ln I-V$ plot [9]. The most common method to calculate $R_{s}$ is the Cheung method [8]. In addition to $R_{s}, n$ and barrier height ( $\Phi_{b}$ ) under a bias can be determined from Cheung's functions. The following equations are known as Cheung's functions:
$\frac{\mathrm{d} V}{\mathrm{~d} \ln I}=I R_{s}+\frac{n k T}{q}$
$H(I)=V-\frac{n k T}{q} \ln \left(\frac{I}{A A^{*} T^{2}}\right)$

$$
\begin{equation*}
H(I)=I R_{s}+n \Phi_{b} \tag{7}
\end{equation*}
$$

Taking into above the equations, $\mathrm{d} V / \mathrm{d} \ln I$ and $H(I)$ versus $I$ plots are drawn by using data determined from the downward curvature region of $\ln I-V$ plot
should be linear. $R_{s}$ is determined from slope of these plots. $n$ and $\Phi_{b}$ values is obtained from the $y$ axis intercept of $\mathrm{d} V / \mathrm{d} \ln I$ and $H(I)$ versus $I$ plots, respectively. Table 2 summarizes the evaluation of the Ni/n-GaAs Schottky diode parameters based on Cheung functions is given in Figure 2.

Table 2. Schottky diode parameters obtained by using Cheung functions

| Temperature (K) | $d V / d \ln I-I$ |  | $H(I)-I$ |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $n$ | $R_{s}(\Omega)$ | $\boldsymbol{R}_{s}(\Omega)$ | $\Phi_{b}(\mathrm{eV})$ |
| 320 | 2.246 | 5.724 | 5.724 | 0.609 |
| 300 | 2.330 | 5.788 | 5.783 | 0.584 |
| 280 | 2.435 | 5.571 | 5.560 | 0.552 |
| 260 | 2.513 | 5.726 | 5.725 | 0.545 |
| 240 | 2.580 | 5.491 | 5.493 | 0.519 |
| 220 | 2.772 | 5.498 | 5.485 | 0.487 |
| 200 | 3.014 | 5.271 | 5.251 | 0.440 |
| 180 | 3.066 | 5.546 | 5.527 | 0.435 |
| 160 | 3.147 | 5.503 | 5.512 | 0.424 |
| 140 | 3.223 | 5.525 | 5.514 | 0.396 |
| 120 | 3.411 | 5.668 | 5.675 | 0.359 |
| 100 | 4.278 | 5.734 | 5.722 | 0.290 |
| 80 | 4.928 | 6.116 | 6.110 | 0.247 |
| 60 | 5.990 | 6.816 | 6.845 | 0.197 |



Figure 2. First and second Cheung functions of $\mathrm{Ni} / n-G a A s$ Schottky diodes

### 3.3. Norde and generalized Norde plot

The generalized Norde function has been proposed by Bohlin to calculate $\Phi_{b}$ and $R_{s}$ values from only one $I-V$ measurement at a fixed temperature. The generalized Norde function is described as [4]:
$F(V, \alpha)=\frac{V}{\alpha}-\frac{1}{\beta} \ln \left(\frac{I(V)}{A A^{*} T^{2}}\right)$
$\alpha$ is defined as an arbitrary constant greater than $n$, $\beta$ is equal to $q / k T, I(V)$ is extracted from the measured $I-V$ curve. $\Phi_{b}$ and $R_{s}$ values can be obtained by determining of $F(V, \alpha)$ against $V$ plot minimum in according to following equations [4]:

$$
\begin{equation*}
\Phi_{b}=F_{\text {min }}(V, \alpha)+\left(\frac{\alpha-n}{n}\right)\left(\frac{V_{\text {min }}}{\alpha}-\frac{1}{\beta}\right) \tag{10}
\end{equation*}
$$

$R_{s}=\frac{(\alpha-n)}{\beta I_{\text {min }}}$
where $n$ value is obtained from the $\ln I-V$ plot. $F_{\text {min }}$ $(V, \alpha)$ is the minimum point $F(V, \alpha)-V$ plot, $V_{\text {min }}$ and $I_{\text {min }}$ are the corresponding voltage and current, respectively. For $n=1$ and $\alpha=2$, Eqs. (9), (10), and (11) are same for normal Norde method are given in Ref. [3]. Figure 3 shows $F(V, 2)$ - $V$ plots for different temperatures. Table 3 contains $\Phi_{b}$ and $R_{s}$ values obtained from Norde and generalized Norde plots.


Figure 3. Norde plots of $\mathrm{Ni} / n-\mathrm{GaAs} /$ In Schottky diodes at different temperatures

Table 3. Schottky diode parameters determined from Norde and generalized Norde method

| Temperature <br> $(\mathbf{K})$ | $\boldsymbol{n}$ | $\boldsymbol{\alpha}$ | (Gen. Norde) <br> $\boldsymbol{R}_{\boldsymbol{s}}(\boldsymbol{\Omega})$ | (Norde) <br> $\boldsymbol{R}_{\boldsymbol{s}}(\boldsymbol{\Omega})$ | (Norde) <br> $\boldsymbol{\Phi}_{\boldsymbol{b}}(\mathbf{e V})$ | (Gen. Norde) <br> $\boldsymbol{\Phi}_{\boldsymbol{b}}(\mathbf{e V})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 320 | 1.015 | 2 | 4.248 | 4.313 | 0.757 | 0.753 |
| 300 | 1.026 | 2 | 5.794 | 5.948 | 0.756 | 0.748 |
| 280 | 1.033 | 2 | 5.004 | 5.175 | 0.770 | 0.758 |
| 260 | 1.056 | 2 | 4.885 | 5.175 | 0.780 | 0.758 |
| 240 | 1.048 | 2 | 4.878 | 5.123 | 0.789 | 0.768 |
| 220 | 1.101 | 2 | 6.553 | 7.289 | 0.791 | 0.747 |
| 200 | 1.121 | 2 | 7.897 | 8.984 | 0.800 | 0.745 |
| 180 | 1.113 | 2 | 6.653 | 7.500 | 0.814 | 0.757 |
| 160 | 1.147 | 2 | 7.154 | 8.388 | 0.819 | 0.743 |
| 140 | 1.192 | 2 | 7.494 | 9.274 | 0.818 | 0.717 |
| 120 | 1.321 | 2 | 12.871 | 18.956 | 0.813 | 0.656 |
| 100 | 1.415 | 2 | 14.280 | 24.410 | 0.811 | 0.615 |
| 80 | 1.461 | 2 | 18.472 | 34.272 | 0.809 | 0.590 |
| 60 | 1.934 | 2 | 18.265 | 276.738 | 0.796 | 0.457 |

### 3.4. The comparison of the methods

Figure 4 shows the temperature dependence of the barrier heights obtained by different methods. As can be seen from the figure, the barrier heights obtained by using the conventional $I$ - $V$ method are in good agreement with those calculated by using the generalized Norde method. There is an inconsistency between conventional $I-V$ method and Norde method because Norde method assumes that the contact between metal semiconductor is ideal (i.e, $n=1$ ). Schottky BH values calculated by Cheung functions exhibit same temperature dependence of Schottky BH values determined by using conventional $I-V$ method.


Figure 4.The temperature dependence of barrier height values obtained by different methods

However, BH values calculated from Cheung functions are smaller than those obtained by using conventional $I-V$ and generalized Norde methods. The traditional $I-V$ and generalized Norde methods use the data on the linear region of the curve. In this part of the $I-V$ characteristics the essential effects are interfacial layer and interface states. In the Cheung method the BH values is obtained by using data taken on the nonlinear region of the $I-V$ characteristics where also $R_{s}$ is effective as well as the interfacial layer and interface states. [9,17]. Since the each method use the different regions of $\ln I-V$ characteristic, the inconsistency between BH values occurs. Same things can be said for ideality factor values obtained by both methods (Figure 5).


Figure 5 . The temperature dependence of ideality factor values obtained by different methods

The temperature dependence of $R_{s}$ values obtained by different methods is given in Figure 6. $R_{s}$ values calculated by Cheung functions are different from those determined from Norde and generalized Norde plots. Cheung functions are only applied to downward curvature region of $\ln I-V$ characteristic. But, Norde and generalized Norde functions are applied to the full forward bias region of the $\ln I-V$ characteristic. Therefore, $R_{s}$ values obtained by Norde plots are larger than those determined from Cheung functions [14].
Furthermore, the difficulty in the determination of minimum points of $F(V, 2)-V$ plots is still a problem can cause this inconsistency [8,9]. The Norde method and generalized Norde method determines series resistance value by using $I-V$ data in the linear region of $\ln I-V$ characteristic in where the current changes as exponential. Therefore, a little mistake in determining the graphical turning points can affect series resistance value, dramatically. As referring to Figure 4 and

Figure 6, all parameters obtained by different methods approach each other at high temperature region. This result is expected since the ideality factor value approaches to the unity as the temperature increases.


Figure 6. The temperature dependence of series resistance values obtained by different methods (The series resistance value at 60 K determined from Norde plot was excluded for a better visuality)

### 3.5. Conclusion

The Schottky BH values obtained from conventional $I-V$ and generalized Norde methods show good agreement with each other. If we assume the conventional $I-V$ method as the reference method, Norde methods lose their reliability at low temperatures. The Schottky BH values calculated by Cheung functions have same temperature dependence of Schottky BH values obtained by the reference method. $R_{s}$ values determined from Cheung and generalized Norde (also, Norde) functions exhibit inconsistency because that the different nature of the methods. For the reliable results, the $I-V$ data must be taken by little voltage steps when the series resistance is calculated by means of Norde and generalized Norde methods.

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