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# Exponential growth of solutions for a parabolic system

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#### ARTICLE INFO

ABSTRACT

| Article history:<br>Received 17 Oct 2019<br>Received in revised form 14 Nov<br>2019 | In this paper, we investigated the initial boundary problem of a class of doubly nonlinear parabolic systems. We prove exponential growth of solution with negative initial energy. |
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#### 1. Introduction

In this paper, we study the following parabolic system

$$\begin{cases} u_t - \Delta u + |u|^{q-2} u_t = f_1(u, v), & x \in \Omega, \ t > 0, \\ v_t - \Delta v + |v|^{q-2} v_t = f_2(u, v), & x \in \Omega, \ t > 0, \\ u(x, t) = v(x, t) = 0, & x \in \partial\Omega, \ t \ge 0, \\ u(x, 0) = u_0(x), & v(x, 0) = v_0(x) & x \in \Omega, \end{cases}$$
(1)

where q > 2 are real numbers and  $\Omega$  is a bounded domain in  $\mathbb{R}^n$   $(n \ge 1)$  with smooth boundary  $\partial \Omega$ .  $f_i(u, v)$  (i = 1, 2) will be given later.

Pang and Qiao [1] studied the blow up properties of the problem (1) with negative and positive initial energy.

In the absence of  $|u|^{q-2}u_t$  and  $|v|^{q-2}v_t$  term become the following problem

$$\begin{cases} u_t - \Delta u = f_1(u, v), \\ v_t - \Delta v = f_2(u, v). \end{cases}$$
(2)

This type of equation aries from a variety of mathematical models in engineering and physical sciences, it appears naturally in the models of physics, chemistry, biology, ecology and so on (see [2-12]). In [13], the authors got the global existence solution, blow-up in finite time solution, and asymptotic behavior of solution for (2).

Currently, in [14] the author also discussed the problem (2). He got global existence of the solutions, the asymptotic stability of solution and the blow up of solution.

In detail, this paper is organized as follows: In the next section, we present some notations and statement of assumptions. In section 3, the growth of solution is given.

# 2. Preliminaries

In this section, we shall give some assumptions for the proof of our results. Let  $\|.\|$ ,  $\|.\|_p$  and  $(u, v) = \int_{\Omega} u(x)v(x) dx$  denote the usual  $L^2(\Omega)$  norm,  $L^p(\Omega)$  norm and inner product of  $L^2(\Omega)$ , respectively. Throughout this paper, C is used to point out general positive constants.

For the numbers m and q, we suppose that

$$\begin{cases} 2 < q < m \le \frac{n+2}{n-2} & \text{if } n > 2, \\ 2 < q < m \le +\infty & \text{if } n = 1, 2. \end{cases}$$
(3)

Regarding the functions  $f_1(u, v)$ ,  $f_2(u, v) \in C^1$  such that

$$f_1(u,v) = \frac{\partial F(u,v)}{\partial u}, f_2(u,v) = \frac{\partial F(u,v)}{\partial v}$$

and

$$\begin{cases} k_0(|u|^m + |v|^m) \le F(u, v) \le k_1(|u|^m + |v|^m), \\ uf_1(u, v) + vf_2(u, v) = (m+1)F(u, v) \end{cases}$$
(4)

where  $k_0, k_1$  are positive constants.

Combining arguments of [15,12,16], u(x,t), v(x,t) are called a solution of problem (1) on  $\Omega \times [0,T)$  if

$$\begin{cases} u, v \in C(0, T; H_0^1(\Omega)) \cap C^1(0, T; L^2(\Omega)), \\ |u|^{q-2}u_t, \ |v|^{q-2}v_t \in L^2(\Omega \times [0, T)) \end{cases}$$
(5)

satisfying the initial condition  $u(x, 0) = u_0(x), v(x, 0) = v_0(x)$  and

$$\int_{0}^{t} \int_{\Omega} \left[ \nabla u(s) \nabla w(s) + u_{t}(s) w(s) + |u|^{q-2} u_{t} w - f_{1}(u, v) w \right] dx \, ds = 0, \, (6)$$

$$\int_{0}^{t} \int_{\Omega} \left[ \nabla v(s) \nabla w(s) + v_{t}(s) w(s) + |v|^{q-2} v_{t} w - f_{2}(u, v) w \right] dx \, ds = 0 \quad (7)$$

for all  $w \in C(0,T; H_0^1(\Omega))$ .

The energy functional associated with problem (1) is

$$E(t) = \frac{1}{2} \|\nabla u\|^2 + \frac{1}{p} \|\nabla v\|^2 - \int_{\Omega} F(u, v) dx,$$
(8)

where  $u, v \in H_0^1(\Omega)$ .

Lemma 2.1 Suppose that (3) and (4) hold. E'(t) is noncreasing function t > 0 and  $E'(t) = -||u_t||^2 - ||v_t||^2 - \int_{\Omega} |u|^{q-2} u_t^2 dx - \int_{\Omega} |v|^{q-2} v_t^2 dx < 0.$  (9)

**Proof.** Multiplying  $Eq. (1)_1$  by  $u_t$  and  $Eq. (1)_2$  by  $v_t$  and integrating over  $\Omega$ , we obtain  $\int_0^t E'(\tau) d\tau = -\left[\int_0^t (||u_t||^2 + ||v_t||^2) d\tau + \int_0^t \int_{\Omega} |u|^{q-2} u_t^2 dx d\tau + \int_0^t \int_{\Omega} |v|^{q-2} v_t^2 dx d\tau\right],$   $E(t) - E(0) = -\left[\int_0^t (||u_t||^2 + ||v_t||^2) d\tau + \int_0^t \int_{\Omega} |u|^{q-2} u_t^2 dx d\tau + \int_0^t \int_{\Omega} |v|^{q-2} v_t^2 dx d\tau\right]$ for t > 0.

### 3. Exponential Growth of Solution

In this section, we state and prove exponential growth result.

**Theorem 3.1** Suppose that (3) holds,  $u_0, v_0 \in H_0^1(\Omega)$  and E(0) < 0. Then, the solution of the system (1) grows exponentially.

Proof. We set H(t) = -E(t).(10)From (10) and (9), we have  $H'(t) = -E'(t) \ge 0.$ (11)Since E(0) < 0, we get H(0) = -E(0) > 0.(12)By the integrate (11), we get  $0 < H(0) \le H(t).$ (13)By using (10) and (8) $H(t) - \int_{0} F(u, v) dx = -\frac{1}{2} \left( \|\nabla u\|^{2} + \|\nabla v\|^{2} \right) < 0.$ (14)Then, by using (4), we have  $0 < H(0) \le H(t) \le \int_{\Omega} F(u, v) dx \le k_1(||u||_m^m + ||v||_m^m).$ (15)We define the functional  $\Phi(t) = H(t) + \frac{\varepsilon}{2} ||u||^2 + \frac{\varepsilon}{2} ||v||^2.$ (16)By differentiating (16) and using Eq.(1), we get  $\Phi'(t) = H'(t) + \varepsilon \left( \int_{\Omega} u u_t dx + \int_{\Omega} v v_t dx \right)$  $= \|u_t\|^2 + \|v_t\|^2 - \varepsilon \|\nabla u\|^2 - \varepsilon \|\nabla v\|^2 + \varepsilon \int_{\Omega} [uf_1(u,v) + vf_2(u,v)] dx$  $+\int_{\Omega}|u|^{q-2}u_t^2dx+\int_{\Omega}|v|^{q-2}v_t^2dx-\varepsilon\int_{\Omega}|u|^{q-2}uu_tdx-\varepsilon\int_{\Omega}|v|^{q-2}vv_tdx$  $= \|u_t\|^2 + \|v_t\|^2 - \varepsilon \|\nabla u\|^2 - \varepsilon \|\nabla v\|^2 + \varepsilon (m+1) \int_{\Omega} F(u,v) dx$  $+ \int_{\Omega} |u|^{q-2} u_t^2 dx + \int_{\Omega} |v|^{q-2} v_t^2 dx - \varepsilon \int_{\Omega} |u|^{q-2} u u_t dx$  $-\varepsilon \int_{\Omega} |v|^{q-2} v v_t dx.$ (17)

In order to estimate the last two terms in the right-hand side of (17), we use the following Young's inequality,

$$ab \leq \delta^{-1}a^2 + \delta b^2,$$

so we have

$$\begin{split} \int_{\Omega} & |u|^{q-2} u u_t dx \leq \int_{\Omega} & |u|^{\frac{q-2}{2}} u_t |u|^{\frac{q-2}{2}} dx \\ & \leq \delta^{-1} \int_{\Omega} & |u|^{q-2} u_t^2 dx + \delta \int_{\Omega} & |u|^q dx. \end{split}$$

Similarly,

$$\int_{\Omega} |v|^{q-2} v v_t dx \leq \delta^{-1} \int_{\Omega} |v|^{q-2} v_t^2 dx + \delta \int_{\Omega} |v|^q dx$$

Then, (17) becomes

$$\begin{split} \Phi'(t) &\geq \|u_t\|^2 + \|v_t\|^2 - \varepsilon \|\nabla u\|^2 - \varepsilon \|\nabla v\|^2 + \varepsilon (m+1)(\|u\|_m^m + \|v\|_m^m) \\ &- \varepsilon \delta (\|u\|_q^q + \|v\|_q^q) + (1 - \varepsilon \delta^{-1}) \int_{\Omega} |u|^{q-2} u_t^{-2} dx \\ &+ (1 - \varepsilon \delta^{-1}) \int_{\Omega} |v|^{q-2} v_t^{-2} dx. \end{split}$$
(18)

By using follows equality that

$$-\|\nabla u\|^2 - \|\nabla v\|^2 = 2H(t) - 2\int_{\Omega} F(u, v)dx.$$

Hence, (18) becomes

$$\Phi'(t) \ge 2\varepsilon H(t) + \|u_t\|^2 + \|v_t\|^2 + \varepsilon (m-1)(\|u\|_m^m + \|v\|_m^m) -\varepsilon \delta(\|u\|_q^q + \|v\|_q^q) + (1 - \varepsilon \delta^{-1}) \int_{\Omega} |u|^{q-2} u_t^2 dx + (1 - \varepsilon \delta^{-1}) \int_{\Omega} |v|^{q-2} v_t^2 dx.$$
(19)

As the embedding  $L^m \hookrightarrow L^q \hookrightarrow L^2$ , m > q > 2, we have

$$\begin{cases} \|u\|_{q}^{q} \leq C \|u\|_{m}^{q} \leq C (\|u\|_{m}^{m})^{\frac{q}{m}}, \\ \|v\|_{q}^{q} \leq C \|v\|_{m}^{q} \leq C (\|v\|_{m}^{m})^{\frac{q}{m}}. \end{cases}$$
(20)

Since  $0 < \frac{q}{m} < 1$ , now applying the following inequality

 $x^{l} \leq (x+1) \leq \left(1 + \frac{1}{z}\right)(x+z),$ 

which holds for all  $x \ge 0$ ,  $0 \le l \le 1$ , z > 0, especially, taking  $x = ||u||_{m}^{m}$ ,  $l = \frac{q}{m}$ , z = H(0), we get

$$C(\|u\|_m^m)^{\frac{q}{m}} \le \left(1 + \frac{1}{H(0)}\right) \left(\|u\|_m^m + H(0)\right),$$

similarly

$$C(\|v\|_m^m)^{\frac{q}{m}} \le \left(1 + \frac{1}{H(0)}\right) \left(\|v\|_m^m + H(0)\right).$$

Then, from (15) and (20), we get

$$\begin{aligned} \|u\|_{q}^{q} + \|v\|_{q}^{q} &\leq C\left(\|u\|_{m}^{q} + \|v\|_{m}^{q}\right) \\ &\leq C_{1}(\|u\|_{m}^{m} + \|u\|_{m}^{m}). \end{aligned}$$
(21)

Then, from (21) we obtain

$$\begin{split} \Phi'(t) &\geq 2\varepsilon H(t) + \|u_t\|^2 + \|v_t\|^2 + \varepsilon a_1(\|u\|_m^m + \|v\|_m^m) \\ &+ (1 - \varepsilon \delta^{-1}) \int_{\Omega} |u|^{q-2} u_t^2 dx + (1 - \varepsilon \delta^{-1}) \int_{\Omega} |v|^{q-2} v_t^2 dx \end{split}$$

where  $\delta$  small enough such that  $a_1 = m - 1 - \delta C_1 > 0$  and taking  $\varepsilon$  and  $\delta$  small enough such that  $1 - \varepsilon \delta^{-1} > 0$ , then

 $\Phi'(t) \ge C(H(t) + ||u_t||^2 + ||v_t||^2 + ||u||_m^m + ||v||_m^m).$ (22)On the other hand, by definition of  $\Phi(t)$  and Poincare's inequality, we get  $\Phi(t) = H(t) + \frac{\varepsilon}{2} \|u\|^2 + \frac{\varepsilon}{2} \|v\|^2$  $\leq C(H(t) + \|\nabla u\|^2 + \|\nabla v\|^2).$ From definition of H(t), we get  $\Phi(t) \le C(H(t) + \|u\|_m^m + \|v\|_m^m)$ (23) $\leq \mathcal{C}(H(t) + \|u\|_m^m + \|v\|_m^m + \|u_t\|^2 + \|v_t\|^2).$ (24)From (22) and (24), we arrive at  $\Phi'(t) \ge r\Phi(t),$ (25)where r is a positive constant. Integration of (25) over (0, t) gives us  $\phi(t) > \phi(0) \exp(rt)$ 

$$\Psi(t) \geq \Psi(0) \exp(rt)$$

From (23) and (15), we get

 $\Phi(t) \le H(t) \le \|u\|_m^m + \|v\|_m^m.$ 

Consequently, we show that the solution in the  $L_m$ -norm growths exponentially.

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